Flavour Symmetry Embedded - GLoBES (FaSE-GLoBES)

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I. INTRODUCTION

The discovery of neutrino oscillations points out the fact that neutrinos have mass, and provides evidence beyond the Standard Model (BSM). This phenomenon is successfully described by a theoretical framework with the help of three neutrino mixing angles (θ_{12} , θ_{13} , θ_{23}), two mass-square splittings (Δm_{21}^2 , Δm_{31}^2), and one Dirac CP phase (δ) [1–4]. Thanks to the great efforts in the past two decades, we almost have a complete understanding of such a neutrino oscillation framework. More data in the neutrino oscillation experiments is needed to determine the sign of Δm_{31}^2 , to measure the value of $\sin \theta_{23}$, to discover the potential CP violation in the leptonic sector and even to constrain the size of δ [4]. For these purposes, the on-going long baseline experiments (LBLs), such as the NuMI Off-axis ν_e Appearance experiment (NO ν A) [5] and the Tokai-to-Kamioka experiment (T2K) [6], can answer these questions with the statistical significance $\gtrsim 3\sigma$ in most of the parameter space. Based on the analysis with their data, the normal mass ordering ($\Delta m_{31}^2 > 0$), the higher θ_{23} octant ($\theta_{23} > 0$) 45°), and $\delta \sim 270^{\circ}$ are preferred so far [4]. The future LBLs, Deep Underground Neutrino Experiment (DUNE) [7], Tokai to Hyper-Kamiokande (T2HK) [8], and the medium baseline reactor experiment, the Jiangmen Underground Neutrino Observatory (JUNO) [9, 10] will further complete our knowledge of neutrino oscillations.

Flavour symmetry models are used to explain the origin of the neutrino mixing, and to predict the value of oscillation parameters (some of useful review articles are [11–17]). These models are motivated by some interesting features, such as $\theta_{12} \sim 33^{\circ}$, and $\theta_{23} \sim 45^{\circ}$. Before the discovery of non-zero θ_{13} measurement by Daya Bay experiment [18], the 'tri-bi-maximal' neutrino mixing (TBM) ansatz, which was proposed in 2002 by Horrison, Perkins, and Scott [19], fitted with the experimental data in a good agreement:

$$U_{\text{TBM}} = \begin{pmatrix} 2/\sqrt{6} & 1/\sqrt{3} & 0\\ -1/\sqrt{6} & 1/\sqrt{3} & 1/\sqrt{2}\\ 1/\sqrt{6} & -1/\sqrt{3} & 1/\sqrt{2} \end{pmatrix}.$$

With the fact that $\theta_{13} \approx 8^{\circ}$, several ways to obtain such non-zero value of θ_{13} are proposed. One of popular proposals is to correct the tri-bi-maximal neutrino mixing such that

$$\sin \theta_{12} = (1+s)/\sqrt{3}$$
, $\sin \theta_{13} = r/\sqrt{2}$, and $\sin \theta_{23} = (1+a)/\sqrt{2}$.

Currently, some works discuss on how the future experiments can be used for testing

these flavour symmetry models in the phenomenological point of view, e.g. Ref. [20–23]. A large number of these works are based on the c-library – General Long Baseline Experiment Simulator (GLoBES) [24, 25], which is a convenient simulation tool to simulate neutrino oscillation experiments via the Abstract Experiment Definition Language (AEDL). Some AEDL files for experiments are also available on GLoBES website, while the working group of DUNE experiment also releases their AEDL files [26]. GLoBES can be taken as one of popular and useful tools in the community of neutrino oscillation physics. However, it has not yet to be extended for the purpose of analysing flavour symmetry models. In this work, we will present our simulation tool Flavour Symmetry Embedded - GLoBES (FaSE-GloBES) that is a simulation code based on GLoBES library for studying the flavour symmetry. FaSE is a supplemental tool for GLoBES, written in c/c++ language, and FASE-GLoBES allows the user to assign the flavour symmetry model and analysis how a flavour symmetry model can be excluded by the simulated neutrino oscillation experiment.

