

INTRODUCTION TO EXPERIMENTAL QUANTUM COMPUTATION

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Summer School on Quantum Computing
Universidad Internacional Menéndez y Pelayo
September 3d, 2020



Institut de Física
d'Altes Energies



INTRODUCTION TO EXPERIMENTAL SUPERCONDUCTING QUANTUM COMPUTATION

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Barcelona Institute of
Science and Technology

OUTLINE

Lecture I

- **Quantum computation**
- Circuit quantization
- Superconducting qubit zoo
- Qubit state control

Lecture II

- Resonators for quantum computation
- Circuit quantum electrodynamics
- Qubit-qubit couplings and 2-qubit gates
- State of the art

QUANTUM COMPUTATION

Why do we need Quantum Computation?



MareNostrum 4, Barcelona Supercomputing Center:

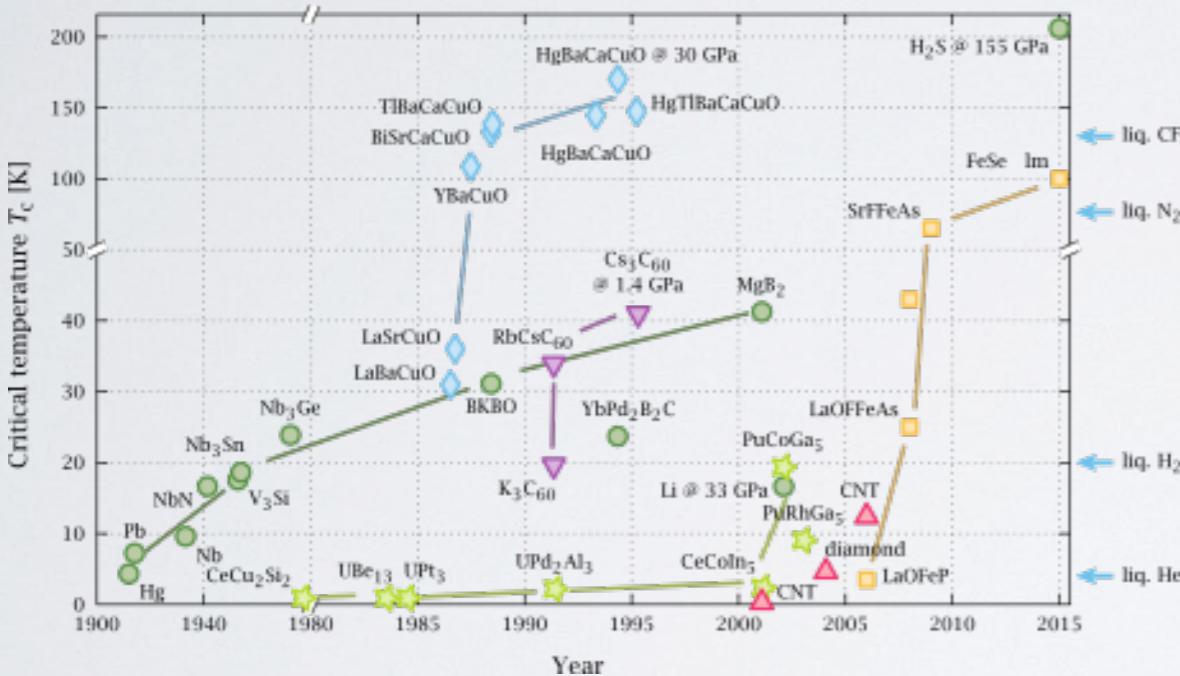
390 TB, 11.5 Pflops \sim 40 qubits

But also energy consumption!
Power consumed 1.3MW
TFlop generation many MWs...

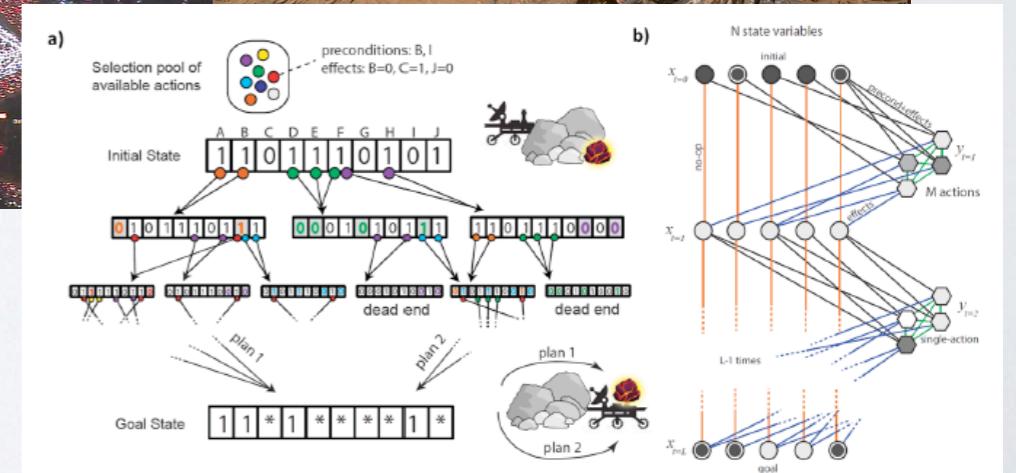
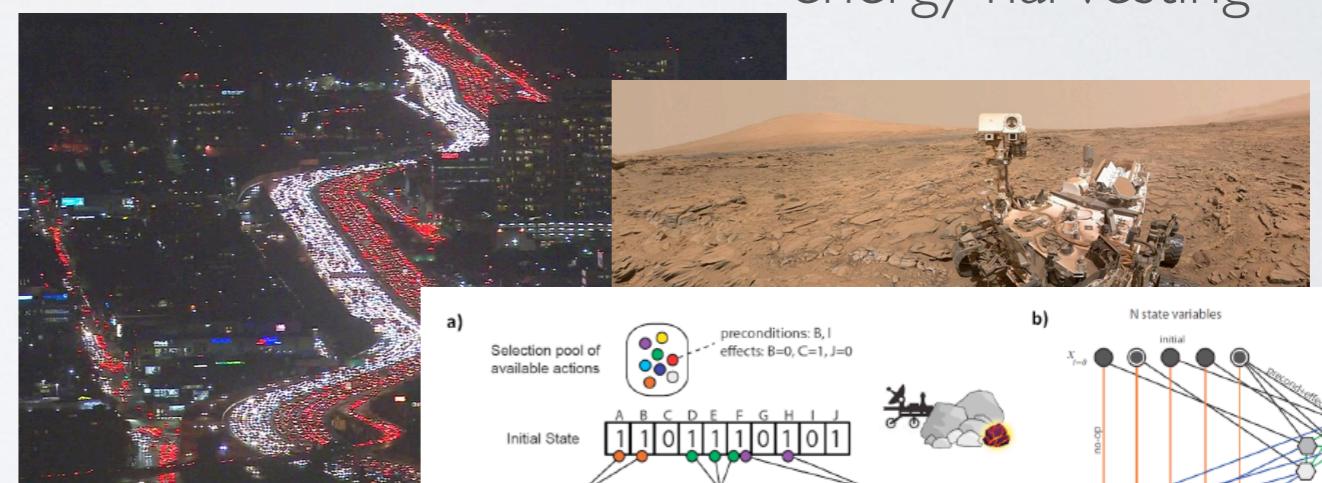
QUANTUM COMPUTATION

Why do we need Quantum Computation?

- Applications where small-sized quantum processors can outperform classical computers

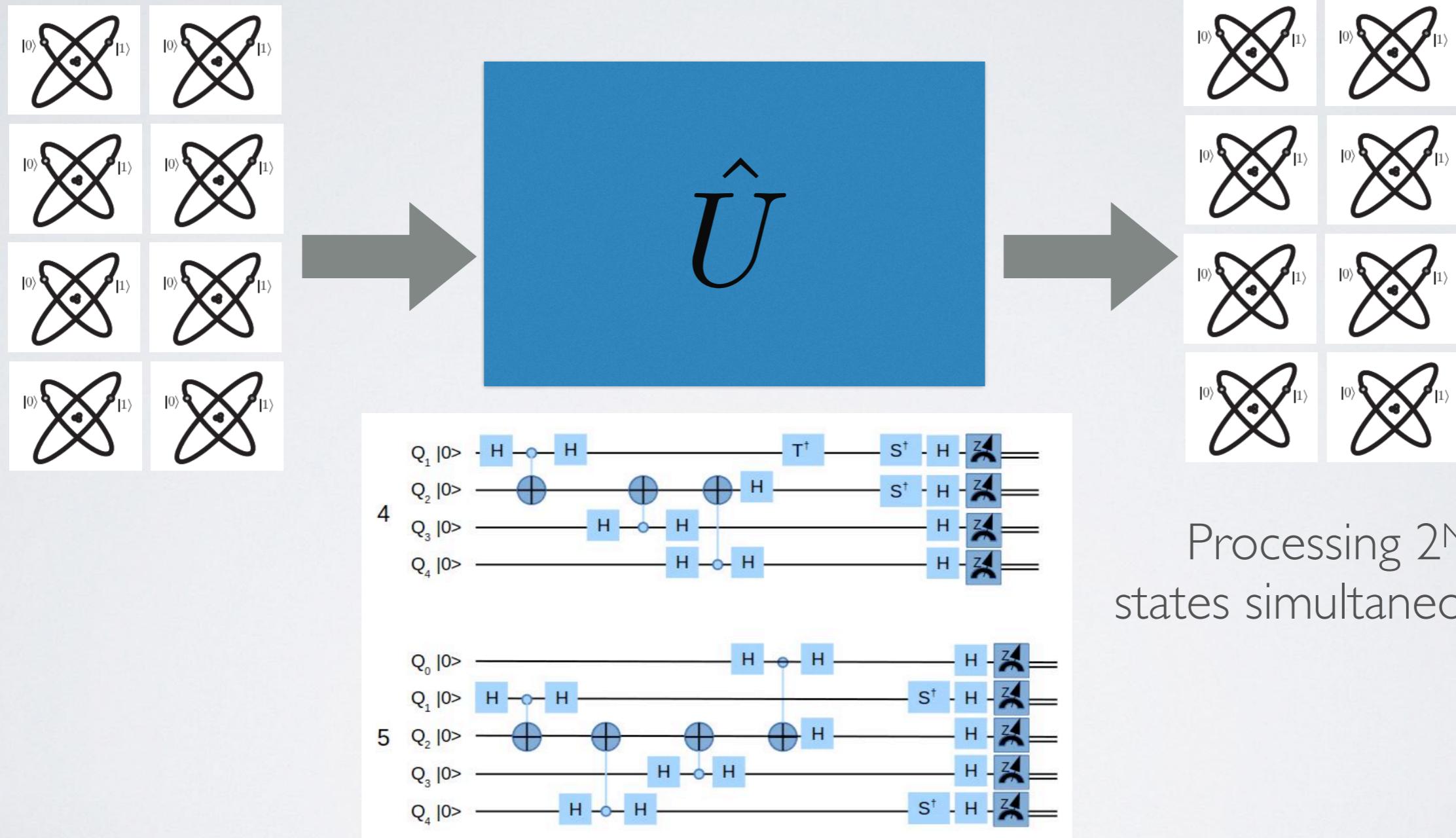


Simulate higher T_c superconductor



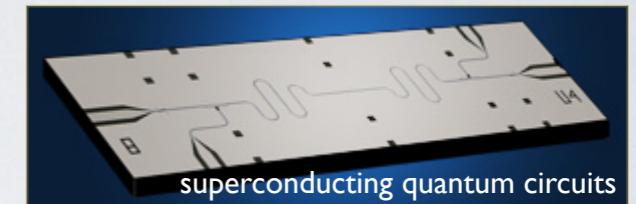
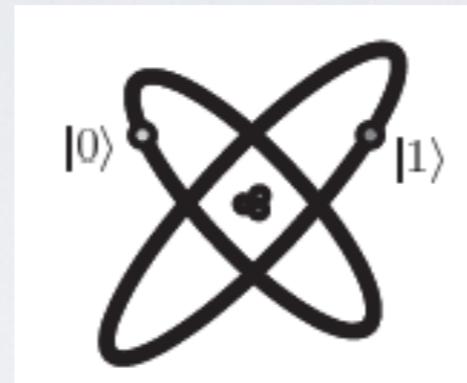
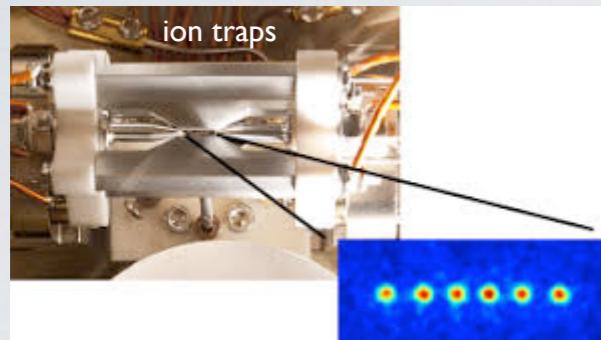
Optimization: traffic, navigation, scheduling, machine learning, etc.

QUANTUM COMPUTATION

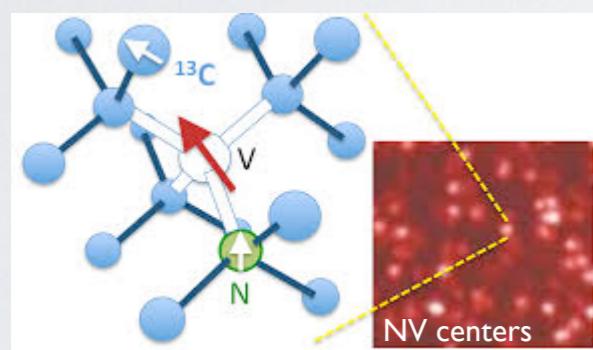


Nielsen and Chuang, Quantum computation and Quantum information

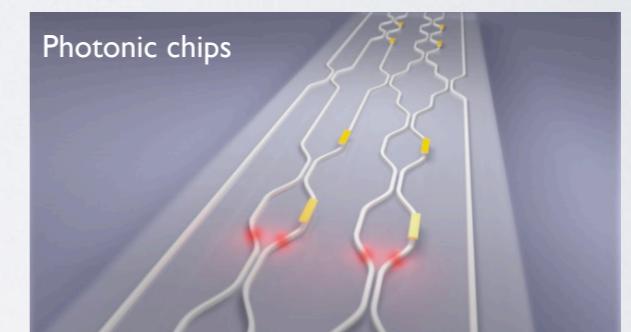
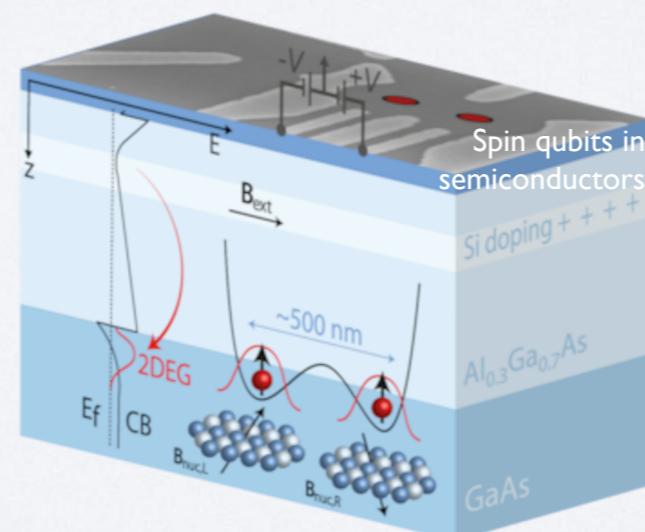
QUANTUM COMPUTATION



Quantum systems store information in their quantum states.
Quantum bit = qubit: minimal quantum information unit with 2 states



$$|\psi\rangle = \frac{|0\rangle + |1\rangle}{\sqrt{2}}$$



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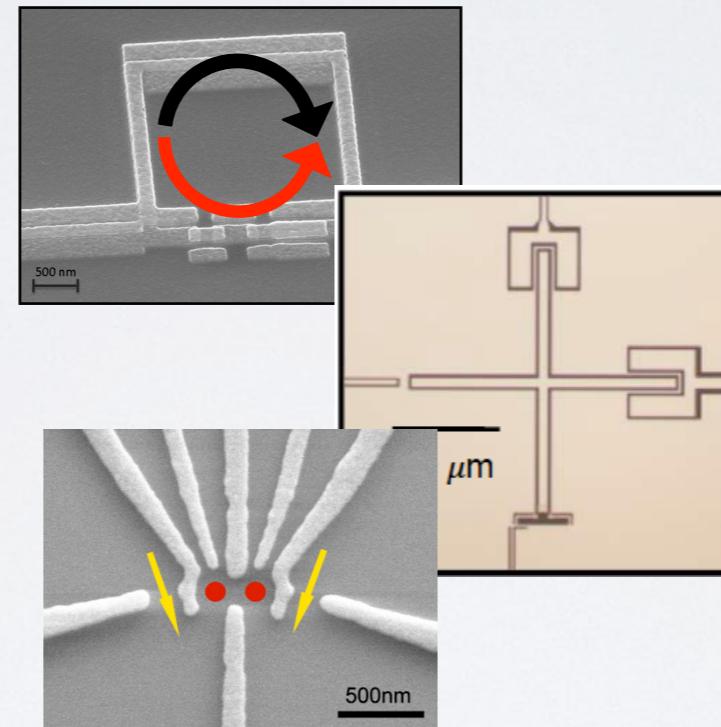
CIRCUIT QUANTIZATION

Microscopic



$$i\hbar \frac{\partial \psi(\vec{r}, t)}{\partial t} = \mathcal{H}(\vec{r}, t)\psi(\vec{r}, t)$$

Mesoscopic



Large number of particles
Artificial, man-made

Quantum collective degrees of freedom

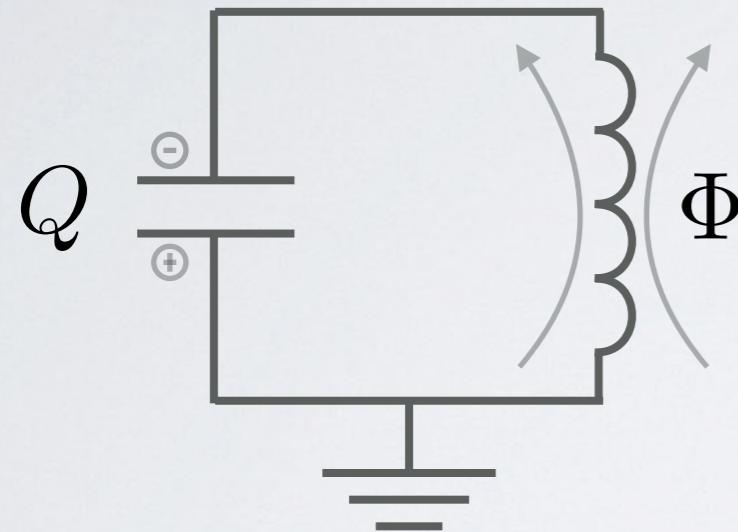
Macroscopic



$$\vec{F} = m\vec{a}$$

CIRCUIT QUANTIZATION

What corresponds to these collective degrees of freedom?



