

# INTEGRATED LOCAL ENERGY DECAY (MORAWETZ) ESTIMATES

OVIDIU-NECULAI AVADANEI

We consider the Cauchy problem for the wave equation

$$(1) \quad \begin{cases} (-\partial_t^2 + \Delta)u = f \\ u(0) = u_0, \\ \partial_t u(0) = u_1 \end{cases}$$

As we have already seen, its solution satisfies the energy estimate

$$\|\nabla_{t,x}u\|_{L_t^\infty L_x^2} \lesssim \|(u_0, u_1)\|_{\dot{H}_x^1 \times L_x^2} + \|f\|_{L_t^1 L_x^2}$$

In Section 1, we shall prove another energy estimate that quantifies the dispersive decay of  $u$  in terms of  $L^2$ -based norms, which is known as the *Morawetz* or the *integrated local energy decay* bound.

One of the reasons why this estimate is useful is that it can be employed as an *intermediary* decay bound, which can enable one to obtain stronger decay bounds under an appropriate asymptotic flatness condition. We shall illustrate this in Section 2 by presenting a result of Rodnianski-Schlag that establishes non-endpoint Strichartz estimates for sub-principal asymptotically flat perturbations of  $\square$  (which involve only first and zeroth order perturbations), assuming the integrated local energy decay condition. We note that this is only the simplest example of a plethora of such results (see [Tataru], [Dafermos-Rodnianski], [Moschidis], [Oliver-Sterbenz], etc.).

In section 3 we give a spectral characterization of the integrated local energy decay. This shows that tools from spectral theory (such as Fredholm theory, resonances, semi-classical analysis, etc.) can be used to obtain the integrated local energy decay estimate, whose usefulness is illustrated in Section 2.

## 1. AN INTEGRATED LOCAL ENERGY DECAY ESTIMATE

We define

$$A_j = \begin{cases} \{x \in \mathbb{R}^d | 2^j \leq |x| < 2^{j+1}\}, j \geq 1 \\ \{x \in \mathbb{R}^d | |x| < 2\}, j = 0 \end{cases}$$

We also define

$$\begin{aligned} \|u\|_{LE} &= \sup_{j \geq 0} \|\langle r \rangle^{-\frac{1}{2}} u\|_{L^2 L^2(\mathbb{R}_t \times A_j)} \\ \|f\|_{LE^*} &= \sum_{j \geq 0} \|\langle r \rangle^{\frac{1}{2}} f\|_{L^2 L^2(\mathbb{R}_t \times A_j)} \end{aligned}$$

We consider the equation (1)

**Theorem 1.1.** *Every solution of (1) satisfies*

$$\|\nabla_{t,x}u\|_{LE} + \|\langle r \rangle^{-1}u\|_{LE} \lesssim \|(u_0, u_1)\|_{\dot{H}^1 \times L^2} + \|f\|_{LE^*}.$$

---

*Date:* 10/20/2021.

*Proof.* Let  $X = \varphi(r)x^j\partial_j$ , where  $\varphi$  is going to be chosen later. We have

$$\begin{aligned}
\langle Xu, \Delta u \rangle &= \int_{\mathbb{R}^d} \varphi(r)x^j\partial_j u \cdot \Delta u \, dx = \int_{\mathbb{R}^d} \partial_j u \cdot \varphi(r)x^j \Delta u \, dx = - \int_{\mathbb{R}^d} u \cdot \partial_j(\varphi(r)x^j \Delta u) \, dx \\
&= - \int_{\mathbb{R}^d} u \cdot \varphi'(r) \frac{x^j x_j}{r} \Delta u \, dx - \int_{\mathbb{R}^d} u \cdot \varphi(r) \delta_j^j \Delta u \, dx - \int_{\mathbb{R}^d} u \cdot \varphi(r)x^j \partial_j \Delta u \, dx \\
&= - \int_{\mathbb{R}^d} u \cdot r\varphi'(r) \Delta u \, dx - d \int_{\mathbb{R}^d} u \cdot \varphi(r) \Delta u \, dx - \int_{\mathbb{R}^d} u \cdot X \Delta u \, dx \\
&= - \langle (r\varphi'(r) + d\varphi(r))u, \Delta u \rangle - \langle X(\Delta u), u \rangle \\
&= - \langle (r\varphi'(r) + d\varphi(r))u, \Delta u \rangle - \langle \Delta(Xu), u \rangle - \langle [X, \Delta]u, u \rangle \\
&= - \langle (r\varphi'(r) + d\varphi(r))u, \Delta u \rangle - \langle Xu, \Delta u \rangle + \langle [\Delta, X]u, u \rangle
\end{aligned}$$

Thus,

$$\langle Xu, \Delta u \rangle = \frac{1}{2} \langle [\Delta, X]u, u \rangle - \frac{1}{2} \langle (r\varphi'(r) + d\varphi(r))u, \Delta u \rangle$$

For every index  $k \in \{1, 2, \dots, d\}$ , we have

$$\begin{aligned}
\partial_k(Xu) &= \partial_k(\varphi(r)x^j\partial_j u) = \varphi'(r) \frac{x_k}{r} x^j \partial_j u + \varphi(r) \partial_k u + \varphi(r)x^j \partial_j \partial_k u \\
\partial_k^2(Xu) &= \varphi''(r) \frac{x_k^2}{r^2} x^j \partial_j u + \varphi'(r) \frac{r - x_k \cdot \frac{x_k}{r}}{r^2} x^j \partial_j u + \varphi'(r) \frac{x_k}{r} \partial_k u + \varphi'(r) \frac{x_k}{r} x^j \partial_j \partial_k u \\
&\quad + \varphi'(r) \frac{x_k}{r} \partial_k u + \varphi(r) \partial_k^2 u \\
&\quad + \varphi'(r) \frac{x_k}{r} x^j \partial_j \partial_k u + \varphi(r) \partial_k^2 u + \varphi(r)x^j \partial_j \partial_k^2 u
\end{aligned}$$

Thus,

$$\begin{aligned}
\Delta(Xu) &= \varphi''(r)x^j\partial_j u + \varphi'(r) \frac{d-1}{r} x^j \partial_j u + \varphi'(r) \frac{2}{r} x^j \partial_j u + 2 \frac{\varphi'(r)}{r} x^k x^j \partial_j \partial_k u + 2\varphi(r) \Delta u + X(\Delta u) \\
&= \left( \varphi''(r) + \frac{\varphi'(r)}{r} (d+1) \right) x^j \partial_j u + 2 \frac{\varphi'(r)}{r} x^k x^j \partial_j \partial_k u + 2\varphi(r) \Delta u + X(\Delta u) \\
[\Delta, X]u &= \left( \varphi''(r) + \frac{\varphi'(r)}{r} (d+1) \right) x^j \partial_j u + 2 \frac{\varphi'(r)}{r} x^k x^j \partial_j \partial_k u + 2\varphi(r) \Delta u
\end{aligned}$$

In this case,

$$\langle [\Delta, X]u, u \rangle = \left\langle \left( \varphi''(r) + \frac{\varphi'(r)}{r} (d+1) \right) x^j \partial_j u, u \right\rangle + 2 \left\langle \frac{\varphi'(r)}{r} x^k x^j \partial_j \partial_k u, u \right\rangle + 2 \langle \varphi(r) \Delta u, u \rangle$$

We also have

$$\begin{aligned}
\left\langle \frac{\varphi'(r)}{r} x_k x_j \partial_j \partial_k u, u \right\rangle &= \int_{\mathbb{R}^d} \partial_j \partial_k u \cdot \frac{\varphi'(r)}{r} x_k x_j u \, dx = - \int_{\mathbb{R}^d} \partial_k u \cdot \partial_j \left( \frac{\varphi'(r)}{r} x_k x_j u \right) \, dx \\
&= - \int_{\mathbb{R}^d} \partial_k u \cdot \frac{\varphi''(r) \frac{x_j}{r} \cdot r - \varphi'(r) \frac{x_j}{r}}{r^2} x_k x_j u \, dx - \int_{\mathbb{R}^d} \partial_k u \cdot \frac{\varphi'(r)}{r} \delta_k^j x_j u \, dx \\
&\quad - \int_{\mathbb{R}^d} \partial_k u \cdot \frac{\varphi'(r)}{r} x_k u \, dx - \int_{\mathbb{R}^d} \partial_k u \cdot \frac{\varphi'(r)}{r} x_k x_j \partial_j u \, dx
\end{aligned}$$

Thus,

$$\begin{aligned}
 \left\langle \frac{\varphi'(r)}{r} x^k x^j \partial_j \partial_k u, u \right\rangle &= - \int_{\mathbb{R}^d} \frac{\varphi''(r)r - \varphi'(r)}{r} x^j \partial_j u \cdot u \, dx - \int_{\mathbb{R}^d} 2 \frac{\varphi'(r)}{r} x^j \partial_j u \cdot u \, dx \\
 &\quad - \int_{\mathbb{R}^d} \frac{\varphi'(r)}{r} (x^j \partial_j u)^2 \, dx \\
 &= - \int_{\mathbb{R}^d} \left( \varphi''(r) + \frac{\varphi'(r)}{r} \right) x^j \partial_j u \cdot u \, dx - \int_{\mathbb{R}^d} \frac{\varphi'(r)}{r} (x^j \partial_j u)^2 \, dx
 \end{aligned}$$

Besides this,

$$\begin{aligned}
 \langle \varphi(r) \partial_j^2 u, u \rangle &= \int_{\mathbb{R}^d} \partial_j^2 u \cdot \varphi(r) u \, dx = - \int_{\mathbb{R}^d} \partial_j u \cdot \partial_j (\varphi(r) u) \, dx \\
 &= - \int_{\mathbb{R}^d} \partial_j u \cdot \frac{\varphi'(r)}{r} x_j u \, dx - \int_{\mathbb{R}^d} \partial_j u \cdot \varphi(r) \partial_j u \, dx
 \end{aligned}$$

This means that

$$\langle \varphi(r) \Delta u, u \rangle = - \int_{\mathbb{R}^d} \frac{\varphi'(r)}{r} x^j \partial_j u \cdot u \, dx - \int_{\mathbb{R}^d} \varphi(r) |\nabla u|^2 \, dx$$

We have

$$\begin{aligned}
 \langle [\Delta, X] u, u \rangle &= \left\langle \left( \varphi''(r) + \frac{\varphi'(r)}{r} (d+1) \right) x^j \partial_j u, u \right\rangle + 2 \left\langle \frac{\varphi'(r)}{r} x^k x^j \partial_j \partial_k u, u \right\rangle + 2 \langle \varphi(r) \Delta u, u \rangle \\
 &= \left\langle \left( \varphi''(r) + \frac{\varphi'(r)}{r} (d+1) \right) x^j \partial_j u, u \right\rangle - 2 \int_{\mathbb{R}^d} \left( \varphi''(r) + \frac{\varphi'(r)}{r} \right) x^j \partial_j u \cdot u \, dx \\
 &\quad - 2 \int_{\mathbb{R}^d} \frac{\varphi'(r)}{r} (x^j \partial_j u)^2 \, dx - 2 \int_{\mathbb{R}^d} \frac{\varphi'(r)}{r} x^j \partial_j u \cdot u \, dx - 2 \int_{\mathbb{R}^d} \varphi(r) |\nabla u|^2 \, dx \\
 &= \left\langle \left( -\varphi''(r) + \frac{\varphi'(r)}{r} (d-3) \right) x^j \partial_j u, u \right\rangle - 2 \int_{\mathbb{R}^d} \frac{\varphi'(r)}{r} (x^j \partial_j u)^2 \, dx - 2 \int_{\mathbb{R}^d} \varphi(r) |\nabla u|^2 \, dx
 \end{aligned}$$

We also note that

$$\begin{aligned}
 \left\langle \left( -\varphi''(r) + \frac{\varphi'(r)}{r} (d-3) \right) x_j \partial_j u, u \right\rangle &= \int_{\mathbb{R}^d} \partial_j u \cdot x_j \psi(r) u \, dx = - \int_{\mathbb{R}^d} u \cdot \partial_j (x_j \psi(r) u) \, dx \\
 &= - \int_{\mathbb{R}^d} u \cdot \psi(r) u \, dx - \int_{\mathbb{R}^d} u \cdot x_j \psi'(r) \frac{x_j}{r} u \, dx \\
 &\quad - \int_{\mathbb{R}^d} u \cdot x_j \psi(r) \partial_j u \, dx
 \end{aligned}$$

Thus,

$$\int_{\mathbb{R}^d} \psi(r) x^j \partial_j u \cdot u \, dx = - \int_{\mathbb{R}^d} d\psi(r) u^2 \, dx - \int_{\mathbb{R}^d} r\psi'(r) u^2 \, dx - \int_{\mathbb{R}^d} \psi(r) x^j \partial_j u \cdot u \, dx,$$

which shows that

$$\int_{\mathbb{R}^d} \psi(r) x^j \partial_j u \cdot u \, dx = - \frac{1}{2} \int_{\mathbb{R}^d} (d\psi(r) + r\psi'(r)) u^2 \, dx$$

