

This notebook derives some manipulations/special cases of formulae from O'Toole & Woodhouse (2011) "Numerically stable computation of complete synthetic seismograms including the static displacement in plane layered media".

```
In[1]:= $Assumptions = v_p > 0 && v_p ∈ Reals && v_s > 0 &&
v_s ∈ Reals && ρ > 0 && ρ ∈ Reals && k ∈ Reals && k ≥ 0 && h ∈ Reals && h > 0
Out[1]= v_p > 0 && v_p ∈ ℝ && v_s > 0 && v_s ∈ ℝ && ρ > 0 && ρ ∈ ℝ && k ∈ ℝ && k ≥ 0 && h ∈ ℝ && h > 0
```

In[2]:=

Equation 14 gives:

```
In[3]:= A := {{0, k λ σ, 1 σ, 0}, {-k, 0, 0, 1 μ}, {-ρ ω², 0, 0, k}, {0, k² (γ + μ) - ρ ω², -k λ σ, 0}}
```

Equation 55 defines Z:

```
In[4]:= Z := {{1 √ρ, 0, 0, 0}, {0, 0, 0, -1 √ρ}, {0, 0, √ρ, -2 μ k √ρ}, {2 μ k √ρ, √ρ, 0, 0}}
```

and we can verify the stated form of the inverse matrix:

```
In[5]:= Inverse[Z] // MatrixForm
```

```
Out[5]/MatrixForm=
{{√ρ, 0, 0, 0},
 {-2 k μ √ρ, 0, 0, 1 √ρ},
 {0, -2 k μ √ρ, 1 √ρ, 0},
 {0, -√ρ, 0, 0}}
```

We can then form A' and verify that it is as given in eq.57:

```
In[6]:= Ap := FullSimplify[Inverse[Z].A.Z, {σ == λ + 2 μ, γ == μ 3 λ + 2 μ / λ + 2 μ}]
```

```
In[7]:= Ap // MatrixForm
```

```
Out[7]/MatrixForm=
{{0, 0, ρ σ, -k},
 {0, 0, -k, ω²},
 {-ω², -k, 0, 0},
 {-k, -ρ μ, 0, 0}}
```

Substituting in for the Lamé parameters we find

```
In[8]:= Ap /. σ → ρ v_p² /. μ → ρ v_s²
```

```
Out[8]= {{0, 0, 1 v_p², -k}, {0, 0, -k, ω²}, {-ω², -k, 0, 0}, {-k, -1 v_s², 0, 0}}
```

Converting to use ζ_p , ζ_s we have

```
In[9]:= Ap := {{0, 0, k² - ξ_p² ω², -k}, {0, 0, -k, ω²}, {-ω², -k, 0, 0}, {-k, ξ_s² - k² ω², 0, 0}}
```

which has the eigendecomposition as given in eqs. 59 & 60 (up to normalisation differences).

In[10]:= **Eigenvalues**[Ap] // FullSimplify

$$\text{Out}[10]= \{-\xi_p, \xi_p, -\xi_s, \xi_s\}$$

In[11]:= **Eigenvectors**[Ap] // FullSimplify

$$\text{Out}[11]= \left\{ \left\{ \frac{\xi_p}{k}, 0, \frac{\omega^2}{k}, 1 \right\}, \left\{ -\frac{\xi_p}{k}, 0, \frac{\omega^2}{k}, 1 \right\}, \left\{ \frac{k}{\xi_s}, -\frac{\omega^2}{\xi_s}, 0, 1 \right\}, \left\{ -\frac{k}{\xi_s}, \frac{\omega^2}{\xi_s}, 0, 1 \right\} \right\}$$