$$\mathcal{H} = \frac{Q^2}{2C} + \frac{\Phi^2}{2L}$$

$$\omega_{\text{LC}} = \frac{1}{\sqrt{LC}}$$

$$k_B T \ll \omega_{\text{LC}}$$

$$P_e = e^{\frac{\hbar\omega_{\text{LC}}}{k_B T}} \approx 6 \times 10^{-6}$$

$$\langle \hat{\Phi}^2 \rangle = \frac{\hbar Z_0}{2} \coth \left(\frac{\hbar\omega_{\text{LC}}}{2k_B T} \right)$$

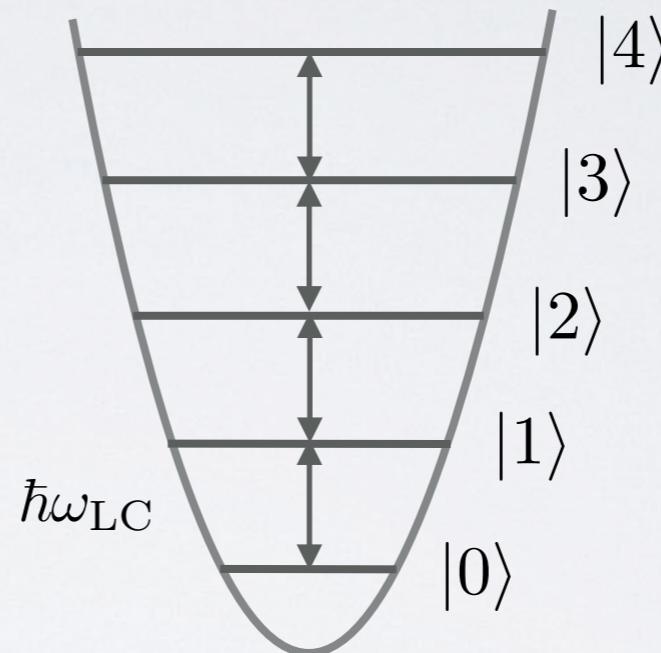
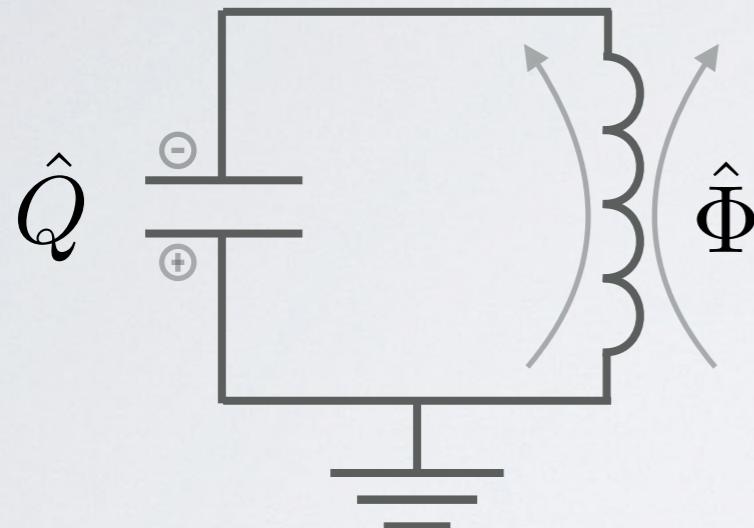
$$\begin{matrix} \Phi \\ Q \end{matrix} \rightarrow \begin{matrix} \hat{\Phi} \\ \hat{Q} \end{matrix}$$

$$[\hat{\Phi}, \hat{Q}] = i\hbar$$

Quantum fluctuations of collective d.o.f!

CIRCUIT QUANTIZATION

Quantum circuit spectrum and energy levels



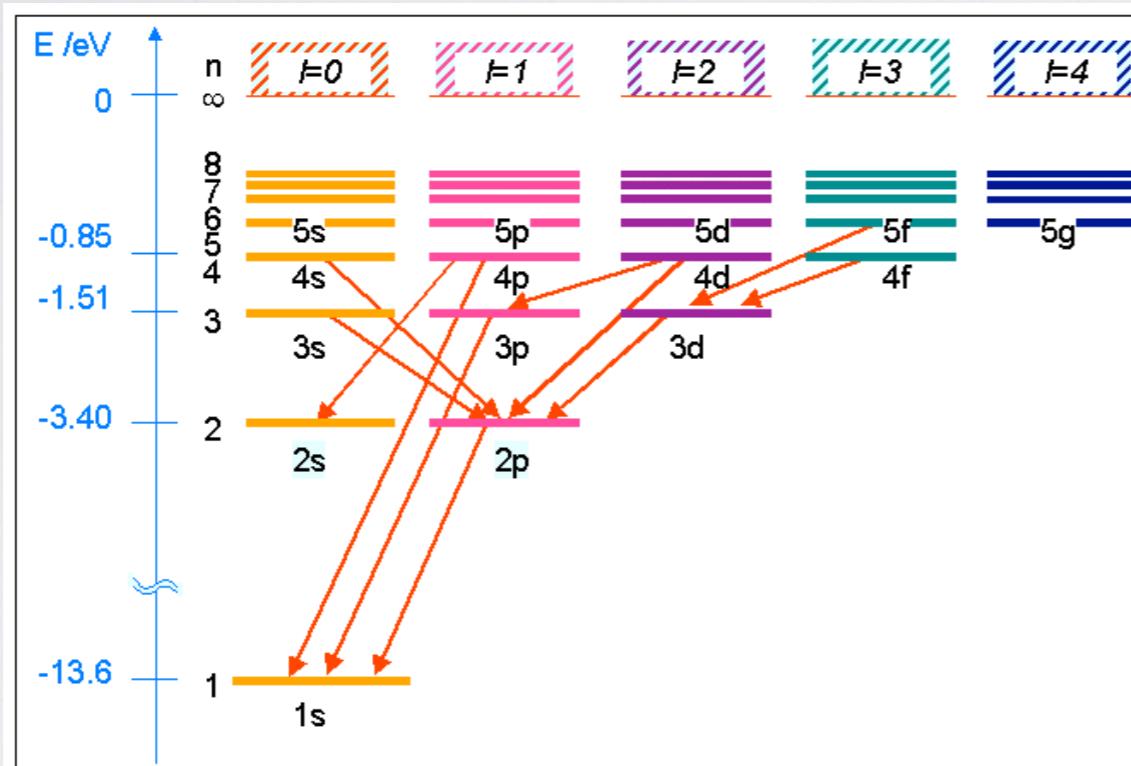
$$Q \gg 1$$

Superconducting
quantum circuit!

Problems:

- We cannot single out a particular transition
- Not a proper quantum system for quantum information processing

CIRCUIT QUANTIZATION



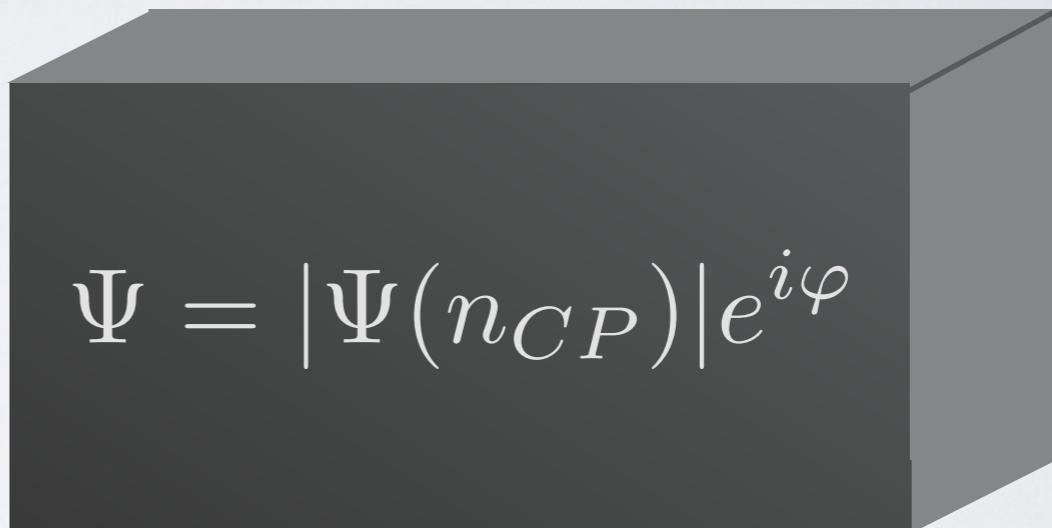
Need nonlinear circuit element to alter circuit spectrum...

- Atomic spectra intrinsically anharmonic
- All transitions addressable by external radiation

CIRCUIT QUANTIZATION

Superconductivity for qubits 101

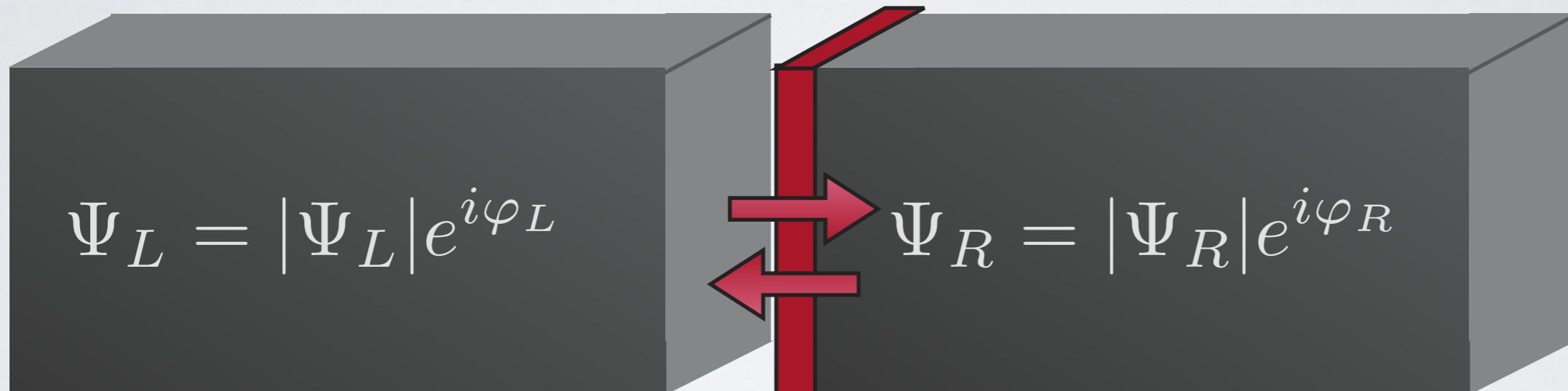
$$T < T_C$$



- Macroscopic quantum system with collective wave function
- No dissipation by charge transport
- Charge and phase conjugate variables: $[\hat{\varphi}, \hat{n}] = i \longleftrightarrow [\hat{x}, \hat{p}] = i\hbar$

CIRCUIT QUANTIZATION

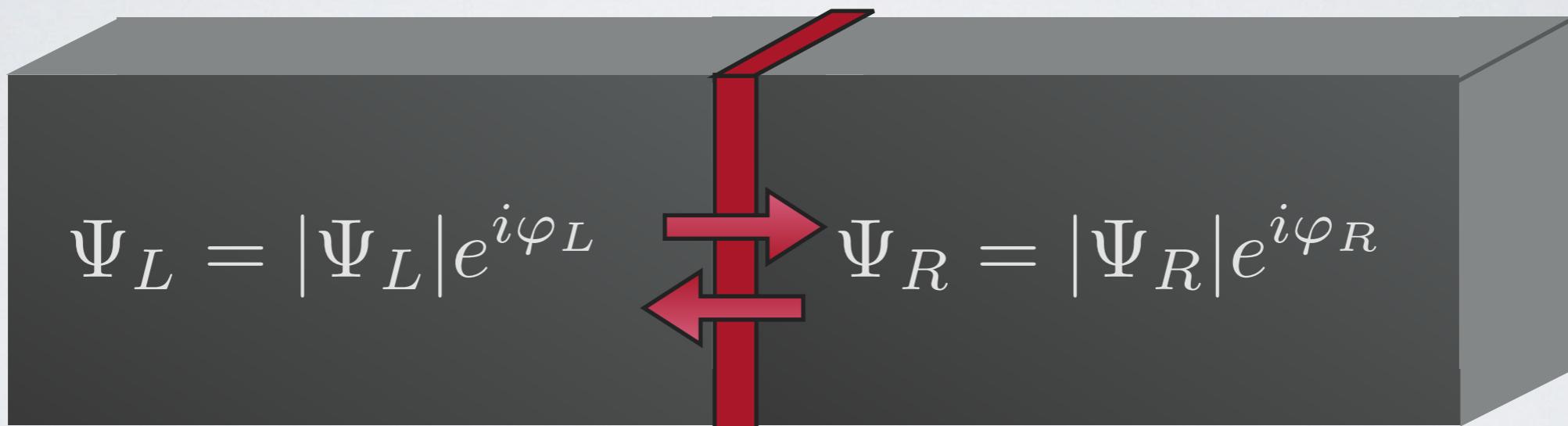
Superconductivity for qubits 101



- Each block has its own wavefunction
- Perfectly isolated systems
- Insulating layer separates superconducting blocks
- Quantum tunneling is established between both sides leading to SUPERCURRENTS. It's a Josephson tunnel junction.

CIRCUIT QUANTIZATION

Josephson junctions: a macroscopic quantum effect



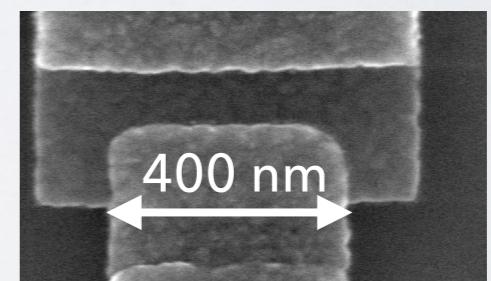
Constituent relation:

$$\frac{dI}{dt} = V \frac{2\pi I_C}{\Phi_0} \sqrt{\frac{I_C^2 - I^2}{I_C^2}}$$

$$\equiv 1/L_J(I)$$

Josephson inductance

Electrical circuits symbol



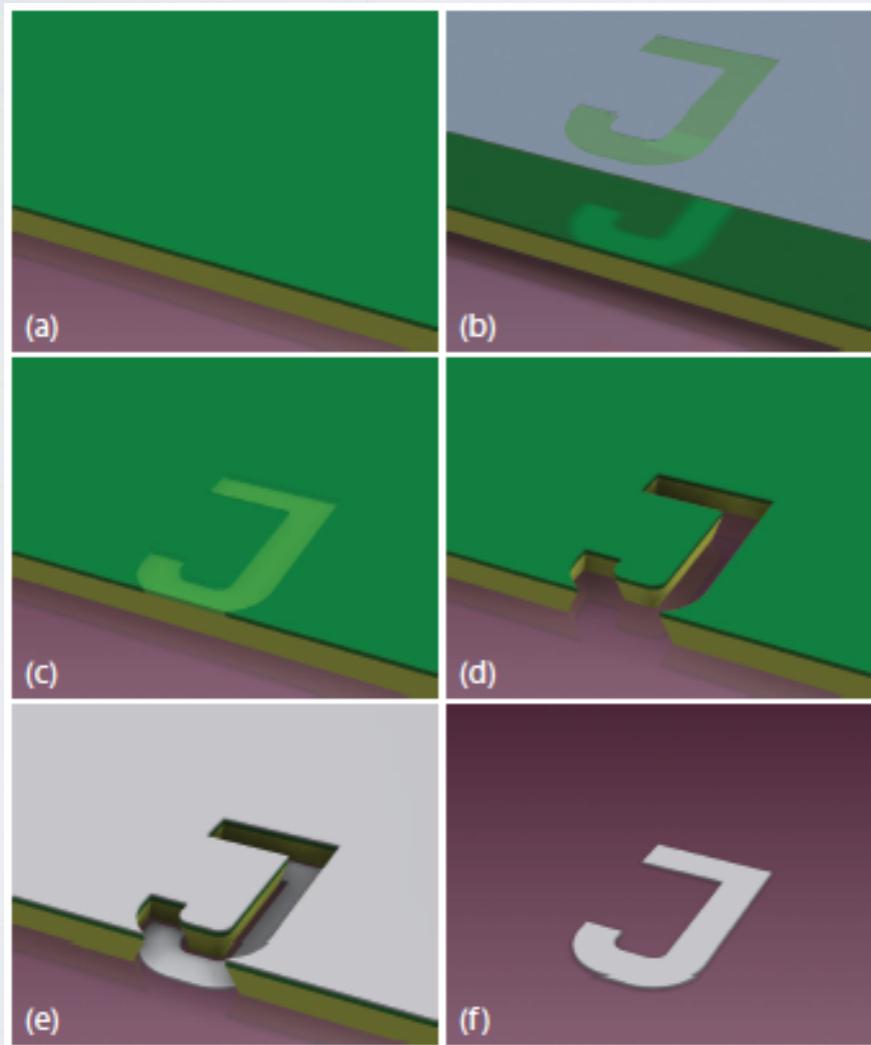
SEM micrograph

- ▶ Nonlinear circuit element
- ▶ Lossless

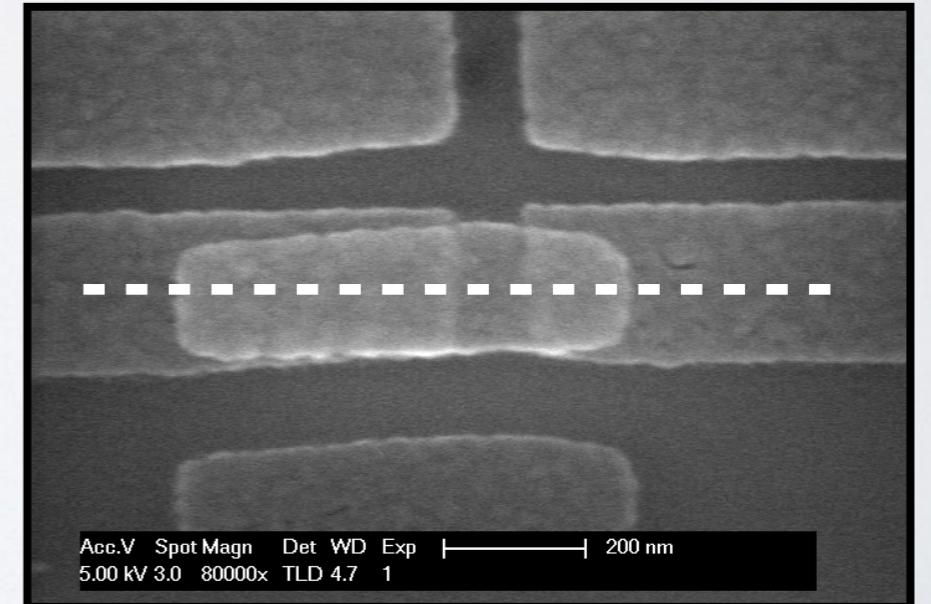
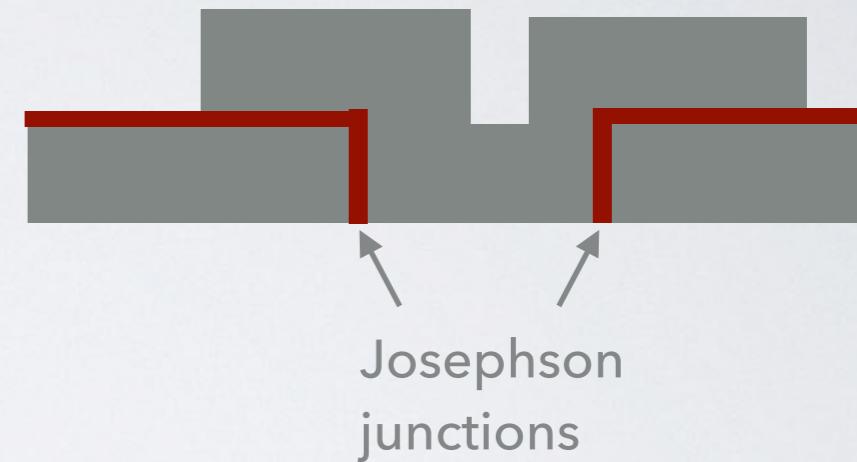
CIRCUIT QUANTIZATION