As  $\psi(r) = -\varphi''(r) + \frac{\varphi'(r)}{r}(d-3)$ , we get that

$$\begin{aligned}\psi'(r) &= -\varphi'''(r) + (d-3)\frac{\varphi''(r)r - \varphi'(r)}{r^2} \\ d\psi(r) + r\psi'(r) &= -d\varphi''(r) + d(d-3)\frac{\varphi'(r)}{r} - r\varphi'''(r) + (d-3)\frac{\varphi''(r)r - \varphi'(r)}{r} \\ &= -r\varphi'''(r) - 3\varphi''(r) + (d-1)(d-3)\frac{\varphi'(r)}{r}\end{aligned}$$

Therefore,

$$\left\langle \left( -\varphi''(r) + \frac{\varphi'(r)}{r}(d-3) \right) x_j \partial_j u, u \right\rangle = -\frac{1}{2} \int_{\mathbb{R}^d} \left( -r\varphi'''(r) - 3\varphi''(r) + (d-1)(d-3)\frac{\varphi'(r)}{r} \right) \cdot u^2 dx$$

In this case,

$$\begin{aligned}\langle [\Delta, X]u, u \rangle &= -2 \int_{\mathbb{R}^d} \frac{\varphi'(r)}{r} (x^j \partial_j u)^2 dx - 2 \int_{\mathbb{R}^d} \varphi(r) |\nabla u|^2 dx \\ &\quad - \frac{1}{2} \int_{\mathbb{R}^d} \left( -r\varphi'''(r) - 3\varphi''(r) + (d-1)(d-3)\frac{\varphi'(r)}{r} \right) \cdot u^2 dx\end{aligned}$$

Therefore,

$$\begin{aligned}\left\langle Xu + \frac{r\varphi'(r) + d\varphi(r)}{2}u, \Delta u \right\rangle &= \frac{1}{2} \langle [\Delta, X]u, u \rangle \\ &= - \int_{\mathbb{R}^d} \frac{\varphi'(r)}{r} (x^j \partial_j u)^2 dx - \int_{\mathbb{R}^d} \varphi(r) |\nabla u|^2 dx \\ &\quad - \frac{1}{4} \int_{\mathbb{R}^d} \left( -r\varphi'''(r) - 3\varphi''(r) + (d-1)(d-3)\frac{\varphi'(r)}{r} \right) \cdot u^2 dx\end{aligned}$$

We note that, in general,

$$\begin{aligned}\int_{\mathbb{R}^d} \partial_j h \cdot \varphi(r) x_j h dx &= - \int_{\mathbb{R}^d} h \cdot \partial_j (\varphi(r) x_j h) dx = - \int_{\mathbb{R}^d} h \cdot \varphi'(r) \frac{x_j}{r} x_j h dx - \int_{\mathbb{R}^d} h \cdot \varphi(r) h dx \\ &\quad - \int_{\mathbb{R}^d} h \cdot \varphi(r) x_j \partial_j h dx \Rightarrow \\ \int_{\mathbb{R}^d} \partial_j h \cdot \varphi(r) x_j h dx &= -\frac{1}{2} \int_{\mathbb{R}^d} h \cdot \varphi'(r) \frac{x_j^2}{r} h dx - \frac{1}{2} \int_{\mathbb{R}^d} h \cdot \varphi(r) h dx \\ \langle Xh, h \rangle &= -\frac{1}{2} \int_{\mathbb{R}^d} r\varphi'(r) h^2 dx - \frac{1}{2} \int_{\mathbb{R}^d} d\varphi(r) h^2 dx\end{aligned}$$

We also have

$$\begin{aligned}\int_0^T \left\langle Xu + \frac{r\varphi'(r) + d\varphi(r)}{2}u, -\partial_t^2 u \right\rangle dt &= \left\langle Xu + \frac{r\varphi'(r) + d\varphi(r)}{2}u, -\partial_t u \right\rangle \Big|_0^T \\ &\quad + \int_0^T \left\langle X\partial_t u + \frac{r\varphi'(r) + d\varphi(r)}{2}\partial_t u, \partial_t u \right\rangle dt \\ &= \left\langle Xu + \frac{r\varphi'(r) + d\varphi(r)}{2}u, -\partial_t u \right\rangle \Big|_0^T \\ \langle X\partial_t u, \partial_t u \rangle &= - \int_{\mathbb{R}^d} \frac{r\varphi'(r) + d\varphi(r)}{2} (\partial_t u)^2\end{aligned}$$

We can now write

$$\begin{aligned}
 \int_0^T \left\langle Xu + \frac{r\varphi'(r) + d\varphi(r)}{2}u, f \right\rangle dt &= \int_0^T \left\langle Xu + \frac{r\varphi'(r) + d\varphi(r)}{2}u, -\partial_t^2 u \right\rangle dt \\
 &\quad + \int_0^T \left\langle Xu + \frac{r\varphi'(r) + d\varphi(r)}{2}u, \Delta u \right\rangle dt \\
 &= \left\langle Xu + \frac{r\varphi'(r) + d\varphi(r)}{2}u, -\partial_t u \right\rangle \Big|_0^T \\
 &\quad - \int_0^T \int_{\mathbb{R}^d} \frac{\varphi'(r)}{r} (x^j \partial_j u)^2 dx dt - \int_0^T \int_{\mathbb{R}^d} \varphi(r) |\nabla u|^2 dx dt \\
 &\quad - \frac{1}{4} \int_0^T \int_{\mathbb{R}^d} \left( -r\varphi'''(r) - 3\varphi''(r) + (d-1)(d-3) \frac{\varphi'(r)}{r} \right) \cdot u^2 dx dt
 \end{aligned}$$

Therefore,

$$\begin{aligned}
 &- \int_0^T \int_{\mathbb{R}^d} \frac{\varphi'(r)}{r} (x^j \partial_j u)^2 dx dt - \int_0^T \int_{\mathbb{R}^d} \varphi(r) |\nabla u|^2 dx dt \\
 &- \frac{1}{4} \int_0^T \int_{\mathbb{R}^d} \left( -r\varphi'''(r) - 3\varphi''(r) + (d-1)(d-3) \frac{\varphi'(r)}{r} \right) \cdot u^2 dx dt \\
 &= \int_0^T \left\langle \varphi(r) x^j \partial_j u + \frac{r\varphi'(r) + d\varphi(r)}{2}u, f \right\rangle dt + \left\langle \varphi(r) x^j \partial_j u + \frac{r\varphi'(r) + d\varphi(r)}{2}u, \partial_t u \right\rangle \Big|_0^T
 \end{aligned}$$

On the other hand, if  $w$  is a smooth function (to be chosen later as well),

$$\begin{aligned}
 \int_0^T \langle wu, f \rangle dt &= \int_0^T \langle wu, -\partial_t^2 u \rangle dt + \int_0^T \langle wu, \Delta u \rangle dt \\
 &= \langle wu, -\partial_t u \rangle \Big|_0^T + \int_0^T \langle w \partial_t u, \partial_t u \rangle dt - \int_0^T \langle \nabla(wu), \nabla u \rangle dt \\
 &= \langle wu, -\partial_t u \rangle \Big|_0^T + \int_0^T \int_{\mathbb{R}^d} w(r) (\partial_t u)^2 dx dt - \int_0^T \int_{\mathbb{R}^d} w(r) |\nabla u|^2 dx dt \\
 &\quad - \int_0^T \int_{\mathbb{R}^d} u \frac{w'(r)}{r} x^j \partial_j u dx dt \\
 &= \langle wu, -\partial_t u \rangle \Big|_0^T + \int_0^T \int_{\mathbb{R}^d} w(r) (\partial_t u)^2 dx dt - \int_0^T \int_{\mathbb{R}^d} w(r) |\nabla u|^2 dx dt \\
 &\quad + \frac{1}{2} \int_0^T \int_{\mathbb{R}^d} \left( w''(r) + \frac{d-1}{r} w'(r) \right) u^2 dx dt
 \end{aligned}$$

Thus,

$$\begin{aligned}
 \int_0^T \langle wu, f \rangle dt + \langle wu, \partial_t u \rangle \Big|_0^T + \int_0^T \int_{\mathbb{R}^d} w(r) |\nabla u|^2 dx dt - \frac{1}{2} \int_0^T \int_{\mathbb{R}^d} \left( w''(r) + \frac{d-1}{r} w'(r) \right) u^2 dx dt \\
 = \int_0^T \int_{\mathbb{R}^d} w(r) (\partial_t u)^2 dx dt
 \end{aligned}$$

In this case,

$$\begin{aligned}
& \int_0^T \int_{\mathbb{R}^d} w(r) (\partial_t u)^2 dx dt - \int_0^T \int_{\mathbb{R}^d} \frac{\varphi'(r)}{r} (x^j \partial_j u)^2 dx dt - \int_0^T \int_{\mathbb{R}^d} \varphi(r) |\nabla u|^2 dx dt \\
& - \frac{1}{4} \int_0^T \int_{\mathbb{R}^d} \left( -r\varphi'''(r) - 3\varphi''(r) + (d-1)(d-3) \frac{\varphi'(r)}{r} \right) \cdot u^2 dx dt \\
& = \int_0^T \left\langle \varphi(r) x^j \partial_j u + \frac{r\varphi'(r) + d\varphi(r) + 2w(r)}{2} u, f \right\rangle dt + \left\langle \varphi(r) x^j \partial_j u + \frac{r\varphi'(r) + d\varphi(r)}{2} u, \partial_t u \right\rangle \Big|_0^T \\
& + \langle wu, \partial_t u \rangle \Big|_0^T + \int_0^T \int_{\mathbb{R}^d} w(r) |\nabla u|^2 dx dt - \frac{1}{2} \int_0^T \int_{\mathbb{R}^d} \left( w''(r) + \frac{d-1}{r} w'(r) \right) u^2 dx dt,
\end{aligned}$$

which means that

$$\begin{aligned}
& \int_0^T \int_{\mathbb{R}^d} w(r) (\partial_t u)^2 dx dt - \int_0^T \int_{\mathbb{R}^d} \frac{\varphi'(r)}{r} (x^j \partial_j u)^2 dx dt - \int_0^T \int_{\mathbb{R}^d} (\varphi(r) - w(r)) |\nabla u|^2 dx dt \\
& - \frac{1}{4} \int_0^T \int_{\mathbb{R}^d} \left( -r\varphi'''(r) - 3\varphi''(r) - 2w''(r) - 2\frac{d-1}{r} w'(r) + (d-1)(d-3) \frac{\varphi'(r)}{r} \right) \cdot u^2 dx dt \\
& = \int_0^T \left\langle \varphi(r) x^j \partial_j u + \frac{r\varphi'(r) + d\varphi(r) + 2w(r)}{2} u, f \right\rangle dt + \left\langle \varphi(r) x^j \partial_j u + \frac{r\varphi'(r) + d\varphi(r)}{2} u, \partial_t u \right\rangle \Big|_0^T \\
& + \langle wu, \partial_t u \rangle \Big|_0^T
\end{aligned}$$

We choose  $\varphi \in \mathcal{S}(\mathbb{R})$ ,  $\varphi \leq -1$  on  $[-2, 2]$ ,  $|\varphi'(r)| \ll r|\varphi(r)|$ , and  $r|\varphi'(r)| \ll |\varphi(r)|$ . Let  $w = -\frac{r\varphi'(r) + d\varphi(r)}{2}$ . In this case,  $-r\varphi'''(r) - 3\varphi''(r) - 2w''(r) - 2\frac{d-1}{r} w'(r) + (d-1)(d-3) \frac{\varphi'(r)}{r} = (2d-2) \left( \varphi''(r) + (d-1) \frac{\varphi'(r)}{r} \right)$ , and we impose  $\varphi''(r) + (d-1) \frac{\varphi'(r)}{r} \leq 0$  on  $[0, \infty)$  and  $\varphi''(r) + (d-1) \frac{\varphi'(r)}{r} \leq -\lambda$  on  $[0, 2]$ , where  $\lambda > 0$  is a fixed small parameter. We thus have

$$\begin{aligned}
& - \int_0^T \int_{\mathbb{R}^d} \frac{r\varphi'(r) + d\varphi(r)}{2} (\partial_t u)^2 dx dt - \int_0^T \int_{\mathbb{R}^d} \frac{\varphi'(r)}{r} (x^j \partial_j u)^2 dx dt \\
& - \int_0^T \int_{\mathbb{R}^d} \frac{r\varphi'(r) + (d+2)\varphi(r)}{2} |\nabla u|^2 dx dt \\
& - \frac{d-1}{2} \int_0^T \int_{\mathbb{R}^d} \left( \varphi''(r) + (d-1) \frac{\varphi'(r)}{r} \right) \cdot u^2 dx dt \\
& = \int_0^T \langle \varphi(r) x^j \partial_j u, f \rangle dt + \langle \varphi(r) x^j \partial_j u, \partial_t u \rangle \Big|_0^T
\end{aligned}$$

