## legend('exp(t)','exp(0)','exp(-t)') grid hold off exp(at) -- a real 8 exp(t) 7 exp(0) exp(-t) 6 exp(t) and exp(-t) 1 0 -1 -0.5 0 0.5 1.5 2 -1 t (s) You can regenerate this image generated with this Matlab script: expon.m. • When a < 0 the response is a decaying exponential (red line in plot) • When a=0 $e^{at}=1$ -- essentially a model of DC • When a > 0 the response is an *unbounded* increasing exponential (blue line in plot) Case when a is imaginary $e^{j\omega t} = \cos \omega t + j\sin \omega t$ Phasor Plot 0.8 0.5 0.6 0.4 cos(omegat) + jsin(omegat) 0.5 Real 0.4 0.6 0.8 omega t (rad) This is the case that helps us simplify the computation of sinusoidal Fourier series. It was Leonhard Euler who discovered the formula visualized above. Some important values of $\omega t$ These are useful when simplifying expressions that result from integrating functions that involve the imaginary exponential Give the following: $e^{j\omega t}$ when $\omega t = 0$ $e^{j\omega t}$ when $\omega t = \pi/2$ $e^{j\omega t}$ when $\omega t = \pi$ $e^{j\omega t}$ when $\omega t = 3\pi/2$ • $e^{j\omega t}$ when $\omega t = 2\pi$ Case where a is complex We shall not say much about this case except to note that the Laplace transform equation includes such a number. The variable *s* in the Laplace Transform $\int_0^\infty f(t)e^{-st}\,dt$ is a complex exponential. The consequences of a complex s have particular significance in the development of system stability theories and in control systems analysis and design. Look out for them in EG-243. **Two Other Important Properties** By use of trig. identities, it is relatively straight forward to show that: $\cos \omega t = \frac{e^{j\omega t} + e^{-j\omega t}}{2}$ and $\sin \omega t = \frac{e^{j\omega t} - e^{-j\omega t}}{j2}$ We can use this result to convert the *Trigonometric Fourier Series* into an *Exponential Fourier Series* which has only one integral term to solve per harmonic. The Exponential Fourier Series As as stated in the notes on the Trigonometric Fourier Series any periodic waveform f(t) can be represented as $f(t) = \frac{1}{2}a_0 + a_1 \cos \Omega_0 t + a_2 \cos 2\Omega_0 t + \cdots$ $+b_1 \sin \Omega_0 t + b_2 \sin 2\Omega_0 t + \cdots$ If we replace the cos and sin terms with their imaginary expontial equivalents: $f(t) = \frac{1}{2}a_0 + a_1 \left(\frac{e^{j\Omega_0 t} + e^{-j\Omega_0 t}}{2}\right) + a_2 \left(\frac{e^{j2\Omega_0 t} + e^{-j2\Omega_0 t}}{2}\right) + \cdots$ $+b_1\left(\frac{e^{j\Omega_0t}-e^{-j\Omega_0t}}{i2}\right)+b_2\left(\frac{e^{j2\Omega_0t}-e^{-j2\Omega_0t}}{i2}\right)+\cdots$ Grouping terms with same exponents $f(t) = \dots + \left(\frac{a_2}{2} - \frac{b_2}{i2}\right) e^{-j2\Omega_0 t} + \left(\frac{a_1}{2} - \frac{b_1}{i2}\right) e^{-j\Omega_0 t} + \frac{1}{2} a_0 + \left(\frac{a_1}{2} + \frac{b_1}{i2}\right) e^{j\Omega_0 t} + \left(\frac{a_2}{2} + \frac{b_2}{i2}\right) e^{j2\Omega_0 t} + \dots$ **New coefficents** The terms in parentheses are usually denoted as $C_{-k} = \frac{1}{2} \left( a_k - \frac{b_k}{i} \right) = \frac{1}{2} (a_k + jb_k)$ $C_k = \frac{1}{2} \left( a_k + \frac{b_k^2}{i} \right) = \frac{1}{2} (a_k - jb_k)$ $C_0 = \frac{1}{2}a_0$

**Exponential Fourier Series** ¶

This page is downloadable as a PDF file.