Junction fabrication

Electron beam lithography



Shadow angle evaporation



Scanning electron microscope (SEM)

OUTLINE

Lecture I

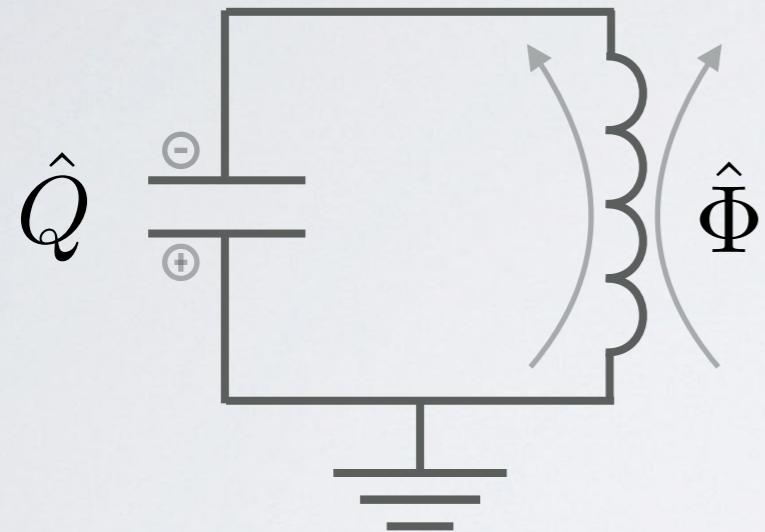
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- Circuit quantization
- **Superconducting qubit zoo**
- Qubit state control

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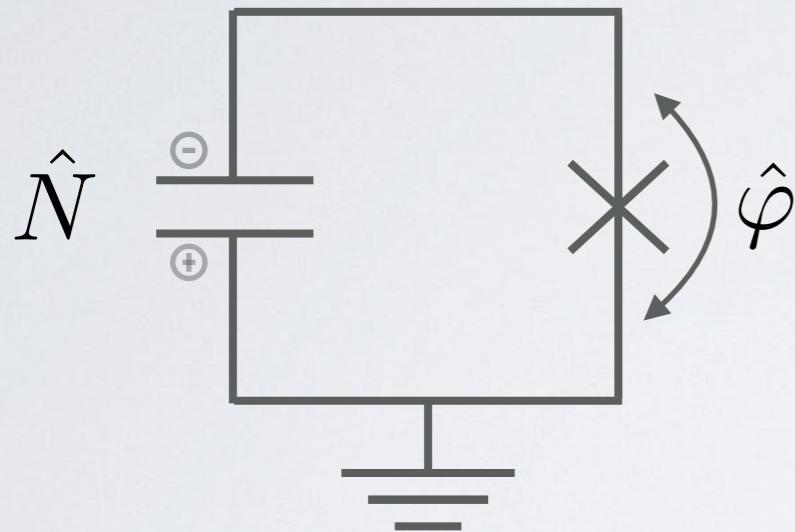
SUPERCONDUCTING QUBITS

The hydrogen atom of superconducting qubits

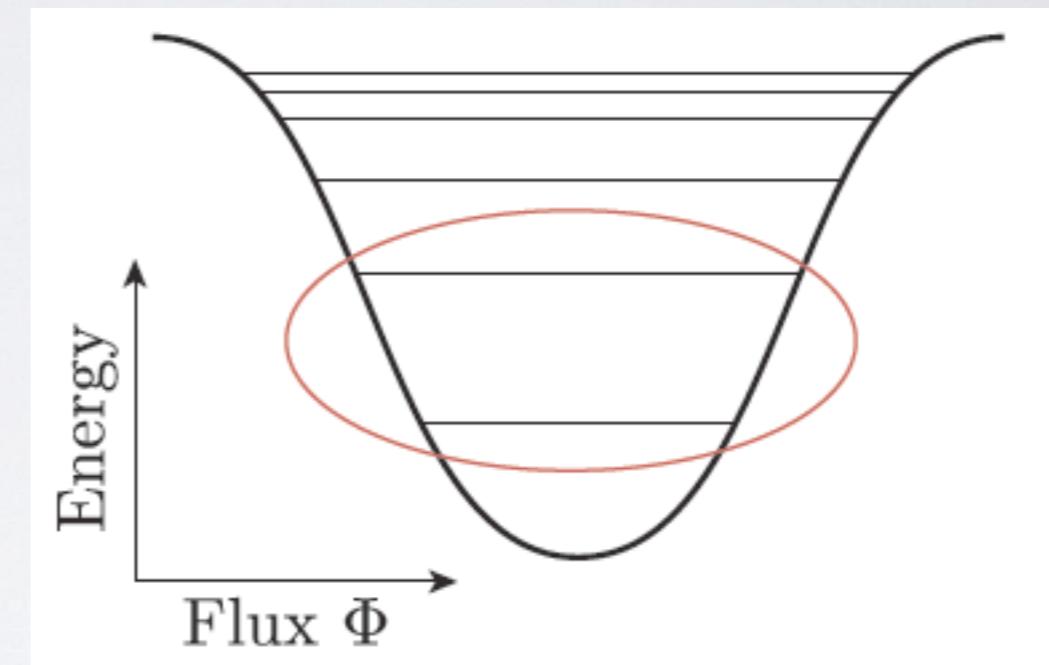


SUPERCONDUCTING QUBITS

The hydrogen atom of superconducting qubits



$$\hat{\mathcal{H}} = E_C \hat{N}^2 - E_J \cos(\hat{\varphi})$$

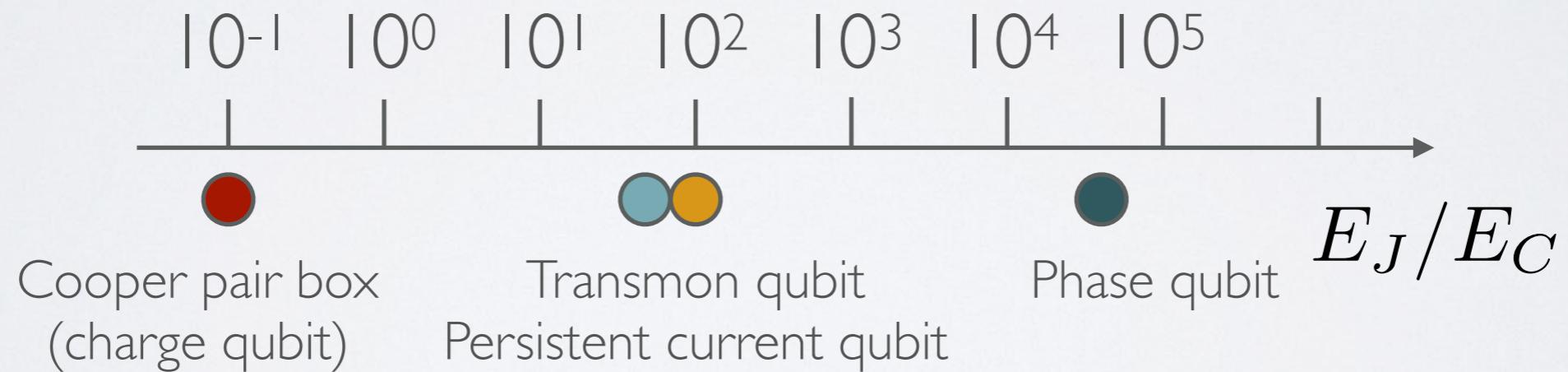


- Anharmonic spectrum
- Energy levels related to charge tunneling
- Lowest two states form a qubit

SUPERCONDUCTING QUBITS

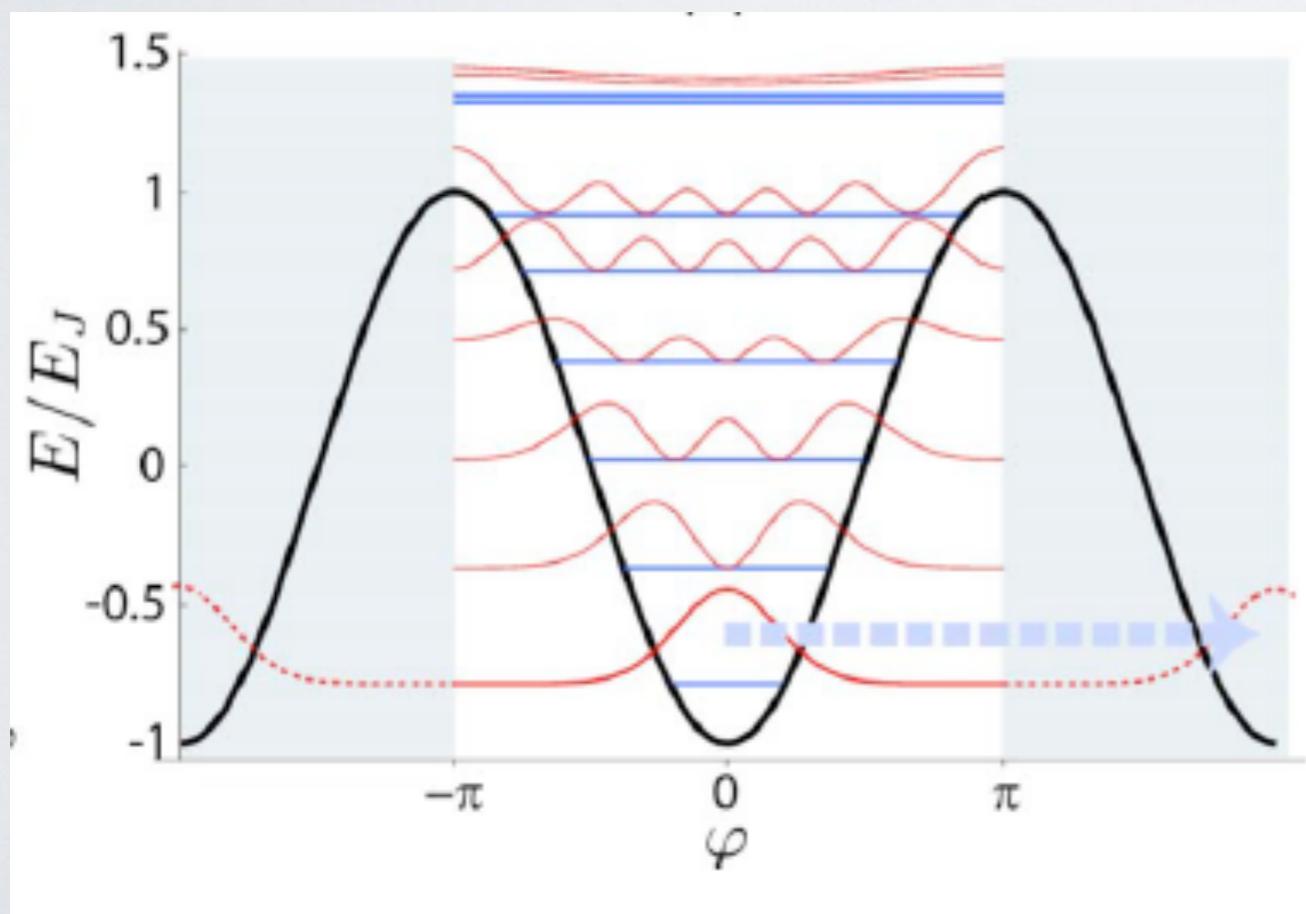
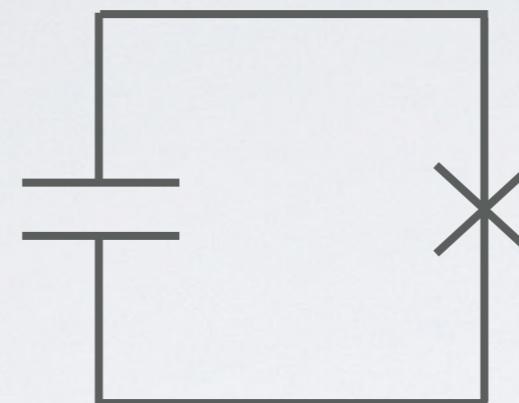
A zoo of superconducting qubits