Therefore,

$$\begin{aligned}
& \|\nabla u\|_{L^2 L^2([0, T] \times B_1)}^2 + \|\partial_t u\|_{L^2 L^2([0, T] \times B_1)}^2 + \|u\|_{L^2 L^2([0, T] \times B_1)}^2 \\
& \lesssim \left| \int_0^T \langle \varphi(r) x^j \partial_j u, f \rangle dt \right| + \|\nabla u(T)\|_{L_x^2}^2 + \|\nabla u(0)\|_{L_x^2}^2 \\
& + \|\partial_t u(T)\|_{L_x^2}^2 + \|\partial_t u(0)\|_{L_x^2}^2
\end{aligned}$$

Hardy's inequality implies that

$$\begin{aligned}
& \|\nabla_{t,x} u\|_{L^2 L^2([0, T] \times B_1)}^2 + \|r^{-1} u\|_{L^2 L^2([0, T] \times B_1)}^2 \lesssim \left| \int_0^T \langle \varphi(r) x^j \partial_j u, f \rangle dt \right| \\
& + \|\nabla_{t,x} u(T)\|_{L_x^2}^2 + \|\nabla_{t,x} u(0)\|_{L_x^2}^2
\end{aligned}$$

By taking  $\beta(r) = r\varphi(r)$ , we can rewrite the previous inequality in the form

$$\begin{aligned} \|\nabla_{t,x}u\|_{L^2L^2([0,T]\times B_1)}^2 + \|r^{-1}u\|_{L^2L^2([0,T]\times B_1)}^2 &\lesssim \left| \int_0^T \langle \beta(r)\partial_r u, f \rangle dt \right| \\ &\quad + \|\nabla_{t,x}u(T)\|_{L_x^2}^2 + \|\nabla_{t,x}u(0)\|_{L_x^2}^2, \end{aligned}$$

with  $\beta$  bounded. We note that for every  $k \geq 1$ , the function  $u^k(t, x) = u(2^k t, 2^k x)$  solves  $(-\partial_t^2 + \Delta)u = 2^{2k}f^k$ , where  $f^k(t, x) = f(2^k t, 2^k x)$ . We have

$$\begin{aligned} \|\nabla_{t,x}u^k\|_{L^2L^2([0,T]\times B_1)}^2 + \|r^{-1}u^k\|_{L^2L^2([0,T]\times B_1)}^2 &\lesssim \left| \int_0^T \langle \beta(r)\partial_r u^k, 2^{2k}f^k \rangle dt \right| \\ &\quad + \|\nabla_{t,x}u^k(T)\|_{L_x^2}^2 + \|\nabla_{t,x}u^k(0)\|_{L_x^2}^2, \end{aligned}$$

As  $2^{-k}A_k \subset B_1$ , we also get that

$$\begin{aligned} \|\nabla_{t,x}u^k\|_{L^2L^2([0,T]\times 2^{-k}A_k)}^2 + \|r^{-1}u^k\|_{L^2L^2([0,T]\times 2^{-k}A_k)}^2 &\lesssim \left| \int_0^T \langle \beta(r)\partial_r u^k, 2^{2k}f^k \rangle dt \right| \\ &\quad + \|\nabla_{t,x}u^k(T)\|_{L_x^2}^2 + \|\nabla_{t,x}u^k(0)\|_{L_x^2}^2, \end{aligned}$$

We can also see that

$$\begin{aligned} \|\nabla_{t,x}u^k\|_{L^2L^2([0,T]\times 2^{-k}A_k)}^2 &= 2^{-k(d-1)}\|\nabla_{t,x}u\|_{L^2L^2([0,2^kT]\times A_k)}^2 \\ \|r^{-1}u^k\|_{L^2L^2([0,T]\times 2^{-k}A_k)}^2 &= 2^{-k(d-1)}\|r^{-1}u\|_{L^2L^2([0,2^kT]\times A_k)}^2 \\ \|\nabla_{t,x}u^k(0)\|_{L_x^2}^2 &= 2^{-k(d-2)}\|\nabla_{t,x}u(0)\|_{L_x^2}^2 \\ \|\nabla_{t,x}u^k(T)\|_{L_x^2}^2 &= 2^{-k(d-2)}\|\nabla_{t,x}u(2^kT)\|_{L_x^2}^2 \\ \int_0^T \langle \beta(r)\partial_r u^k, 2^{2k}f^k \rangle dt &= 2^{-k(d-2)} \int_0^{2^kT} \langle \beta\left(\frac{r}{2^k}\right)\partial_r u, f \rangle dt \end{aligned}$$

Therefore,

$$\begin{aligned} 2^{-k}\|\nabla_{t,x}u\|_{L^2L^2([0,2^kT]\times A_k)}^2 + 2^{-k}\|r^{-1}u\|_{L^2L^2([0,2^kT]\times A_k)}^2 &\lesssim \left| \int_0^{2^kT} \langle \beta\left(\frac{r}{2^k}\right)\partial_r u, f \rangle dt \right| \\ &\quad + \|\nabla_{t,x}u(2^kT)\|_{L_x^2}^2 + \|\nabla_{t,x}u(0)\|_{L_x^2}^2 \end{aligned}$$

From the energy identity

$$\|\nabla_{t,x}u^k(2^kT)\|_{L_x^2}^2 = \|\nabla_{t,x}u^k(0)\|_{L_x^2}^2 + 2 \int_0^{2^kT} \langle \partial_t u, f \rangle dt$$

we get that

$$\begin{aligned} 2^{-k}\|\nabla_{t,x}u\|_{L^2L^2([0,2^kT]\times A_k)}^2 + 2^{-k}\|r^{-1}u\|_{L^2L^2([0,2^kT]\times A_k)}^2 &\lesssim \left| \int_0^{2^kT} \langle \beta\left(\frac{r}{2^k}\right)\partial_r u, f \rangle dt \right| \\ &\quad + \|\nabla_{t,x}u(0)\|_{L_x^2}^2 + 2 \int_0^{2^kT} |\langle \partial_t u, f \rangle| dt \\ &\lesssim \int_{\mathbb{R}} \int_{\mathbb{R}^d} (|\nabla_{t,x}u|) |f| dx dt \\ &\quad + \|\nabla_{t,x}u(0)\|_{L_x^2}^2 \end{aligned}$$

It is also clear that we have the same inequality for  $k = 0$ . By taking the supremum with respect to  $k \geq 0$ , we deduce that for every  $\delta > 0$ ,

$$\begin{aligned} \|\nabla_{t,x}u\|_{LE}^2 + \|r^{-1}u\|_{LE}^2 &\lesssim \int_{\mathbb{R}} \int_{\mathbb{R}^d} |\nabla_{t,x}u| |f| dx dt + \|\nabla_{t,x}u(0)\|_{L_x^2}^2 \\ &\lesssim \|\nabla_{t,x}u(0)\|_{L_x^2}^2 + \delta \|\nabla_{t,x}u\|_{LE}^2 + \frac{1}{\delta} \|f\|_{LE^*}^2 \end{aligned}$$

By choosing  $\delta > 0$  sufficiently small, we deduce that

$$\|\nabla_{t,x}u\|_{LE}^2 + \|r^{-1}u\|_{LE}^2 \lesssim \|(u_0, u_1)\|_{\dot{H}^1 \times L^2}^2 + \|f\|_{LE^*}^2$$

We can now immediately see that

$$\|\nabla_{t,x}u\|_{LE} + \|\langle r \rangle^{-1}u\|_{LE} \lesssim \|(u_0, u_1)\|_{\dot{H}^1 \times L^2} + \|f\|_{LE^*}.$$

□

**Corollary 1.2.** *Every solution of (1) satisfies*

$$\|\nabla_{t,x}u\|_{L_t^\infty L_x^2 \cap LE} + \|\langle r \rangle^{-1}u\|_{L_t^\infty L_x^2 \cap LE} \lesssim \|\nabla_{t,x}u(0)\|_{L_x^2} + \|f\|_{L_t^1 L_x^2 + LE^*}.$$

*Proof.* As we have already seen,

$$\|\nabla_{t,x}u\|_{LE}^2 + \|\langle r \rangle^{-1}u\|_{LE}^2 \lesssim \int_{\mathbb{R}} \int_{\mathbb{R}^d} |\nabla_{t,x}u| |f| dx dt + \|\nabla_{t,x}u(0)\|_{L_x^2}^2$$

From the energy identity

$$\|\nabla_{t,x}u(T)\|_{L_x^2}^2 = \|\nabla_{t,x}u(0)\|_{L_x^2}^2 + 2 \int_0^T \langle \partial_t u, f \rangle dt$$

we immediately get that

$$\|\nabla_{t,x}u\|_{L_t^\infty L_x^2}^2 \lesssim \|\nabla_{t,x}u(0)\|_{L_x^2}^2 + \int_{\mathbb{R}} |\partial_t u| |f| dt$$

Thus,

$$\|\nabla_{t,x}u\|_{L_t^\infty L_x^2 \cap LE}^2 + \|\langle r \rangle^{-1}u\|_{LE}^2 \lesssim \int_{\mathbb{R}} \int_{\mathbb{R}^d} |\nabla_{t,x}u| |f| dx dt + \|\nabla_{t,x}u(0)\|_{L_x^2}^2$$

From Hardy's inequality ( $d \geq 3$ ), we get that

$$\|\nabla_{t,x}u\|_{L_t^\infty L_x^2 \cap LE}^2 + \|\langle r \rangle^{-1}u\|_{L_t^\infty L_x^2 \cap LE}^2 \lesssim \int_{\mathbb{R}} \int_{\mathbb{R}^d} |\nabla_{t,x}u| |f| dx dt + \|\nabla_{t,x}u(0)\|_{L_x^2}^2$$

Thus, for every  $\delta > 0$

$$\begin{aligned} \|\nabla_{t,x}u\|_{L_t^\infty L_x^2 \cap LE}^2 + \|\langle r \rangle^{-1}u\|_{L_t^\infty L_x^2 \cap LE}^2 &\lesssim \delta \|\nabla_{t,x}u\|_{L_t^\infty L_x^2 \cap LE}^2 + \frac{1}{\delta} \|f\|_{L_t^1 L_x^2 + LE^*}^2 \\ &\quad + \|\nabla_{t,x}u(0)\|_{L_x^2}^2 \end{aligned}$$

By choosing  $\delta > 0$  small enough, we get that

$$\|\nabla_{t,x}u\|_{L_t^\infty L_x^2 \cap LE}^2 + \|\langle r \rangle^{-1}u\|_{L_t^\infty L_x^2 \cap LE}^2 \lesssim \|\nabla_{t,x}u(0)\|_{L_x^2}^2 + \|f\|_{L_t^1 L_x^2 + LE^*}^2$$

Therefore,

$$\|\nabla_{t,x}u\|_{L_t^\infty L_x^2 \cap LE} + \|\langle r \rangle^{-1}u\|_{L_t^\infty L_x^2 \cap LE} \lesssim \|\nabla_{t,x}u(0)\|_{L_x^2} + \|f\|_{L_t^1 L_x^2 + LE^*}.$$

□



We now consider the perturbed equation

$$(2) \quad \begin{cases} (-\partial_t^2 + L)u = 0 \\ u(0) = u_0, \partial_t u(0) = u_1 \end{cases}$$

where  $L = -\Delta + b^k \partial_k + c$  satisfies the decay condition

$$(3) \quad \sum_{j=0}^{\infty} \sup_{\mathbb{R}_t \times A_j} \langle x \rangle |b| + \langle x \rangle^2 |\partial_l b^l| + \langle x \rangle^2 |c| < K,$$

where  $K > 0$  is a positive constant.