The result is called the *Exponential Fourier Series*.

Exponents and Euler's Equation

Symmetry in Exponential Fourier Series

The Exponential Function  $e^{at}$ 

The Exponential Fourier series

Case when a is real.

%% The decaying exponential

title('exp(at) -- a real')

ylabel('exp(t) and exp(-t)')

plot(t, exp(t), t, exp(0.\*t), t, exp(-t))

t=linspace(-1,2,1000);

axis([-1,2,-1,8])

xlabel('t (s)')

figure

In [1]:

An annotatable worksheet for this presentation is available as Worksheet 10.

• The source code for this page is <u>content/fourier\_series/2/exp\_fs1.ipynb</u>.

• You can view the notes for this presentation as a webpage (HTML).

This section builds on our Revision of the to Trigonometrical Fourier Series.

concentrated on the properties and left the computation to a computer.

Trigonometric Fourier series uses integration of a periodic signal multiplied by sines and cosines at the fundamental and

harmonic frequencies. If performed by hand, this can a painstaking process. Even with the simplifications made possible by exploiting waveform symmetries, there is still a need to integrate cosine and sine terms, be aware of and able to exploit the

trigonometric identities, and the properties of orthogonal functions before we can arrive at the simplified solutions. This is why I

However, by exploiting the exponential function  $e^{at}$ , we can derive a method for calculating the coefficients of the harmonics

that is much easier to calculate by hand and convert into an algorithm that can be executed by computer.

• You should already be familiar with  $e^{at}$  because it appears in the solution of differential equations.

It is also a function that appears in the definition of the Laplace and Inverse Laplace Transform.

It pops up again and again in tables and properies of the Laplace Transform.

When a is real the function  $e^{at}$  will take one of the two forms illustrated below:

Colophon

**Agenda** 

Example

From knowledge of the trig. fourier series, even functions have no sine terms so the  $b_k$  coefficients are 0. Therefore both  $C_{-k}$ and  $C_k$  are real. **Odd Functions** 

imaginary.

**Proof** 

Half-wave symmetry

For proof see notes

 $C_{-k}=C_k^st$  always

**Example 1** 

Solved in in Class

Some questions for you

Hence

•  $C_0 = [?]$ 

Solution to example 1

For n odd\*,  $e^{-jk\pi} = -1$ . Therefore

 $^{^{*}}$  You may want to verify that  $C_0=0$  and

From the definition of the exponential Fourier series

E.g. since  $C_3 = 2A/j3\pi$ ,  $C_{-3} = C_3^* = -2A/j3\pi$ 

gathering terms at each harmonic frequency gives

Since

In [2]:

Define f(t)

X =

Plot

In [6]:

In [7]:

/(5\*pi)]

-5

Plot the numerical results from MATLAB calculation.

title('Exponential Fourier Series for Square Waveform with Odd Symmetry')

Exponential Fourier Series for Square Waveform with Odd Symmetry

0

 $\Omega_0$  (rad/sec)

3

3

2

Convert symbolic to numeric result

Xw = subs(X,A,1);

stem(w,abs(Xw), 'o-');

stem(w,angle(Xw), 'o-');

xlabel('\Omega\_0 (rad/sec)');

xlabel('\Omega\_0 (rad/sec)'); ylabel('\angle c\_k [radians]');

subplot(211)

subplot(212)

0.6

0.4

0.2

-2

**Summary** 

Example

Hence

•  $C_0 = 0$ 

-5

<u>ပ</u>

ylabel('|c\_k|');

the exponential Fourier series for the square wave with odd symmetry is

Note sign change in first two terms. This is due to the fact that  $C_{-k}=C_k^st.$ 

Trig. Fourier Series from Exponential Fourier Series

• Square wave is an [odd/even/neither] function?

• Coefficients  $C_k$  are [real/imaginary/complex]?

• What is the integral that needs to be solved for  $C_k$ ?

Square wave [has/does not have] half-wave symmetry?

• Subscripts k are [odd only/even only/both odd and even]?

• DC component is [zero/non-zero]?