$$\hat{\mathcal{H}} = E_C \hat{N}^2 - E_J \cos(\hat{\varphi})$$

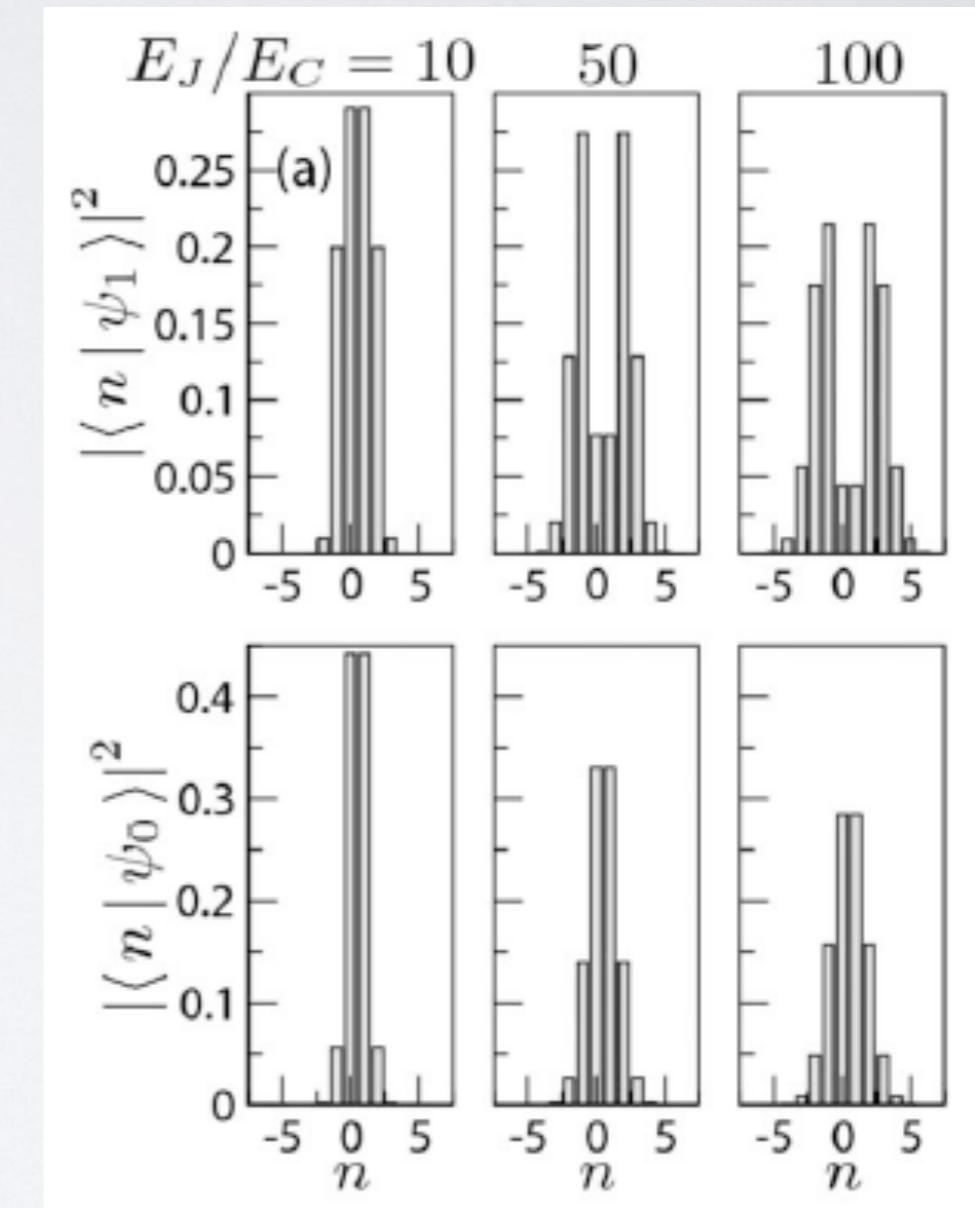


SUPERCONDUCTING QUBITS

Transmon qubit



Eigenstates: charge superposition



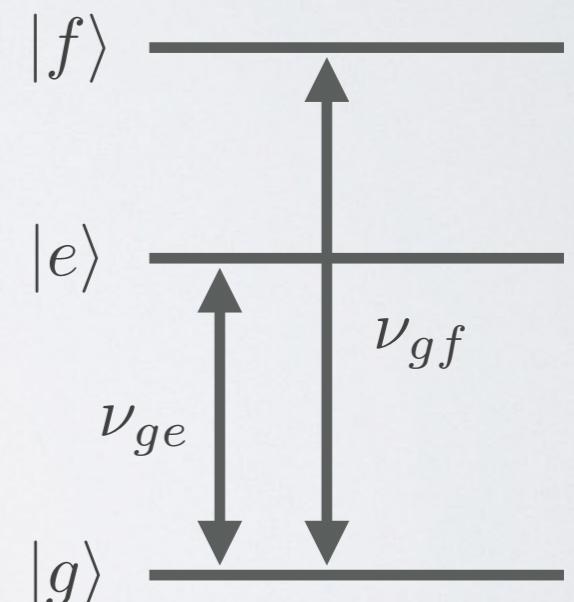
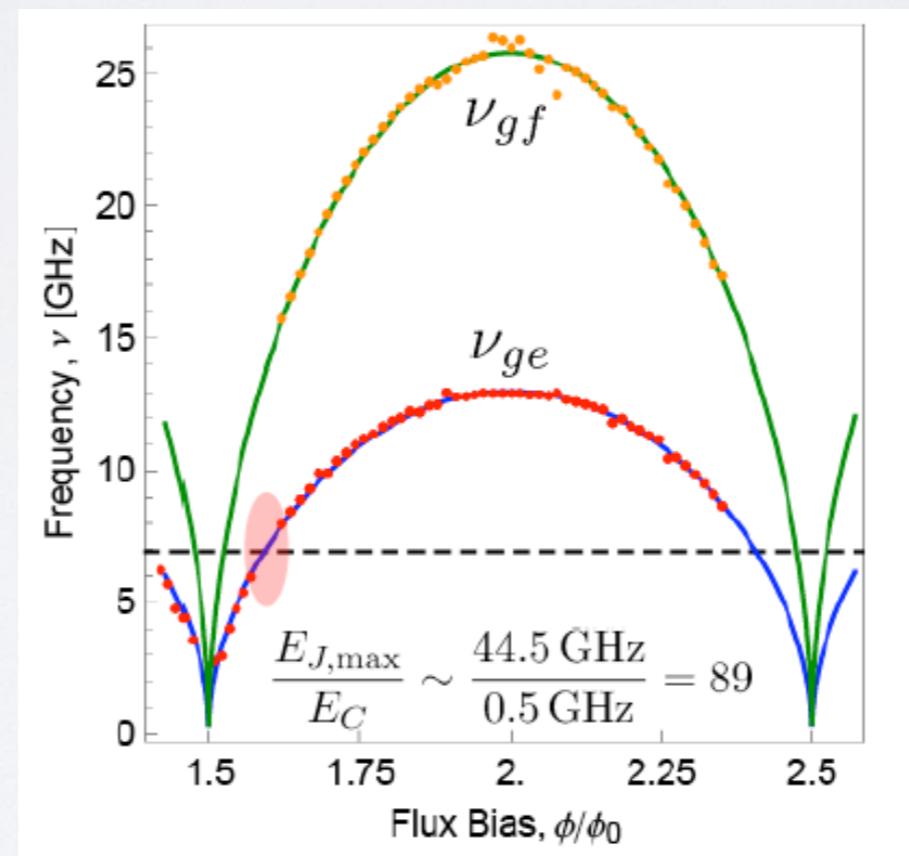
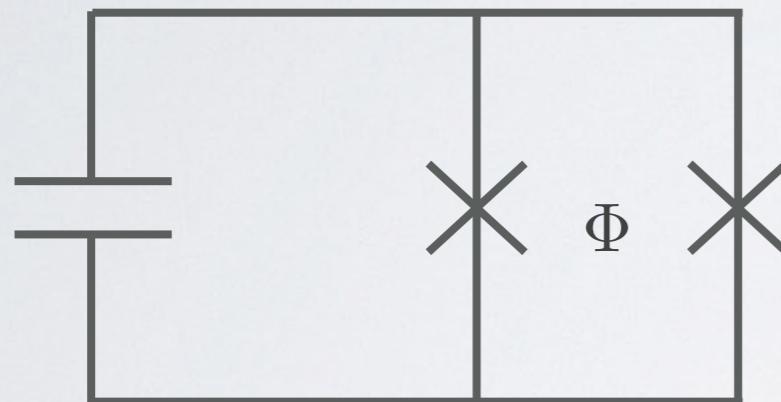
Koch et al. PRA 76, 042319 (2007)

SUPERCONDUCTING QUBITS

Transmon qubit



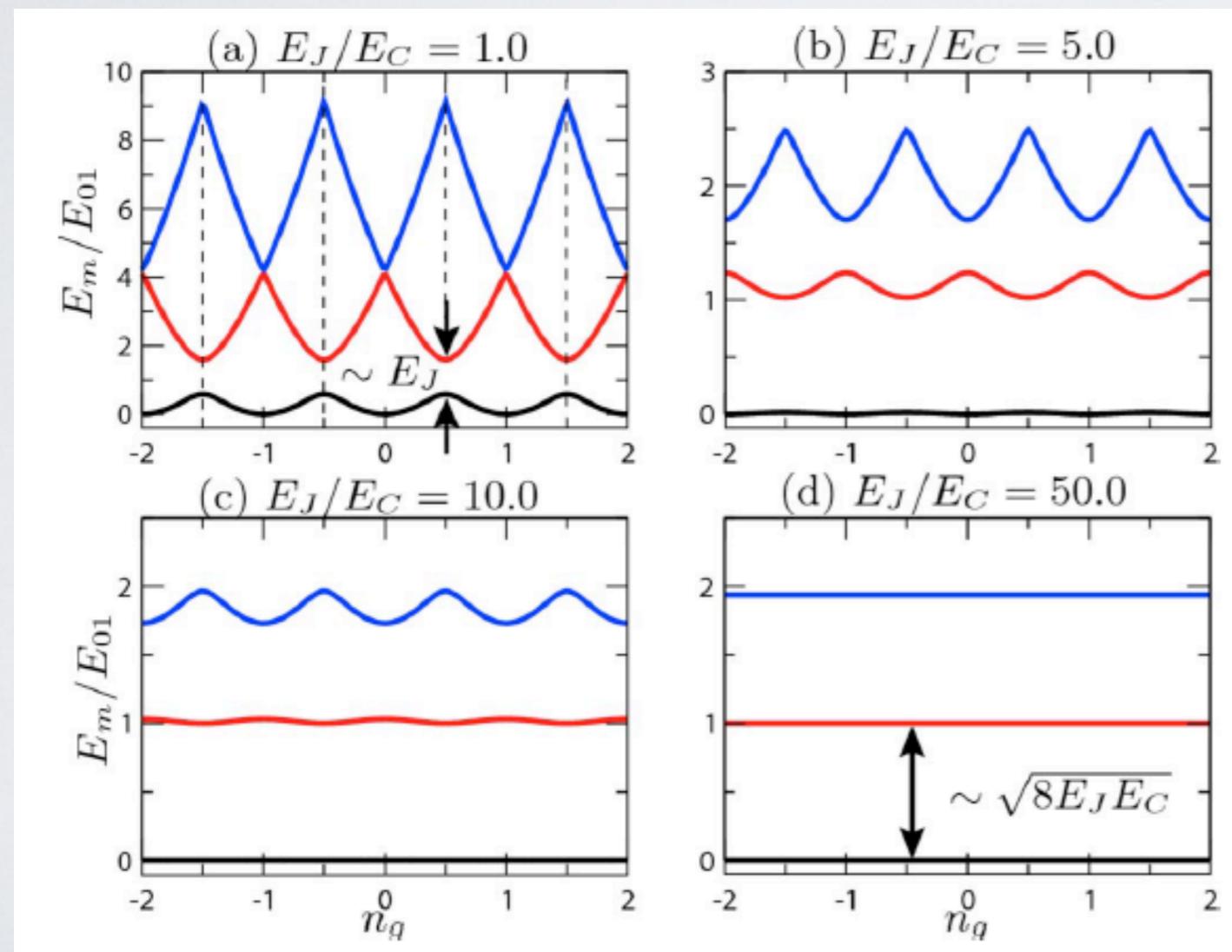
Flux dependence of energy levels



SUPERCONDUCTING QUBITS

Transmon qubit

Charge insensitivity of transmon from circuit design



Flat bands mean
larger insensitivity

System remains slightly
anharmonic

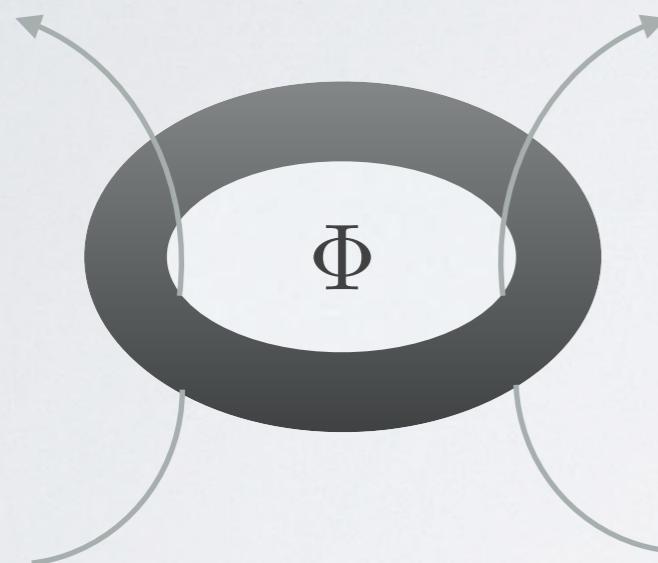
$$\omega_{01} - \omega_{12} = -\alpha = -E_C$$

anharmonicity

SUPERCONDUCTING QUBITS

Flux qubit

Superconducting ring: flux quantized

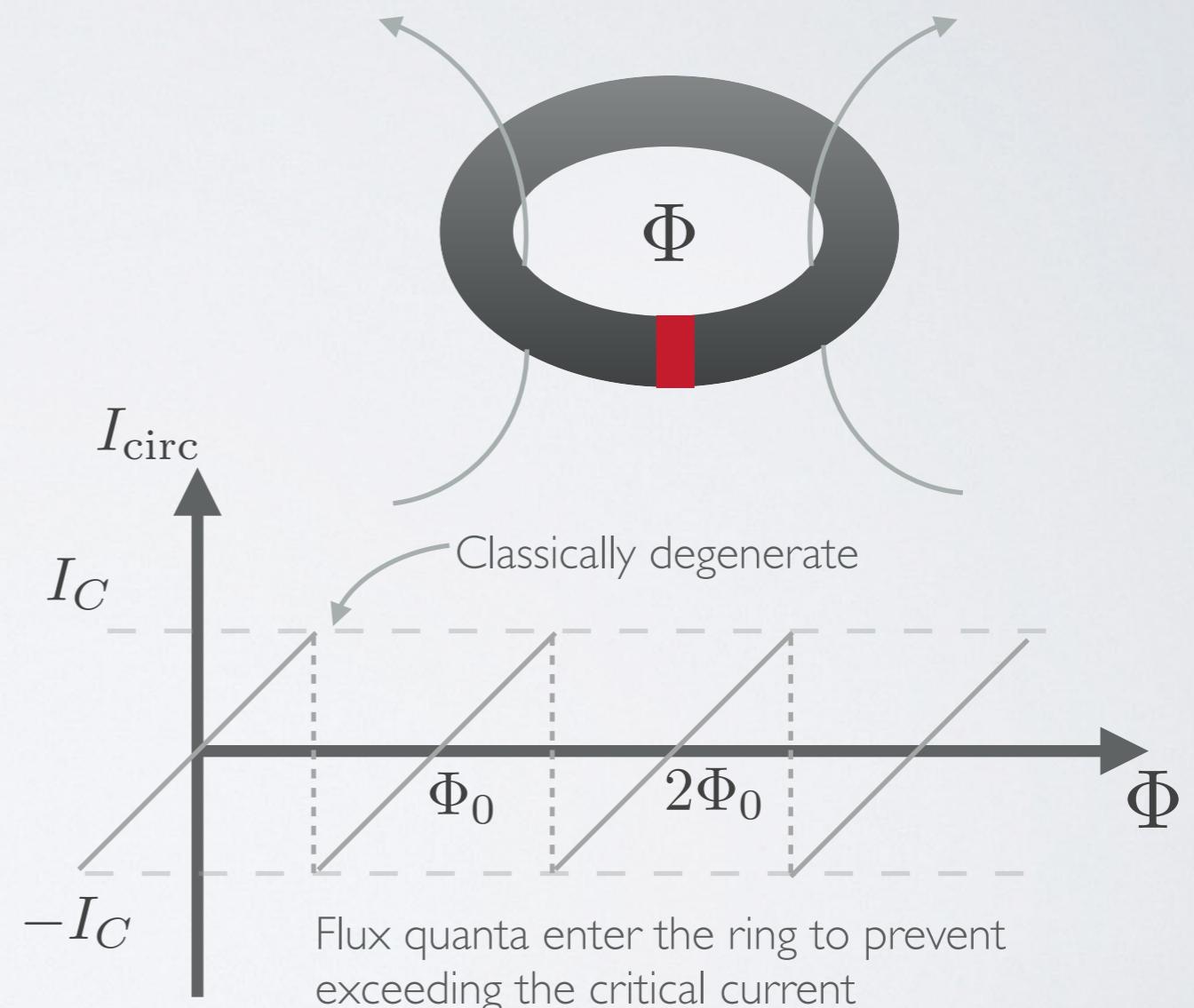


$$\Psi = |\Psi(n_{CP})| e^{i\varphi}$$

$$\Delta\varphi = \varphi + 2\pi N$$

fixed upon cooldown:
Number of flux quanta trapped

Single-junction loop: rf-SQUID / Flux qubit

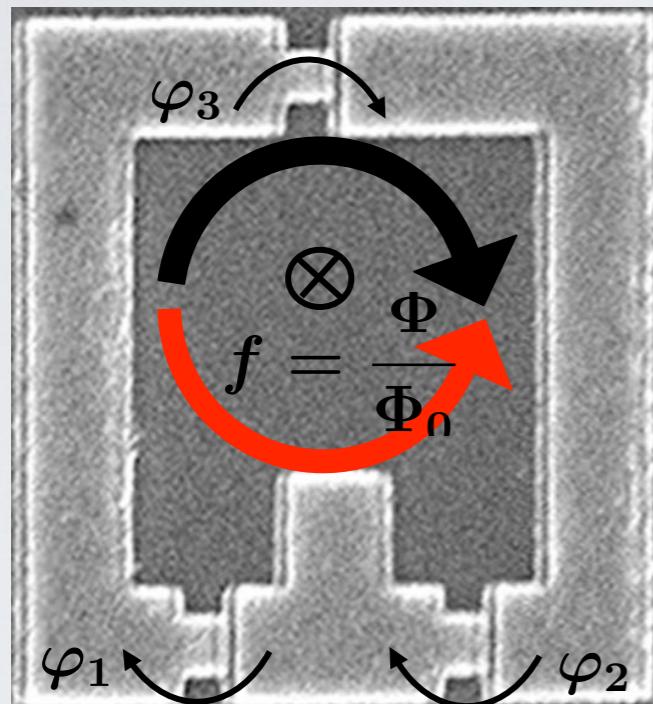


But: large loop, too sensitive to flux noise...

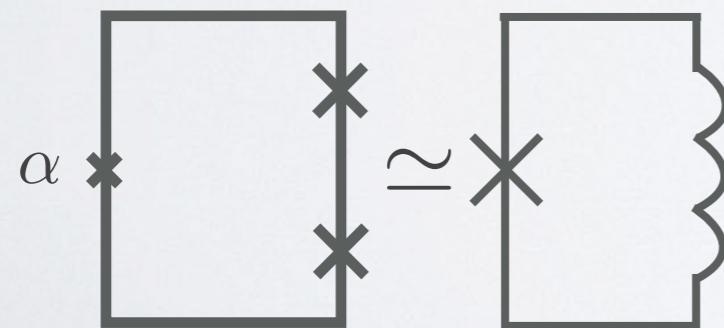
SUPERCONDUCTING QUBITS

Persistent current qubit

Aluminum loop with 3 Josephson junctions



Circuit diagram:

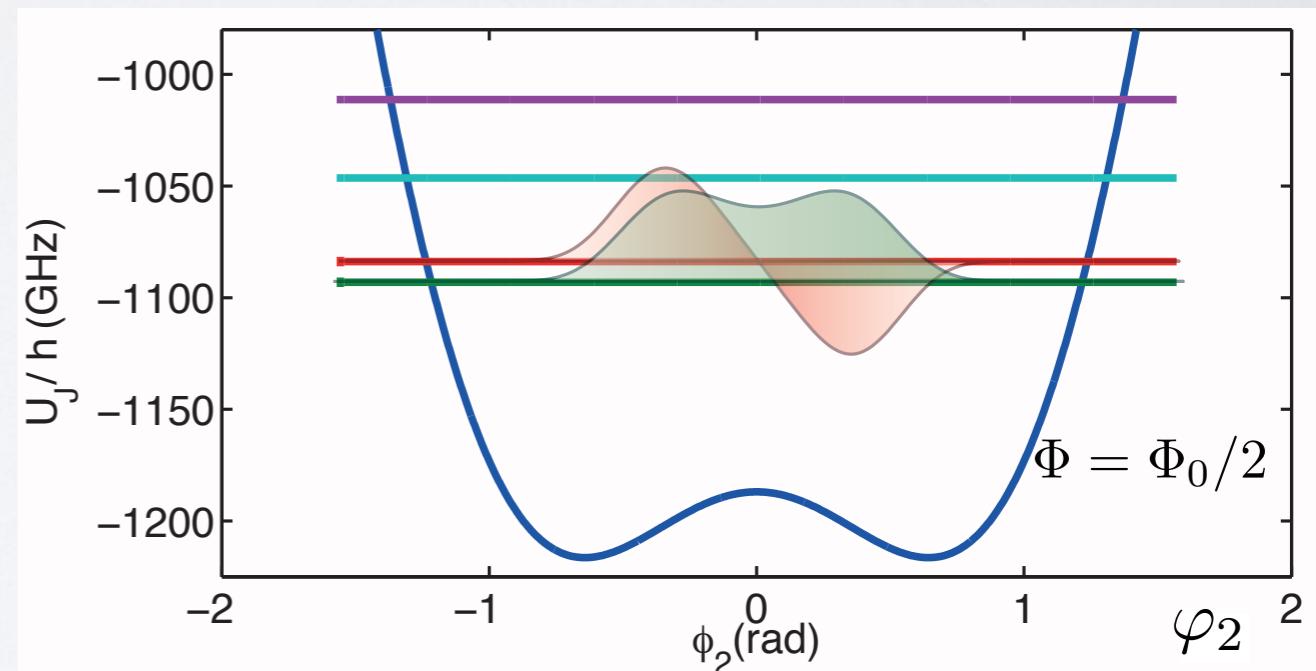


T. P. Orlando, et al. PRB 60 15398 (1999)

J. E. Mooij, et al. Science 285 1036 (1998)

Circuit Hamiltonian:

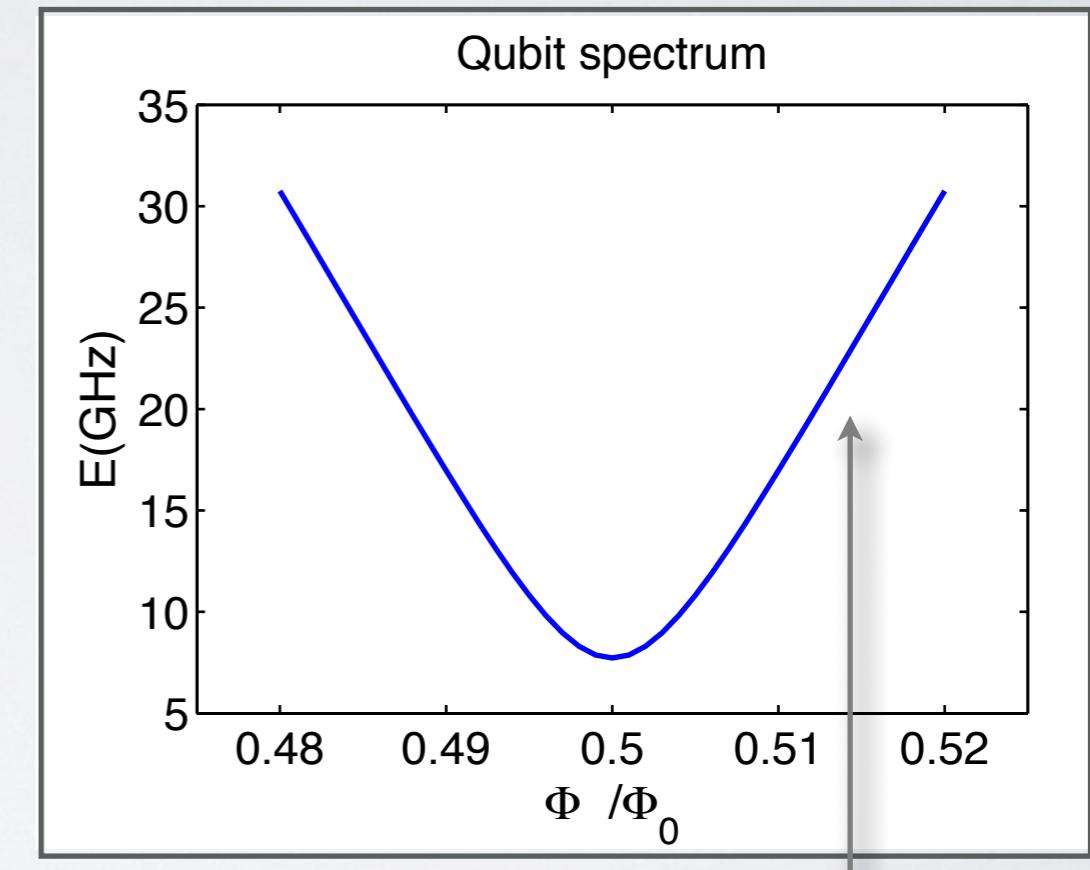
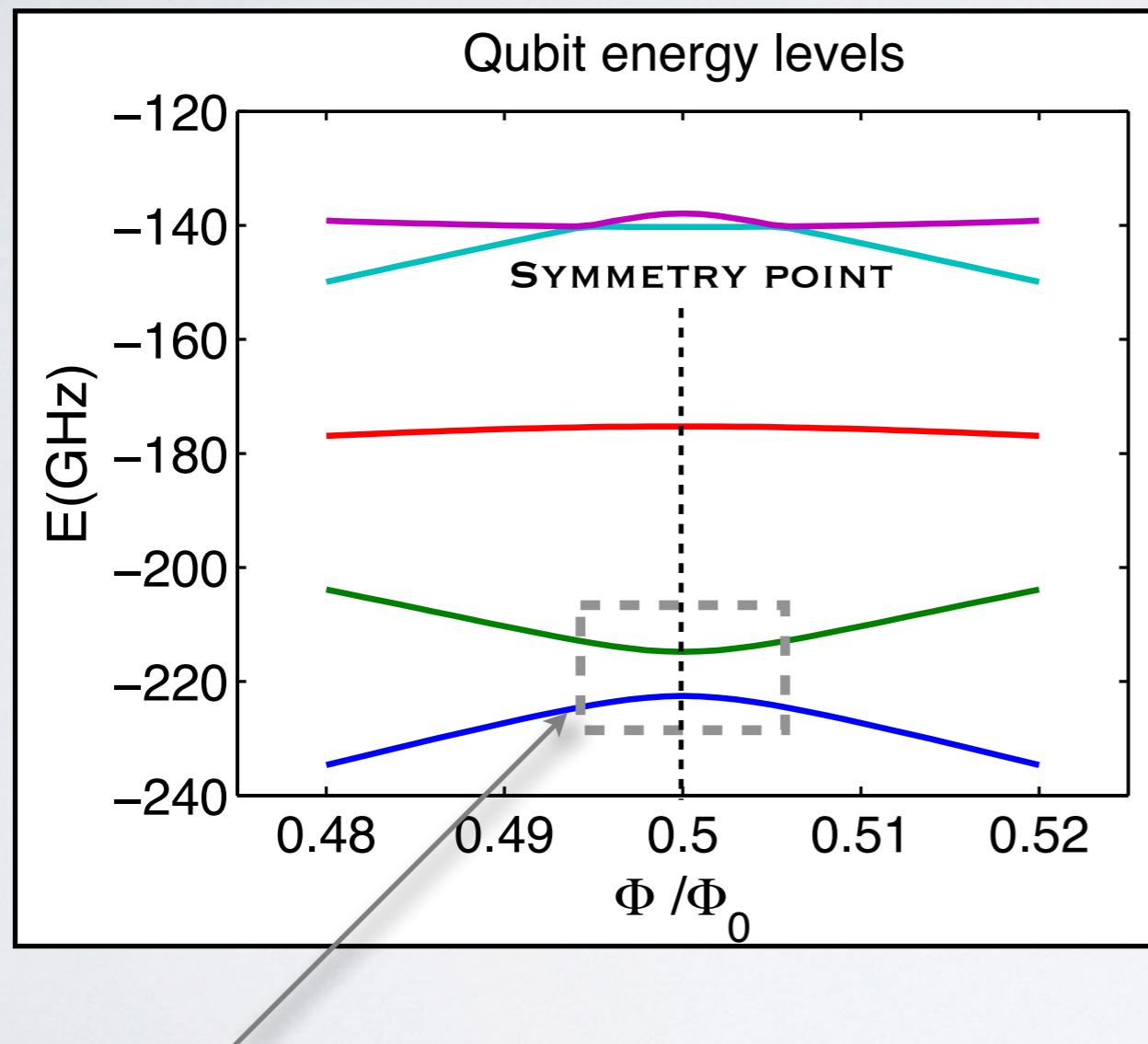
$$\mathcal{H}_{\text{PCQ}} = \frac{E_C}{1+2\alpha} [n_1^2 - 2\alpha n_1 n_2 + (1+\alpha)n_2^2] + E_J [\cos \varphi_1 + \cos \varphi_2 + \alpha \cos(\varphi_1 + \varphi_2 + 2\pi\Phi/\Phi_0)]$$



The tunneling through or over the barrier couples states into superposition of phase, hence current states

SUPERCONDUCTING QUBITS

Persistent current qubit



High sensitivity to
magnetic flux*

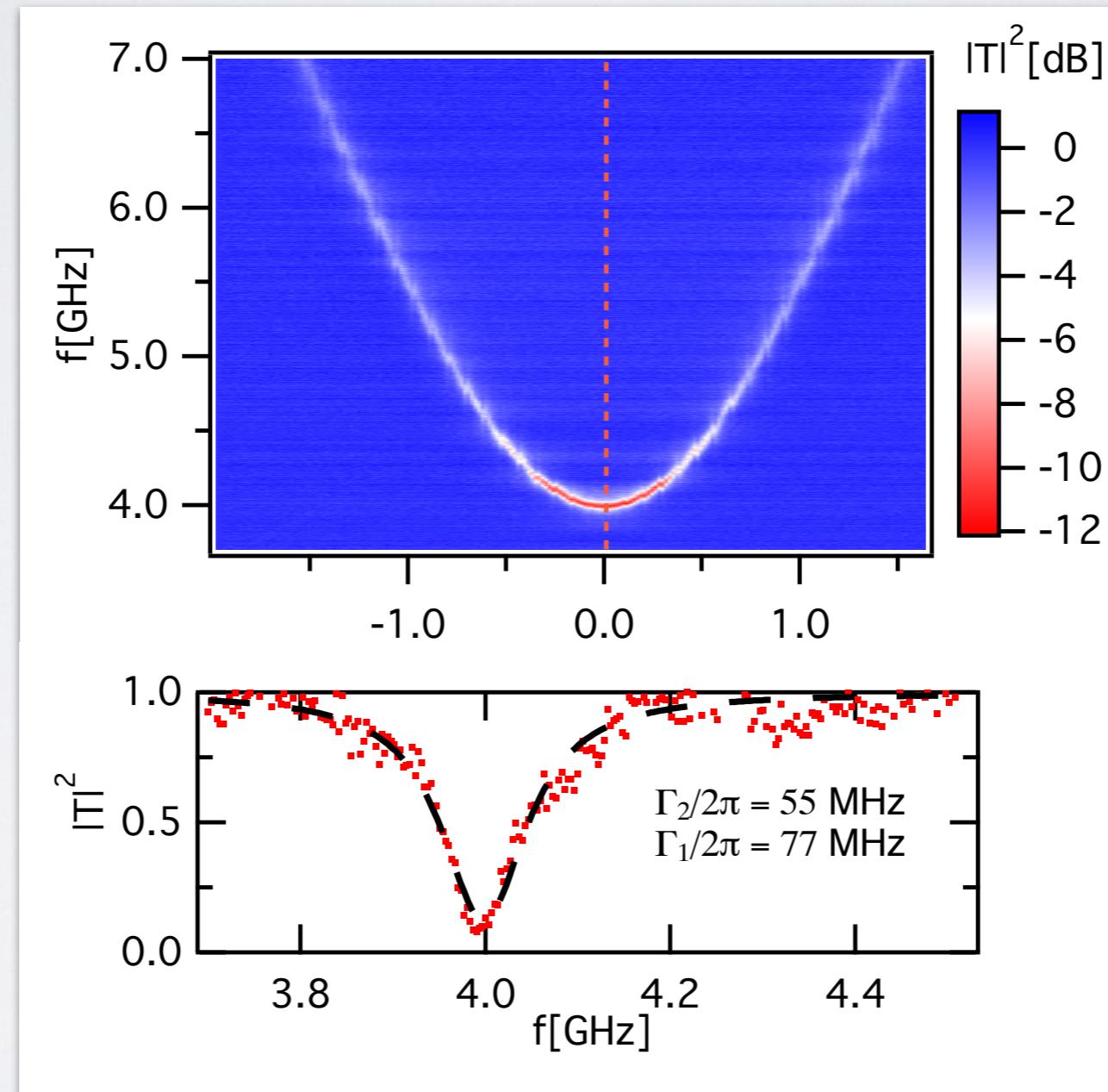
Persistent current states are degenerate in energy

*M. Bal, et al., Nat. Comm. 3, 1324 (2012)

SUPERCONDUCTING QUBITS

Persistent current qubit

Experimental
signature?



OUTLINE

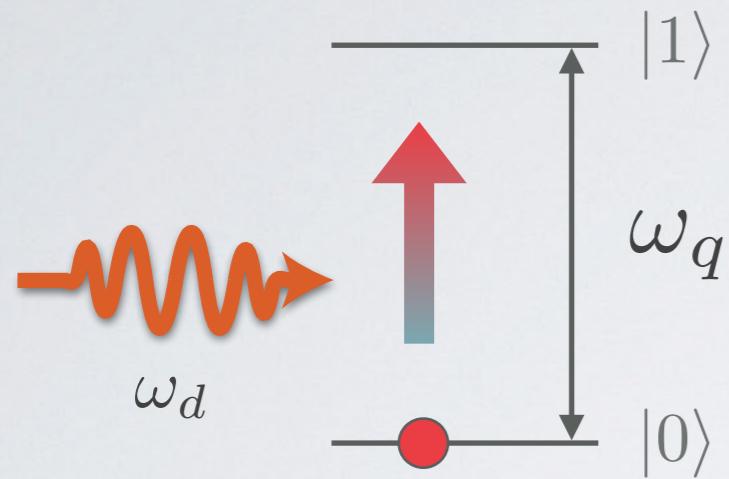
Lecture I

- Quantum computation
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- **Qubit state control**

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QUBIT STATE CONTROL



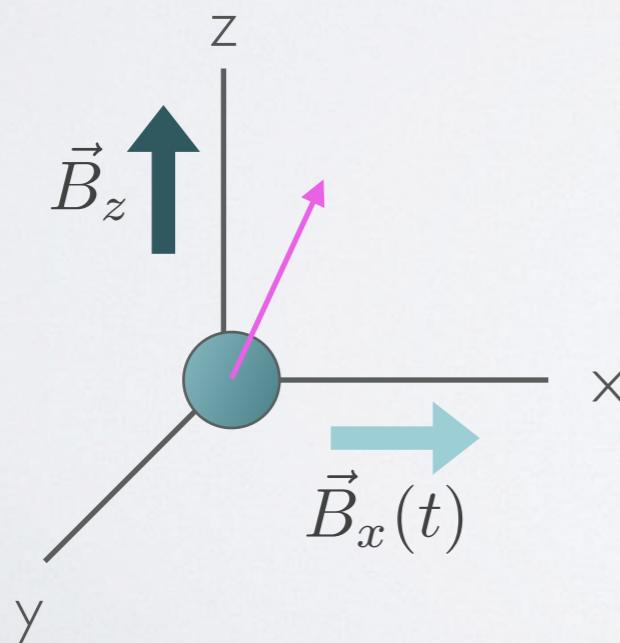
Analogy with spin-1/2

Rabi formula

$$P_{0 \rightarrow 1} = \frac{\Omega_R^2}{\Delta^2 + \Omega_R^2} \sin^2 \left(\sqrt{(\Delta^2 + \Omega_R^2)} t / \hbar \right)$$

Ω_R Rabi frequency “bare”

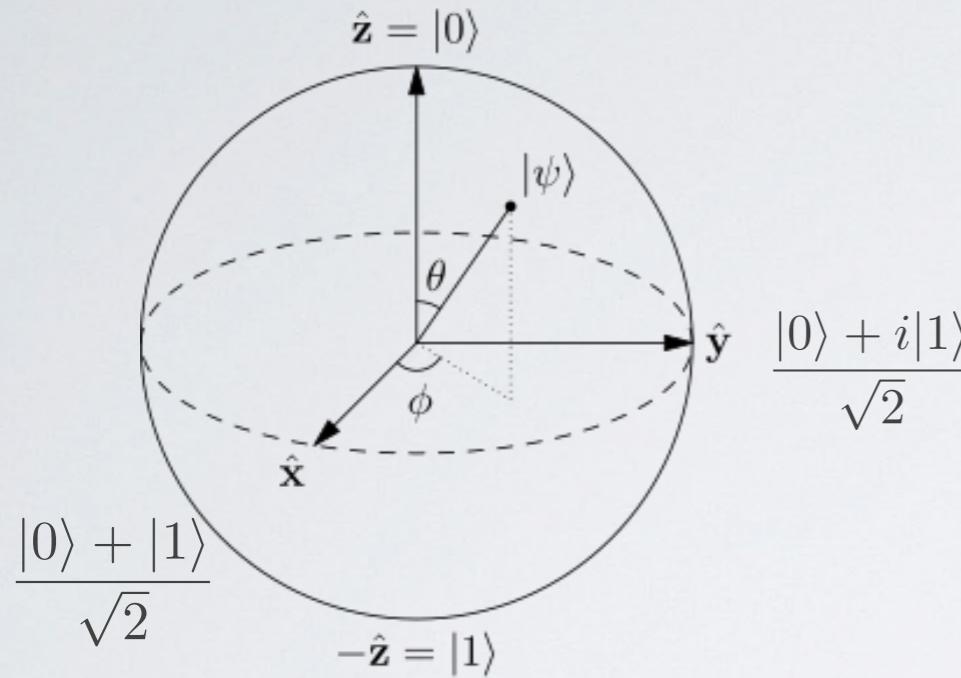
$\Delta \equiv \omega_q - \omega_d$ Detuning



Import all NMR techniques for
qubit state control...

Quantum Mechanics, C.Tannoudji
Nielsen and Chuang, Quantum computation and Quantum information

QUBIT STATE CONTROL



Rabi formula

$$P_{0 \rightarrow 1} = \frac{\Omega_R^2}{\Delta^2 + \Omega_R^2} \sin^2 \left(\sqrt{(\Delta^2 + \Omega_R^2)} t / \hbar \right)$$

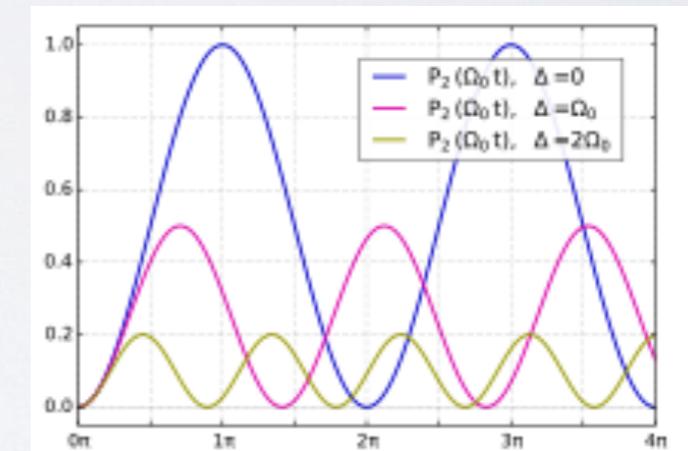
Ω_R Rabi frequency “bare”

$\Delta \equiv \omega_q - \omega_d$ Detuning

$$|\Psi\rangle = \cos \theta |0\rangle + \sin \theta e^{i\phi} e^{-iE_{01}t/\hbar} |1\rangle$$

