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Fluid Simulation Using Smoothed Particle Hydrodynamics

Abstract

We explore the method of Smoothed Particle Hydrodynamics for the simulation of water and other interesting liquids, such as liquids with high viscosity and high surface tension. We use two different but related methods for surface reconstruction of the isosurface generated by the density field, Marching Cubes and Ray Marching. For the former we utilized a library and for the latter we reused a previously made ray tracer for Assignment 2 of CS184. We discovered that although ray marching provides more interesting results, along with them comes a significant increase in computation time.

Introduction

The idea behind SPH is that for each particle we can determine any desired scalar quantity by weighting that quantity from neighboring particles with an appropriate kernel function. The formula states that for any value of interest A_i at particle i:

$$A_i = \sum_j A_j \frac{m_j}{\rho_j} W(r_{ij}, h)$$

where m_j and ρ_j are the mass and density of particle j, W is a kernel function, h is the support radius of W, and r_{ij} is the vector from particle j to particle i.

For example, to get the density of particle i, we just need to plug in ρ_j for the quantity of interest A_j. This tells us that:

$$\rho_i = \sum_j m_j W(r_{ij}, h)$$

Solving Navier-Stokes Equation

For each particle we need to update its position and velocity based on the current arrangement of particles. The equation we're solving to get the change in velocity of each particle is:

$$\rho_i \frac{dv_i}{dt} = - \bigtriangledown p_i + \mu \bigtriangledown^2 \cdot v_i + f_{external}$$

Thus, to get the values necessary to solve the Navier-Stokes equation, we need to know the density, pressure gradient, viscosity, and external forces acting on the particle, such as gravitational force and surface tension. Computing most of the desired quantities relies on knowing the density at each particle. Thus, we first loop through all particles and compute density, and then loop again to compute the remaining quantities.

The values we need in terms of our kernel method, in addition to the density calculation above, are then just:

Pressure Gradient

$$-\bigtriangledown p_i = -\sum_j m_j \frac{p_i + p_j}{2\rho_j} \bigtriangledown W(r_{ij}, h)$$

where the p_i and p_j are pressures given by $p_i = k(\rho_i - \rho_0)$ with ρ_0 the ideal density of the fluid and k a stiffness constant.

Viscosity

$$\mu \bigtriangledown^2 \cdot v_i = \mu \sum_j m_j \frac{v_j - v_i}{\rho_j} \bigtriangledown^2 W(r_{ij}, h)$$

where μ is the viscosity coefficient of the fluid.

Surface Tension

$$f_{tension} = -\sigma \bigtriangledown^2 c_i \mathbf{n_i}$$

where c_i is the color field calculated by plugging in 1 for A_j in the generic kernel form, and $\mathbf{n_i}$ is the normalized gradient of the color field.

Gravitational Force

$$f_{gravity} = \rho_i \mathbf{g}$$

We add up all these quanities in the Navier-Stokes equations and solve the equation for the change in velocity of particle i. Then we update the current velocity and position for each particle. To do this, we use Euler's method with a step size of .01. We tried the method of Verlet as well, but experienced no change in the simulation, so we went back to using the simpler method. This perhaps would have not been the case had we used a larger timestep.

Kernels

For our computations we use the following kernels as suggested from [PRESENTATION CITATION]. In all of the following, h is the support radius of the kernel, $\mathbf{r_{ij}}$ is the vector from point r_j to r_i, and r is the length of this vector. For all the kernels, we use the below formulas provided $r < h$; otherwise, we return 0.

Default

$$W(r, h) = \frac{315}{64\pi h^9}(h^2 - r^2)^3$$

The default kernel is used for all quantities except pressure and viscosity.

Pressure

$$W(r, h) = \frac{15}{\pi h^6}(h - r)^3$$

We use this kernel only for calculating pressure gradient, so really we need the gradient of this kernel, which is:

$$W(\mathbf{r_{ij}}, h) = \frac{-45}{\pi h^6} \frac{\mathbf{r_{ij}}}{r} (h - r)^2$$

We use this kernel for pressure because the gradient is large as r approaches 0.