We can then take the matrix exponential and obtain eq. 62

In[12]:= **MatrixExp**[h Ap] // FullSimplify // MatrixForm

Out[12]//MatrixForm=

$$\begin{pmatrix} \cosh[h\xi_p] & \frac{k(\cosh[h\xi_p]-\cosh[h\xi_s])}{\omega^2} & \frac{-\sinh[h\xi_p]\xi_p + \frac{k^2 \sinh[h\xi_s]}{\xi_s}}{\omega^2} & -\frac{k \sinh[h\xi_s]}{\xi_s} \\ 0 & \cosh[h\xi_s] & -\frac{k \sinh[h\xi_s]}{\xi_s} & \frac{\omega^2 \sinh[h\xi_s]}{\xi_s} \\ -\frac{\omega^2 \sinh[h\xi_p]}{\xi_p} & -\frac{k \sinh[h\xi_p]}{\xi_p} & \cosh[h\xi_p] & 0 \\ -\frac{k \sinh[h\xi_p]}{\xi_p} & -\frac{k^2 \sinh[h\xi_p]}{\xi_p} + \sinh[h\xi_s] \xi_s & \frac{k(\cosh[h\xi_p]-\cosh[h\xi_s])}{\omega^2} & \cosh[h\xi_s] \end{pmatrix}$$

and taking the second-order minors yields eq. 66

In[13]:= **Minors**[**MatrixExp**[h Ap], 2] // FullSimplify // MatrixForm

Out[13]//MatrixForm=

$$\begin{pmatrix} \cosh[h\xi_p] \cosh[h\xi_s] & -\frac{k \cosh[h\xi_p] \sinh[h\xi_s]}{\xi_s} & \frac{\omega^2 \cosh[h\xi_p]}{\xi} \\ -\frac{k \cosh[h\xi_s] \sinh[h\xi_p]}{\xi_p} & 1 + \frac{k^2 \sinh[h\xi_p] \sinh[h\xi_s]}{\xi_p \xi_s} & -\frac{k \omega^2 \sinh[h\xi_s]}{\xi} \\ -\frac{k^2 \cosh[h\xi_s] \sinh[h\xi_p]}{\xi_p} + \cosh[h\xi_p] \sinh[h\xi_s] \xi_s & \frac{k \left(1 - \cosh[h\xi_p] \cosh[h\xi_s] + \frac{k^2 \sinh[h\xi_p] \sinh[h\xi_s]}{\xi_p \xi_s}\right)}{\omega^2} & \cosh[h\xi_p] \cosh[h\xi_s] \\ \frac{\omega^2 \cosh[h\xi_s] \sinh[h\xi_p]}{\xi_p} & -\frac{k \omega^2 \sinh[h\xi_p] \sinh[h\xi_s]}{\xi_p \xi_s} & \frac{\omega^4 \sinh[h\xi_p]}{\xi_p} \\ \frac{k \cosh[h\xi_s] \sinh[h\xi_p]}{\xi_p} & -\frac{k^2 \sinh[h\xi_p] \sinh[h\xi_s]}{\xi_p \xi_s} & \frac{k \omega^2 \sinh[h\xi_s]}{\xi_p} \\ -\frac{\sinh[h\xi_p] \sinh[h\xi_s] \xi_s}{\xi_p} & \frac{k \cosh[h\xi_s] \sinh[h\xi_p]}{\xi_p} & -\frac{\omega^2 \cosh[h\xi_s]}{\xi} \end{pmatrix}$$

A naïve implementation of the above expression is not numerically stable as $\omega \rightarrow 0$. To obtain a version useful in this case we begin by finding the zero-th order term in the Taylor expansion:

In[14]:= **ZeroFreq** =

$$\begin{aligned} & \text{Series}\left[\text{Minors}[\text{MatrixExp}[h A p], 2] /. \xi_p \rightarrow \sqrt{k^2 - \frac{\omega^2}{v_p^2}} / . \xi_s \rightarrow \sqrt{k^2 - \frac{\omega^2}{v_s^2}}, \{\omega, \theta, \theta\}\right] // \\ & \text{Normal} \\ \text{Out}[14]= & \left\{ \left\{ \frac{1}{4} e^{-2 h k} (1 + e^{2 h k})^2, -\frac{1}{4} e^{-2 h k} (-1 + e^{2 h k}) (1 + e^{2 h k}), \right. \right. \\ & \theta, -\frac{e^{-2 h k} (-v_p^2 + e^{4 h k} v_p^2 - 4 e^{2 h k} h k v_p^2 - v_s^2 + e^{4 h k} v_s^2 + 4 e^{2 h k} h k v_s^2)}{8 k v_p^2 v_s^2}, \\ & \left. \left. \frac{1}{4} e^{-2 h k} (-1 + e^{4 h k}), -\frac{1}{4} e^{-2 h k} (-1 + e^{2 h k})^2 \right\}, \right. \\ & \left\{ -\frac{1}{4} e^{-2 h k} (-1 + e^{4 h k}), \frac{1}{4} e^{-2 h k} (1 + e^{2 h k})^2, \theta, \frac{e^{-2 h k} (-1 + e^{2 h k})^2 (v_p^2 + v_s^2)}{8 k v_p^2 v_s^2}, \right. \\ & -\frac{1}{4} e^{-2 h k} (-1 + e^{2 h k})^2, \frac{1}{4} e^{-2 h k} (-1 + e^{2 h k}) (1 + e^{2 h k}) \}, \\ & \left\{ -\frac{e^{-2 h k} (-v_p^2 + e^{4 h k} v_p^2 + 4 e^{2 h k} h k v_p^2 - v_s^2 + e^{4 h k} v_s^2 - 4 e^{2 h k} h k v_s^2)}{8 k v_p^2 v_s^2}, \right. \\ & \frac{e^{-2 h k} (-1 + e^{2 h k})^2 (v_p^2 + v_s^2)}{8 k v_p^2 v_s^2}, 1, \frac{1}{16 k^2 v_p^4 v_s^4} \\ & e^{-2 h k} (v_p^4 - 2 e^{2 h k} v_p^4 + e^{4 h k} v_p^4 - 4 e^{2 h k} h^2 k^2 v_p^4 + 2 v_p^2 v_s^2 - 4 e^{2 h k} v_p^2 v_s^2 + 2 e^{4 h k} v_p^2 v_s^2 + 8 e^{2 h k} \\ & h^2 k^2 v_p^2 v_s^2 + v_s^4 - 2 e^{2 h k} v_s^4 + e^{4 h k} v_s^4 - 4 e^{2 h k} h^2 k^2 v_s^4), -\frac{e^{-2 h k} (-1 + e^{2 h k})^2 (v_p^2 + v_s^2)}{8 k v_p^2 v_s^2}, \\ & \left. \frac{e^{-2 h k} (-v_p^2 + e^{4 h k} v_p^2 - 4 e^{2 h k} h k v_p^2 - v_s^2 + e^{4 h k} v_s^2 + 4 e^{2 h k} h k v_s^2)}{8 k v_p^2 v_s^2} \right\}, \\ & \{0, 0, 0, 1, 0, 0\}, \left\{ \frac{1}{4} e^{-2 h k} (-1 + e^{2 h k}) (1 + e^{2 h k}), -\frac{1}{4} e^{-2 h k} (-1 + e^{2 h k})^2, \theta, \right. \\ & -\frac{e^{-2 h k} (-1 + e^{2 h k})^2 (v_p^2 + v_s^2)}{8 k v_p^2 v_s^2}, \frac{1}{4} e^{-2 h k} (1 + e^{2 h k})^2, -\frac{1}{4} e^{-2 h k} (-1 + e^{4 h k}) \}, \\ & \left\{ -\frac{1}{4} e^{-2 h k} (-1 + e^{2 h k})^2, \frac{1}{4} e^{-2 h k} (-1 + e^{4 h k}), \theta, \right. \\ & \left. \frac{e^{-2 h k} (-v_p^2 + e^{4 h k} v_p^2 + 4 e^{2 h k} h k v_p^2 - v_s^2 + e^{4 h k} v_s^2 - 4 e^{2 h k} h k v_s^2)}{8 k v_p^2 v_s^2}, \right. \\ & \left. -\frac{1}{4} e^{-2 h k} (-1 + e^{2 h k}) (1 + e^{2 h k}), \frac{1}{4} e^{-2 h k} (1 + e^{2 h k})^2 \right\} \} \end{aligned}$$