**Corollary 1.3.** *Let  $u$  be a solution of (2). If  $K$  is small enough, then  $u$  satisfies*

$$\|\nabla_{t,x} u\|_{L_t^\infty L_x^2 \cap LE} + \|\langle r \rangle^{-1} u\|_{L_t^\infty L_x^2 \cap LE} \lesssim \|\nabla_{t,x} u(0)\|_{L_x^2} + \|f\|_{L_t^1 L_x^2 + LE^*}.$$

*Proof.* We write  $Bu = b^l \partial_l u + cu$ . We rewrite the equation as  $(-\partial_t^2 + \Delta)u = f - Bu$ . Thus,

$$\begin{aligned} \|\nabla_{t,x} u\|_{L_t^\infty L_x^2 \cap LE} + \|\langle r \rangle^{-1} u\|_{L_t^\infty L_x^2 \cap LE} &\lesssim \|(u_0, u_1)\|_{\dot{H}^1 \times L^2} + \|f - Bu\|_{LE^*} \\ &\lesssim \|(u_0, u_1)\|_{\dot{H}^1 \times L^2} + \|f\|_{LE^*} + \|Bu\|_{LE^*} \end{aligned}$$

For every  $k \geq 0$ , we have

$$\begin{aligned} \|Bu\|_{LE^*} &= \sum_{k=0}^{\infty} 2^{\frac{k}{2}} \|Bu\|_{L^2 L^2(\mathbb{R}_t \times A_k)} = \sum_{k=0}^{\infty} 2^{\frac{k}{2}} \|b^l \partial_l u\|_{L^2 L^2(\mathbb{R}_t \times A_k)} + 2^{\frac{k}{2}} \|cu\|_{L^2 L^2(\mathbb{R}_t \times A_k)} \\ &\lesssim K(\|\nabla_{t,x} u\|_{LE} + \|\langle r \rangle^{-1} u\|_{LE}) \end{aligned}$$

Thus,

$$\begin{aligned} \|\nabla_{t,x} u\|_{L_t^\infty L_x^2 \cap LE} + \|\langle r \rangle^{-1} u\|_{L_t^\infty L_x^2 \cap LE} &\lesssim \|(u_0, u_1)\|_{\dot{H}^1 \times L^2} + \|f - Bu\|_{LE^*} \\ &\lesssim \|(u_0, u_1)\|_{\dot{H}^1 \times L^2} + \|f\|_{LE^*} + \|Bu\|_{LE^*} \\ &\lesssim \|(u_0, u_1)\|_{\dot{H}^1 \times L^2} + \|f\|_{LE^*} \\ &\quad + K(\|\nabla_{t,x} u\|_{LE} + \|\langle r \rangle^{-1} u\|_{LE}) \end{aligned}$$

If  $K > 0$  is sufficiently small, we deduce that

$$\|\nabla_{t,x} u\|_{L_t^\infty L_x^2 \cap LE} + \|\langle r \rangle^{-1} u\|_{L_t^\infty L_x^2 \cap LE} \lesssim \|(u_0, u_1)\|_{\dot{H}^1 \times L^2} + \|f\|_{LE^*},$$

□

**Remark 1.4.** *We note that the integrated local energy decay estimated also takes place under perturbations of the metric (see [Metcalf-Tataru]).*

## 2. STRICHARTZ ESTIMATES

**Definition 2.1.** *A pair  $(p, q)$  is said to be **wave-admissible** in dimension  $d + 1$  if*

$$p \in [2, \infty], \frac{1}{p} + \frac{d-1}{2q} \leq \frac{d-1}{4}, (p, q, d) \neq (2, \infty, 3)$$

We recall the following result:

**Theorem 2.2.** *Let  $u_0, u_1 \in \mathcal{S}(\mathbb{R}^d)$ , and let  $u$  be the solution to (1) with this initial data. Let  $(p, q)$  and  $(\tilde{p}, \tilde{q})$  be pairs of wave-admissible exponents, which also obey the scaling conditions*

$$\frac{1}{p} + \frac{d}{q} = -2 + \frac{1}{\tilde{p}'} + \frac{d}{\tilde{q}'} = -1 + \frac{d}{2},$$

where  $\tilde{p}'$  and  $\tilde{q}'$  are the Lebesgue duals to  $\tilde{p}$  and  $\tilde{q}$ , i.e.  $\frac{1}{\tilde{p}'} + \frac{1}{\tilde{p}} = \frac{1}{\tilde{q}'} + \frac{1}{\tilde{q}} = 1$ , and  $(\hat{p}, \hat{q})$  another pair of wave-admissible exponents satisfying

$$\frac{1}{\hat{p}} + \frac{d}{\hat{q}} = \frac{d}{2}.$$

In addition, we assume that they also satisfy the **non-endpoint** condition  $p, \tilde{p}, \hat{p} > 2$ . Then,

$$\|\nabla_{t,x} u\|_{L^\infty L^2} + \|u\|_{L^p L^q} + \|\nabla_{t,x} u\|_{L^{\tilde{p}} L^{\tilde{q}}} \lesssim_{p,q,\tilde{p},\tilde{q},\hat{p},\hat{q}} \|(u_0, u_1)\|_{\dot{H}^1 \times L^2} + \|f\|_{L^{\tilde{p}'} L^{\tilde{q}'}}$$

We are also going to need the following theorem

**Theorem 2.3.** *Let  $X$  and  $Y$  be Banach spaces, and let  $T : L^p(\mathbb{R}; X) \rightarrow L^q(\mathbb{R}; Y)$  ( $1 \leq p, q \leq \infty$ ) be of the form*

$$Tf(t) = \int_{-\infty}^{\infty} K(t, s) f(s) ds$$

for some kernel  $K : \mathbb{R} \times \mathbb{R} \rightarrow \mathcal{B}(X \rightarrow Y)$ . Provided that  $p < q$ , the truncated operator

$$\tilde{T}f(t) = \int_{-\infty}^t K(t, s) f(s) ds$$

defines a bounded operator from  $L^p(\mathbb{R}; X)$  to  $L^q(\mathbb{R}; Y)$ .

We are now going to prove the following result:

**Theorem 2.4.** (Rodnianski-Schlag) *We assume that the coefficients of (2) satisfy (3) for some  $K > 0$  (in particular,  $K$  can be large). We also assume that (ILED) holds for (2).*

*Let  $u_0, u_1 \in \mathcal{S}(\mathbb{R}^d)$ , and let  $u$  be the solution to (2) with this initial data. Let  $(p, q)$  and  $(\hat{p}, \hat{q})$ , be pairs of wave-admissible exponents, which also obey the scaling conditions*

$$\frac{1}{p} + \frac{d}{q} = \frac{1}{\hat{p}} + \frac{d}{\hat{q}} - 1 = -1 + \frac{d}{2}.$$

*In addition, we assume that they also satisfy the **non-endpoint** condition  $p, \hat{p} > 2$ . Then,*

$$\|\nabla_{t,x} u\|_{L^\infty L^2} + \|u\|_{L^p L^q} + \|\nabla_{t,x} u\|_{L^{\tilde{p}} L^{\tilde{q}}} \lesssim_{p,q,\tilde{p},\tilde{q}} \|(u_0, u_1)\|_{\dot{H}^1 \times L^2} + \|f\|_{L^{-1} L_x^2 + L E^*}$$

*Proof.* We recall some results concerning the linear wave equation. The first one says that the solution of the homogeneous problem

$$\begin{aligned} (-\partial_t^2 + \Delta)u &= 0 \\ u(0) &= u_0, \\ \partial_t u(0) &= u_1 \end{aligned}$$

is given by the formula

$$u(t) = \cos(t\sqrt{-\Delta})u_0 + \frac{\sin(t\sqrt{-\Delta})}{\sqrt{-\Delta}}u_1$$

Besides this, the purely inhomogeneous problem

$$\begin{aligned} (-\partial_t^2 + \Delta)u &= F \\ u(0) &= 0, \\ \partial_t u(0) &= 0 \end{aligned}$$

has a solution given by the formula

$$u(t) = - \int_{-\infty}^t \frac{\sin(t-s)\sqrt{-\Delta}}{\sqrt{-\Delta}} F(s) ds$$

We also note that when  $u$  is a solution of the homogeneous problem,

$$\|\nabla_{t,x} u\|_{L_t^\infty L_x^2 \cap LE} \lesssim \|(u_0, u_1)\|_{\dot{H}_x^1 \times L_x^2}$$

By duality, we get that

$$\begin{aligned} \left\| \int_{-\infty}^{\infty} \cos(t\sqrt{-\Delta}) F(t) dt \right\|_{L_x^2} &\lesssim \|F\|_{L_t^1 L_x^2 + LE^*} \\ \left\| \int_{-\infty}^{\infty} \sin(t\sqrt{-\Delta}) F(t) dt \right\|_{L_x^2} &\lesssim \|F\|_{L_t^1 L_x^2 + LE^*} \end{aligned}$$

By applying the homogeneous Strichartz estimate, we deduce that

$$\begin{aligned} &\left\| \frac{\sin(t\sqrt{-\Delta})}{\sqrt{-\Delta}} \int_{-\infty}^{\infty} \cos(s\sqrt{-\Delta}) F(s) ds \right\|_{L_t^p L_x^q} + \left\| \nabla_{t,x} \frac{\sin(t\sqrt{-\Delta})}{\sqrt{-\Delta}} \int_{-\infty}^{\infty} \cos(s\sqrt{-\Delta}) F(s) ds \right\|_{L_t^{\hat{p}} L_x^{\hat{q}}} \\ &\lesssim \left\| \int_{-\infty}^{\infty} \cos(s\sqrt{-\Delta}) F(s) ds \right\|_{L_x^2} \\ &\left\| \cos(t\sqrt{-\Delta}) \int_{-\infty}^{\infty} \frac{\sin(s\sqrt{-\Delta})}{\sqrt{-\Delta}} F(s) dt \right\|_{L_t^p L_x^q} + \left\| \nabla_{t,x} \cos(t\sqrt{-\Delta}) \int_{-\infty}^{\infty} \frac{\sin(s\sqrt{-\Delta})}{\sqrt{-\Delta}} F(s) dt \right\|_{L_t^{\hat{p}} L_x^{\hat{q}}} \\ &\lesssim \left\| \int_{-\infty}^{\infty} \sin(s\sqrt{-\Delta}) F(s) ds \right\|_{L_x^2} \end{aligned}$$

Thus,

$$\left\| \int_{-\infty}^{\infty} \frac{\sin(t-s)\sqrt{-\Delta}}{\sqrt{-\Delta}} F(s) ds \right\|_{L_t^p L_x^q} + \left\| \nabla_{t,x} \int_{-\infty}^{\infty} \frac{\sin(t-s)\sqrt{-\Delta}}{\sqrt{-\Delta}} F(s) ds \right\|_{L_t^{\hat{p}} L_x^{\hat{q}}} \lesssim \|F\|_{L_t^1 L_x^2 + LE^*}$$

As  $p > 2$ , the Christ-Kiselev lemma enables us to deduce that

$$\left\| \int_{-\infty}^t \frac{\sin(t-s)\sqrt{-\Delta}}{\sqrt{-\Delta}} F(s) ds \right\|_{L_t^p L_x^q} + \left\| \nabla_{t,x} \int_{-\infty}^t \frac{\sin(t-s)\sqrt{-\Delta}}{\sqrt{-\Delta}} F(s) ds \right\|_{L_t^{\hat{p}} L_x^{\hat{q}}} \lesssim \|F\|_{L_t^1 L_x^2 + LE^*}$$

Along with the homogeneous Strichartz estimate, this immediately implies that

$$\|u\|_{L_t^p L_x^q} + \|\nabla_{t,x} u\|_{L_t^{\hat{p}} L_x^{\hat{q}}} \lesssim \|(u_0, u_1)\|_{\dot{H}_x^1 \times L_x^2} + \|F\|_{L_t^1 L_x^2 + LE^*}$$

We now return to the problem

$$\begin{aligned} (-\partial_t^2 - L)u &= f \\ u(0) &= u_0, \\ \partial_t u(0) &= u_1 \end{aligned}$$

This can be rewritten as

$$\begin{aligned} (-\partial_t^2 + \Delta)u &= f + Bu \\ u(0) &= u_0, \\ \partial_t u(0) &= u_1 \end{aligned}$$

As in the proof of Corollary 1.3, we have

$$\|Bu\|_{LE^*} \lesssim K(\|\nabla_{t,x} u\|_{LE} + \|\langle r \rangle^{-1} u\|_{LE}) \lesssim \|(u_0, u_1)\|_{\dot{H}_x^1 \times L_x^2} + \|f\|_{L_t^1 L_x^2 + LE^*}$$

We deduce that

$$\begin{aligned}
\|u\|_{L^p L^q} + \|\nabla_{t,x} u\|_{L_t^{\hat{p}} L_x^{\hat{q}}} &\lesssim \|(u_0, u_1)\|_{\dot{H}_x^1 \times L_x^2} + \|f + Bu\|_{L_t^1 L_x^2 + LE^*} \\
&\lesssim \|(u_0, u_1)\|_{\dot{H}_x^1 \times L_x^2} + \|f\|_{L_t^1 L_x^2 + LE^*} + \|Bu\|_{L_t^1 L_x^2 + LE^*} \\
&\lesssim \|(u_0, u_1)\|_{\dot{H}_x^1 \times L_x^2} + \|f\|_{L_t^1 L_x^2 + LE^*}
\end{aligned}$$

This finishes the proof.  $\square$

**Remark 2.5.** We note that we can further generalize the previous Strichartz estimates to general non-endpoint wave-admissible pairs  $(\tilde{p}, \tilde{q})$  satisfying the scaling condition  $\frac{1}{\tilde{p}} + \frac{d}{\tilde{q}} = \frac{d}{2}$  if we replace the integrated local energy decay estimate by the condition

$$\|\nabla_{t,x} u\|_{LE} + \|\langle r \rangle^{-1} u\|_{LE} \lesssim \|(u_0, u_1)\|_{\dot{H}_x^1 \times L_x^2} + \|f\|_{L_t^{\tilde{p}'} L_x^{\tilde{q}'} + LE^*}.$$

**Remark 2.6.** For the endpoint case see the argument of [Keel-Tao] or [Ionescu-Kenig].

**Remark 2.7.** If  $(b, c)$  don't satisfy (3), then the previous result can fail in some instances. For example, if  $c = \alpha r^{-2}$ , then the wave admissible pairs  $(p, q)$ ,  $(\hat{p}, \hat{q})$ , and  $(\tilde{p}, \tilde{q})$  for which Strichartz estimates hold depend sensitively on  $\alpha$  (see [Burq Planchon Stalker Tahvildar-Zadeh]).