II. OVERVIEW FASE-GLOBES

Flavor Symmetry Embedded (FaSE) is a supplemental tool for General Long Baseline Experiment Simulator (GLoBES) in order to analysis how a leptonic flavour symmetry model can be tested in neutrino oscillation experiments. FaSE is written in the c/c++ language, and consist with three codes FASE_GLoBES.c and model-input.c. The user defines the model in to model-input.c, while FASE_GLoBES.c does not need to be touched.

The concept of FaSE-GLoBES is shown in Fig. 1, in which three parts are shown: 1. the parameter translation (the blue box), 2. giving oscillation-parameter values (the green box), and 3. the χ^2 -value calculation (the orange box). The idea behind this flow chart Fig. 1 is that given a set of value for model parameter, the corresponding values for oscillation parameters are obtained by a translation, which is assigned by the user in model-input.c. And then, through FASE_GLoBES.c, these oscillation-parameter values are passed in to GLoBES library to simulate the event spectrum for evaluating the χ^2 value.

API functions to use **FaSE** are listed:

1. $MODEL_init(N_{para})$,

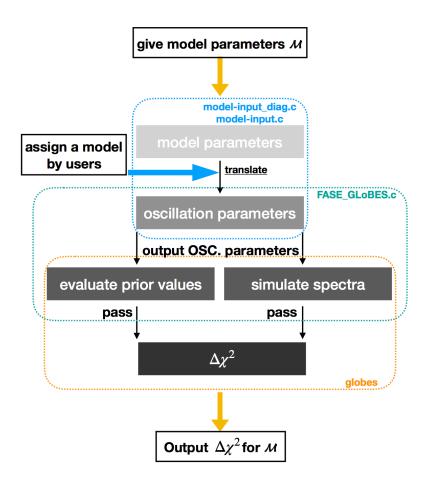


Figure 1: A scheme to correlate the model parameters with standard neutrino oscillation parameters. The error propagation is implemented in the simulation code up to the spectra analysis.

- 2. FASE_glb_probability_matrix,
- 3. FASE_glb_set_oscillation_parameters,
- 4. FASE_glb_get_oscillation_parameters,
- 5. FASE_prior_OSC,
- 6. FASE_prior_model.

The first one is to initialise **FaSE** with the parameter N_{para} the number of input parameters. The next three functions need to be included to replace the default GLoBES probability engine by the one that can read the output from **model-input.c**, as follows glbRegisterProbabilityEngine(6,

```
&FASE_glb_probability_matrix,
&FASE_glb_set_oscillation_parameters,
&FASE_glb_get_oscillation_parameters,
NULL);
```

This probability engine can work with the oscillation or model parameter. It can be set by the user with the parameter PARA. If PARA=STAN (PARA=MODEL) the probability engine works with oscillation (model) parameters. The final two functions are prior functions. Once the user gives the prior in oscillation (model) parameters, the user needs to call FASE_prior_OSC (FASE_prior_model) as follows.

```
glbRegisterPriorFunction(FASE_prior_OSC, NULL, NULL, NULL);
or
glbRegisterPriorFunction(FASE_prior_model, NULL, NULL, NULL);
```

We note that except for setting the probability engine and the prior function, the other parts in the main code should follow with the GLoBES manual.

III. MODEL SETTING

In the function MtoS in model-input.c, the user can assign the relation between the oscillation and model parameter sets, or define the mass matrix in model parameters, which will be diagonalised by the function ModelTO to obtain the corresponding oscillation-parameter values. The user needs give the input of MtoS – a set of model parameters ($\vec{\theta}_{Model}$), and the output is the corresponding oscillation parameter $\vec{\theta}_{OSC}$, of which components are θ_{12} , θ_{13} , θ_{23} , δ , Δm_{21}^2 , and Δm_{31}^2 . These values will be passed in to FaSE_GLoBES to simulate the experimental spectra and compute the prior value.