In lab frame, spin precesses (Larmor frequency),
z-axis control

Rabi oscillations



Wikipedia

x-y axes control

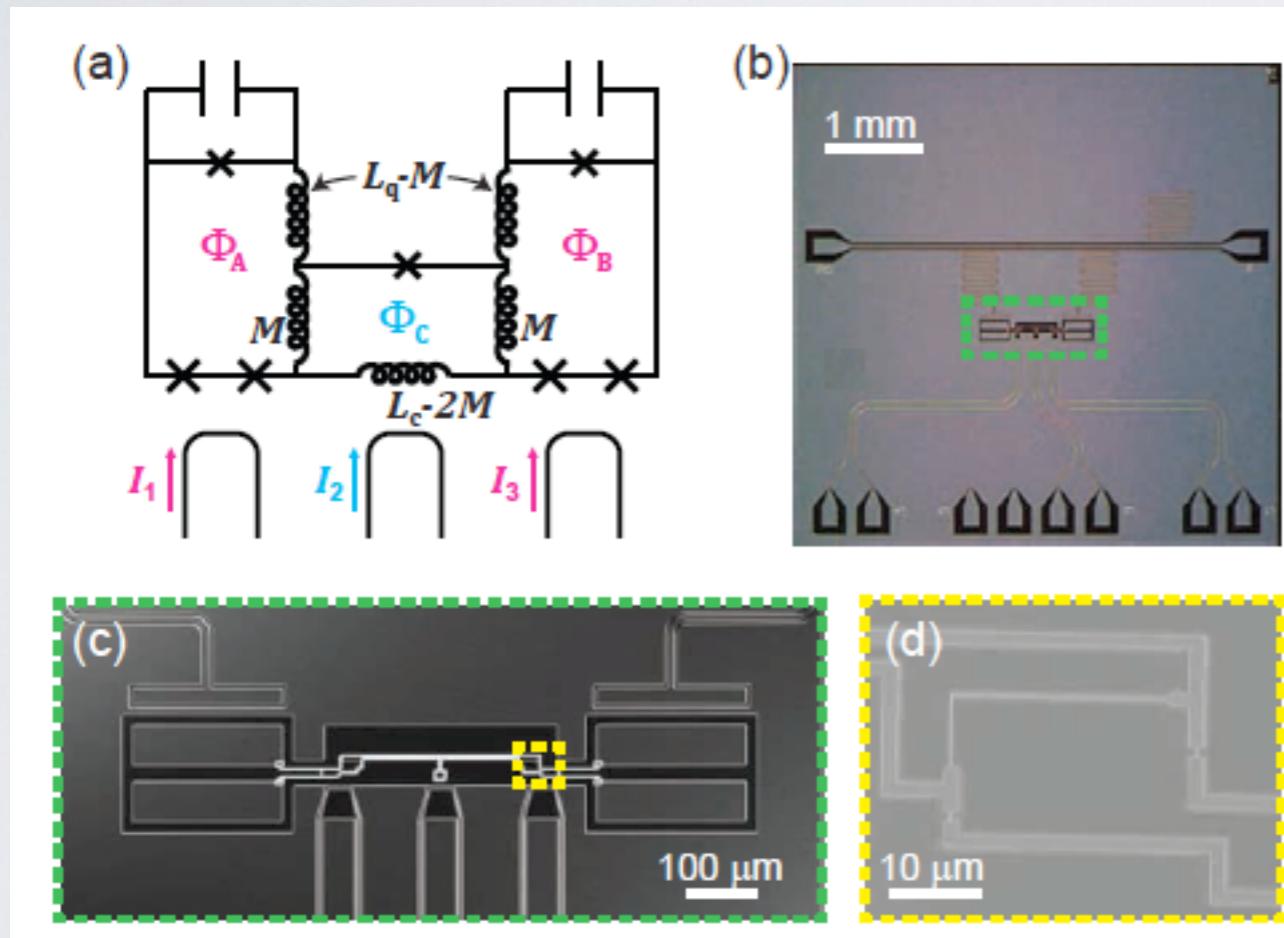
Quantum Mechanics, C.Tannoudji

Nielsen and Chuang, Quantum computation and Quantum information

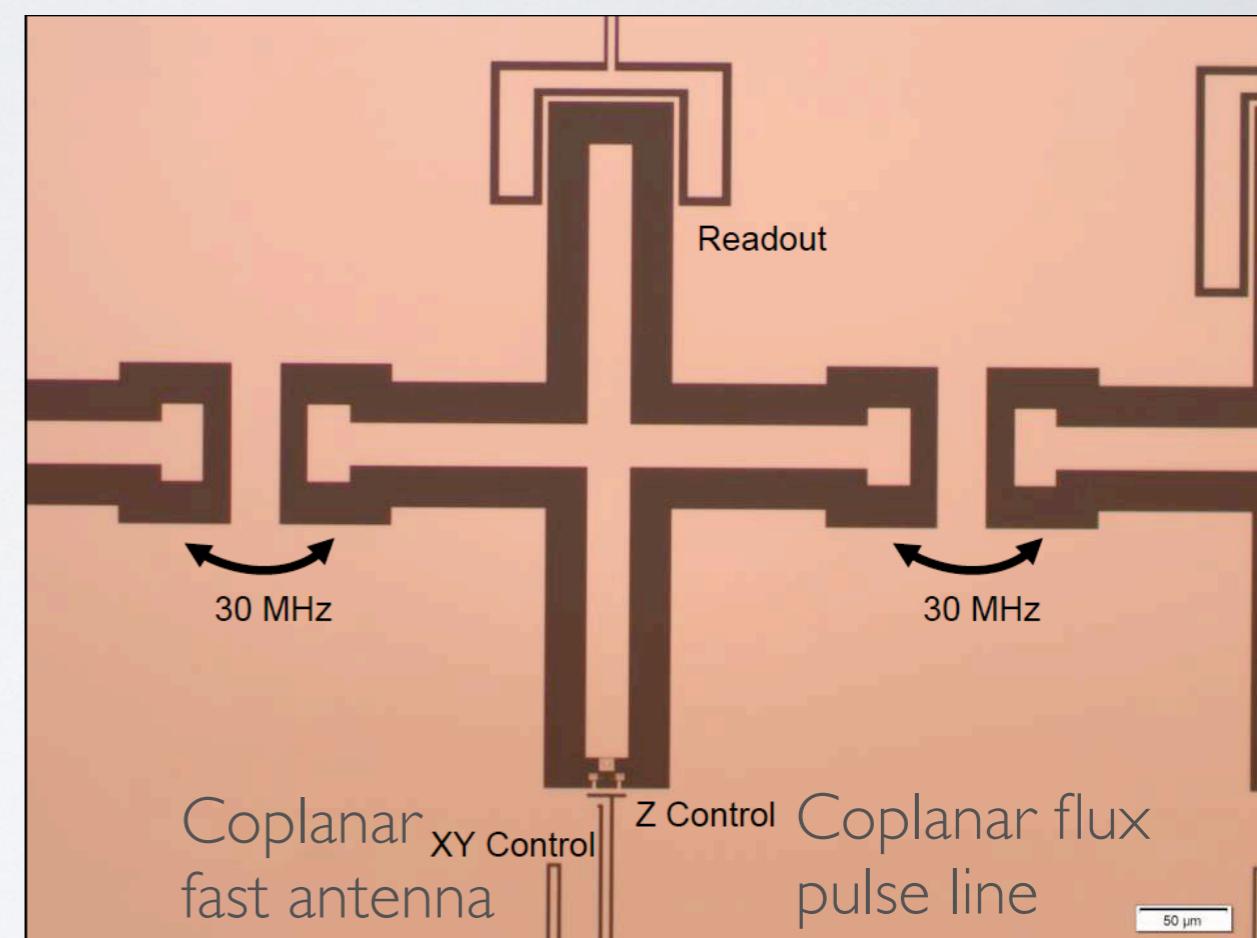
QUBIT STATE CONTROL

Local control lines

Flux qubit



Transmon qubit



Coplanar flux bias lines x-y-z control

Steven J. Weber, Phys. Rev. Applied 8, 014004 (2017)

Barends et al., PRL 111, 080502 (2013)

QUBIT STATE CONTROL

Setup

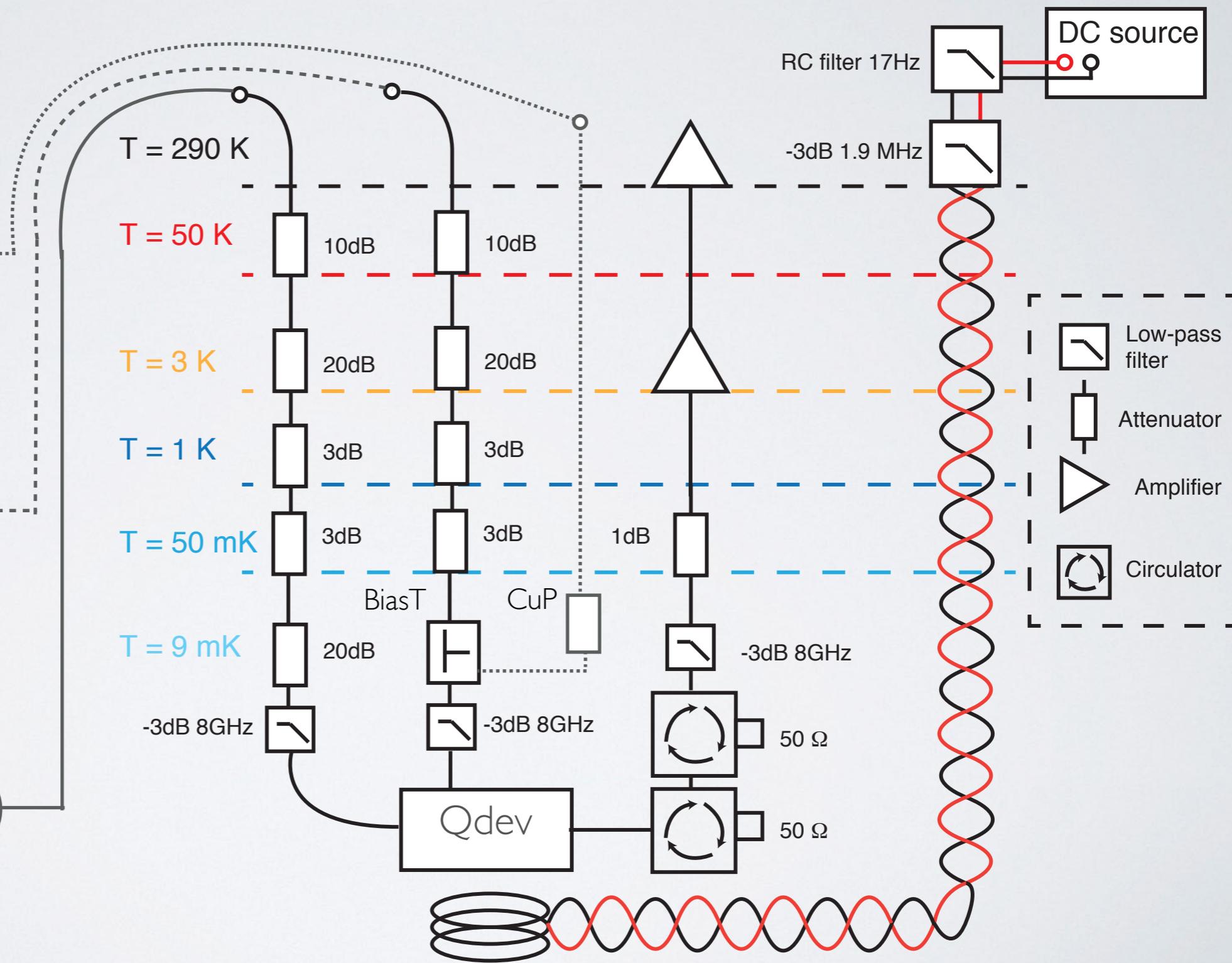
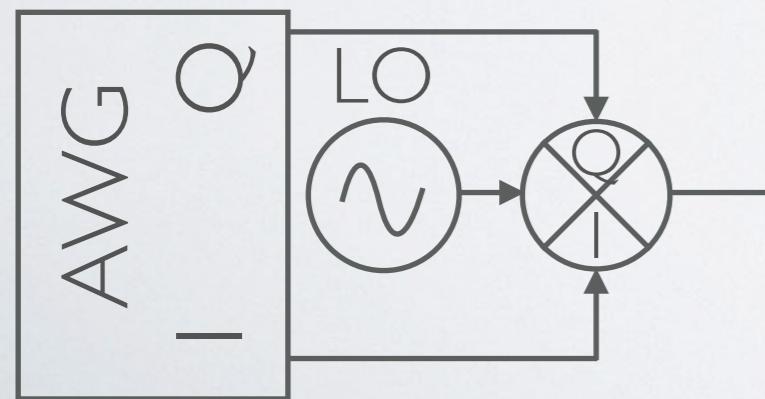
Per qubit:



z control

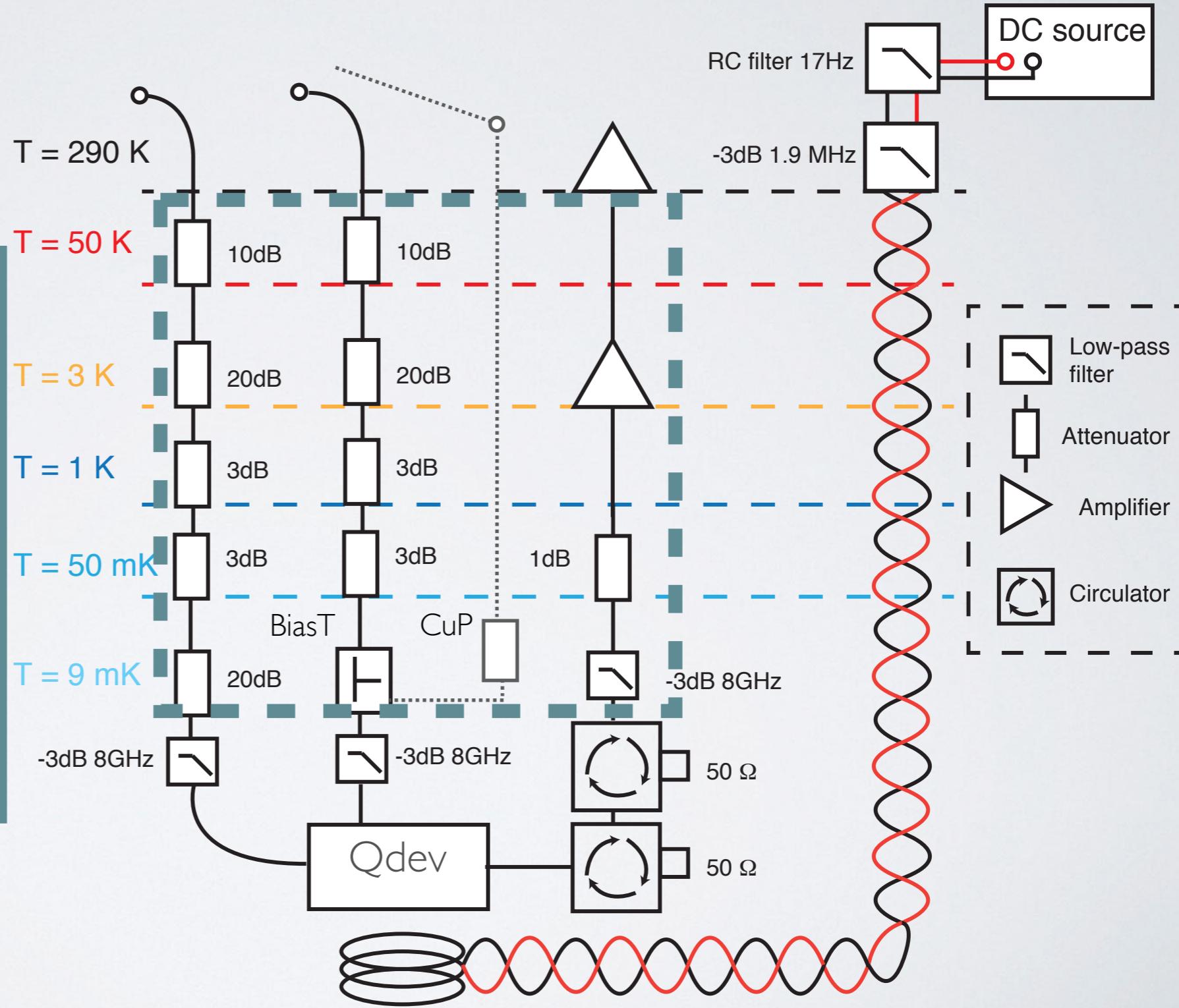
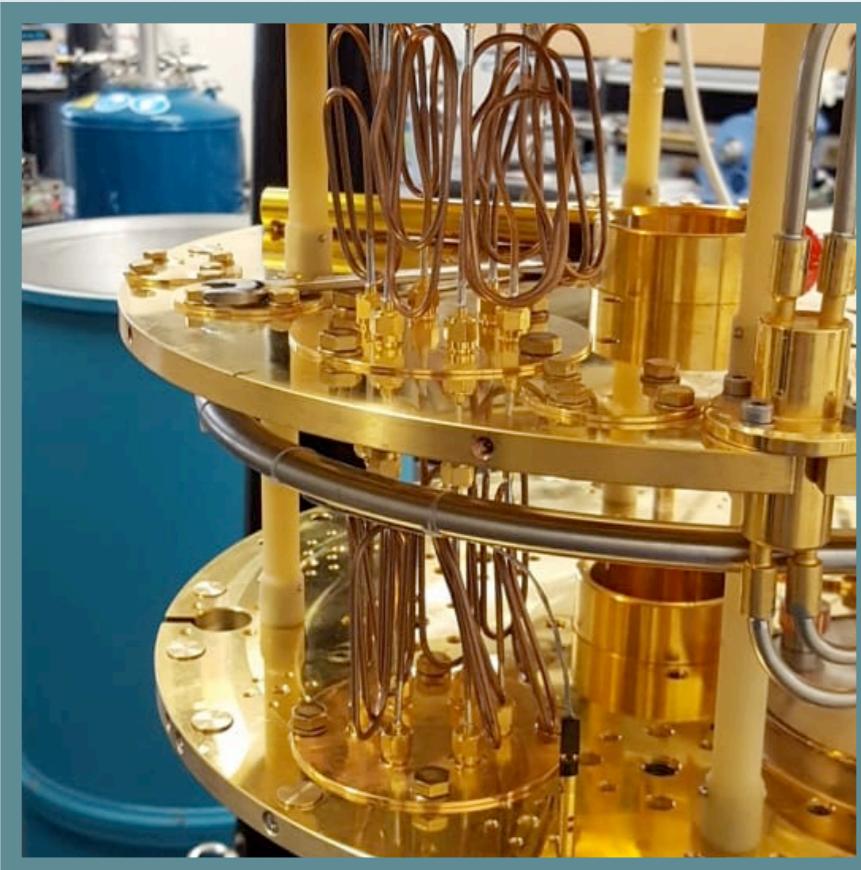


