

Moment (physics)

From Wikipedia, the free encyclopedia

In physics, **moment** is an expression involving the product of a distance and physical quantity, and in this way it accounts for how the physical quantity is located or arranged. Moments are usually defined with respect to a fixed reference point; they deal with physical quantities located at some distance relative to that reference point. For example, **the moment of force**, often called **torque**, is the **product of a force on an object and the distance from the reference point to the object**. In principle, any physical quantity can be multiplied by a distance to produce a moment. Commonly used quantities include forces, masses, and electric charge distributions.

- Examples** [edit]
- The **moment of force**, or **torque**, is a first moment.

τ
=

r

P
¯

{\displaystyle \tau =r{\bar {P}}}

 or, more generally,

τ
=

r
×

F
¯

{\displaystyle \tau =r\times {\bar {F}}}
 - Similarly, **angular momentum** is the 1st moment of momentum.

L
¯
=

r
×

p
¯

{\displaystyle {\bar {L}}=r\times {\bar {p}}}

. Note that momentum itself is not a moment.
 - The electric **dipole moment** is also a 1st moment

p
¯
=

q

d
¯

{\displaystyle {\bar {p}}=q{\bar {d}}}

 for two opposite point charges or

∫

r

ρ
(

r
¯
)

d

3

r
¯

{\displaystyle \int {\bar {r}}\rho ({\bar {r}})d^{3}{\bar {r}}}

 for a distributed charge with charge density

ρ
(

r
¯
)

{\displaystyle \rho ({\bar {r}})}

- Moments of mass:
- The total mass is the zeroth moment of mass
 - The **center of mass** is the 1st moment of mass normalized by total mass:

R
¯
=

1
M

∑

i

r

i

m

i

{\displaystyle {\bar {R}}={\frac {1}{M}}\sum _{i}{\bar {r}}_{i}m_{i}}

 for a collection of point masses, or

1
M

∫

r

ρ
(

r
¯
)

d

3

r
¯

{\displaystyle {\frac {1}{M}}\int {\bar {r}}\rho ({\bar {r}})d^{3}{\bar {r}}}

 for an object with mass distribution

ρ
(

r
¯
)

{\displaystyle \rho ({\bar {r}})}
 - The **moment of inertia** is the **2nd moment of mass**.

I
¯
=

r

2

m

{\displaystyle {\bar {I}}=r^{2}m}

 for a point mass

∑

i

r

i

2

m

i

{\displaystyle \sum _{i}r_{i}^{2}m_{i}}

 for a collection of point masses, or

∫

r

2

ρ
(

r
¯
)

d

3

r
¯

{\displaystyle \int r^{2}\rho ({\bar {r}})d^{3}{\bar {r}}}

 for an object with mass distribution

ρ
(

r
¯
)

{\displaystyle \rho ({\bar {r}})}

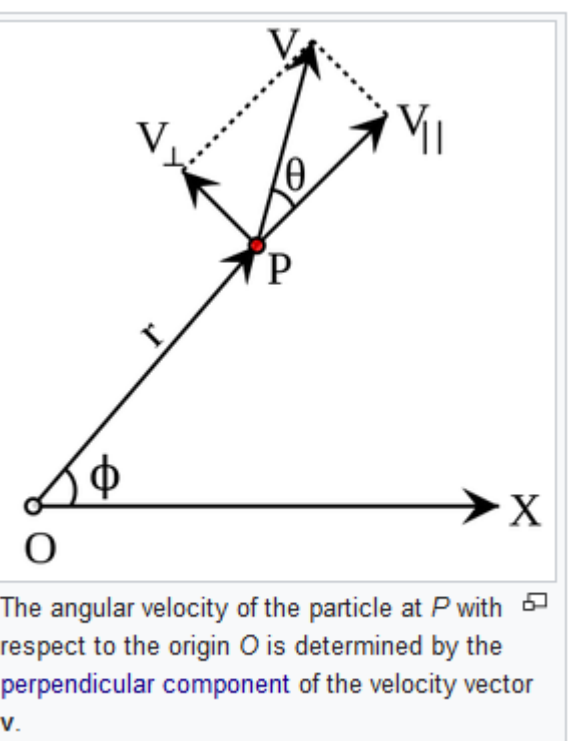
. Note that the center of mass is often (but not always) taken as the reference point.