### 3. A SPECTRAL CHARACTERIZATION

We assume that  $L = -\Delta + b^k \partial_k + c$ , where  $b$  is purely imaginary and  $c$  is real, and both of them are time-independent and satisfy the decay condition (3) for some (possibly large)  $K \in (0, \infty)$ . These conditions imply that  $L$  is self-adjoint. It is clear that  $\langle Lu, u \rangle$  is bounded on  $\dot{H}^1$ . We also assume that  $L$  is coercive, in the sense that  $\langle Lu, u \rangle \gtrsim \|u\|_{\dot{H}^1}^2$ . We define the wave *resolvents* of  $L$  as

$$\mathbf{R}_z = (z^2 - L)^{-1},$$

which are well-defined as bounded operators from  $\dot{H}^1$  to  $D(L)$  as long as  $z^2 \notin \sigma(L)$ . The spectral theorem implies the following bound for  $\mathbf{R}_z$ :

**Lemma 3.1.** For any  $\tau \in \mathbb{R}$  and  $\varepsilon > 0$ , we have

$$|\tau| \|\mathbf{R}_{\tau \pm i\varepsilon} g\|_{L^2} \lesssim \varepsilon^{-1} \|g\|_{L^2}$$

*Proof.* From the spectral theorem, we have (we keep in mind that  $\sigma(L) \subset \mathbb{R}$ ) (for every  $\tau \neq 0$ )

$$\begin{aligned}
\|\mathbf{R}_{\tau \pm i\varepsilon} g\|_{L^2} &\lesssim d((\tau \pm i\varepsilon)^2, \sigma(L))^{-1} \|g\|_{L^2} = d(\tau^2 - \varepsilon^2 \pm 2\tau i\varepsilon, \sigma(L))^{-1} \|g\|_{L^2} \\
&\lesssim |\tau|^{-1} \varepsilon^{-1} \|g\|_{L^2}
\end{aligned}$$

Thus,

$$|\tau| \|\mathbf{R}_{\tau \pm i\varepsilon} g\|_{L^2} \lesssim \varepsilon^{-1} \|g\|_{L^2}$$

Here we have used the bound

$$d((\tau \pm i\varepsilon)^2, \sigma(L)) \geq d((\tau \pm i\varepsilon)^2, \mathbb{R}) = d(\tau^2 - \varepsilon^2 \pm 2\tau i\varepsilon, \mathbb{R}) = 2|\tau|\varepsilon,$$

which follows from the fact that  $\sigma(L) \subset \mathbb{R}$ .

The same conclusion is clearly true for  $\tau = 0$ .  $\square$

**Lemma 3.2.** For any  $\tau \in \mathbb{R}$  and  $\varepsilon > 0$ , we have

$$\|\nabla_x \mathbf{R}_{\tau \pm i\varepsilon} g\|_{L^2} + \varepsilon \|\mathbf{R}_{\tau \pm i\varepsilon} g\|_{L^2} \lesssim \varepsilon^{-1} \|g\|_{L^2}$$

*Proof.* Let  $u = \mathbf{R}_{\tau \pm i\varepsilon} g$ . By definition,  $((\tau \pm i\varepsilon)^2 - L)u = g$ . We have

$$\Re \langle u, g \rangle = \Re \langle u, (\tau^2 - \varepsilon^2)u \pm 2i\tau\varepsilon u - Lu \rangle = \langle u, (\tau^2 - \varepsilon^2)u - Lu \rangle$$

Thus, for every  $\delta > 0$ ,

$$\begin{aligned} \langle Lu, u \rangle + \varepsilon^2 \|u\|_{L^2}^2 &= \tau^2 \|u\|_{L^2}^2 - \Re \langle u, g \rangle \lesssim \tau^2 \|u\|_{L^2}^2 + |\langle u, g \rangle| \\ &\lesssim \tau^2 \|u\|_{L^2}^2 + \|u\|_{L^2} \|g\|_{L^2} \lesssim \varepsilon^{-2} \|g\|_{L^2}^2 + \delta^{-2} \varepsilon^{-2} \|g\|_{L^2}^2 + \delta^2 \varepsilon^2 \|u\|_{L^2}^2 \end{aligned}$$

By choosing  $\delta > 0$  sufficiently small, we deduce that

$$\langle Lu, u \rangle + \varepsilon^2 \|u\|_{L^2}^2 \lesssim \varepsilon^{-2} \|g\|_{L^2}^2$$

As  $\langle Lu, u \rangle \gtrsim \|u\|_{\dot{H}^1}^2$ , we get that

$$\|u\|_{\dot{H}^1}^2 + \varepsilon^2 \|u\|_{L^2}^2 \lesssim \varepsilon^{-2} \|g\|_{L^2}^2,$$

hence

$$\|u\|_{\dot{H}^1} + \varepsilon \|u\|_{L^2} \lesssim \varepsilon^{-1} \|g\|_{L^2},$$

The conclusion immediately follows.  $\square$

We are also going to define the spatial counterparts of  $LE$  and  $LE^*$ :

$$\begin{aligned} \|u\|_{\mathcal{LE}} &= \sup_{j \geq 0} \|\langle r \rangle^{-\frac{1}{2}} u\|_{L^2(A_j)} \\ \|f\|_{\mathcal{LE}^*} &= \sum_{j \geq 0} \|\langle r \rangle^{\frac{1}{2}} f\|_{L^2(A_j)} \end{aligned}$$

We shall prove the following result (for further reference, see [Metcalf-Sterbenz-Tataru]):

**Theorem 3.3.** *The following are equivalent:*

- (1) *Every solution  $u$  to (2) satisfies ILED*
- (2) *For every  $\tau \in \mathbb{R}$  and  $\varepsilon > 0$ , we have*

$$|\tau \pm i\varepsilon| \|\mathbf{R}_{\tau \pm i\varepsilon} g\|_{\mathcal{LE}} + \|\nabla_x \mathbf{R}_{\tau \pm i\varepsilon} g\|_{\mathcal{LE}} + \|\langle r \rangle^{-1} \mathbf{R}_{\tau \pm i\varepsilon} g\|_{\mathcal{LE}} \leq C \|g\|_{\mathcal{LE}^*},$$

where  $C > 0$  is an universal constant.

*Proof. Step 1: Reduction to forward solutions.* We first prove that every solution of  $(-\partial_t^2 - L)u = f$  with  $u(0) = u_0$  and  $u_t(0) = u_1$  satisfies

$$\|\nabla u\|_{LE} + \|\langle r \rangle^{-1} u\|_{LE} \lesssim \|(u_0, u_1)\|_{\dot{H}^1 \times L^2} + \|f\|_{LE^*}$$

if and only if every solution of the forward problem  $(-\partial_t^2 - L)u = f$ , with  $f$  supported away from  $\{t = -\infty\}$  satisfies

$$\|\nabla u\|_{LE} + \|\langle r \rangle^{-1} u\|_{LE} \lesssim \|f\|_{LE^*}$$

When  $f$  is supported away from  $\{t = -\infty\}$ , and  $u(t) = 0$  for  $t$  sufficiently negative, we deduce that there exists  $t_0$  such that  $u(t) = 0, \forall t \leq t_0$ . Thus,  $\partial_t u(t_0) = 0$ . By applying ILED to  $\tilde{u}(t) = u(t + t_0)$  (which satisfies  $\tilde{u}(0) = 0, \partial_t \tilde{u}(0) = 0$ ), and by using the time-invariance of the  $LE$  and  $LE^*$  norms, we deduce that

$$\|\nabla u\|_{LE} + \|\langle r \rangle^{-1} u\|_{LE} \lesssim \|f\|_{LE^*}$$

For the converse, we use a method that is similar to the one employed by Rodnianski and Schlag. We consider  $v$  to be the solution of the problem

$$\begin{aligned} (-\partial_t^2 + \Delta)v &= f \\ v(0) &= u_0 \\ \partial_t v(0) &= u_1 \end{aligned}$$

By 1.2, we have

$$\|\nabla_{t,x}v\|_{LE} + \|\langle r \rangle^{-1}v\|_{LE} \lesssim \|(u_0, u_1)\|_{\dot{H}_x^1 \times L_x^2} + \|f\|_{LE^*}$$

We can now immediately see that  $u - v$  satisfies

$$\begin{aligned} (-\partial_t^2 + L)(u - v) &= -Bv \\ (u - v)(0) &= 0 \\ \partial_t(u - v)(0) &= 0 \end{aligned}$$

Let  $v_+$  be the forward solution of

$$\begin{aligned} (-\partial_t^2 + L)(v_+) &= \mathbf{1}_{[0,\infty)}(-Bv) \\ (v_+)(0) &= 0 \\ \partial_t(v_+)(0) &= 0 \end{aligned}$$

and  $v_-$  be the backward solution of

$$\begin{aligned} (-\partial_t^2 + L)(v_-) &= \mathbf{1}_{(-\infty,0)}(-Bv) \\ (v_-)(0) &= 0 \\ \partial_t(v_-)(0) &= 0 \end{aligned}$$

Thus,

$$\|\nabla_{t,x}v_+\|_{LE} + \|\langle r \rangle^{-1}v_+\|_{LE} \lesssim \|\mathbf{1}_{[0,\infty)}(-Bv)\|_{LE^*} \lesssim \|Bv\|_{LE^*},$$

and by using the time symmetry of the  $LE$  and  $LE^*$  norms,

$$\|\nabla_{t,x}v_-\|_{LE} + \|\langle r \rangle^{-1}v_-\|_{LE} \lesssim \|\mathbf{1}_{(-\infty,0)}(-Bv)\|_{LE^*} \lesssim \|Bv\|_{LE^*},$$

As in the proof of Corollary 1.3, we get that

$$\|Bv\|_{LE^*} \lesssim K(\|\nabla_{t,x}v\|_{LE} + \|\langle r \rangle^{-1}v\|_{LE}) \lesssim \|(u_0, u_1)\|_{\dot{H}_x^1 \times L_x^2} + \|f\|_{LE^*}$$

We note that  $u - v - v_+ - v_-$  is a finite energy solution of

$$\begin{aligned} (-\partial_t^2 + L)(u - v - v_+ - v_-) &= 0 \\ (u - v - v_+ - v_-)(0) &= 0 \\ \partial_t(u - v - v_+ - v_-)(0) &= 0 \end{aligned}$$

Thus,  $u = v + v_+ + v_-$ . This means that

$$\begin{aligned} \|\nabla_{t,x}u\|_{LE} + \|\langle r \rangle^{-1}u\|_{LE} &\lesssim \|\nabla_{t,x}v\|_{LE} + \|\langle r \rangle^{-1}v\|_{LE} + \|\nabla_{t,x}v_+\|_{LE} + \|\langle r \rangle^{-1}v_+\|_{LE} \\ &\quad + \|\nabla_{t,x}v_-\|_{LE} + \|\langle r \rangle^{-1}v_-\|_{LE} \\ &\lesssim \|(u_0, u_1)\|_{\dot{H}_x^1 \times L_x^2} + \|f\|_{LE^*} + \|Bv\|_{LE^*} \\ &\lesssim \|(u_0, u_1)\|_{\dot{H}_x^1 \times L_x^2} + \|f\|_{LE^*}, \end{aligned}$$

as desired.

*Step 2: Reduction to damped forward solutions.* Now we prove that the condition

$$\|\nabla_{t,x}u\|_{LE} + \|\langle r \rangle^{-1}u\|_{LE} \lesssim \|f\|_{LE^*}$$

for forward solutions is equivalent to the condition

$$\|e^{-\varepsilon t}\nabla_{t,x}u\|_{LE} + \|\langle r \rangle^{-1}e^{-\varepsilon t}u\|_{LE} \lesssim \|e^{-\varepsilon t}f\|_{LE^*}, \forall \varepsilon > 0$$

We first show that the latter implies the former.

