

Computational Physics - Project 3

Johannes Scheller, Vincent Noculak, Lukas Powalla

October 23, 2015

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1 Introduction to Project 3

1.1 analytical derivation of the integral

We want to calculate the analytical solution to the integral:

$$I = \int d\vec{r}_1 \int d\vec{r}_2 \frac{1}{|\vec{r}_1 - \vec{r}_2|} \cdot e^{-4(r_1 + r_2)}$$

Therefore, we first transform to spherical Coordinates for each variable. In addition to that, we choose while calculating the integral over r_2 the axis of r_1 as z-axis. With simple scalar product ($\vec{r}_1 \cdot \vec{r}_2 = r_1 \cdot r_2 \cdot \cos(\theta)$) we get the following expression: (don't forget about Jacobi-determinant)

$$I = \int_0^{2\pi} d\phi_1 \int_0^\pi \sin(\theta_1) d\theta_1 \int_0^\infty r_1^2 \cdot e^{-4 \cdot r_1} dr_1 \cdot \int_0^{2\pi} d\phi_2 \int_0^\pi \sin(\theta_2) d\theta_2 \int_0^\infty r_2^2 dr_2 \frac{1}{\sqrt{r_1^2 + r_2^2 - 2 \cdot r_1 \cdot r_2}} \cdot e^{-4 \cdot r_2} \quad (1)$$

$$= 4\pi^2 \cdot \int_0^\pi \sin(\theta_1) d\theta_1 \int_0^\infty r_1^2 \cdot e^{-4 \cdot r_1} dr_1 \cdot \int_0^\infty dr_2 \int_0^\pi d\theta_2 \cdot \frac{r_2^2 \cdot \sin(\theta_2)}{\sqrt{r_1^2 + r_2^2 - 2 \cdot r_1 r_2 \cos(\theta_2)}} \quad (2)$$

$$= 4\pi^2 \cdot \int_0^\pi \sin(\theta_1) d\theta_1 \int_0^\infty r_1^2 \cdot e^{-4 \cdot r_1} dr_1 \cdot I(\vec{r}_1)_2 \quad (3)$$

We now calculate the integral $I(\vec{r}_1)_2$:

$$I(\vec{r}_1)_2 = \int_0^\infty dr_2 \int_0^\pi d\theta_2 \cdot \frac{r_2^2 \cdot \sin(\theta_2)}{\sqrt{r_1^2 + r_2^2 - 2 \cdot r_1 r_2 \cos(\theta_2)}} \quad (4)$$

$$= \int_0^\infty dr_2 \cdot r_2^2 \left[\frac{1}{r_1 r_2} \sqrt{r_1^2 + r_2^2 - 2 r_1 r_2 \cos(\theta_2)} \right]_0^\pi \quad (5)$$

$$= \int_0^\infty dr_2 \cdot r_2^2 \left(\sqrt{r_1^2 + r_2^2 + 2 \cdot r_1 \cdot r_2} - \sqrt{r_1^2 + r_2^2 - 2 \cdot r_1 \cdot r_2} \right) \cdot e^{-4 r_2} \quad (6)$$

$$= \int_0^\infty dr_2 \frac{r_2}{r_1} (r_1 + r_2 - |r_1 - r_2|) \quad (7)$$

Now, we split up the integral in two parts:

$$I(\vec{r}_1)_2 = \int_0^{r_1} dr_2 2 \cdot \frac{r_2^2}{r_1} e^{-4 r_2} + 2 \cdot \int_{r_1}^\infty dr_2 e^{-4 r_2} \quad (8)$$

$$= \frac{2}{r_1} \cdot \hat{I}_1 + \hat{I}_2 \quad (9)$$

Through partial integration, we can determine the following integral:

$$\hat{I}_1 = \int_0^{r_1} dr_2 \cdot r_2^2 e^{-4 r_2} \quad (10)$$

$$= -\frac{1}{4} r_1^2 e^{-4 r_2} \Big|_0^{r_1} + \int_0^{r_1} \frac{1}{2} r_2 e^{-4 r_2} dr_2 \quad (11)$$

$$= -\frac{1}{4} r_1^2 e^{-4 r_1} - \frac{1}{8} r_2 e^{-4 r_2} \Big|_0^{r_1} + \int_0^{r_1} \frac{1}{8} e^{-4 r_2} dr_2 \quad (12)$$