Angular velocity

From Wikipedia, the free encyclopedia

See also: *Angular frequency*

In physics, **angular velocity** refers to **how fast an object rotates or revolves relative to another point**, i.e. how fast the angular position or orientation of an object changes with time. There are two types of angular



The **angular velocity**

ω

{\displaystyle \omega }

 is the rate of change of angular position with respect to time, which can be computed from the cross-radial velocity as:

ω
=

d
θ

d
t

=

v

⊥

r

{\displaystyle \omega ={\frac {d\theta }{dt}}={\frac {v_{\perp }}{r}}}

Here the cross-radial speed

v

⊥

{\displaystyle v_{\perp }}

 is the signed magnitude of

v

⊥

{\displaystyle v_{\perp }}

, positive for counter-clockwise motion, negative for clockwise. Taking polar coordinates for the linear velocity

v
¯

{\displaystyle {\bar {v}}}

 gives magnitude

v

{\displaystyle v}

 (linear speed) and angle

θ

{\displaystyle \theta }

 relative to the radius vector, in these terms,

v

⊥

=
v
sin
⁡
(
θ
)

{\displaystyle v_{\perp }=v\sin(\theta)}

, so that

ω
=

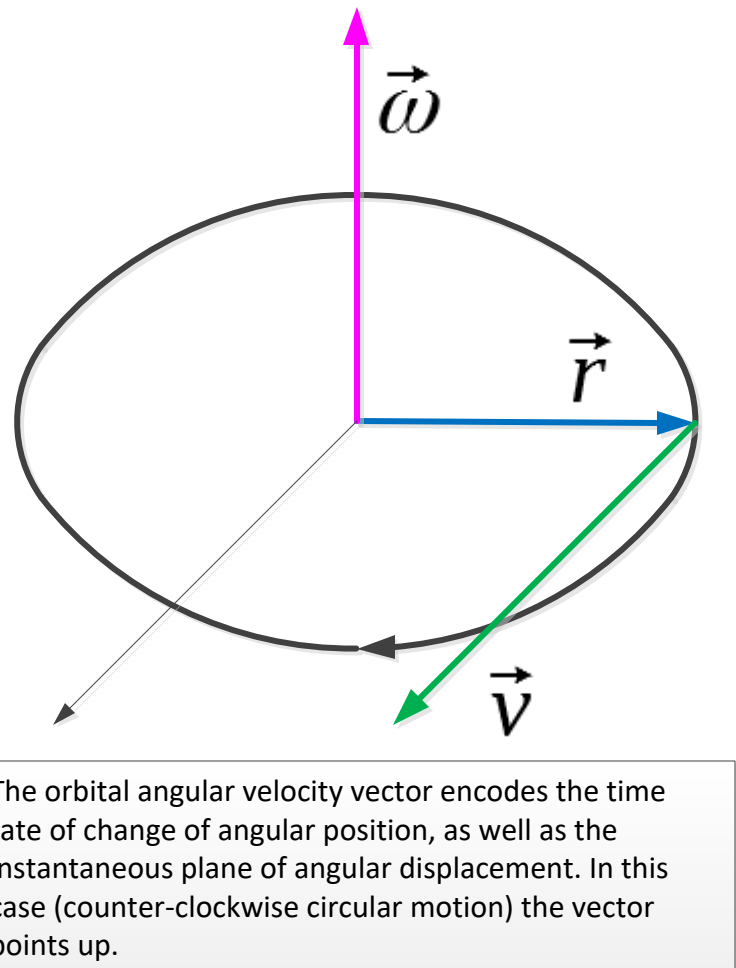
v
sin
⁡
(
θ
)

r

{\displaystyle \omega ={\frac {v\sin(\theta)}{r}}}

Angular Velocity in 3D

In three dimensions, angular velocity is a **pseudovector**, with its magnitude measuring the rate at which an object rotates or revolves, **and its direction pointing perpendicular to the instantaneous plane of rotation or angular displacement**. The orientation of angular velocity is conventionally specified by the right-hand rule.^[1]