x-y control

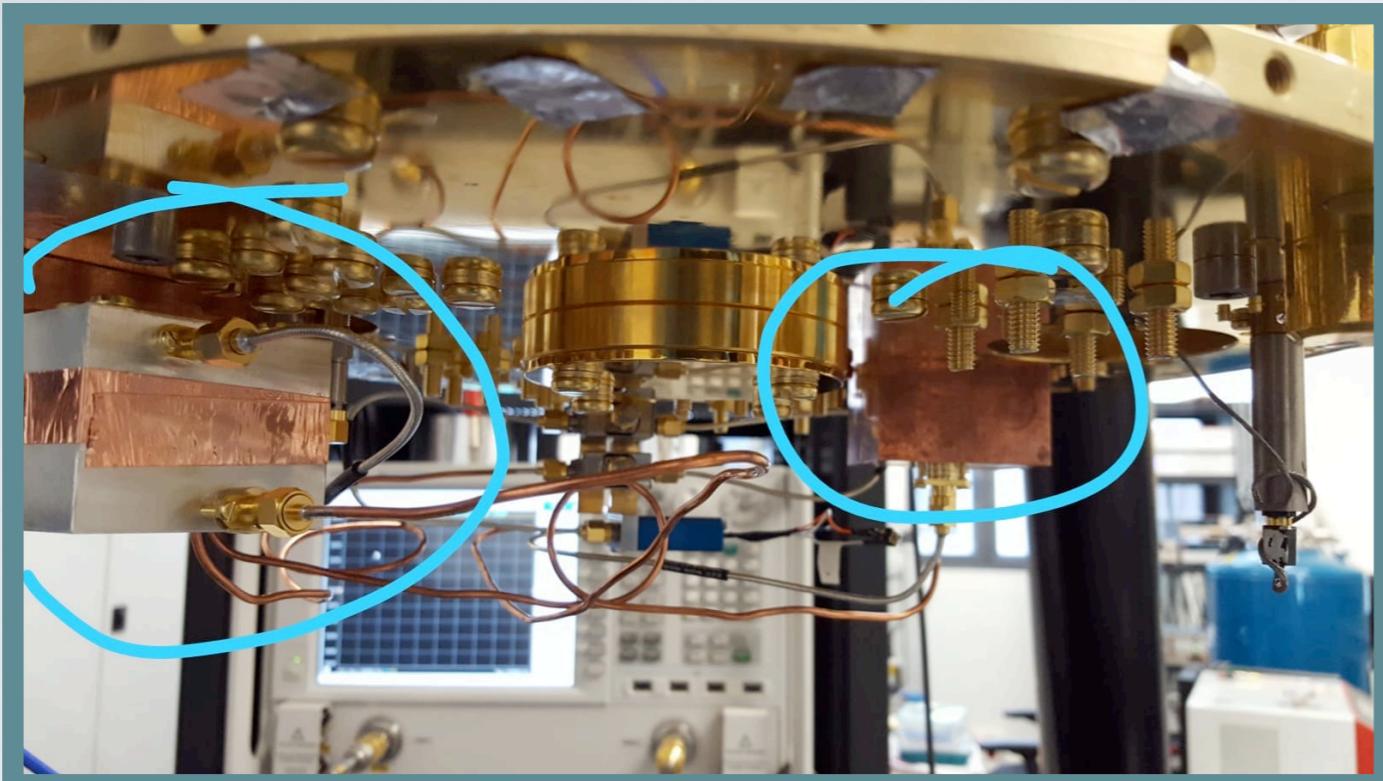


QUBIT STATE CONTROL

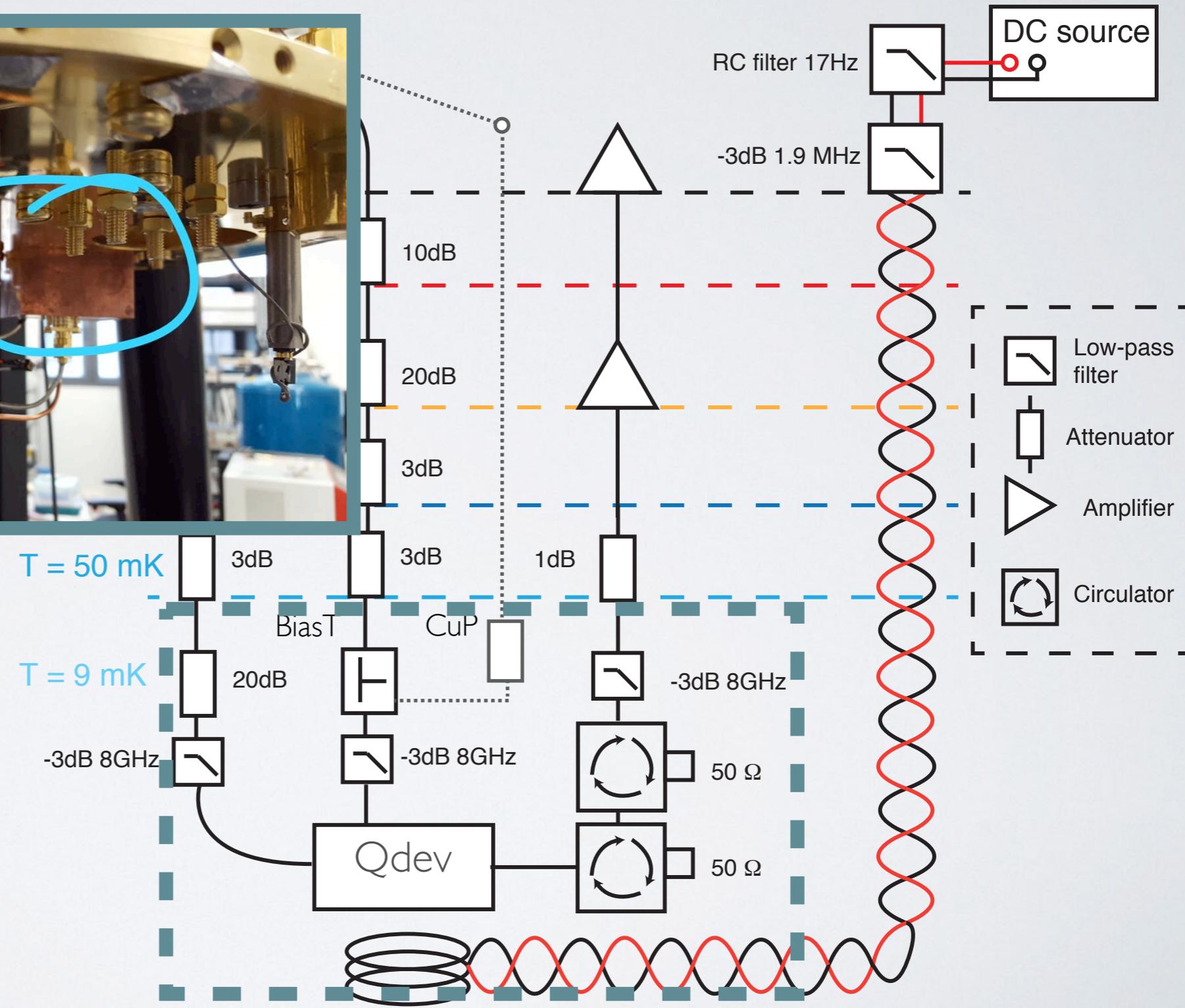
Coaxial wiring
dilution fridge



QUBIT STATE CONTROL

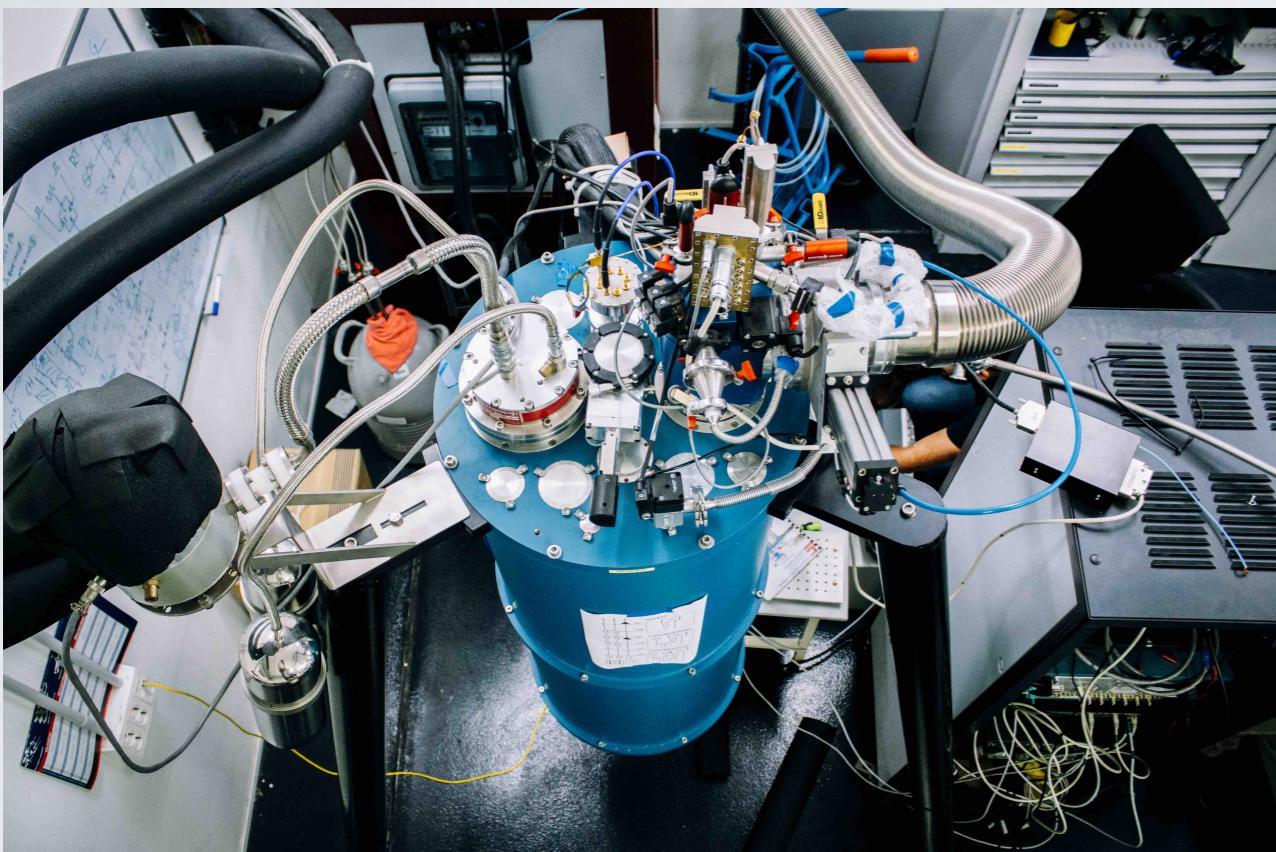


Superconducting
cavities hosting qubits



QUBIT STATE CONTROL

In a real experiment...



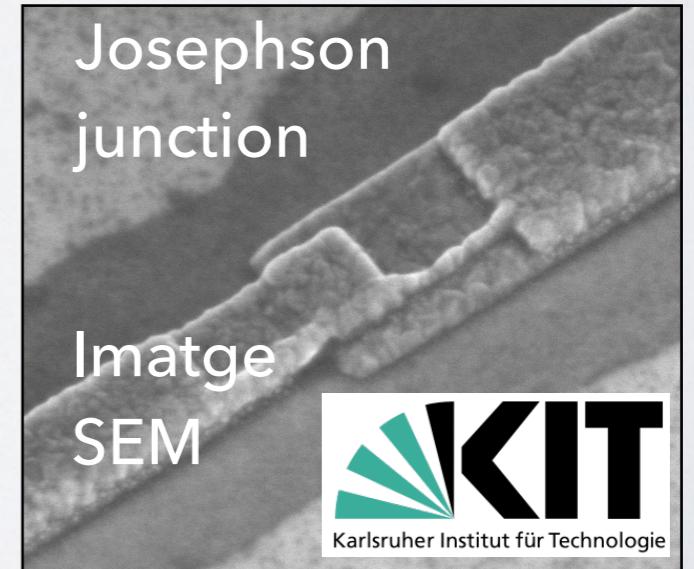
T = 0.01 K (-273.13 C)
Power consumption: 10kW

Currently at:



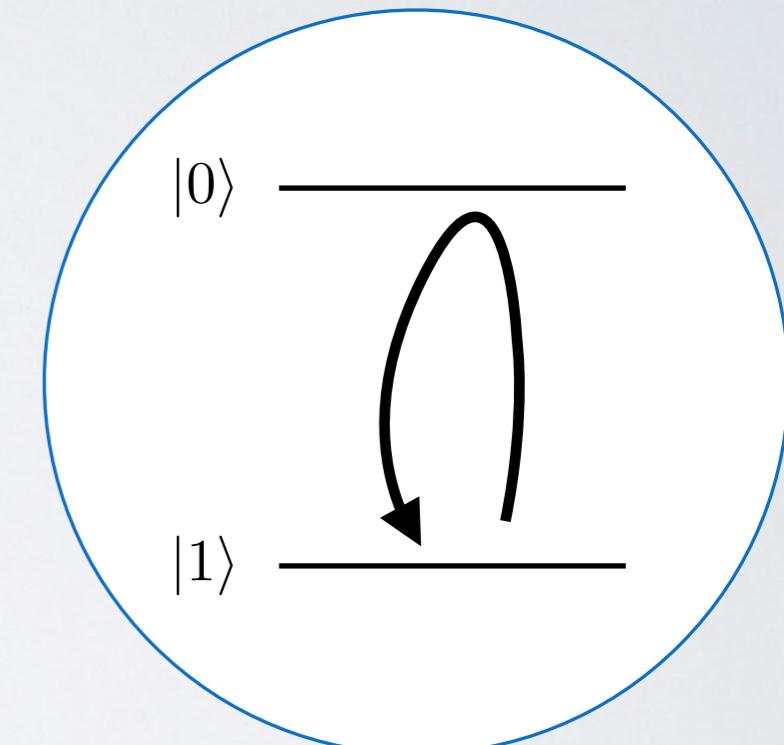
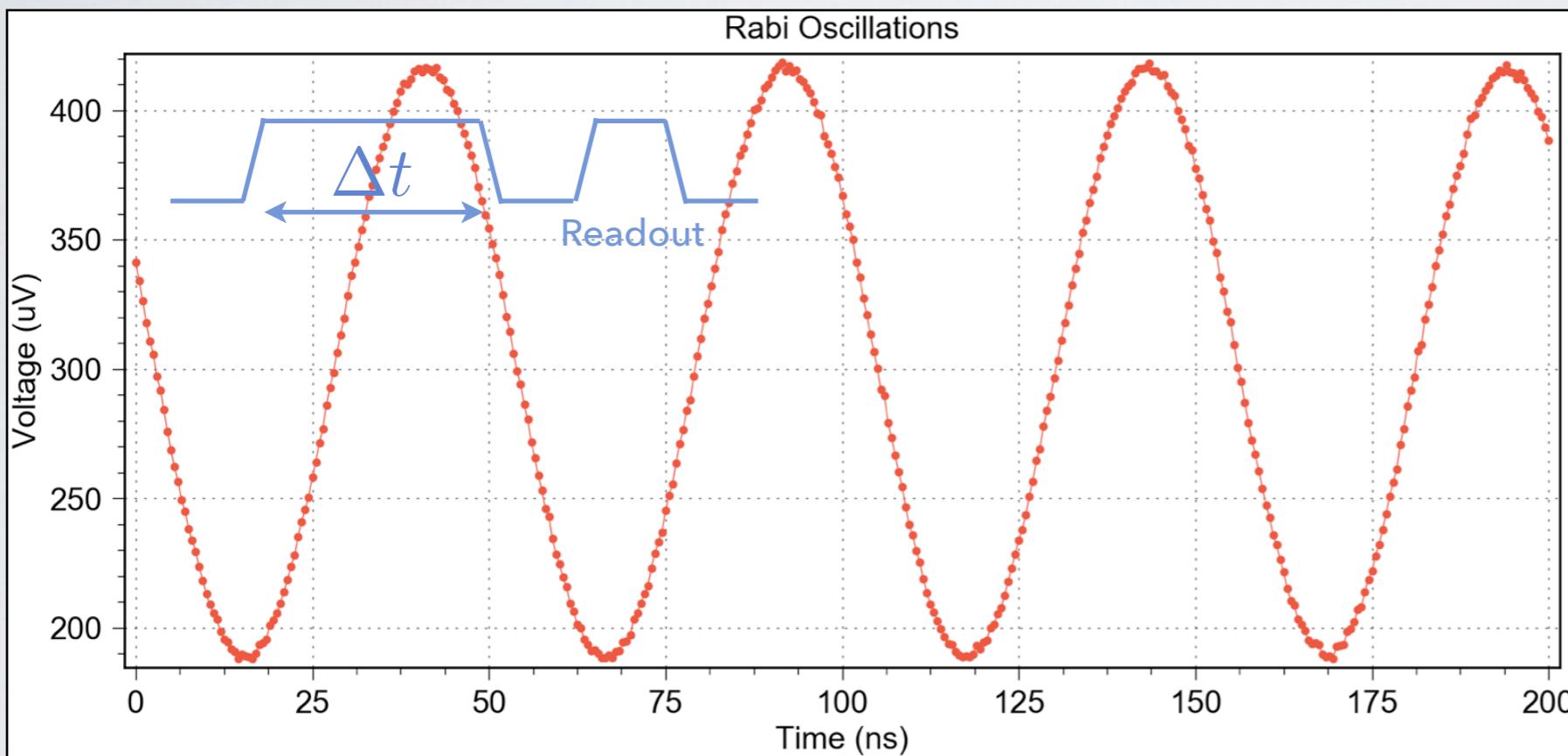
Josephson
junction