$$= -\frac{1}{4} r_1^2 e^{-4 r_1} - \frac{1}{8} r_1 e^{-4 r_1} + \frac{1}{32} - \frac{1}{32} \cdot e^{-4 r_1} \quad (13)$$

$$(14)$$

similarly, we determine the second integral:

$$\hat{I}_2 = 2 \cdot \int_{r_1}^\infty dr_2 e^{-4 r_2} \quad (15)$$

$$= -\frac{1}{2} r_2 e^{-4 r_2} \Big|_{r_1}^\infty + \int_{r_1}^\infty \frac{1}{2} e^{-4 r_2} dr_2 \quad (16)$$

$$= \frac{1}{2} r_1 e^{-4 r_1} + \frac{1}{8} e^{-4 r_1} \quad (17)$$

In total, we can now determine the integral I_2 :

$$I_2 = -\frac{2}{r_1} \left(\frac{1}{4} r_1^2 e^{-4 r_1} - \frac{1}{8} r_1 e^{-4 r_1} + \frac{1}{32} - \frac{1}{32} \cdot e^{-4 r_1} \right) + \frac{1}{2} r_1 e^{-4 r_1} + \frac{1}{8} e^{-4 r_1} \quad (18)$$

$$= -\frac{1}{16} \frac{(2 e^{-4 r_1} r_1 + e^{-4 r_1} - 1)}{r_1} \quad (19)$$

Finally, we can calculate the integral:

$$I = 4\pi^2 \cdot \int_0^\pi \sin(\theta_1) d\theta_1 \int_0^\infty r_1^2 \cdot e^{-4 \cdot r_1} dr_1 \left(-\frac{1}{16} \frac{(2e^{-4r_1} r_1 + e^{-4r_1} - 1)}{r_1} \right) \quad (20)$$

$$= 8\pi^2 \cdot \int_0^\infty r_1^2 \cdot e^{-4 \cdot r_1} dr_1 \left(-\frac{1}{16} \frac{(2e^{-4r_1} r_1 + e^{-4r_1} - 1)}{r_1} \right) \quad (21)$$

$$= 8\pi^2 \cdot \int_0^\infty r_1^2 \cdot e^{-4 \cdot r_1} dr_1 \left(-\frac{1}{16} \frac{(2e^{-4r_1} r_1 + e^{-4r_1} - 1)}{r_1} \right) \quad (22)$$

$$= -\frac{1}{2}\pi^2 \cdot \int_0^\infty e^{-4 \cdot r_1} dr_1 (2e^{-4r_1} \cdot r_1^2 + e^{-4r_1} \cdot r_1 - r_1) \quad (23)$$

$$= -\frac{1}{2}\pi^2 \cdot \int_0^\infty dr_1 (2e^{-8r_1} \cdot r_1^2 + e^{-8r_1} \cdot r_1 - r_1 \cdot e^{-4 \cdot r_1}) \quad (24)$$

$$(25)$$

We can derive with partial integration the expression:

$$\int_0^\infty dx \cdot x^n \cdot e^{-\beta x} = \frac{n!}{\beta^{n+1}} \quad (26)$$

Then, we can calculate the integral as follows:

$$I = -\frac{1}{2}\pi^2 \cdot \left[\frac{2 \cdot 2!}{8^{2+1}} + \frac{1!}{8^{1+1}} - \frac{1!}{4^{1+1}} \right] \quad (27)$$

$$= -\frac{1}{2}\pi^2 \cdot \left[\frac{4}{8 \cdot 8 \cdot 8} + \frac{8}{8 \cdot 8 \cdot 8} - \frac{1}{8 \cdot 2} \right] \quad (28)$$

$$= -\frac{1}{2}\pi^2 \cdot \left[\frac{12}{8 \cdot 8 \cdot 8} - \frac{32}{8 \cdot 8 \cdot 8} \right] \quad (29)$$

$$= \pi^2 \cdot \left[\frac{16}{8 \cdot 8 \cdot 8} - \frac{6}{8 \cdot 8 \cdot 8} \right] = \pi^2 \frac{10}{8 \cdot 8 \cdot 8} = \pi^2 \frac{5}{16^2} \approx 0.19277 \quad (30)$$