Viscosity

$$W(r, h) = \frac{15}{2\pi h^3} \left(-\frac{r^3}{2h^3} + \frac{r^2}{h^2} + \frac{h}{2r} - 1\right)$$

Similar to pressure, we really need the laplacian of this kernel for the viscosity calculations. The laplacian is:

$$\bigtriangledown^2 W(r, h) = \frac{45}{\pi h^6} (h - r)$$

Improving Computation Time

A naïve implementation of a fluid simulation using SPH takes into account interaction between all pairs of particles each time step, producing a runtime of $O(n^2)$, where n is the number of particles. To speed up the simulation, we used the nearest neighbor algorithm which assumes that particles will only interact if they are within a certain distance of each other. This distance, called the radius of support, can be adjusted so that the fluid simulation produced looks identical to the simulation produced by the naïve implementation. Using this algorithm, the runtime is reduced to an expected runtime of $O(nm)$, where m is the average number of neighbors of each particle.

Explicit Constants

In the following table we list all the constants we used for simulating water. To simulate other more viscous fluids we just vary the values in the table below to preference. In addition to those listed, we use a length tolerance of 6.0 for the surface tension calculation; ie, if the length of the surface normal is less than the tolerance we do not include it in the calculation.

Mass	.02
Stiffness	3.5
Ideal Density	1000
Viscosity	3.5
Surface Tension	.07
Support Radius	.0625

The majority of these constants come from [KELAGER], which also has reccomended constants for other fluids, such as mucous. One could consult those constants if interested in replicating those results.

Surface Reconstruction

To reconstruct the surface we use an isosurface level of 500. Although 1000 would be more physically accurate for water, we found the results for 1000 to make the water appear significantly more globular and compressible, and less like actual water. For this reason it would seem the method of SPH is more suited for compressible fluids such as gasses or smoke.

For the first method of surface reconstruction with Marching Cubes we used a standard implementation from [INSERT WEBISTE AND NAME]. The algorithm breaks up the container of the fluid into a grid and considers the density value at each of the eight vertices of a cube in the grid. Depending on which vertex densities are above a certain threshold we have different cases to consider for how to draw the triangles. We implemented the marching squares version ourselves for 2D simulation (which requires considering 16 different cases when drawing triangles) but relied on the implementation from [WEBSITE AND NAME] to perform marching cubes (which requires 256 cases).

For the second method, we reused the ray tracer we had made for a previous assignment. To do this, we extended a super class of Shape, which we had previously used to extend to Ellipsoids and Triangles, to a subclass of ParticleBlob and implemented a new intersection checking method using ray marching. To determine the density and normal vector at a point we reuse the kernel functions from the SPH simulation. In addition, we reuse the neighbor structure to accelerate the processing of the particles within the ray tracer. With this ray tracer we were able to generate much more interesting images, which used reflections and refractions. But the computation cost of using ray marching for the surface reconstruction was significantly higher than that of Marching Cubes.

Results

Implementing the nearest neighbor algorithm allowed us to run real-time simulations with approximately 1200 particles at an average of 15 frames per second, rendering particles as spheres.

Conclusion & Future Consideration

References

[1] Kelager, M., "Lagrangian Fluid Dynamics Using Smoothed Particle Hydrodynamics," *Department of Computer Science, University of Copenhagen*, January 9th, 2006.

[2] Fisher, M., Stanford, Matt's Webcorner,
http://graphics.stanford.edu/~mdfisher/MarchingCubes.html (accessed April 20th, 2013).