We can rewrite this substituting c,s for the hyperbolic functions as in the paper - some effort seems to be required to force Mathematica to do this completely.

```
In[15]:= ZeroFreqCS = Simplify [
  Simplify[FullSimplify[ZeroFreq /. vp → Sqrt[σ/ρ] /. vs → Sqrt[μ/ρ], {Cosh[h k]^2 == cc, Cosh[h k] == c, Sinh[h k]^2 == ss, Sinh[h k] == s, Cosh[h k] Sinh[h k] == cs}] // TrigExpand, {Cosh[h k]^2 == cc, Cosh[h k] == c, Sinh[h k]^2 == ss, Sinh[h k] == s, Cosh[h k] Sinh[h k] == cs}, Trig → False] /.
  cc → c^2 /. ss → s^2 /. cs → c s, c^2 + s^2 == 1 + 2 s^2]

Out[15]= {{c^2, -c s, 0, -ρ (h k (μ - σ) + c s (μ + σ)) / (2 k μ σ), c s, -s^2}, {-c s, c^2, 0, s^2 ρ (μ + σ) / (2 k μ σ), -s^2, c s}, {-ρ (h k (-μ + σ) + c s (μ + σ)) / (2 k μ σ), s^2 ρ (μ + σ) / (2 k μ σ), 1, -ρ^2 (h^2 k^2 (μ - σ)^2 - s^2 (μ + σ)^2) / (4 k^2 μ^2 σ^2), -s^2 ρ (μ + σ) / (2 k μ σ), ρ (h k (μ - σ) + c s (μ + σ)) / (2 k μ σ)}, {0, 0, 0, 1, 0, 0}, {c s, -s^2, 0, -s^2 ρ (μ + σ) / (2 k μ σ), c^2, -c s}, {-s^2, c s, 0, ρ (h k (-μ + σ) + c s (μ + σ)) / (2 k μ σ), -c s, c^2}}
```

Most terms in this expression feature a product of two hyperbolic functions, but not all do. We want to separate the two:

```
In[16]:= ZFCS0 = ZeroFreqCS /. c → 0 /. s → 0 // FullSimplify
Out[16]= {{0, 0, 0, h ρ (-μ + σ) / (2 μ σ), 0, 0}, {0, 0, 0, 0, 0, 0}, {h ρ (μ - σ) / (2 μ σ), 0, 1, -h^2 ρ^2 (μ - σ)^2 / (4 μ^2 σ^2), 0, h ρ (μ - σ) / (2 μ σ)}, {0, 0, 0, 1, 0, 0}, {0, 0, 0, 0, 0, 0}, {0, 0, 0, h ρ (-μ + σ) / (2 μ σ), 0, 0}}
```

```
In[17]:= ZFCS1 = (ZeroFreqCS - ZFCS0) // FullSimplify
Out[17]= {{c^2, -c s, 0, -c s ρ (μ + σ) / (2 k μ σ), c s, -s^2}, {-c s, c^2, 0, s^2 ρ (μ + σ) / (2 k μ σ), -s^2, c s}, {-c s ρ (μ + σ) / (2 k μ σ), s^2 ρ (μ + σ) / (2 k μ σ), 0, s^2 ρ^2 (μ + σ)^2 / (4 k^2 μ^2 σ^2), -s^2 ρ (μ + σ) / (2 k μ σ), c s ρ (μ + σ) / (2 k μ σ)}, {0, 0, 0, 0, 0, 0}, {c s, -s^2, 0, -s^2 ρ (μ + σ) / (2 k μ σ), c^2, -c s}, {-s^2, c s, 0, c s ρ (μ + σ) / (2 k μ σ), -c s, c^2}}
```

In[18]:=

At this point we have $\text{ZeroFreqCS} = \text{ZFCS0} + \text{ZFCS1}$. We do some rescaling of rows and columns to

get things into a nice format:

In[19]:=

```
SL = DiagonalMatrix[{1, 1, 2 k μ σ, 1, 1, 1}]
```

```
Out[19]= {{1, 0, 0, 0, 0, 0}, {0, 1, 0, 0, 0, 0}, {0, 0, 2 k μ σ, 0, 0, 0}, {0, 0, 0, 1, 0, 0}, {0, 0, 0, 0, 1, 0}, {0, 0, 0, 0, 0, 1}}
```

In[20]:= SR = DiagonalMatrix[{1, 1, 1, 2 k μ σ, 1, 1}]

```
Out[20]= {{1, 0, 0, 0, 0, 0}, {0, 1, 0, 0, 0, 0}, {0, 0, 1, 0, 0, 0}, {0, 0, 0, 2 k μ σ, 0, 0}, {0, 0, 0, 0, 1, 0}, {0, 0, 0, 0, 0, 1}}
```