Let  $u_k$  be the forward solution corresponding to  $\mathbf{1}_{[\frac{k}{\varepsilon}, \frac{k+1}{\varepsilon})}(t)f$ . As the operator  $L$  is elliptic, we have

$$u_k(t) = - \int_{-\infty}^t \frac{\sin(t-s)L}{L} (\mathbf{1}_{[\frac{k}{\varepsilon}, \frac{k+1}{\varepsilon})} f)(s) ds,$$

which shows that  $u_k$  is supported in  $\{t \in [k, \infty)\}$ . We also have  $u = \sum_{k \in \mathbb{Z}} u_k$ .

We note that

$$\begin{aligned} \|e^{-\varepsilon t} \nabla_{t,x} u\|_{LE} &\simeq \sup_{j \geq 0} 2^{-\frac{j}{2}} \|e^{-\varepsilon t} \nabla_{t,x} u\|_{L^2 L^2(\mathbb{R} \times A_j)} \\ &\simeq \sup_{j \geq 0} \left( \sum_{k \in \mathbb{Z}} 2^{-j} \|e^{-\varepsilon t} \mathbf{1}_{[\frac{k}{\varepsilon}, \frac{k+1}{\varepsilon})}(t) \nabla_{t,x} u\|_{L^2 L^2(\mathbb{R} \times A_j)}^2 \right)^{\frac{1}{2}} \\ &\simeq \sup_{j \geq 0} \left( \sum_{k \in \mathbb{Z}} 2^{-j} e^{-2k} \|\mathbf{1}_{[\frac{k}{\varepsilon}, \frac{k+1}{\varepsilon})}(t) \nabla_{t,x} u\|_{L^2 L^2(\mathbb{R} \times A_j)}^2 \right)^{\frac{1}{2}} \\ &\simeq \|2^{-\frac{j}{2}} e^{-k} \|\mathbf{1}_{[\frac{k}{\varepsilon}, \frac{k+1}{\varepsilon})}(t) \nabla_{t,x} u\|_{L^2 L^2(\mathbb{R} \times A_j)}\|_{l_j^\infty l_k^2} \\ &\lesssim \|2^{-\frac{j}{2}} e^{-k} \|\mathbf{1}_{[\frac{k}{\varepsilon}, \frac{k+1}{\varepsilon})}(t) \nabla_{t,x} u\|_{L^2 L^2(\mathbb{R} \times A_j)}\|_{l_k^2 l_j^\infty} \end{aligned}$$

Similarly,

$$\begin{aligned} \|\langle r \rangle^{-1} e^{-\varepsilon t} u\|_{LE} &\simeq \sup_{j \geq 0} 2^{-\frac{3j}{2}} \|e^{-\varepsilon t} u\|_{L^2 L^2(\mathbb{R} \times A_j)} \\ &\simeq \|2^{-\frac{3j}{2}} e^{-k} \|\mathbf{1}_{[\frac{k}{\varepsilon}, \frac{k+1}{\varepsilon})}(t) u\|_{L^2 L^2(\mathbb{R} \times A_j)}\|_{l_j^\infty l_k^2} \\ &\lesssim \|2^{-\frac{3j}{2}} e^{-k} \|\mathbf{1}_{[\frac{k}{\varepsilon}, \frac{k+1}{\varepsilon})}(t) u\|_{L^2 L^2(\mathbb{R} \times A_j)}\|_{l_k^2 l_j^\infty} \end{aligned}$$

We also note that

$$\begin{aligned} \|f\|_{LE^*} &\simeq \sum_{l \geq 0} 2^{\frac{l}{2}} \|f\|_{L^2 L^2(\mathbb{R} \times A_l)} \simeq \sum_{l \geq 0} 2^{\frac{l}{2}} \left( \|\mathbf{1}_{[\frac{k}{\varepsilon}, \frac{k+1}{\varepsilon})}(t) f\|_{L^2 L^2(\mathbb{R} \times A_l)}^2 \right)^{\frac{1}{2}} \\ &\simeq \|2^{\frac{l}{2}} \|\mathbf{1}_{[\frac{k}{\varepsilon}, \frac{k+1}{\varepsilon})}(t) f\|_{L^2 L^2(\mathbb{R} \times A_l)}\|_{l_l^2 l_k^1} \gtrsim \|2^{\frac{l}{2}} \|\mathbf{1}_{[\frac{k}{\varepsilon}, \frac{k+1}{\varepsilon})}(t) f\|_{L^2 L^2(\mathbb{R} \times A_l)}\|_{l_k^2 l_l^1} \end{aligned}$$

This means that it suffices to prove the inequality

$$\begin{aligned} \|2^{-\frac{j}{2}} e^{-k} \|\mathbf{1}_{[\frac{k}{\varepsilon}, \frac{k+1}{\varepsilon})}(t) \nabla_{t,x} u\|_{L^2 L^2(\mathbb{R} \times A_j)}\|_{l_k^2 l_j^\infty} + \|2^{-\frac{3j}{2}} e^{-k} \|\mathbf{1}_{[\frac{k}{\varepsilon}, \frac{k+1}{\varepsilon})}(t) u\|_{L^2 L^2(\mathbb{R} \times A_j)}\|_{l_k^2 l_j^\infty} \\ \lesssim \|2^{\frac{l}{2}} \|\mathbf{1}_{[\frac{k}{\varepsilon}, \frac{k+1}{\varepsilon})}(t) f\|_{L^2 L^2(\mathbb{R} \times A_l)}\|_{l_k^2 l_l^1} \end{aligned}$$

As  $u_k$  is supported in  $[k, \infty)$ , for every  $j \geq 0$  we have

$$\begin{aligned} \|\mathbf{1}_{[\frac{k}{\varepsilon}, \frac{k+1}{\varepsilon})}(t) \nabla_{t,x} u\|_{L^2 L^2(\mathbb{R} \times A_j)} &= \|\mathbf{1}_{[\frac{k}{\varepsilon}, \frac{k+1}{\varepsilon})}(t) \sum_{k' \in \mathbb{Z}} \nabla_{t,x} u_{k'}\|_{L^2 L^2(\mathbb{R} \times A_j)} \\ &= \|\mathbf{1}_{[\frac{k}{\varepsilon}, \frac{k+1}{\varepsilon})}(t) \sum_{k' \leq k} \nabla_{t,x} u_{k'}\|_{L^2 L^2(\mathbb{R} \times A_j)} \\ \|\mathbf{1}_{[\frac{k}{\varepsilon}, \frac{k+1}{\varepsilon})}(t) u\|_{L^2 L^2(\mathbb{R} \times A_j)} &= \|\mathbf{1}_{[\frac{k}{\varepsilon}, \frac{k+1}{\varepsilon})}(t) \sum_{k' \in \mathbb{Z}} u_{k'}\|_{L^2 L^2(\mathbb{R} \times A_j)} \\ &= \|\mathbf{1}_{[\frac{k}{\varepsilon}, \frac{k+1}{\varepsilon})}(t) \sum_{k' \leq k} u_{k'}\|_{L^2 L^2(\mathbb{R} \times A_j)} \end{aligned}$$

We now note that

$$\begin{aligned}
& \|2^{-\frac{j}{2}} e^{-k} \mathbf{1}_{[\frac{k}{\varepsilon}, \frac{k+1}{\varepsilon})}(t) \nabla_{t,x} u\|_{L^2 L^2(\mathbb{R} \times A_j)}\|_{l_k^2 l_j^\infty} = \left( \sum_{k \in \mathbb{Z}} e^{-2k} \sup_{j \geq 0} 2^{-j} \|\mathbf{1}_{[\frac{k}{\varepsilon}, \frac{k+1}{\varepsilon})}(t) \nabla_{t,x} u\|_{L^2 L^2(\mathbb{R} \times A_j)}^2 \right)^{\frac{1}{2}} \\
& = \left( \sum_{k \in \mathbb{Z}} e^{-2k} \sup_{j \geq 0} 2^{-j} \|\mathbf{1}_{[\frac{k}{\varepsilon}, \frac{k+1}{\varepsilon})}(t) \sum_{k' \leq k} \nabla_{t,x} u_{k'}\|_{L^2 L^2(\mathbb{R} \times A_j)}^2 \right)^{\frac{1}{2}} \\
& \lesssim \left( \sum_{k \in \mathbb{Z}} e^{-2k} \left( \sum_{k' \leq k} \sup_{j \geq 0} 2^{-\frac{j}{2}} \|\mathbf{1}_{[\frac{k'}{\varepsilon}, \frac{k'+1}{\varepsilon})}(t) \nabla_{t,x} u_{k'}\|_{L^2 L^2(\mathbb{R} \times A_j)} \right)^2 \right)^{\frac{1}{2}} \\
& \lesssim \left( \sum_{k \in \mathbb{Z}} e^{-2k} \left( \sum_{k' \leq k} \sup_{j \geq 0} 2^{-\frac{j}{2}} \|\nabla_{t,x} u_{k'}\|_{L^2 L^2(\mathbb{R} \times A_j)} \right)^2 \right)^{\frac{1}{2}} \\
& \lesssim \left( \sum_{k \in \mathbb{Z}} e^{-2k} \left( \sum_{k' \leq k} \sum_{l \geq 0} 2^{\frac{l}{2}} \|\mathbf{1}_{[\frac{k'}{\varepsilon}, \frac{k'+1}{\varepsilon})}(t) f\|_{L^2 L^2(\mathbb{R} \times A_l)} \right)^2 \right)^{\frac{1}{2}} \\
& \lesssim \left( \sum_{k \in \mathbb{Z}} \left( \sum_{k' \leq k} e^{-(k-k')} \sum_{l \geq 0} 2^{\frac{l}{2}} \|e^{-\varepsilon t} \mathbf{1}_{[\frac{k'}{\varepsilon}, \frac{k'+1}{\varepsilon})}(t) f\|_{L^2 L^2(\mathbb{R} \times A_l)} \right)^2 \right)^{\frac{1}{2}}
\end{aligned}$$

Similarly,

$$\begin{aligned}
& \|2^{-\frac{3j}{2}} e^{-k} \mathbf{1}_{[\frac{k}{\varepsilon}, \frac{k+1}{\varepsilon})}(t) u\|_{L^2 L^2(\mathbb{R} \times A_j)}\|_{l_k^2 l_j^\infty} \\
& \lesssim \left( \sum_{k \in \mathbb{Z}} \left( \sum_{k' \leq k} e^{-(k-k')} \sum_{l \geq 0} 2^{\frac{l}{2}} \|e^{-\varepsilon t} \mathbf{1}_{[\frac{k'}{\varepsilon}, \frac{k'+1}{\varepsilon})}(t) f\|_{L^2 L^2(\mathbb{R} \times A_l)} \right)^2 \right)^{\frac{1}{2}}
\end{aligned}$$

Thus,

$$\begin{aligned}
& \|2^{-\frac{j}{2}} e^{-k} \mathbf{1}_{[\frac{k}{\varepsilon}, \frac{k+1}{\varepsilon})}(t) \nabla_{t,x} u\|_{L^2 L^2(\mathbb{R} \times A_j)}\|_{l_k^2 l_j^\infty} + \|2^{-\frac{3j}{2}} e^{-k} \mathbf{1}_{[\frac{k}{\varepsilon}, \frac{k+1}{\varepsilon})}(t) u\|_{L^2 L^2(\mathbb{R} \times A_j)}\|_{l_k^2 l_j^\infty} \\
& \lesssim \left( \sum_{k \in \mathbb{Z}} \left( \sum_{k' \leq k} e^{-(k-k')} \sum_{l \geq 0} 2^{\frac{l}{2}} \|e^{-\varepsilon t} \mathbf{1}_{[\frac{k'}{\varepsilon}, \frac{k'+1}{\varepsilon})}(t) f\|_{L^2 L^2(\mathbb{R} \times A_l)} \right)^2 \right)^{\frac{1}{2}}
\end{aligned}$$



However, the sequence  $(c_k)_{k \geq 0}$  given by  $c_k = e^{-k}$  is  $l_k^1$ , and  $\|c_k\|_{l_k^1} = (1 - e^{-1})^{-1}$ . By Young's inequality, we get

$$\begin{aligned} & \|2^{-\frac{j}{2}} e^{-k} \mathbf{1}_{[\frac{k}{\varepsilon}, \frac{k+1}{\varepsilon})}(t) \nabla_{t,x} u\|_{L^2 L^2(\mathbb{R} \times A_j)} \|l_k^2 l_j^\infty\| + \|2^{-\frac{3j}{2}} e^{-k} \mathbf{1}_{[\frac{k}{\varepsilon}, \frac{k+1}{\varepsilon})}(t) u\|_{L^2 L^2(\mathbb{R} \times A_j)} \|l_k^2 l_j^\infty\| \\ & \lesssim \left( \sum_{k \in \mathbb{Z}} \left( \sum_{k' \leq k} e^{-(k-k')} \sum_{l \geq 0} 2^{\frac{l}{2}} \|e^{-\varepsilon t} \mathbf{1}_{[\frac{k'}{\varepsilon}, \frac{k'+1}{\varepsilon})}(t) f\|_{L^2 L^2(\mathbb{R} \times A_l)} \right)^2 \right)^{\frac{1}{2}} \\ & \lesssim \|c_k\|_{l_k^1} \left( \sum_{k \in \mathbb{Z}} \left( \sum_{l \geq 0} 2^{\frac{l}{2}} \|e^{-\varepsilon t} \mathbf{1}_{[\frac{k}{\varepsilon}, \frac{k+1}{\varepsilon})}(t) f\|_{L^2 L^2(\mathbb{R} \times A_l)} \right)^2 \right)^{\frac{1}{2}} \\ & \lesssim \left( \sum_{k \in \mathbb{Z}} \left( \sum_{l \geq 0} 2^{\frac{l}{2}} \|e^{-\varepsilon t} \mathbf{1}_{[\frac{k}{\varepsilon}, \frac{k+1}{\varepsilon})}(t) f\|_{L^2 L^2(\mathbb{R} \times A_l)} \right)^2 \right)^{\frac{1}{2}} \end{aligned}$$