[3] Auer, S., "Realtime Particle-Based Fluid Simulation", Computer Graphics and Visualization, October 4th, 2008.
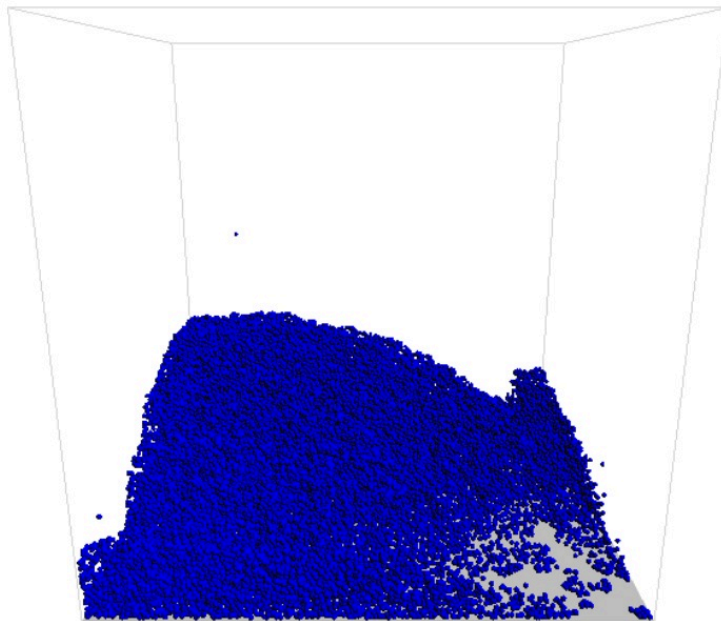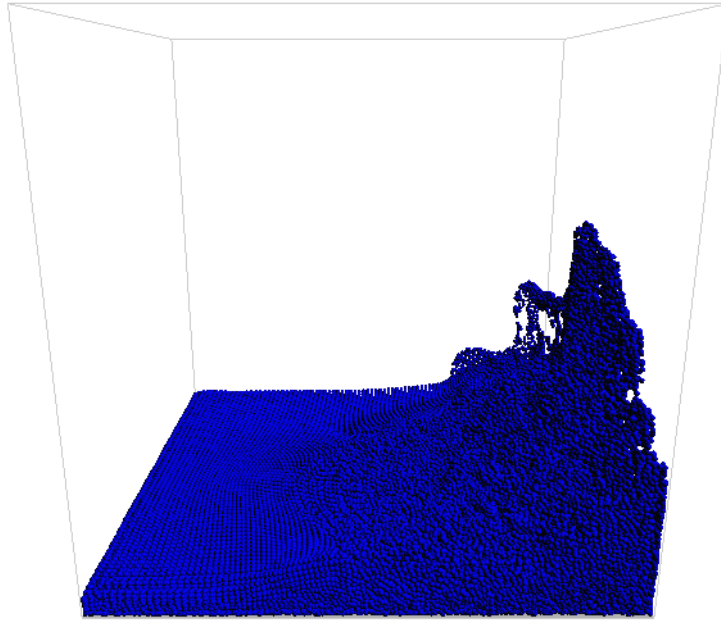
Figure 1: Approximately 100k particles being sloshed to the right. Approximately 80k particles released from the corner as a dam break.
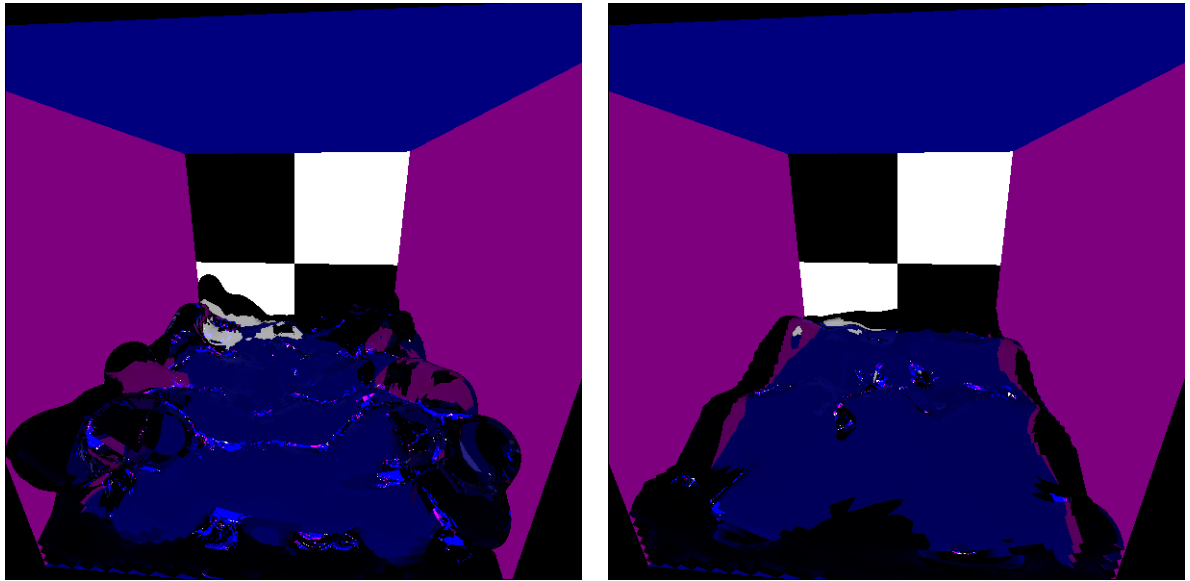
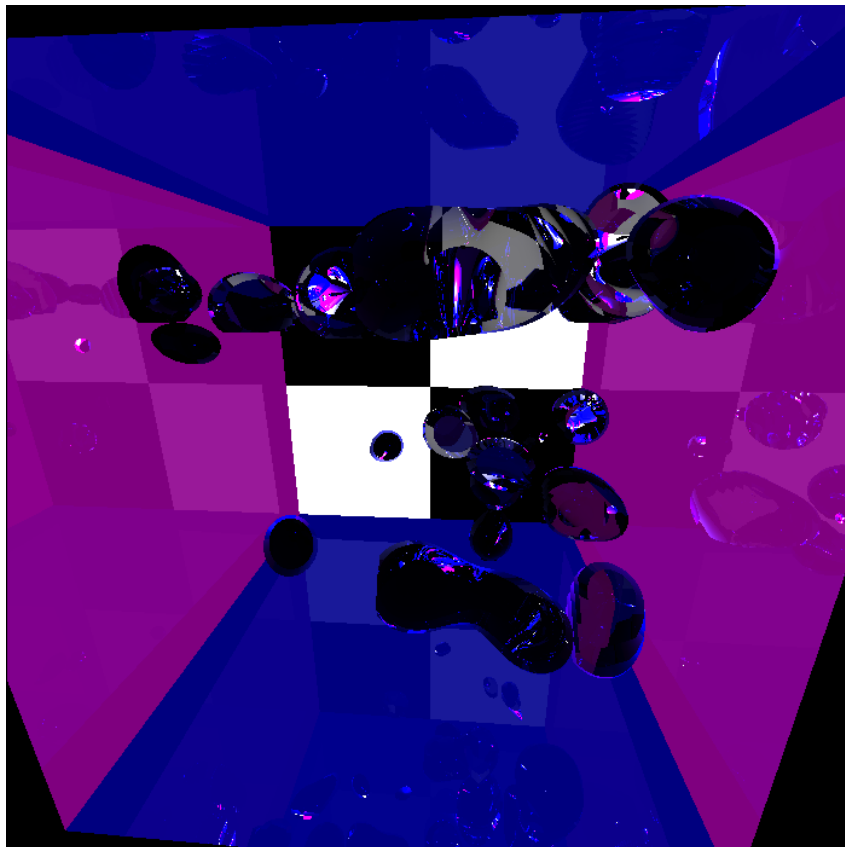Figure 2: Two ray traced images a few frames apart in the scene. The image shows refraction and reflection.