In[21]:= SL.ZFCS0.SR // MatrixForm

Out[21]//MatrixForm=

$$\begin{pmatrix} 0 & 0 & 0 & h k \rho (-\mu + \sigma) & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 \\ h k \rho (\mu - \sigma) & 0 & 2 k \mu \sigma & -h^2 k^2 \rho^2 (\mu - \sigma)^2 & 0 & h k \rho (\mu - \sigma) \\ 0 & 0 & 0 & 2 k \mu \sigma & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & h k \rho (-\mu + \sigma) & 0 & 0 \end{pmatrix}$$

In[22]:= SL.ZFCS1.SR // MatrixForm

Out[22]//MatrixForm=

$$\begin{pmatrix} c^2 & -c s & 0 & -c s \rho (\mu + \sigma) & c s & -s^2 \\ -c s & c^2 & 0 & s^2 \rho (\mu + \sigma) & -s^2 & c s \\ -c s \rho (\mu + \sigma) & s^2 \rho (\mu + \sigma) & 0 & s^2 \rho^2 (\mu + \sigma)^2 & -s^2 \rho (\mu + \sigma) & c s \rho (\mu + \sigma) \\ 0 & 0 & 0 & 0 & 0 & 0 \\ c s & -s^2 & 0 & -s^2 \rho (\mu + \sigma) & c^2 & -c s \\ -s^2 & c s & 0 & c s \rho (\mu + \sigma) & -c s & c^2 \end{pmatrix}$$

In[23]:=

The above two matrices represent S0L(exp(h A'))S0R. We want to form Z(exp(hA'))Z^{-1}. We therefore need to roll the inverse scalings into our definitions of Z,Z^{-1}.

In[24]:= Minors[Z, 2].Inverse[SL] // FullSimplify

```
Out[24]= {{0, 0, -1/(2 k μ ρ σ), 0, 0, 0}, {0, 1, -1/(ρ σ), 0, 0, 0}, {1, 0, 0, 0, 0, 0}, {0, 0, 0, 0, 0, 1}, {0, 0, 1/(ρ σ), 0, 1, 0}, {0, -2 k μ, 2 k μ/(ρ σ), -ρ, 2 k μ, 0}}
```

In[25]:=

Inverse[SR].Minors[Inverse[Z], 2] // FullSimplify

```
Out[25]= {{0, 0, 1, 0, 0, 0}, {-2 k μ, 1, 0, 0, 0, 0}, {-ρ, 0, 0, 0, 0, 0}, {2 k μ/(ρ σ), -1/(ρ σ), 0, 0, 1/(ρ σ), -1/(2 k μ ρ σ)}, {2 k μ, 0, 0, 0, 1, 0}, {0, 0, 0, 1, 0, 0}}
```

For the 4D system at zero frequency things are simpler:

```
In[26]:= FullSimplify[
  Series[MatrixExp[h A p] /. {ξp → Sqrt[k^2 - ω^2/vp^2], ξs → Sqrt[k^2 - ω^2/vs^2]}, {ω, 0, 0}] //.
  Normal] /. {vp → Sqrt[σ/ρ], vs → Sqrt[μ/ρ]}] // MatrixForm
```

Out[26]//MatrixForm=

$$\begin{pmatrix} \cosh[hk] & \frac{h\rho(-\mu+\sigma)\sinh[hk]}{2\mu\sigma} & \frac{h k \rho (\mu-\sigma) \cosh[hk]+\rho(\mu+\sigma)\sinh[hk]}{2 k \mu \sigma} & -\sinh[hk] \\ 0 & \cosh[hk] & -\sinh[hk] & 0 \\ 0 & -\sinh[hk] & \cosh[hk] & 0 \\ -\sinh[hk] & \frac{h k \rho (\mu-\sigma) \cosh[hk]-\rho(\mu+\sigma)\sinh[hk]}{2 k \mu \sigma} & \frac{h \rho (-\mu+\sigma)\sinh[hk]}{2 \mu \sigma} & \cosh[hk] \end{pmatrix}$$