We have now derived a analytical expression for the integral. The integral has the value:

$$I = \frac{5\pi^2}{16^2} \quad (31)$$

2 Theoretical background for numerical integration

2.1 Gaussian quadrature

2.2 Montecarlo integration

3 Execution

3.1 Gaussian quadrature

3.2 Montecarlo integration

We calculated the same integral with montecarlo method. First, we calculated the integral in a brute force way. This means that we calculate the integral using (pseudo) random numbers, which obey uniform distribution functions. The random numbers for each of the six dimensional integral are uniform distributed in a chosen interval (-a to a). Furthermore, we don't transform the integral, but we calculate it in Cartesian coordinates. We got the values in table 1. (We used the interval for a=2) In one dimension, the montecarlo method can be described by formula 31.

$$I = \int_a^b f(x) dx \approx \langle f(x) \rangle \cdot (b-a) = \frac{1}{n} \sum_{i=1}^n f(x_i) \cdot (b-a) = \hat{I} \quad (32)$$

In addition to that, we tried to improve our calculations. First, we transformed the integral into spherical coordinates. In spherical coordinates, we use the variables θ (from 0 to π), ϕ (from 0 to 2π) and r (from 0 to infinity) instead of using Cartesian coordinates $x_{i,k}$ ($i=1,2,3; k=1,2$). We also used a distribution function in order to get appropriate values of the random numbers. Formula 32 to 34 describe the general one dimensional reformulation if you want to use a other particle distribution function.

$$P(x) = \int_0^x p(x) dx \quad (33)$$

$$I = \int_a^b \frac{f(x)}{p(x)} \cdot p(x) dx = \int_a^b \hat{f}(x) \cdot p(x) dx \approx \frac{1}{n} \sum_{i=1}^n \hat{f}(y_i) \cdot (b-a) = \hat{I} \quad (34)$$

$$y_i(x_i) = P^{-1}(p(y(x_i))) = P^{-1}(x_i) \quad (35)$$

In order to improve the precision of the integral, we used a not uniform distribution function, which can be found in formula 35ff.

$$P(x) = \int_0^x 4 \cdot e^{-4x} = 1 - e^{-4x} \quad (36)$$

$$y_i(x_i) = -\frac{1}{4} \ln(1 - x_i) \quad (37)$$

We transformate the integral to spherical Coordinates and use the distribution function for r_1 and r_2 . Finally, the integral can be calculated through formula 38. The results are in table 2.

$$f(r_{1,i}, r_{2,i}, \theta_{1,i}, \theta_{2,i}, \phi_{1,i}, \phi_{2,i}) = \frac{r_{1,i}^2 \cdot r_{2,i}^2 \cdot \sin(\theta_{1,i}) \sin(\theta_{2,i})}{\sqrt{r_{1,i}^2 + r_{2,i}^2 - 2 \cdot r_{1,i} r_{2,i} \cos(\theta_{1,i}) \cos(\theta_{2,i}) + \sin(\theta_{1,i}) \sin(\theta_{2,i}) \cdot \cos(\phi_{1,i} - \phi_{2,i})} \cdot 4^2} \quad (38)$$

$$\hat{I} = \frac{1}{n} \sum_{i=1}^n f(r_{1,i}, r_{2,i}, \theta_{1,i}, \theta_{2,i}, \phi_{1,i}, \phi_{2,i}) \cdot (2\pi - 0)^2 \cdot (\pi - 0)^2 \quad (39)$$

Table 1: Data from the brute force montecarlo algorithm (part c))

n	Integral	standart deviation	time in s
100	0.0563094	0.0362644	0
1000	0.226437	0.159513	0.001
10000	0.0986332	0.0263534	0.006
100000	0.155652	0.01757	0.062
1000000	0.178452	0.00754909	0.667
10000000	0.188683	0.00290854	6.645
100000000	0.19169	0.000942637	70.881

Table 2: Data from the montecarlo algorithm with distribution function (in spherical coordinates) (part d))

n	Integral	standart deviation	time in s
100	0.216734	0.0769193	0
1000	0.173694	0.0241284	0.002
10000	0.186069	0.00885492	0.021
100000	0.19571	0.00343622	0.21
1000000	0.193259	0.00101078	2.07

4 Comparison and discussion of the results

5 Source-code