**Proof** 

Recall

and

**Even Functions** 

For even functions, all coefficients  $C_k$  are real.

For odd functions, all coefficients  $C_k$  are imaginary.

If there is *half-wave symmetry*,  $C_k = 0$  for k even.

k even. Hence  $C_{-k}$  and  $C_k$  are also zero when k is even.

The Exponential Fourier Series is

or more compactly

**Important** 

or

SO

SO

Similarly

 $f(t) = \dots + C_{-2}e^{-j2\Omega_0t} + C_{-1}e^{-j\Omega_0t} + C_0 + C_1e^{j\Omega_0t} + C_2e^{j2\Omega_0t} + \dots$ 

 $f(t) = \sum_{k=-n}^{n} C_k e^{jk\Omega_0 t}$ 

 $C_{-k} = C_k^*$ 

 $C_k = \frac{1}{2\pi} \int_0^{2\pi} f(\Omega_0 t) e^{-jk(\Omega_0 t)} d(\Omega_0 t)$ 

 $C_k = \frac{1}{T} \int_0^T f(t) e^{-jk\Omega_0 t} dt$ 

\* The analysis that leads to this result is provided between pages 7-31 and 7-32 of the text book. It is not a difficult proof, but

 $C_k + C_{-k} = \frac{1}{2}(a_k - jb_k + a_k + jb_k)$ 

 $a_k = C_k + C_{-k}$ 

 $C_k - C_{-k} = \frac{1}{2}(a_k - jb_k - a_k - jb_k)$ 

 $b_k = j \left( C_k - C_{-k} \right)$ 

Since the coefficients of the Exponential Fourier Series are complex numbers, we can use symmetry to determine the form of

 $C_{-k} = \frac{1}{2} \left( a_k - \frac{b_k}{i} \right) = \frac{1}{2} (a_k + jb_k)$ 

 $C_k = \frac{1}{2} \left( a_k + \frac{b_k}{i} \right) = \frac{1}{2} (a_k - jb_k)$ 

By a similar argument, all odd functions have no cosine terms so the  $a_k$  coefficients are 0. Therefore both  $C_{-k}$  and  $C_k$  are

From Trigonometric Fourier Series, if there is half-wave symmetry, all even harnonics are zero, thus both  $a_k$  and  $b_k$  are zero for

T

 $2\pi$ 

 $\omega t$ 

A

0

-A

 $\frac{1}{2\pi} \left[ \int_{0}^{\pi} A e^{-jk(\Omega_{0}t)} d(\Omega_{0}t) + \int_{0}^{2\pi} (-A)e^{-jk(\Omega_{0}t)} d(\Omega_{0}t) \right] = \frac{1}{2\pi} \left[ \frac{A}{-ik} e^{-jk(\Omega_{0}t)} \Big|_{0}^{\pi} + \frac{-A}{-ik} e^{-jk(\Omega_{0}t)} \Big|_{0}^{2\pi} \right]$ 

 $= \frac{1}{2\pi} \left[ \frac{A}{-ik} \left( e^{-jk\pi} - 1 \right) + \frac{A}{ik} \left( e^{-j2k\pi} - e^{-jk\pi} \right) \right] = \frac{A}{2i\pi k} \left( 1 - e^{-jk\pi} + e^{-j2k\pi} - e^{-jk\pi} \right)$ 

 $\frac{A}{2i\pi k} \left( e^{-j2k\pi} - 2e^{-jk\pi} - 1 \right) = \frac{A}{2i\pi k} \left( e^{-jk\pi} - 1 \right)^2$ 

 $\frac{C_k}{k = \text{odd}} = \frac{A}{2i\pi k} \left( e^{-jk\pi} - 1 \right)^2 = \frac{A}{2i\pi k} (-1 - 1)^2 = \frac{A}{2i\pi k} (-2)^2 = \frac{2A}{i\pi k}$ 

 $\frac{C_k}{k = \text{even}} = 0.$ 

 $f(t) = \dots + C_{-2}e^{-j2\Omega_0 t} + C_{-1}e^{-j\Omega_0 t} + C_0 + C_1e^{j\Omega_0 t} + C_2e^{j2\Omega_0 t} + \dots$ 