The relation between the oscillation and model parameter sets can be derived out by the user,

$$\vec{\theta}_{Model} = \vec{f}(\vec{\theta}_{OSC.}) \tag{1}$$

and give in the function MtoS. The oscillation parameters are based on

$$U^{\dagger} \mathcal{M} \mathcal{M}^{\dagger} U = \mathbf{M}^2$$
, where $\mathbf{M}_{\alpha\beta}^2 = m_{\alpha}^2 \delta_{\alpha\beta}$, (2)

where \mathcal{M} (M) is the mass matrix in the flavour (mass) state. The matrix \mathcal{M} is given by user with model parameters ($\vec{\theta}_{Model}$). The matrix U is the mixing matrix, and can be used

for getting mixing angles, and the diagonal elements of \mathbf{M} are for the corresponding mass-squared differences. The diagnolisation will be done in the function ModelTO, which needs to be called in MtoS. The output is the vector of oscillation parameters, of which components are θ_{12} , θ_{13} , θ_{23} , δ , Δm_{21}^2 , and Δm_{31}^2 . For the first four components, values are given in the unit of \mathbf{rad} , while the other two are in \mathbf{eV}^2 .

IV. PRIOR SETTING

Given a set of values for model parameters, FASE_GLoBES.c will obtain the corresponding oscillation-parameter values from model-input.c, and will pass these values to simulate the event spectrum and to give the prior value. Two gaussian prior functions are provided in FASE – FASE_prior_OSC and FASE_prior_model. These two functions are for different purposes. If the user give the prior in oscillation (model) parameters, the user should register FASE_prior_OSC (FASE_prior_model) for the prior with the GLoBES function glbRegisterPriorFunction, as we introduced in the beginning of this section. The Gaussian prior is

$$\chi_{prior}^2 = \sum_i \frac{(\theta_i - \theta_i^c)^2}{\sigma_i^2},\tag{3}$$

where θ_i is one of parameters constrained by prior, θ_i^c (σ_c) is the central value (Gaussian width) of the prior for θ_i . We note that θ_i can be either model ($\vec{\theta}_{Model}$) or oscillation parameters ($\vec{\theta}_{OSC.}$). But, oscillation and model parameters can be mixed up together in the using of prior.

The values of θ_i^c and σ_i need to be given by the user through three arrays: Central_prior, UPPER_prior, and LOWER_prior, in which there are six components. To treat asymmetry of width for the upper and lower Gaussian widths, we give values in two arrays UPPER_prior, and LOWER_prior, respectively. If the user gives the prior in model parameters, the order of each component follow with the setup of input of the probability engine. While the user gives the prior in oscillation parameters, the six components in order are θ_{12} , θ_{13} , θ_{23} , δ , Δm_{21}^2 , and Δm_{31}^2 . The first four parameter are in rad, and the final two are in eV².

Finally, some restrictions are imposed by the studied flavour symmetry model. We set up these restrictions in the function model_restriction, which is in model-input_diga.c and model-input.c. In the function model_restriction, the user needs to return 0 once the restriction is broken. For example, if the normal ordering is imposed, we give "if

(DMS31<0) { return 1;} "in model_restriction, where DMS31 is the variable for Δm_{31}^2 . Then, when the restriction is broken, model_restriction returns the value 1 to the prior function FASE_prior_OSC or FASE_prior_model. In the following, the prior function will give 10^6 for the χ^2 value, and it will be selected out when the user studies the statistically reasonable region in the parameter space.

V. CONSTRAINT OF MODEL PARAMETERS

The user can use **FaSE-GLoBES** to study the constraint of model parameters. Here we take tri-direct littlest seesaw as an example to present how we can adopt **FaSE-GLoBES** to constrain two model parameters – x and η . The χ^2 function is constructed based on a log-likelihood ratio,

$$\chi^{2}(\vec{\theta}, \xi_{s}, \xi_{b}) = 2 \sum_{i} \left(\eta_{i}(\vec{\theta}, \xi_{s}, \xi_{b}) - n_{i} + n_{i} \ln \frac{n_{i}}{\eta_{i}(\vec{\theta}, \xi_{s}, \xi_{b})} \right) + p(\xi_{s}, \sigma_{s}) + p(\xi_{b}, \sigma_{b}) + \chi^{2}_{prior},$$
 (4)