Momentum

In **Newtonian mechanics**, **linear momentum**, **translational momentum**, or simply **momentum** (pl. momenta) is the **product of the mass and velocity of an object**. It is a vector quantity, possessing a magnitude and a direction. If it is an object's mass and

v
¯

{\displaystyle {\bar {v}}}

 is its velocity (also a vector quantity), then the object's momentum is

p
¯
=
m

v
¯

{\displaystyle {\bar {p}}=m{\bar {v}}}

In SI units, momentum is measured in **kilogram meters per second** (kg·m/s).

Angular momentum

In physics, **angular momentum** (rarely, **moment of momentum** or **rotational momentum**) is the **rotational equivalent of linear momentum**. It is an important quantity in physics because it is a **conserved quantity**—the total angular momentum of a closed system remains constant.

In three dimensions, the angular momentum for a point particle is a **pseudovector**

r
×

p
¯

{\displaystyle {\bar {r}}\times {\bar {p}}}

, the **cross product of the particle's position vector

r
¯

{\displaystyle {\bar {r}}}

 (relative to some origin) and its momentum vector**; the latter is

p
¯
=
m

v
¯

{\displaystyle {\bar {p}}=m{\bar {v}}}

 in Newtonian mechanics. This definition can be applied to each point in

Common symbols	 L ¯<!-- ¯ --> {\displaystyle {\bar {L}}}
In SI base units	kg m² s ^{−1}
Conserved?	yes
Derivation from other quantities	 L ¯<!-- ¯ --> = l a o = r ¯<!-- ¯ --> ×<!-- × --> p ¯<!-- ¯ --> {\displaystyle {\bar {L}}=l_{a}o={\bar {r}}\times {\bar {p}}}
Dimension	M L ² T ^{−1}

and treat it as a scalar (more precisely, a **pseudoscalar**).^[2] Angular momentum can be considered a **rotational analog of linear momentum**. Thus, where linear momentum

p
¯

{\displaystyle {\bar {p}}}

 is proportional to mass

m

{\displaystyle m}

 and linear speed

v

{\displaystyle v}

,

p
¯
=
m

v
¯

{\displaystyle {\bar {p}}=m{\bar {v}}}

**angular momentum

L
¯

{\displaystyle {\bar {L}}}

 is proportional to moment of inertia

I
¯

{\displaystyle {\bar {I}}}

 and angular speed

ω

{\displaystyle \omega }

 measured in radians per second**^[1]