Along with the previous discussion, this implies that

$$\|e^{-\varepsilon t} \nabla_{t,x} u\|_{LE} + \|\langle r \rangle^{-1} e^{-\varepsilon t} u\|_{LE} \lesssim \|e^{-\varepsilon t} f\|_{LE^*}, \forall \varepsilon > 0$$

We now prove the converse. As  $f$  is supported away from  $\{t = -\infty\}$ , there exists  $m \in \mathbb{Z}$  such that  $\text{supp } f \subseteq \{t \in [m, \infty)\}$ . For every  $\varepsilon > 0$ , and  $j \geq 0$ , we have

$$2^{-\frac{j}{2}} \|e^{-\varepsilon t} \nabla_{t,x} u\|_{L^2 L^2(\mathbb{R} \times A_j)} + 2^{-\frac{j}{2}} \|\langle r \rangle^{-1} e^{-\varepsilon t} u\|_{L^2 L^2(\mathbb{R} \times A_j)} \lesssim \|e^{-\varepsilon t} f\|_{LE^*} \lesssim e^{-\varepsilon m} \|f\|_{LE^*}$$

For every  $n \in \mathbb{Z}$ , we get that

$$\begin{aligned} & 2^{-\frac{j}{2}} e^{-\varepsilon n} \|\nabla_{t,x} u\|_{L^2 L^2((-\infty, n] \times A_j)} + 2^{-\frac{j}{2}} e^{-\varepsilon n} \|\langle r \rangle^{-1} u\|_{L^2 L^2((-\infty, n] \times A_j)} \\ & \lesssim 2^{-\frac{j}{2}} \|e^{-\varepsilon t} \nabla_{t,x} u\|_{L^2 L^2((-\infty, n] \times A_j)} + 2^{-\frac{j}{2}} \|\langle r \rangle^{-1} e^{-\varepsilon t} u\|_{L^2 L^2((-\infty, n] \times A_j)} \\ & \lesssim 2^{-\frac{j}{2}} \|e^{-\varepsilon t} \nabla_{t,x} u\|_{L^2 L^2(\mathbb{R} \times A_j)} + 2^{-\frac{j}{2}} \|\langle r \rangle^{-1} e^{-\varepsilon t} u\|_{L^2 L^2(\mathbb{R} \times A_j)} \\ & \lesssim e^{-\varepsilon m} \|f\|_{LE^*} \end{aligned}$$

Thus,

$$\begin{aligned} & 2^{-\frac{j}{2}} e^{-\varepsilon n} \|\nabla_{t,x} u\|_{L^2 L^2((-\infty, n] \times A_j)} + 2^{-\frac{j}{2}} e^{-\varepsilon n} \|\langle r \rangle^{-1} u\|_{L^2 L^2((-\infty, n] \times A_j)} \\ & \lesssim e^{-\varepsilon m} \|f\|_{LE^*} \end{aligned}$$

We now take the limit at 0 with respect to  $\varepsilon$ , and we get that

$$2^{-\frac{j}{2}} \|\nabla_{t,x} u\|_{L^2 L^2((-\infty, n] \times A_j)} + 2^{-\frac{j}{2}} \|\langle r \rangle^{-1} u\|_{L^2 L^2((-\infty, n] \times A_j)} \lesssim \|f\|_{LE^*}$$

By taking the limit at  $\infty$  with respect to  $n$ , we deduce that

$$2^{-\frac{j}{2}} \|\nabla_{t,x} u\|_{L^2 L^2(\mathbb{R} \times A_j)} + 2^{-\frac{j}{2}} \|\langle r \rangle^{-1} u\|_{L^2 L^2(\mathbb{R} \times A_j)} \lesssim \|f\|_{LE^*}$$

We now immediately get that

$$\|\nabla_{t,x} u\|_{LE} + \|\langle r \rangle^{-1} u\|_{LE} \lesssim \|f\|_{LE^*},$$

as claimed.

*Step 3: Reduction to a form for which Plancherel's Theorem can be applied.* We define  $A'_0 = A_0$ , and  $A'_j = \{x \in \mathbb{R}^d | 2^{j-1} < |x| < 2^{j+1}\}$ . We note that  $A'_k \subset A_{k-1} \cup A_k$  for every  $k \geq 1$ . We also consider a smooth partition of unity  $(\chi_k)_{k \geq 0}$ , with  $0 \leq \chi_k \leq 1$ ,  $\sum_{k=0}^{\infty} \chi_k = 1$ , and  $\text{supp } \chi_k \subseteq A'_k$  for every  $k \geq 0$ .

We claim that the condition

$$\|e^{-\varepsilon t} \nabla_{t,x} u\|_{LE} + \|\langle r \rangle^{-1} e^{-\varepsilon t} u\|_{LE} \lesssim \|e^{-\varepsilon t} f\|_{LE^*}$$

is equivalent to

$$2^{-\frac{j}{2}} \|e^{-\varepsilon t} \nabla_{t,x} u\|_{L^2 L^2(\mathbb{R}_t \times A'_j)} + 2^{-\frac{3j}{2}} \|e^{-\varepsilon t} u\|_{L^2 L^2(\mathbb{R}_t \times A'_j)} \lesssim 2^{\frac{k}{2}} \|e^{-\varepsilon t} f\|_{L^2 L^2(\mathbb{R}_t \times A'_k)}$$

For one implication, we note that if we fix  $j, k \geq 0$ , and if  $f$  is supported on  $A'_k$ , we have (when  $j \geq 1$ )

$$\begin{aligned} & 2^{-\frac{j}{2}} \|e^{-\varepsilon t} \nabla_{t,x} u\|_{L^2 L^2(\mathbb{R}_t \times A'_j)} + 2^{-\frac{3j}{2}} \|e^{-\varepsilon t} u\|_{L^2 L^2(\mathbb{R}_t \times A'_j)} \\ & \lesssim 2^{-\frac{j}{2}} \|e^{-\varepsilon t} \nabla_{t,x} u\|_{L^2 L^2(\mathbb{R}_t \times A_{j-1})} + 2^{-\frac{3j}{2}} \|e^{-\varepsilon t} u\|_{L^2 L^2(\mathbb{R}_t \times A_{j-1})} \\ & + 2^{-\frac{j}{2}} \|e^{-\varepsilon t} \nabla_{t,x} u\|_{L^2 L^2(\mathbb{R}_t \times A_j)} + 2^{-\frac{3j}{2}} \|e^{-\varepsilon t} u\|_{L^2 L^2(\mathbb{R}_t \times A_j)} \\ & \lesssim 2^{\frac{k}{2}} \|e^{-\varepsilon t} f\|_{L^2 L^2(\mathbb{R}_t \times A_k)} \lesssim 2^{\frac{k}{2}} \|e^{-\varepsilon t} f\|_{L^2 L^2(\mathbb{R}_t \times A'_k)} \end{aligned}$$

and

$$\begin{aligned} & \|e^{-\varepsilon t} \nabla_{t,x} u\|_{L^2 L^2(\mathbb{R}_t \times A'_0)} + 2^{-\frac{3j}{2}} \|e^{-\varepsilon t} u\|_{L^2 L^2(\mathbb{R}_t \times A'_0)} \\ & \lesssim 2^{\frac{k}{2}} \|e^{-\varepsilon t} f\|_{L^2 L^2(\mathbb{R}_t \times A_k)} \lesssim 2^{\frac{k}{2}} \|e^{-\varepsilon t} f\|_{L^2 L^2(\mathbb{R}_t \times A'_k)} \end{aligned}$$

when  $j = 0$ .

For the converse, we denote by  $u^k$  be the forward solution corresponding to  $\chi_k f$ . Thus,  $u = \sum_k u^k$ . For every  $j \geq 0$  we have

$$\begin{aligned} & 2^{-\frac{j}{2}} \|e^{-\varepsilon t} \nabla_{t,x} u\|_{L^2 L^2(\mathbb{R}_t \times A_j)} + 2^{-\frac{j}{2}} \|\langle r \rangle^{-1} e^{-\varepsilon t} u\|_{L^2 L^2(\mathbb{R}_t \times A_j)} \\ & \lesssim 2^{-\frac{j}{2}} \|e^{-\varepsilon t} \nabla_{t,x} u\|_{L^2 L^2(\mathbb{R}_t \times A_j)} + 2^{-\frac{3j}{2}} \|e^{-\varepsilon t} u\|_{L^2 L^2(\mathbb{R}_t \times A_j)} \\ & \lesssim 2^{-\frac{j}{2}} \|e^{-\varepsilon t} \nabla_{t,x} u\|_{L^2 L^2(\mathbb{R}_t \times A'_j)} + 2^{-\frac{3j}{2}} \|e^{-\varepsilon t} u\|_{L^2 L^2(\mathbb{R}_t \times A'_j)} \\ & \lesssim \sum_{k \geq 0} 2^{-\frac{j}{2}} \|e^{-\varepsilon t} \nabla_{t,x} u^k\|_{L^2 L^2(\mathbb{R}_t \times A'_j)} + 2^{-\frac{3j}{2}} \|e^{-\varepsilon t} u^k\|_{L^2 L^2(\mathbb{R}_t \times A'_j)} \\ & \lesssim \sum_{k \geq 0} 2^{\frac{k}{2}} \|e^{-\varepsilon t} \chi_k f\|_{L^2 L^2(\mathbb{R}_t \times A'_k)} \\ & \lesssim \|e^{-\varepsilon t} \chi_k f\|_{L^2 L^2(\mathbb{R}_t \times A'_0)} + \sum_{k \geq 1} 2^{\frac{k}{2}} \|e^{-\varepsilon t} \chi_k f\|_{L^2 L^2(\mathbb{R}_t \times A'_k)} \\ & \lesssim \|e^{-\varepsilon t} \chi_k f\|_{L^2 L^2(\mathbb{R}_t \times A_0)} + \sum_{k \geq 1} 2^{\frac{k}{2}} (\|e^{-\varepsilon t} \chi_k f\|_{L^2 L^2(\mathbb{R}_t \times A_k)} + \|e^{-\varepsilon t} \chi_k f\|_{L^2 L^2(\mathbb{R}_t \times A_{k-1})}) \\ & \lesssim \|e^{-\varepsilon t} f\|_{LE^*} \end{aligned}$$

By taking the supremum in  $j \geq 0$ , we immediately deduce that

$$\|e^{-\varepsilon t} \nabla_{t,x} u\|_{LE} + \|\langle r \rangle^{-1} e^{-\varepsilon t} u\|_{LE} \lesssim \|e^{-\varepsilon t} f\|_{LE^*},$$

as claimed.

*Step 4: Concluding the proof.* We now prove that the condition

$$2^{-\frac{j}{2}} \|e^{-\varepsilon t} \nabla_{t,x} u\|_{L^2 L^2(\mathbb{R}_t \times A'_j)} + 2^{-\frac{3j}{2}} \|e^{-\varepsilon t} u\|_{L^2 L^2(\mathbb{R}_t \times A'_j)} \lesssim 2^{\frac{k}{2}} \|e^{-\varepsilon t} f\|_{L^2 L^2(\mathbb{R}_t \times A'_k)}, \forall j, k \geq 0$$

is equivalent to the spectral characterization.