Figure 3: Another ray traced image with more more noticable refraction and higher reflections.
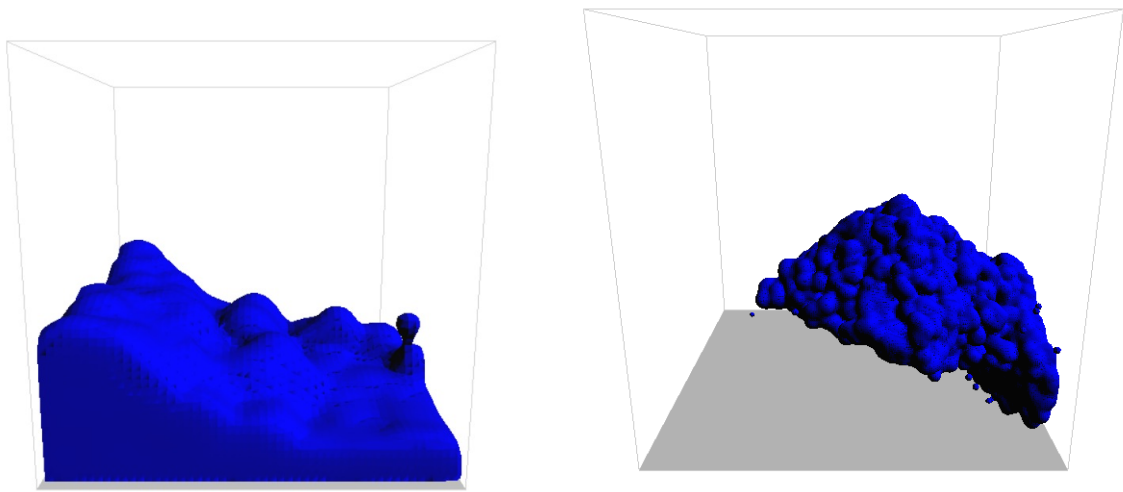
Figure 4: The image on the left uses 1.8k particles, while the image on the right uses 21k. Note the reconstructed surface appears globular, particularly for the image on the right.