L
¯
=

I
¯
ω

{\displaystyle {\bar {L}}={\bar {I}}\omega }

Because

I
¯
=

r

2

m

{\displaystyle {\bar {I}}=r^{2}m}

 for a single particle and

ω
=

v
¯

r

{\displaystyle \omega ={\frac {\bar {v}}{r}}}

 for circular motion, angular momentum can be expanded,

L
¯
=

r

2

m

⋅

v
¯

r

=
r
m

v
¯

{\displaystyle {\bar {L}}=r^{2}m\cdot {\frac {\bar {v}}{r}}=r\,m{\bar {v}}}

the product of the radius of rotation

r

{\displaystyle r}

 and the linear momentum of the particle

p
¯
=
m

v
¯

{\displaystyle {\bar {p}}=m{\bar {v}}}

, where

v

{\displaystyle v}

 in this case is the equivalent linear (tangential) speed at the radius (

r
=

r

u

{\displaystyle r=r_{u}}

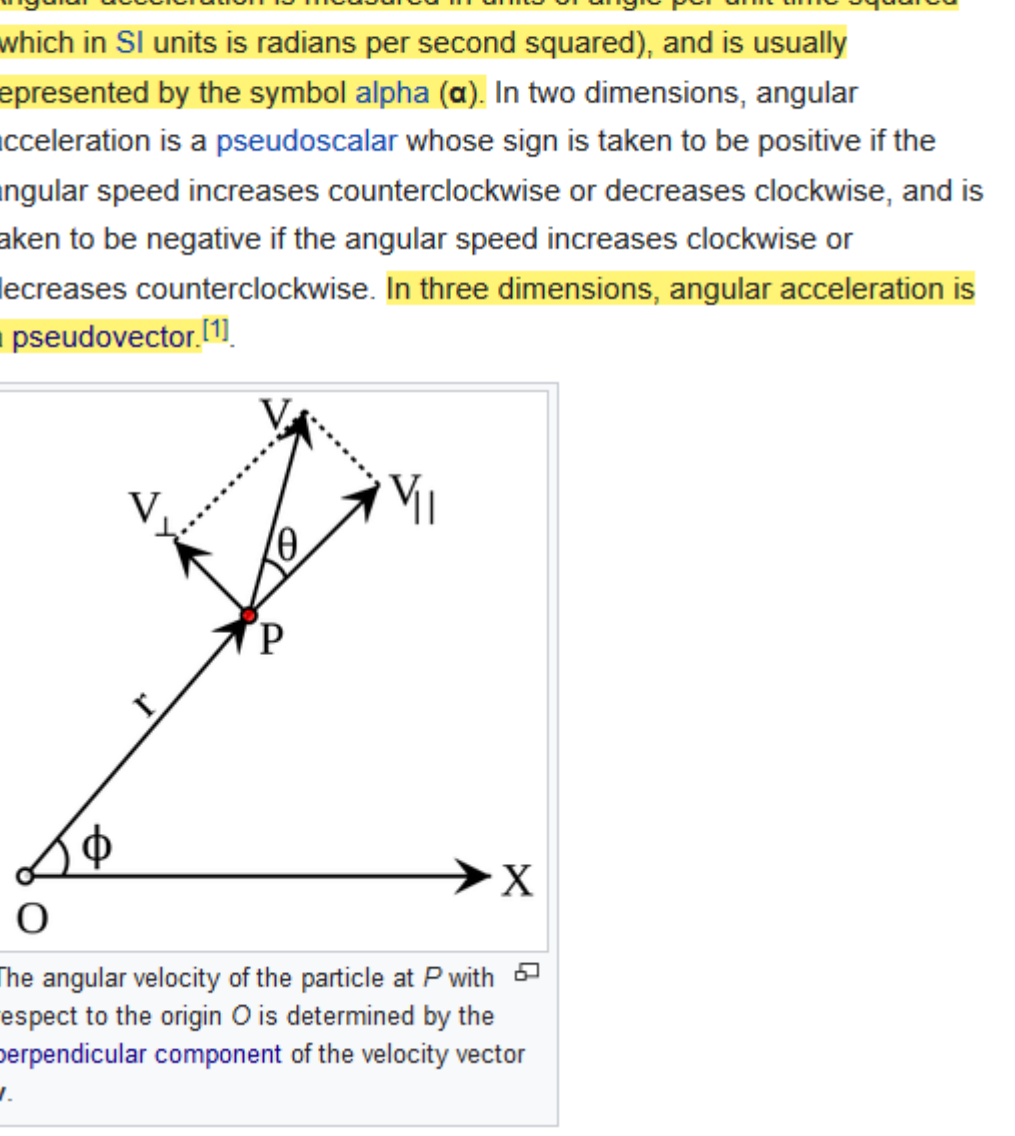
).

Angular acceleration

From Wikipedia, the free encyclopedia

In physics, **angular acceleration** refers to the **time rate of change of angular velocity**. As there are two types of angular velocity, namely spin angular velocity and orbital angular velocity, there are naturally also two types of angular acceleration, called spin angular acceleration and orbital angular acceleration respectively. Spin angular acceleration refers to the angular acceleration of a rigid body about its centre of rotation, and orbital angular acceleration refers to the angular acceleration of a point particle about a fixed origin.