where i runs over the number of bins, $\eta_i \vec{\theta}, \xi_s, \xi_b$ is the hypothesis event rate for bin i and E_i is the central bin energy. The vector $\vec{\theta}$ consist test model or oscillation parameters. Here, $\vec{\theta} = (x, \eta, r, m_a)$. The parameters ξ_s and ξ_b are introduced to account for the systematic uncertainty of normalisation for the signal (subscript $_s$) and background (subscript $_b$) components for the event rate, and are allowed to bary in the fit as nuisance parameters. For a given hypothesized set of parameters $\vec{\theta}$, the event rate for bin i is calculated as

$$\eta_i(\vec{\theta}, \xi_s, \xi_b) = (1 + \xi_s) \times n_i + (1 + \xi_b) \times b_i, \tag{5}$$

where n_i and b_i are the expected number of signal and background events in bin i, respectively. The nuisance parameters are constrained by the Gaussian prior $p(\xi, \sigma) = \xi^2/\sigma^2$ with corresponding uncertainties σ_s and σ_b for the signal and background, respectively. Finally, χ^2_{prior} is a set of Gaussian priors for hypothesis, and is expressed as Eq. 3. To get the constraint of x and η , we obtain the minimum of χ^2 value,

glbSetProjection(free);

float res0=glbChiNP(true_values,NULL,GLB_ALL);

Then, we set two loops to get the $\Delta \chi^2$ value for different x and η

```
float x,eta,r,ma,dx,deta;
float lower_x,upper_x,lower_eta,upper_eta;
FILE* File=fopen("data/constraint_x_eta.dat", "w");
lower_x=-9; upper_x=-3; lower_eta=0.8*M_PI; upper_eta=2*M_PI;
dx=(upper_x-lower_x)/100; deta=(upper_eta-lower_eta)/100;
glbSetProjection(projection);
for (x=lower_x;x<=upper_x;x=x+dx){
  for (eta=lower_eta;eta<=upper_eta;eta=eta+deta){
    glbSetOscParams(test_values,x,0); glbSetOscParams(test_values,eta,1);
    float res=glbChiNP(test_values,NULL,GLB_ALL);
  fprintf(File,"%f %f %f \n",x,eta/M_PI,res-res0);
} fprintf(File,"\n");}</pre>
```

We use the same code to obtain the constraint of any other combination of two parameters. Then, we can have the result as in Fig. 2.

VI. MODEL TESTING

We can also study on excluding the model, assuming different true values for oscillation parameters. In this example, we present testing the tri-driect littlest seesaw model in various θ_{23} and δ . To do so, we set the true value in oscillation parameters. We change values of θ_{23} and δ of the true theory, and compute the minimal χ^2 value for the tested model with all four model parameters, free to be varied. And, the studied statistics function is exactly Eq. 4, but the true event rate n_i is predicted by a set of oscillation parameters, which will be varied in the code. And all four model parameter can be varied with the prior given by Eq. 3.

The example code is given below.

```
{ ... initalize the code ...}
{ ... register the probability engine ...}
float degree = M_PI/180;
float x_true,eta_true,r_true,ma_true;
```