 $f(t) = \frac{2A}{i\pi} \left( \dots - \frac{1}{3} e^{-j3\Omega_0 t} - e^{-j\Omega_0 t} + e^{j\Omega_0 t} + \frac{1}{3} e^{j3\Omega_0 t} + \dots \right) = \frac{2A}{i\pi} \sum_{i=1}^{n} \frac{1}{k} e^{jk\Omega_0 t}$ 

 $f(t) = \frac{2A}{i\pi} \left( \dots - \frac{1}{3} e^{-j3\Omega_0 t} - e^{-j\Omega_0 t} + e^{j\Omega_0 t} + \frac{1}{3} e^{j3\Omega_0 t} + \dots \right)$ 

 $f(t) = \frac{4A}{\pi} \left( \dots + \left( \frac{e^{j\Omega_0 t} - e^{-j\Omega_0 t}}{2i} \right) + \frac{1}{3} \left( \frac{e^{j3\Omega_0 t} - e^{-j3\Omega_0 t}}{2i} \right) + \dots \right)$ 

 $= \frac{4A}{\pi} \left( \sin \Omega_0 t + \frac{1}{3} \sin 3\Omega_0 t + \cdots \right)$ 

Computing coefficients of Exponential Fourier Series in MATLAB

 $= \frac{4A}{\pi} \sum_{k=\text{odd}} \frac{1}{k} \sin k\Omega_0 t.$ 

**Exponential Fourier series for the square wave with odd symmetry** 

the coefficients and thereby simplify the computation of series for wave forms that have symmetry.

These are much easier to derive and compute than the equivalent Trigonemetric Fourier Series coefficients.

The  $C_k$  coefficients, except for  $C_0$  are *complex* and appear in conjugate pairs so

**Trigonometric Fourier Series from Exponential Fourier Series** 

**Evaluation of the complex coefficients** 

we are more interested in the result.

The coefficients are obtained from the following expressions\*:

By substituting  $C_{-k}$  and  $C_k$  back into the original expansion

Thus we can easily go back to the Trigonetric Fourier series if we want to.

Symmetry in Exponential Fourier Series

No symmetry If there is no symmetry the Exponential Fourier Series of f(t) is complex. Relation of  $C_{-k}$  to  $C_k$ 

Compute the Exponential Fourier Series for the square wave shown below assuming that  $\omega=1$ 



5

-2 0 2 -3 -1  $\Omega_0$  (rad/sec) ∠ c<sub>k</sub> [radians]

-3

Exponents and Euler's Equation

Symmetry in Exponential Fourier Series

Answers to in-class problems

• When  $\omega t = 3\pi/2$ :  $e^{j\omega t} = e^{j3\pi/2} = -j$ 

• When  $\omega t = 2\pi$ :  $e^{j\omega t} = e^{j2\pi}e^{j0} = 1$ 

above. For example see  $e^{j2\pi}$  above.

The exponential Fourier series

-2

Some important values of  $\omega t$  - Solution • When  $\omega t = 0$ :  $e^{j\omega t} = e^{j0} = 1$ • When  $\omega t = \pi/2$ :  $e^{j\omega t} = e^{j\pi/2} = j$ • When  $\omega t = \pi$ :  $e^{j\omega t} = e^{j\pi} = -1$ 

Some answers for you Square wave is an odd function! DC component is zero! Square wave has half-wave symmetry!

 $C_k = \frac{1}{2\pi} \int_0^{2\pi} f(\Omega_0 t) e^{-jk(\Omega_0 t)} d(\Omega_0 t) = \frac{1}{2\pi} \left[ \int_0^{\pi} A e^{-jk(\Omega_0 t)} d(\Omega_0 t) + \int_{\pi}^{2\pi} (-A) e^{-jk(\Omega_0 t)} d(\Omega_0 t) \right]$ 

• Coefficients  $C_k$  are imaginary! Subscripts k are odd only! • What is the integral that needs to be solved for  $C_k$ ?

It is also worth being aware that  $n\omega t$ , when n is an integer, produces rotations that often map back to the simpler cases given