Angular acceleration is measured in units of angle per unit time squared (which in SI units is radians per second squared), and is usually represented by the symbol alpha (**α**). In two dimensions, angular acceleration is a pseudoscalar whose sign is taken to be positive if the angular speed increases counterclockwise or decreases clockwise, and is taken to be negative if the angular speed increases clockwise or decreases counterclockwise. **In three dimensions, angular acceleration is a pseudovector**.^[1]



Particle in two dimensions [edit]

In two dimensions, the orbital angular acceleration is the rate at which the two-dimensional orbital angular velocity of the particle about the origin changes. The instantaneous **angular velocity**

ω

{\displaystyle \omega }

 at any point in time is given by

ω
=

v

⊥

r

{\displaystyle \omega ={\frac {v_{\perp }}{r}}}

where

r

{\displaystyle r}

 is the distance from the origin and

v

⊥

{\displaystyle v_{\perp }}

 is the cross-radial component of the instantaneous velocity (i.e. the component perpendicular to the position vector), which by convention is positive for counter-clockwise motion and negative for clockwise motion.

Therefore, the instantaneous angular acceleration

α

{\displaystyle \alpha }

 of the particle is given by

α
=

d

d
t

v

⊥

r

{\displaystyle \alpha ={\frac {d}{dt}}{\frac {v_{\perp }}{r}}}

Expanding the right-hand-side using the product rule from differential calculus, this becomes

α
=

1
r

d

v

⊥

d
t

−

v

⊥

r

2

d
r

d
t

{\displaystyle \alpha ={\frac {1}{r}}{\frac {dv_{\perp }}{dt}}-{\frac {v_{\perp }}{r^{2}}}{\frac {dr}{dt}}}

In the special case where the particle undergoes circular motion about the origin,

d

v

⊥

d
t

{\displaystyle {\frac {dv_{\perp }}{dt}}}

 becomes just the tangential acceleration

a

⊥

{\displaystyle a_{\perp }}

, and

d
r

d
t

{\displaystyle {\frac {dr}{dt}}}

 vanishes (since the distance from the origin stays constant), so the above equation simplifies to

α
=

a

⊥

r

{\displaystyle \alpha ={\frac {a_{\perp }}{r}}}

Particle in three dimensions [edit]

In three dimensions, the orbital angular acceleration is the rate at which three-dimensional orbital angular velocity vector changes with time. The instantaneous angular velocity vector

ω

{\displaystyle \omega }

 at any point in time is given by

ω
=

r
×

v
¯

r

2

{\displaystyle \omega ={\frac {\mathbf {r}\times {\bar {v}}}{r^{2}}}}

where

r

{\displaystyle r}

 is the particle's position vector and

v
¯

{\displaystyle {\bar {v}}}

 is its velocity vector.^[2]

Therefore, the orbital angular acceleration is the vector

α

{\displaystyle \alpha }

 defined by

α
=

d

d
t

r
×

v
¯

r

2

{\displaystyle \alpha ={\frac {d}{dt}}{\frac {\mathbf {r}\times {\bar {v}}}{r^{2}}}}

Expanding this derivative using the **product rule for cross-products** and the **ordinary quotient rule**, one gets:

α
=

1

r

2

(

r
×

d

v
¯

d
t

+

d

r

d
t

×

v
¯
)
−

2
d
r

d
t

r

(

r
×

v
¯
)

{\displaystyle \alpha ={\frac {1}{r^{2}}}\left(\mathbf {r} \times {\frac {d{\bar {v}}}{dt}}+{\frac {d\mathbf {r} }{dt}}\times {\bar {v}}\right)-{\frac {2}{r^{2}}}{\frac {dr}{dt}}\left(\mathbf {r} \times {\bar {v}}\right)}