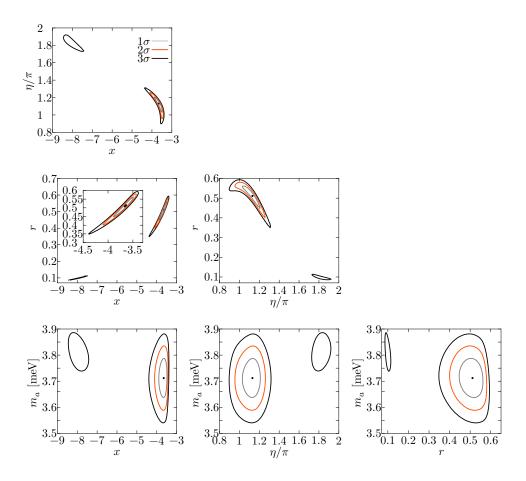


Figure 2: Precision measurements of any two model parameters at 3σ confidence level in the framework of three neutrino oscillations taking uncertainties of the current global fit results, for MOMENT, at 1σ (gray), 2σ (orange), 3σ (black). True values for the model parameters are used $(x, \eta, r, M_a) = (-3.65, 1.13\pi, 0.511, 3.71 \text{ meV})$.

```
x_true=-3.65029; eta_true=1.13067*M_PI; r_true=0.511325;
ma_true=3.71199e-3;
double M_para[4],OSC_para[6];
M_para[0]=x_true; M_para[1]=eta_true; M_para[2]=r_true; M_para[3]=ma_true;
MtoS(OSC_para, M_para);
double th12_true=OSC_para[0];
double th13_true=OSC_para[1];
double th23_true=OSC_para[2];
double dCP_true=OSC_para[3];
double DM21_true=OSC_para[4];
```

```
double DM31_true=OSC_para[5];
glb_params true_values = glbAllocParams();
glb_params test_values = glbAllocParams();
glb_params input_errors = glbAllocParams();
glb_params centers = glbAllocParams();
glbDefineParams(test_values,x_true,eta_true,r_true,ma_true,0,0);
glbSetDensityParams(test_values,1.0,GLB_ALL);
glbDefineParams(true_values,th12_true,th13_true,th23_true,dCP_true,DM21_true,DM31_true);
glbSetDensityParams(true_values,1.0,GLB_ALL);
{ ... set up for prior ...}
float th23,dCP,dth23,ddCP,lower_th23,upper_th23,lower_dCP,upper_dCP;
FILE* File=fopen("data/ModelTest_th23_dCP_test.dat", "w");
lower_th23=40; upper_th23=50; lower_dCP=125; upper_dCP=392;
dth23=(upper_th23-lower_th23)/100; ddCP=(upper_dCP-lower_dCP)/100;
glbSetProjection(free);
for (th23=lower_th23;th23<=upper_th23;th23=th23+dth23){
for (dCP=lower_dCP;dCP<=upper_dCP;dCP=dCP+ddCP){</pre>
glbSetOscParams(true_values,th23*degree,GLB_THETA_23);
glbSetOscParams(true_values,dCP*degree,GLB_DELTA_CP);
PARA=STAN;
glbSetOscillationParameters(true_values);
glbSetRates();
PARA=MODEL;
float res=glbChiNP(test_values,NULL,GLB_ALL);
fprintf(File, "%f %f %f \n", th23, dCP, gsl_cdf_chisq_Qinv(gsl_cdf_chisq_Q(fabs(res), dof), 1)
} fprintf(File,"\n");
}
{ ... destroy the pointer ...}
Finally, in the code we adopt Wilk's theorem [27]. When comparing nested models,
```

the $\Delta \chi^2$ test statistics is a random variable asymptotically distributed according to the

 χ^2 -distribution with the number of degrees of freedom, which is equal to the difference in the number of free model parameters dof. Here the number dof is 2. And that is our output: gsl_cdf_chisq_Qinv(gsl_cdf_chisq_Q(fabs(res),dof),1).