Let now  $f$  be a function supported in  $A'_k$ , and  $u$  the associated forward solution. For every  $j \geq 0$ , we have

$$\begin{aligned} & 2^{-\frac{j}{2}} \|e^{-\varepsilon t} \nabla_{t,x} u\|_{L^2 L^2(\mathbb{R}_t \times A'_j)} + 2^{-\frac{j}{2}} \|e^{-\varepsilon t} \langle r \rangle^{-1} u\|_{L^2 L^2(\mathbb{R}_t \times A'_j)} \\ & \lesssim 2^{-\frac{j}{2}} \|e^{-\varepsilon t} \nabla_{t,x} u\|_{L^2 L^2(\mathbb{R}_t \times A'_j)} + 2^{-\frac{3j}{2}} \|e^{-\varepsilon t} u\|_{L^2 L^2(\mathbb{R}_t \times A'_j)} \\ & \lesssim 2^{\frac{k}{2}} \|e^{-\varepsilon t} f\|_{L^2 L^2(\mathbb{R}_t \times A'_k)} \end{aligned}$$

We take the Fourier transform in time and obtain (via Plancherel's Theorem) that

$$\begin{aligned} & 2^{-\frac{j}{2}} \|(|\tau - i\varepsilon|, \nabla_x) \widehat{u}(\tau - i\varepsilon)\|_{L^2 L^2(\mathbb{R}_\tau \times A'_j)} + 2^{-\frac{j}{2}} \|\langle r \rangle^{-1} \widehat{u}(\tau - i\varepsilon)\|_{L^2 L^2(\mathbb{R}_\tau \times A'_j)} \\ & \lesssim 2^{\frac{k}{2}} \|\widehat{f}(\tau - i\varepsilon)\|_{L^2 L^2(\mathbb{R}_\tau \times A'_k)} \end{aligned}$$

In particular, this holds for every function  $f$  of the form  $f(t, x) = \phi(t)g(x)$ , with  $\phi$  supported away from  $\{t = -\infty\}$ , and  $g$  supported in  $A'_k$ . In this case,  $\widehat{u}(\tau - i\varepsilon) = \mathbf{R}_{\tau - i\varepsilon}(\widehat{\phi}(\tau - i\varepsilon)g(x)) = \widehat{\phi}(\tau - i\varepsilon)\mathbf{R}_{\tau - i\varepsilon}(g(x))$ . Thus, for every  $j, k \geq 0$ ,

$$\begin{aligned} & 2^{-\frac{j}{2}} \|\widehat{\phi}(\tau - i\varepsilon)\|_{L^2(\mathbb{R}_\tau)} \|(|\tau - i\varepsilon|, \nabla_x) \mathbf{R}_{\tau - i\varepsilon} g(x)\|_{L^2(A'_j)} \\ & + 2^{-\frac{j}{2}} \|\widehat{\phi}(\tau - i\varepsilon)\|_{L^2(\mathbb{R}_\tau)} \|\langle r \rangle^{-1} \mathbf{R}_{\tau - i\varepsilon} g(x)\|_{L^2(A'_j)} \\ & \lesssim 2^{\frac{k}{2}} \|\widehat{\phi}(\tau - i\varepsilon)\|_{L^2(\mathbb{R}_\tau)} \|g(x)\|_{L^2(A'_k)} \end{aligned}$$

This shows that

$$2^{-\frac{j}{2}} \|(|\tau - i\varepsilon|, \nabla_x) \mathbf{R}_{\tau - i\varepsilon} g(x)\|_{L^2(A'_j)} + 2^{-\frac{j}{2}} \|\langle r \rangle^{-1} \mathbf{R}_{\tau - i\varepsilon} g(x)\|_{L^2(A'_j)} \lesssim 2^{\frac{k}{2}} \|g(x)\|_{L^2(A'_k)}$$

For arbitrary  $g$ , we take a partition of unity  $(\chi_k)_{k \geq 0}$ , with  $0 \leq \chi_k \leq 1$ ,  $\sum_{k=0}^{\infty} \chi_k = 1$ , and  $\text{supp } \chi_k \subseteq A_k$  for every  $k \geq 0$ . The previous discussion implies that for every  $k \geq 0$ ,

$$\|(|\tau - i\varepsilon|, \nabla_x) \mathbf{R}_{\tau - i\varepsilon}(\chi_k g)\|_{\mathcal{LE}} + \|\langle r \rangle^{-1} \mathbf{R}_{\tau - i\varepsilon}(\chi_k g)\|_{\mathcal{LE}} \lesssim 2^{\frac{k}{2}} \|\chi_k g\|_{L^2 L^2(\mathbb{R}_t \times A'_k)}.$$

Therefore,

$$\begin{aligned} & \|(|\tau - i\varepsilon|, \nabla_x) \mathbf{R}_{\tau - i\varepsilon} g\|_{\mathcal{LE}} + \|\langle r \rangle^{-1} \mathbf{R}_{\tau - i\varepsilon} g\|_{\mathcal{LE}} \leq \sum_{k=0}^{\infty} \|(|\tau - i\varepsilon|, \nabla_x) \mathbf{R}_{\tau - i\varepsilon}(\chi_k g)\|_{\mathcal{LE}} \\ & + \sum_{k=0}^{\infty} \|\langle r \rangle^{-1} \mathbf{R}_{\tau - i\varepsilon}(\chi_k g)\|_{\mathcal{LE}} \lesssim \sum_{k=0}^{\infty} 2^{\frac{k}{2}} \|\chi_k g\|_{L^2 L^2(\mathbb{R}_t \times A'_k)} \\ & \lesssim \|\chi_0 g\|_{L^2 L^2(\mathbb{R}_t \times A'_0)} + \sum_{k=1}^{\infty} 2^{\frac{k}{2}} (\|\chi_k g\|_{L^2 L^2(\mathbb{R}_t \times A_{k-1})} + \|\chi_k g\|_{L^2 L^2(\mathbb{R}_t \times A_k)}) \\ & \lesssim \|g\|_{L^2 L^2(\mathbb{R}_t \times A_0)} + \sum_{k=1}^{\infty} 2^{\frac{k}{2}} (\|g\|_{L^2 L^2(\mathbb{R}_t \times A_{k-1})} + \|g\|_{L^2 L^2(\mathbb{R}_t \times A_k)}) \\ & \lesssim \|g\|_{\mathcal{LE}^*} \end{aligned}$$

Conversely, when  $f$  is a function that is supported away from  $\{t = -\infty\}$  and in  $A'_k$ , then  $\widehat{f}(\tau - i\varepsilon)$  is supported in  $A'_k$  as well, so the previous inequality implies that

$$\begin{aligned} & \|(|\tau - i\varepsilon|, \nabla_x) \mathbf{R}_{\tau - i\varepsilon} \widehat{f}(\tau - i\varepsilon)\|_{\mathcal{LE}} + \|\langle r \rangle^{-1} \mathbf{R}_{\tau - i\varepsilon} \widehat{f}(\tau - i\varepsilon)\|_{\mathcal{LE}} \lesssim \|\widehat{f}(\tau - i\varepsilon)\|_{\mathcal{LE}^*} \\ & = 2^{\frac{k}{2}} \|\widehat{f}(\tau - i\varepsilon)\|_{L^2(A'_k)} \end{aligned}$$

Thus, for every  $j \geq 1$ ,

$$\begin{aligned}
& 2^{-\frac{j}{2}} \|(|\tau - i\varepsilon|, \nabla_x) \mathbf{R}_{\tau - i\varepsilon} \widehat{f}(\tau - i\varepsilon)\|_{L^2(A'_j)} + 2^{-\frac{j}{2}} \|\langle r \rangle^{-1} \mathbf{R}_{\tau - i\varepsilon} \widehat{f}(\tau - i\varepsilon)\|_{L^2(A'_j)} \\
& \lesssim 2^{-\frac{j}{2}} \|(|\tau - i\varepsilon|, \nabla_x) \mathbf{R}_{\tau - i\varepsilon} \widehat{f}(\tau - i\varepsilon)\|_{L^2(A_j)} + 2^{-\frac{j}{2}} \|\langle r \rangle^{-1} \mathbf{R}_{\tau - i\varepsilon} \widehat{f}(\tau - i\varepsilon)\|_{L^2(A_j)} \\
& + 2^{-\frac{j}{2}} \|(|\tau - i\varepsilon|, \nabla_x) \mathbf{R}_{\tau - i\varepsilon} \widehat{f}(\tau - i\varepsilon)\|_{L^2(A_{j-1})} + 2^{-\frac{j}{2}} \|\langle r \rangle^{-1} \mathbf{R}_{\tau - i\varepsilon} \widehat{f}(\tau - i\varepsilon)\|_{L^2(A_{j-1})} \\
& \lesssim 2^{\frac{k}{2}} \|\widehat{f}(\tau - i\varepsilon)\|_{L^2(A_k)} \lesssim 2^{\frac{k}{2}} \|\widehat{f}(\tau - i\varepsilon)\|_{L^2(A'_k)},
\end{aligned}$$

while for  $j = 0$ ,

$$\begin{aligned}
& \|(|\tau - i\varepsilon|, \nabla_x) \mathbf{R}_{\tau - i\varepsilon} \widehat{f}(\tau - i\varepsilon)\|_{L^2(A'_0)} + \|\langle r \rangle^{-1} \mathbf{R}_{\tau - i\varepsilon} \widehat{f}(\tau - i\varepsilon)\|_{L^2(A'_0)} \\
& \lesssim \|(|\tau - i\varepsilon|, \nabla_x) \mathbf{R}_{\tau - i\varepsilon} \widehat{f}(\tau - i\varepsilon)\|_{L^2(A_0)} + \|\langle r \rangle^{-1} \mathbf{R}_{\tau - i\varepsilon} \widehat{f}(\tau - i\varepsilon)\|_{L^2(A_0)} \\
& \lesssim 2^{\frac{k}{2}} \|\widehat{f}(\tau - i\varepsilon)\|_{L^2(A_k)} \lesssim 2^{\frac{k}{2}} \|\widehat{f}(\tau - i\varepsilon)\|_{L^2(A'_k)},
\end{aligned}$$

Thus, for every  $j, k \geq 0$

$$\begin{aligned}
& 2^{-\frac{j}{2}} \|(|\tau - i\varepsilon|, \nabla_x) \mathbf{R}_{\tau - i\varepsilon} \widehat{f}(\tau - i\varepsilon)\|_{L^2(A'_j)} + 2^{-\frac{j}{2}} \|\langle r \rangle^{-1} \mathbf{R}_{\tau - i\varepsilon} \widehat{f}(\tau - i\varepsilon)\|_{L^2(A'_j)} \\
& \lesssim 2^{\frac{k}{2}} \|\widehat{f}(\tau - i\varepsilon)\|_{L^2(A'_k)}
\end{aligned}$$

By taking the  $L^2$ -norm in  $\tau$  and using Plancherel's identity, we deduce that

$$2^{-\frac{j}{2}} \|e^{-\varepsilon t} \nabla u\|_{L^2 L^2(\mathbb{R}_t \times A'_j)} + 2^{-\frac{j}{2}} \|e^{-\varepsilon t} \langle r \rangle^{-1} u\|_{L^2 L^2(\mathbb{R}_t \times A'_j)} \lesssim 2^{\frac{k}{2}} \|e^{-\varepsilon t} f\|_{L^2 L^2(\mathbb{R}_t \times A'_k)}$$

This finishes the proof.  $\square$

## REFERENCES

- [Metcalfe-Tataru] Global parametrices and dispersive estimates for variable coefficient wave equations. *Math. Ann.* 353, 1183–1237 (2012)
- [Keel-Tao] Keel, Markus, and Terence Tao. “Endpoint Strichartz Estimates.” *American Journal of Mathematics* 120, no. 5 (1998): 955–80
- [Ionescu-Kenig] A.D. Ionescu, C. Kenig, Well-posedness and local smoothing of solutions of Schrödinger equations, *Math. Res. Lett.* 12 (2005) 193–205
- [Burq Planchon Stalker Tahvildar-Zadeh] Burq, Nicolas, Fabrice Planchon, John G. Stalker, and A. Shadi Tahvildar-Zadeh. “Strichartz Estimates for the Wave and Schrödinger Equations with Potentials of Critical Decay.” *Indiana University Mathematics Journal* 53, no. 6 (2004): 1665–80
- [Metcalfe-Sterbenz-Tataru] Metcalfe, Jason, Jacob Sterbenz, and Daniel Tataru. “Local energy decay for scalar fields on time dependent non-trapping backgrounds.” *American Journal of Mathematics* 142, no. 3 (2020): 821–883
- [Tataru] Tataru, Daniel. “LOCAL DECAY OF WAVES ON ASYMPTOTICALLY FLAT STATIONARY SPACE-TIMES.” *American Journal of Mathematics* 135, no. 2 (2013): 361–401
- [Dafermos-Rodnianski] Dafermos, Mihalis; Rodnianski, Igor. A new physical-space approach to decay for the wave equation with applications to black hole spacetimes. *XVIth International Congress on Mathematical Physics*. World Scientific Publishing Co., 2010. p. 421–432
- [Moschidis] Moschidis, G. The  $r^p$ -Weighted Energy Method of Dafermos and Rodnianski in General Asymptotically Flat Spacetimes and Applications. *Ann. PDE* 2, 6 (2016)
- [Oliver-Sterbenz] Jesús Oliver. Jacob Sterbenz. “A vector field method for radiating black hole spacetimes.” *Anal. PDE* 13 (1) 29 - 92, 2020