Since

r
×

v
¯

{\displaystyle \mathbf {r} \times {\bar {v}}}

 is just

r

2

ω

{\displaystyle r^{2}\omega }

, the second term may be rewritten as

2
d
r

d
t

ω

{\displaystyle {\frac {2}{r}}{\frac {dr}{dt}}\omega }

. In the case where the distance

r

{\displaystyle r}

 of the particle from the origin does not change with time (which includes circular motion as a subcase), the second term vanishes and the above formula simplifies to

α
=

r
×

a
¯

r

2

{\displaystyle \alpha ={\frac {\mathbf {r}\times {\bar {a}}}{r^{2}}}}

From the above equation, one can recover the cross-radial acceleration in this special case as:

a

⊥

=
α
×

r
¯

{\displaystyle a_{\perp }=\alpha \times {\bar {r}}}

Relation to Torque [edit]

The net torque on a point particle is defined to be the pseudovector

τ
¯
=

r
¯
×

F
¯

{\displaystyle {\bar {\tau }}={\bar {r}}\times {\bar {F}}}

where

F
¯

{\displaystyle {\bar {F}}}

 is the net force on the particle.^[3]

Torque is the rotational analogue of force: it induces change in the rotational state of a system, just as force induces change in the translational state of a system. Since the net force on a particle may be connected to the acceleration of the particle by the equation

F
¯
=
m

a
¯

{\displaystyle {\bar {F}}=m{\bar {a}}}

, one may hope to construct a similar relation connecting the net torque on a particle to the angular acceleration of the particle. That may be done as follows:

First, substituting

F
¯
=
m

a
¯

{\displaystyle {\bar {F}}=m{\bar {a}}}

 into the above equation for torque, one gets

τ
¯
=
m
(

r
×

a
¯
)
=
m

r

2

r
×

a
¯

r

2

{\displaystyle {\bar {\tau }}=m\left(\mathbf {r} \times {\bar {a}}\right)=mr^{2}{\frac {\mathbf {r}\times {\bar {a}}}{r^{2}}}}

But from the previous section, it was derived that

α
=

r
×

a
¯

r

2

−

2
d
r

d
t

a
¯

{\displaystyle \alpha ={\frac {\mathbf {r}\times {\bar {a}}}{r^{2}}}-{\frac {2}{r^{2}}}{\frac {dr}{dt}}{\bar {a}}}

where

α

{\displaystyle \alpha }

 is the orbital angular acceleration of the particle and

ω

{\displaystyle \omega }

 is the orbital angular velocity of the particle. Therefore, it follows that

τ
¯
=
m

r

2

(
α
+

2
d
r

d
t

ω
)

{\displaystyle {\bar {\tau }}=mr^{2}\left(\alpha +{\frac {2}{r}}{\frac {dr}{dt}}\omega \right)}

In the special case where the distance

r

{\displaystyle r}

 of the particle from the origin does not change with time, the second term in the above equation vanishes and the above equation simplifies to

τ
¯
=
m

r

2

α

{\displaystyle {\bar {\tau }}=mr^{2}\alpha }

If the shape of the body does not change, then its moment of inertia appears in Newton's law of motion as the ratio of an applied torque

τ
¯

{\displaystyle {\bar {\tau }}}

 on a body to the angular acceleration

α

{\displaystyle \alpha }

 around a principal axis, that is

τ
¯
=
I
α

{\displaystyle {\bar {\tau }}=I\alpha }

Torque

τ
¯
=

r
¯
×

F
¯

{\displaystyle {\bar {\tau }}={\bar {r}}\times {\bar {F}}}

τ
¯
=
|

r
¯

|

|

F
¯

|

sin
⁡
θ

{\displaystyle {\bar {\tau }}=|{\bar {r}}||{\bar {F}}|\sin \theta }

Newton's second law

The second law states that the **rate of change of momentum of a body is directly proportional to the force applied**, and this change in momentum takes place in the direction of the applied force.