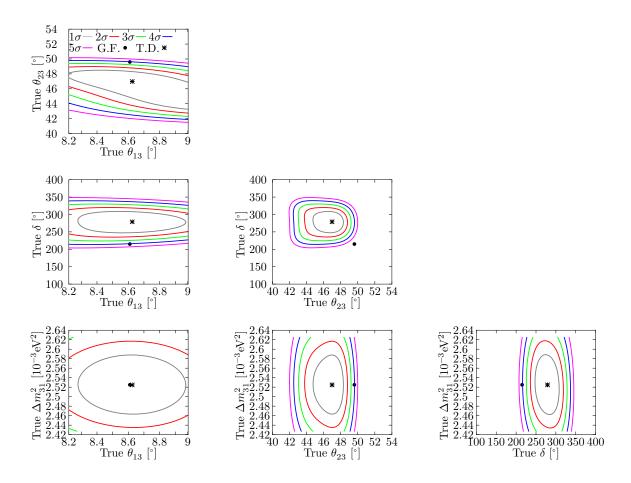


Figure 3: The 2-D exclusion contour for tri-direct littlest seesaw model on the plane of any two true standard parameters, from 1σ to 5σ . The range for each parameter is taken according to the 3σ uncertainty in NuFit4.0 results. The black dot denotes the best fit of NuFit4.0 results $((\theta_{12}, \theta_{13}, \theta_{23}, \delta, \Delta m_{21}^2, \Delta m_{31}^2) = (33.82^{\circ}, 8.61^{\circ}, 49.6^{\circ}, 215^{\circ}, 7.39 \times 10^{-5} \text{ eV}^2, 2.525 \times 10^{-3} \text{ eV}^2)),$ while the star is the prediction by the tri-direct littlest seesaw model with NuFit4.0 results $((\theta_{12}, \theta_{13}, \theta_{23}, \delta, \Delta m_{21}^2, \Delta m_{31}^2) \sim (36.25^{\circ}, 8.63^{\circ}, 47^{\circ}, 279^{\circ}, 7.39 \times 10^{-5} \text{ eV}^2, 2.525 \times 10^{-3} \text{ eV}^2)).$

We use the same code to analysis the exclusion ability for other oscillation-parameter combinations. Then, we have the results shown in Fig. 3.

VII. SUMMARY AND CONCLUSIONS

We have presented FaSE, which is a supplemental simulation tool for GLoBES to study the flavour symmetry with neutrino oscillation experiments. FaSE provides c-codes: model-input.c and FASE_GLoBES.c. Shown in Fig. 1, FASE_GLoBES, which calls functions in model-input.c, plays a role as a bridge between FaSE and GLoBES to simulate the expected spectra and compute the prior value. It can be left to be untouched by users. However, all inputs of the user needs to be given in the code model-input.c. Given a set of model parameters \vec{M} , with GLoBES, the output can be the χ^2 value for the hypothesis \vec{M} . Makefile is easy to include these two binary files (model-input and FASE GLoBES) in the makefile script for GLoBES.

We also present two examples for **FaSE-GLoBES** with the flavour symmetry model – tri-direct littlest seesaw (TDLS) – and the future neutrino oscillation experiment – MO-MENT. We show the model can be assign in two ways: the relation between model and oscillation parameters or the form of mass matrix in model parameters. The input of the oscillation probability and prior value can be in model or oscillation parameters. We further demonstrate how to use **FaSE-GLoBES** to obtain the constraint of any two of model parameters, and to study the ability to TDLS by MOMENT experiment.

Finally, **GLoBES** is a popular and powerful simulation tool to analysis the neutrino oscillation experiments in a simple language (AEDL), without losing too much detail. Considering the success of the flavour symmetry theory to explain the neutrino oscillations, **FaSE-GLoBES** should benefit model builders of leptonic flavour symmetry and phenomenologists for neutrino oscillation physics. We leave the flexibility for the user, and some other improvements and extensions might be done in the future.

Bibliography

- [1] B. Pontecorvo, Neutrino Experiments and the Problem of Conservation of Leptonic Charge, Sov. Phys. JETP 26 (1968) 984.
- [2] Z. Maki, M. Nakagawa and S. Sakata, Remarks on the unified model of elementary particles, Prog. Theor. Phys. 28 (1962) 870.