Image
SEM



QUBIT STATE CONTROL

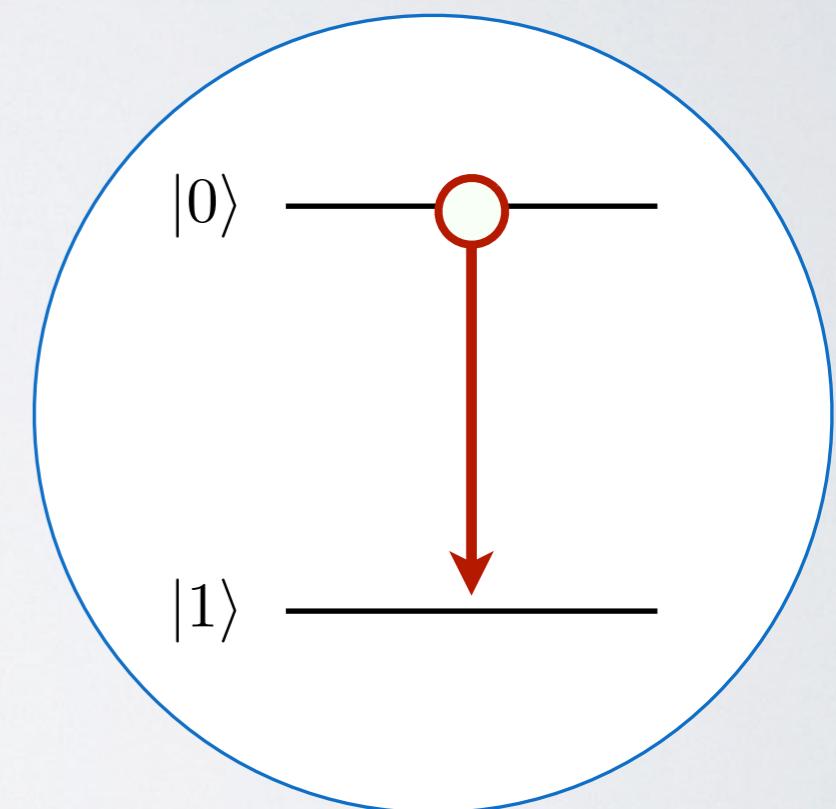
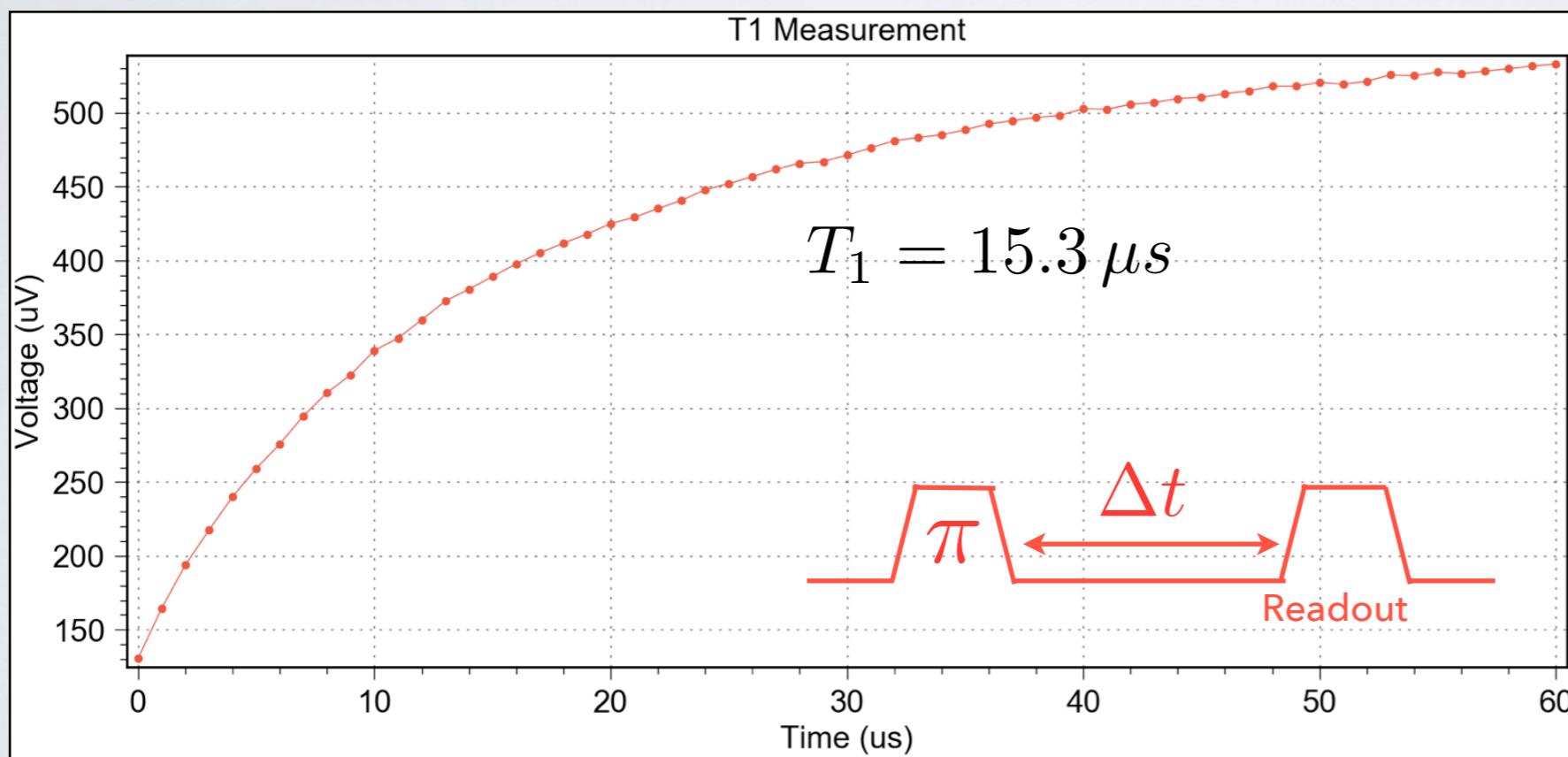
Driven, coherent Rabi oscillations of a superconducting qubit



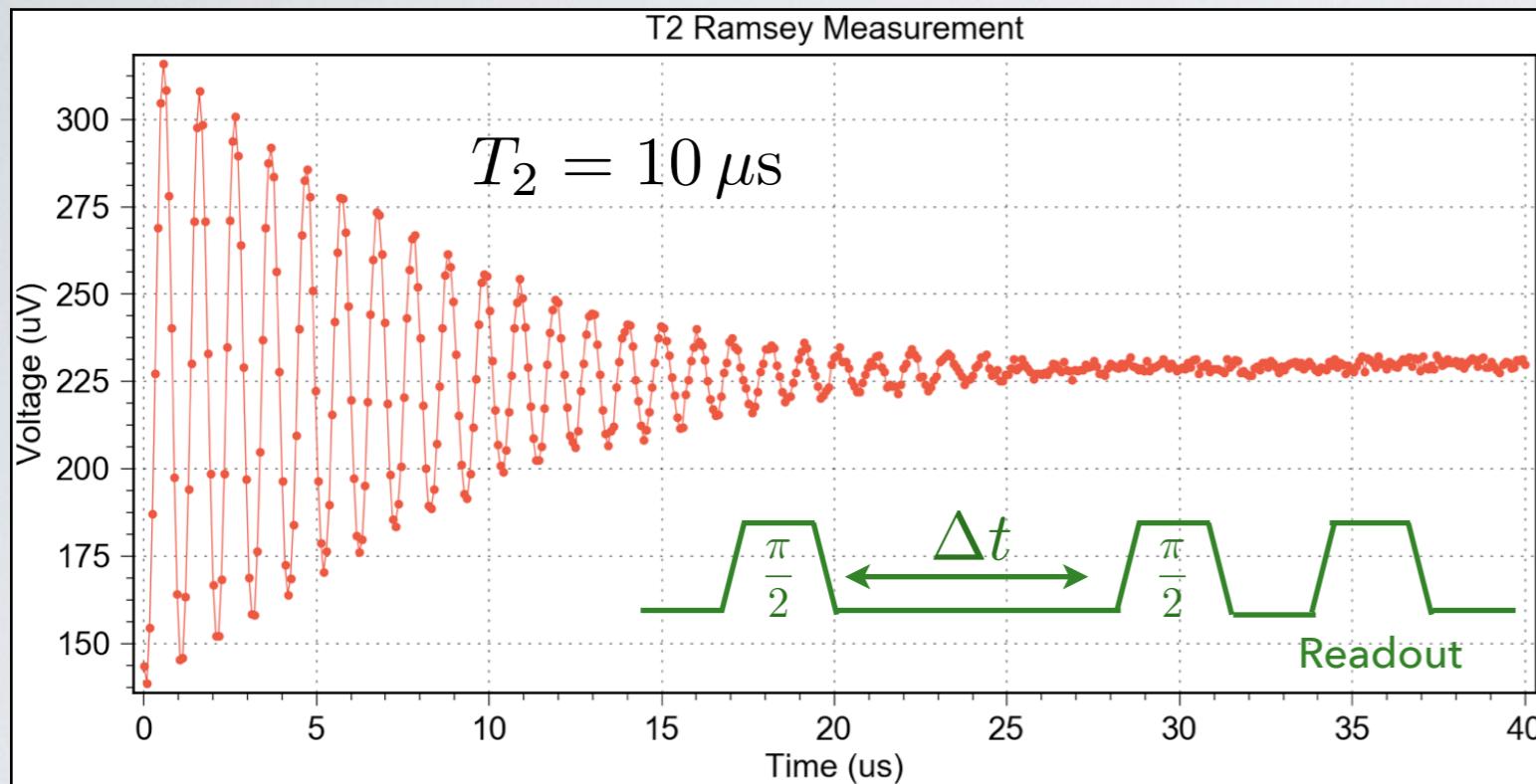
Spin-1/2-like control of qubits allows us to characterize qubit properties borrowing techniques from NMR

QUBIT STATE CONTROL

Energy decay of a superconducting qubit

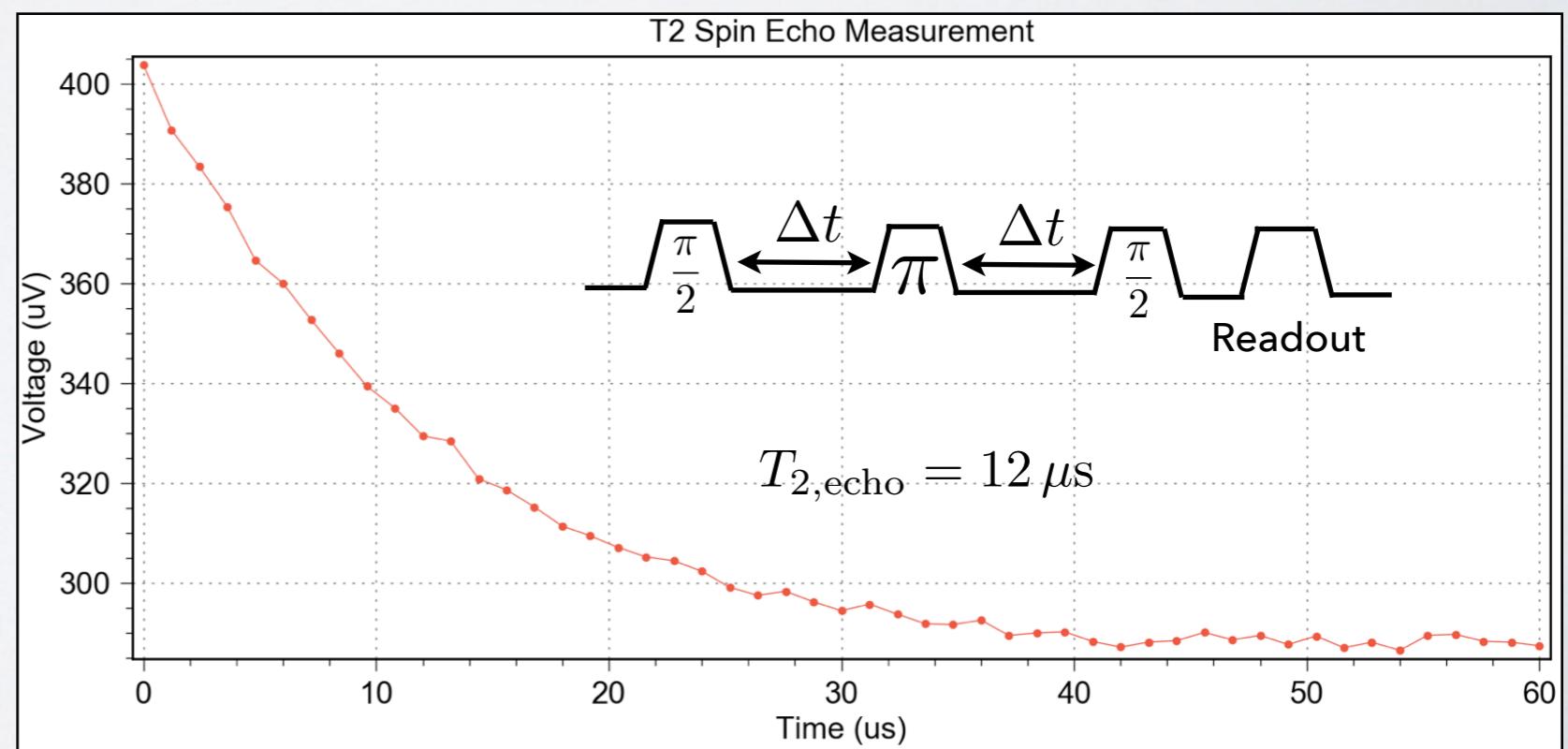


QUBIT STATE CONTROL



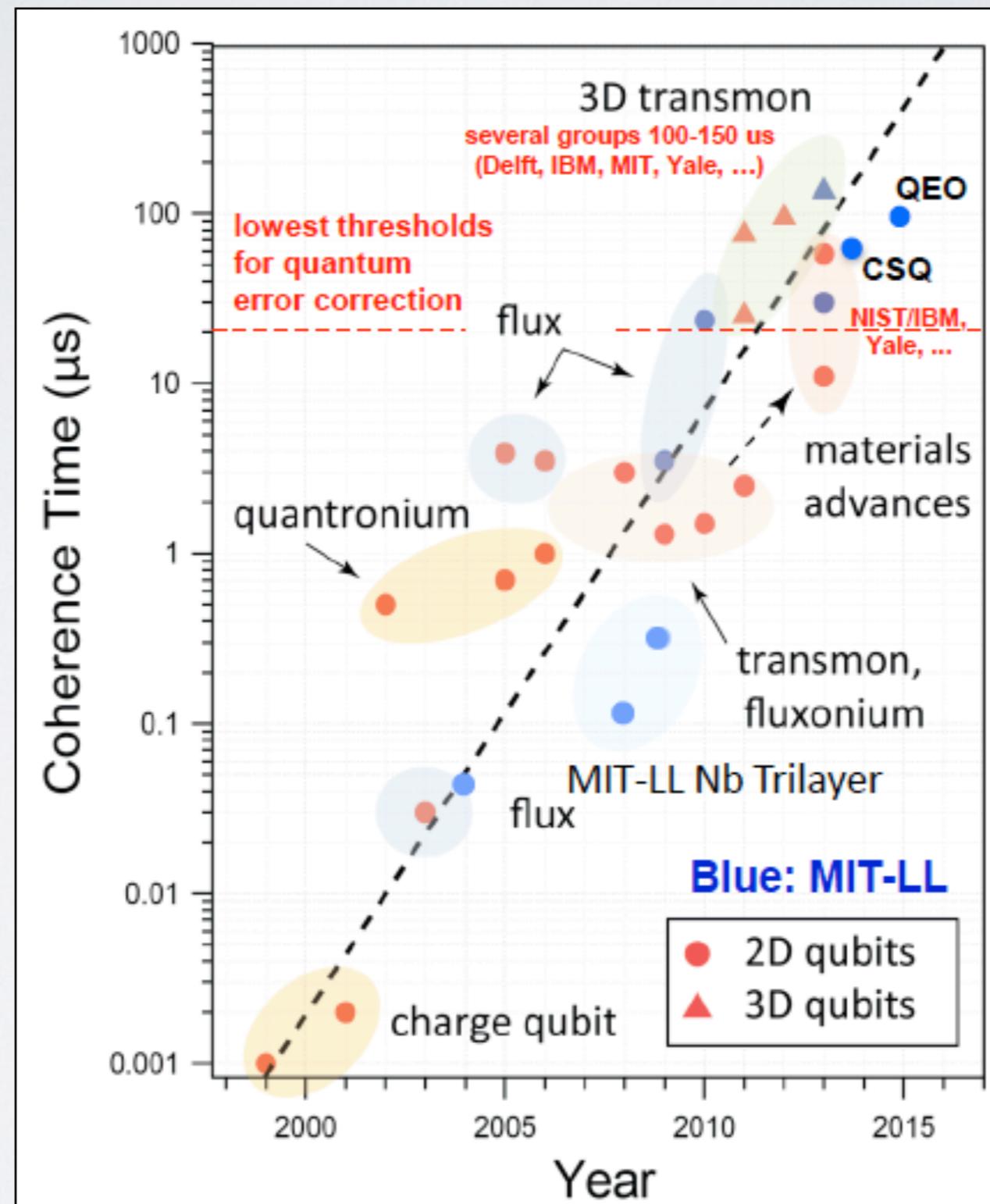
Decay of a superposition state:
Ramsey fringes

Dynamical correction:
Hahn spin-echo



QUBIT STATE CONTROL

A Moore's-like law



OUTLINE

Lecture I

- Quantum computation
- Circuit quantization
- Superconducting qubit zoo
- Qubit state control

End of lecture I!

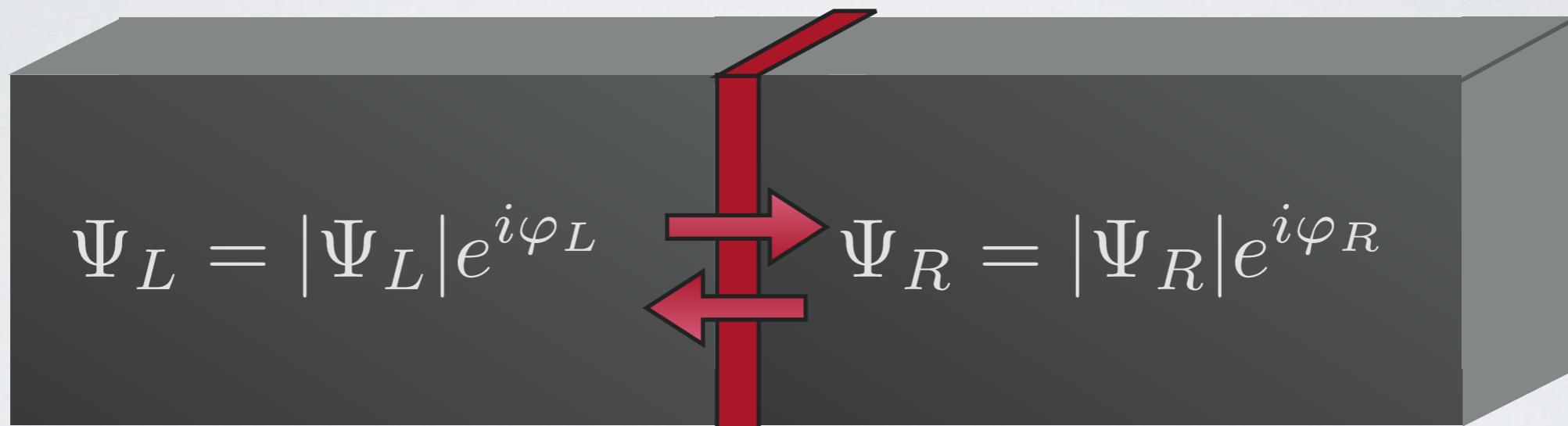
Lecture II

- Resonators for quantum computation
- Circuit quantum electrodynamics
- Qubit-qubit couplings and 2-qubit gates
- State of the art

BACKUP SLIDES

CIRCUIT QUANTIZATION

Superconductivity for qubits 101



First Josephson relation:

$$I_{L \rightarrow R} = I_C \sin(\varphi_L - \varphi_R)$$

Critical current

Second Josephson relation:

$$V = \frac{\Phi_0}{2\pi} \frac{d(\varphi_L - \varphi_R)}{dt}$$

Flux quantum:

$$\Phi_0 = \frac{h}{2e} \sim 2.07 \times 10^{-15} \text{ Wb}$$

Energy stored:

$$U = -E_J \cos(\varphi_L - \varphi_R)$$

Josephson energy:

$$E_J = I_C \Phi_0 / 2\pi$$