F
¯
=

d

(

m

v
¯
)

d
t

=
d
(

m

v
¯)

d
t

{\displaystyle {\bar {F}}={\frac {d\left(m{\bar {v}}\right)}{dt}}={\frac {d\left(m{\bar {v}}\right)}{dt}}}

The second law can also be stated in terms of an object's acceleration. Since Newton's second law is valid only for constant-mass systems,^{[3][132]}

m

{\displaystyle m}

 can be taken outside the differentiation operator by the **constant factor rule** in differentiation. Thus,

F
¯
=
m

d

v
¯

d
t

=
m

a
¯

{\displaystyle {\bar {F}}=m{\frac {d{\bar {v}}}{dt}}=m{\bar {a}}}

where

F
¯

{\displaystyle {\bar {F}}}

 is the net force applied,

m

{\displaystyle m}

 is the mass of the body, and

a
¯

{\displaystyle {\bar {a}}}

 is the body's acceleration. Thus, the net force applied to a body produces a proportional acceleration. In other words, if a body is accelerating, then there is a force on it.

Differentiation rules

The product rule [edit]

Main article: *Product rule*

For the functions

f

{\displaystyle f}

 and

g

{\displaystyle g}

, the derivative of the function

h
(
x
)
=
f
(
x
)
g
(
x
)

{\displaystyle h(x)=f(x)g(x)}

 with respect to

x

{\displaystyle x}

 is

h
′
(
x
)
=
(
f
g
)
′
(
x
)
=
f
′
(
x
)
g
(
x
)
+
f
(
x
)
g
′
(
x
)
.

{\displaystyle h'(x)=(fg)'(x)=f'(x)g(x)+f(x)g'(x).}

In Leibniz's notation this is written

d
(
f
g
)

d
x

=

d
f

d
x

g
+
f

d
g

d
x

{\displaystyle {\frac {d(fg)}{dx}}={\frac {df}{dx}}g+f{\frac {dg}{dx}}}

Proof of the equivalence of definitions [edit]

The definition of angular momentum for a single point particle is:

L
¯
=

r
¯
×

p
¯

{\displaystyle {\bar {L}}={\bar {r}}\times {\bar {p}}}

where

p
¯

{\displaystyle {\bar {p}}}

 is the particle's linear momentum and

r

{\displaystyle r}

 is the position vector from the origin. The time-derivative of this is:

d
L
¯

d
t

=

r
¯
×

d

p
¯

d
t

+

d

r
¯

d
t

×

p
¯

{\displaystyle {\frac {d{\bar {L}}}{dt}}={\bar {r}}\times {\frac {d{\bar {p}}}{dt}}+{\frac {d{\bar {r}}}{dt}}\times {\bar {p}}}

This result can easily be proven by splitting the vectors into components and applying the product rule. Now using the definition of force

F
¯
=

d

p
¯

d
t

{\displaystyle {\bar {F}}={\frac {d{\bar {p}}}{dt}}}

 (whether or not mass is constant) and the definition of velocity

d

r
¯

d
t

=

v
¯

{\displaystyle {\frac {d{\bar {r}}}{dt}}={\bar {v}}}

,

d
L
¯

d
t

=

r
¯
×

F
¯
+

v
¯
×

p
¯

{\displaystyle {\frac {d{\bar {L}}}{dt}}={\bar {r}}\times {\bar {F}}+{\bar {v}}\times {\bar {p}}}

The cross product of momentum

p
¯

{\displaystyle {\bar {p}}}

 with its associated velocity

v
¯

{\displaystyle {\bar {v}}}

 is zero because velocity and momentum are parallel, so the second term vanishes.

By definition, torque

τ
¯
=

r
¯
×

F
¯

{\displaystyle {\bar {\tau }}={\bar {r}}\times {\bar {F}}}

, therefore, torque on a particle is **equal to the first derivative** of its angular momentum with respect to time.