- [3] B. Pontecorvo, Inverse beta processes and nonconservation of lepton charge, Sov. Phys. JETP 7 (1958) 172.
- [4] I. Esteban, M. C. Gonzalez-Garcia, A. Hernandez-Cabezudo, M. Maltoni and T. Schwetz, Global analysis of three-flavour neutrino oscillations: synergies and tensions in the determination of θ₂3, δ_CP, and the mass ordering, JHEP 01 (2019) 106 [1811.05487].
- [5] NOvA collaboration, The NOvA Technical Design Report, .
- [6] T2K collaboration, The T2K Experiment, Nucl. Instrum. Meth. A659 (2011) 106[1106.1238].
- [7] DUNE collaboration, Long-Baseline Neutrino Facility (LBNF) and Deep Underground Neutrino Experiment (DUNE), 1512.06148.
- [8] HYPER-KAMIOKANDE WORKING GROUP collaboration, A Long Baseline Neutrino Oscillation Experiment Using J-PARC Neutrino Beam and Hyper-Kamiokande, 2014, 1412.4673.
- [9] JUNO collaboration, JUNO Conceptual Design Report, 1508.07166.
- [10] JUNO collaboration, Neutrino Physics with JUNO, J. Phys. G43 (2016) 030401[1507.05613].
- [11] G. Altarelli and F. Feruglio, Discrete Flavor Symmetries and Models of Neutrino Mixing, Rev. Mod. Phys. 82 (2010) 2701 [1002.0211].
- [12] H. Ishimori, T. Kobayashi, H. Ohki, Y. Shimizu, H. Okada and M. Tanimoto, Non-Abelian Discrete Symmetries in Particle Physics, Prog. Theor. Phys. Suppl. 183 (2010) 1 [1003.3552].
- [13] S. F. King and C. Luhn, Neutrino Mass and Mixing with Discrete Symmetry, Rept. Prog. Phys. 76 (2013) 056201 [1301.1340].
- [14] S. F. King, A. Merle, S. Morisi, Y. Shimizu and M. Tanimoto, Neutrino Mass and Mixing: from Theory to Experiment, New J. Phys. 16 (2014) 045018 [1402.4271].
- [15] S. F. King, Models of Neutrino Mass, Mixing and CP Violation, J. Phys. G42 (2015) 123001
 [1510.02091].
- [16] S. F. King, Neutrino Mixing: from experiment to theory, Nucl. Part. Phys. Proc. 265-266 (2015) 288.
- [17] S. F. King, Unified Models of Neutrinos, Flavour and CP Violation, Prog. Part. Nucl. Phys. 94 (2017) 217 [1701.04413].
- [18] Daya Bay collaboration, Spectral measurement of electron antineutrino oscillation amplitude

- and frequency at Daya Bay, Phys. Rev. Lett. 112 (2014) 061801 [1310.6732].
- [19] P. F. Harrison, D. H. Perkins and W. G. Scott, Tri-bimaximal mixing and the neutrino oscillation data, Phys. Lett. B530 (2002) 167 [hep-ph/0202074].
- [20] P. Ballett, S. F. King, S. Pascoli, N. W. Prouse and T. Wang, Precision neutrino experiments vs the Littlest Seesaw, JHEP 03 (2017) 110 [1612.01999].
- [21] S. S. Chatterjee, P. Pasquini and J. Valle, Probing atmospheric mixing and leptonic CP violation in current and future long baseline oscillation experiments, Phys. Lett. B 771 (2017) 524 [1702.03160].
- [22] G.-J. Ding, Y.-F. Li, J. Tang and T.-C. Wang, Confronting Tri-direct CP-symmetry models to neutrino oscillation experiments, Phys. Rev. D100 (2019) 055022 [1905.12939].
- [23] J. Tang and T. Wang, Study of a tri-direct littlest seesaw model at MOMENT, Nucl. Phys. B 952 (2020) 114915 [1907.01371].
- [24] P. Huber, M. Lindner and W. Winter, Simulation of long-baseline neutrino oscillation experiments with GLoBES (General Long Baseline Experiment Simulator), Comput. Phys. Commun. 167 (2005) 195 [hep-ph/0407333].
- [25] P. Huber, J. Kopp, M. Lindner, M. Rolinec and W. Winter, New features in the simulation of neutrino oscillation experiments with GLoBES 3.0: General Long Baseline Experiment Simulator, Comput. Phys. Commun. 177 (2007) 432 [hep-ph/0701187].
- [26] DUNE collaboration, Experiment Simulation Configurations Used in DUNE CDR, 1606.09550.
- [27] S. S. Wilks, The Large-Sample Distribution of the Likelihood Ratio for Testing Composite Hypotheses, Annals Math. Statist. 9 (1938) 60.