If multiple forces are applied, Newton's second law instead reads

F

net

¯
=
m

a
¯

{\displaystyle {\bar {\mathbf {F} _{\mathrm {net} }}=m{\bar {\mathbf {a} }}}

, and it follows that

d
L
¯

d
t

=

r
¯
×

F

net

¯

{\displaystyle {\frac {d{\bar {L}}}{dt}}={\bar {r}}\times {\bar {\mathbf {F} _{\mathrm {net} }}}=m{\bar {\mathbf {a} }}}

This is a general proof for point particles.

Torque

From Wikipedia, the free encyclopedia

In three dimensions, the torque is a pseudovector; for point particles, it is given by the cross product of the position vector (distance vector) and the force vector. The magnitude of torque of a rigid body depends on three quantities: the force applied, the lever arm vector^[2] connecting the point about which the torque is being measured to the point of force application, and the angle between the force and lever arm vectors. In symbols:

τ
¯
=

r
¯
×

F
¯

{\displaystyle {\bar {\tau }}={\bar {r}}\times {\bar {F}}}

τ
¯
=
|

r
¯

|

|

F
¯

|

sin
⁡
θ

{\displaystyle {\bar {\tau }}=|{\bar {r}}||{\bar {F}}|\sin \theta }

where

- τ
¯

{\displaystyle {\bar {\tau }}}

 is the torque vector and

r

{\displaystyle r}

 is the magnitude of the torque,
- r

{\displaystyle r}

 is the position vector (a vector from the point about which the torque is being measured to the point where the force is applied)
- F
¯

{\displaystyle {\bar {F}}}

 is the force vector,
- ×

{\displaystyle \times }

 denotes the cross product, which produces a vector that is perpendicular to both

r

{\displaystyle r}

 and

F
¯

{\displaystyle {\bar {F}}}

 following the right-hand rule,
- θ

{\displaystyle \theta }

 is the angle between the force vector and the lever arm vector.

The **net torque** on a body determines the rate of change of the body's angular momentum,

τ
¯
=

d
L
¯

d
t

{\displaystyle {\bar {\tau }}={\frac {d{\bar {L}}}{dt}}}

where

L
¯

{\displaystyle {\bar {L}}}

 is the angular momentum vector and

t

{\displaystyle t}

 is time.

For the motion of a point particle,

L
¯
=
I
ω

{\displaystyle {\bar {L}}=I\omega }

where

I

{\displaystyle I}

 is the moment of inertia and

ω

{\displaystyle \omega }

 is the orbital angular velocity pseudovector. It follows that

τ

net

¯
=

d
L
¯

d
t

=

d
(
I
ω
)

d
t

=
I

d
ω

d
t

+

d
I

d
t

ω
=
I
α
+

d
I

d
t

ω
=
I
α
+
2

r

p

⊥

ω

{\displaystyle {\bar {\tau }}_{\mathrm {net} }={\frac {d{\bar {L}}}{dt}}={\frac {d\left(I\omega \right)}{dt}}=I{\frac {d\omega }{dt}}+{\frac {dI}{dt}}\omega =I\alpha +{\frac {dI}{dt}}\omega =I\alpha +2r_{\perp }\omega }

where

α

{\displaystyle \alpha }

 is the angular acceleration of the particle, and

p

⊥

{\displaystyle p_{\perp }}

 is the radial component of its linear momentum. This equation

Moment of inertia

- The **moment of inertia** is the **2nd moment of mass**.

I
¯
=

r

2

m

{\displaystyle {\bar {I}}=r^{2}m}

 for a point mass,

Torque

τ
¯
=

r
¯
×

F
¯

{\displaystyle {\bar {\tau }}={\bar {r}}\times {\bar {F}}}

τ
¯
=
|

r
¯

|

|

F
¯

|

sin
⁡
θ

{\displaystyle {\bar {\tau }}=|{\bar {r}}||{\bar {F}}|\sin \theta }

```
// test 01
Vector3 force = new Vector3(0, 0, 1) * magnitude;
Vector3 torque = Vector3.Cross(Vector3.right, force);
_rigidbody.AddTorque(torque);
_timer = 5.0f;
```