

Predictive Engineering and Computational Sciences

Locally Conservative Discontinuous Petrov-Galerkin for Fluid Problems

Truman E. Ellis

Institute for Computational and Engineering Sciences
The University of Texas at Austin

July 25, 2013

Truman. E. Ellis Locally Conservative DPG July 25, 2013

1/20

A Summary of DPG

Overview of Features

- Robust for singularly perturbed problems
- Stable in the preasymptotic regime
- Designed for adaptive mesh refinement

DPG is a minimum residual method:

$$u_h = \underset{w_h \in U_h}{\arg \min} \frac{1}{2} \|Bw_h - I\|_{V'}^2$$

$$\downarrow \downarrow$$

$$b(u_h, R_V^{-1} B \delta u_h) = I(R_V^{-1} B \delta u_h) \quad \forall \delta u_h \in U_h$$

where $v_{\delta u_h} := R_V^{-1} B \delta u_h$ are the optimal test functions.

DPG for Convection-Diffusion

Start with the strong-form PDE.

$$\nabla \cdot (\beta u) - \epsilon \Delta u = g$$

Rewrite as a system of first-order equations.

$$abla \cdot (eta u - oldsymbol{\sigma}) = g$$

$$\frac{1}{\epsilon} oldsymbol{\sigma} - \nabla u = \mathbf{0}$$

Multiply by test functions and integrate by parts over each element, K.

$$-(\beta u - \sigma, \nabla v)_{K} + ((\beta u - \sigma) \cdot \mathbf{n}, v)_{\partial K} = (g, v)_{K}$$
$$\frac{1}{\epsilon} (\sigma, \tau)_{K} + (u, \nabla \cdot \tau)_{K} - (u, \tau_{n})_{\partial K} = 0$$

Use the ultraweak (DPG) formulation to obtain bilinear form b(u, v) = I(v).

$$-(\beta u - \sigma, \nabla v)_{K} + (\hat{f}, v)_{\partial K} + \frac{1}{\epsilon} (\sigma, \tau)_{K} + (u, \nabla \cdot \tau)_{K} - (\hat{u}, \tau_{n})_{\partial K} = (g, v)_{K}$$

Local Conservation

The local conservation law in convection diffusion is

$$\int_{\partial K} \hat{f} = \int_K g,$$

which is equivalent to having $\mathbf{v}_{\mathcal{K}} := \{v, \tau\} = \{\mathbf{1}_{\mathcal{K}}, \mathbf{0}\}$ in the test space. In general, this is not satisfied by the optimal test functions. Following Moro et al^[5] (also Chang and Nelson^[2]), we can enforce this condition with Lagrange multipliers:

$$L(u_h, \lambda) = \frac{1}{2} \left\| R_V^{-1}(Bu_h - I) \right\|_V^2 - \sum_K \lambda_K \underbrace{\langle Bu_h - I, \mathbf{v}_K \rangle}_{\langle \hat{f}, \mathbf{1}_K \rangle_{\partial K} - \langle g, \mathbf{1}_K \rangle_K},$$

where $\lambda = \{\lambda_1, \cdots, \lambda_N\}.$

Local Conservation

Finding the critical points of $L(u, \lambda)$, we get the following equations^[4].

$$\frac{\partial L(u_h, \lambda)}{\partial u_h} = b(u_h, R_V^{-1} B \delta u_h) - I(R_V^{-1} B \delta u_h) - \sum_K \lambda_K b(\delta u_h, \mathbf{v}_K) = 0 \quad \forall \delta u_h \in U_h$$

$$\frac{\partial L(u_h, \boldsymbol{\lambda})}{\partial \lambda_K} = -b(u_h, \boldsymbol{v}_K) + I(\boldsymbol{v}_K) = 0 \quad \forall K$$

A few consequences:

- We've turned our minimization problem into a saddlepoint problem.
- Only need to find the optimal test function in the orthogonal complement of constants.

Optimal Test Functions

For each $\mathbf{u} = \{u, \sigma, \hat{u}, \hat{t}\} \in \mathbf{U}_h$, find $\mathbf{v_u} = \{v_u, \tau_u\} \in \mathbf{V}$ such that

$$(\mathbf{v}_{\mathbf{u}}, \mathbf{w})_{\mathbf{V}} = b(\mathbf{u}, \mathbf{w}) \quad \forall \mathbf{w} \in \mathbf{V}$$

where **V** becomes $\mathbf{V}_{p+\Delta p}$ in order to make this computationally tractable. We recently developed this modification to the *robust test norm* ^[1] which behaves better in the presence of singularities.

$$\begin{aligned} \|(\boldsymbol{v}, \boldsymbol{\tau})\|_{\mathbf{V}, \Omega_h}^2 &= \left\|\min\left\{\frac{1}{\sqrt{\epsilon}}, \frac{1}{\sqrt{|K|}}\right\} \boldsymbol{\tau}\right\|^2 + \|\nabla \cdot \boldsymbol{\tau} - \boldsymbol{\beta} \cdot \nabla \boldsymbol{v}\|^2 \\ &+ \|\boldsymbol{\beta} \cdot \nabla \boldsymbol{v}\|^2 + \epsilon \|\nabla \boldsymbol{v}\|^2 \underbrace{\qquad \qquad + \|\boldsymbol{v}\|^2}_{\text{No longer necessary}} \end{aligned}$$

Optimal Test Functions

For each $\mathbf{u} = \{u, \sigma, \hat{u}, \hat{f}\} \in \mathbf{U}_h$, find $\mathbf{v}_{\mathbf{u}} = \{v_{\mathbf{u}}, \tau_{\mathbf{u}}\} \in \mathbf{V}$ such that

$$(\mathbf{v}_{\mathbf{u}}, \mathbf{w})_{\mathbf{V}} = b(\mathbf{u}, \mathbf{w}) \quad \forall \mathbf{w} \in \mathbf{V}$$

where **V** becomes $\mathbf{V}_{p+\Delta p}$ in order to make this computationally tractable. We recently developed this modification to the *robust test norm* ^[1] which behaves better in the presence of singularities.

$$\begin{aligned} \|(\boldsymbol{v}, \boldsymbol{\tau})\|_{\mathbf{V}, \Omega_h}^2 &= \left\| \min \left\{ \frac{1}{\sqrt{\epsilon}}, \frac{1}{\sqrt{|K|}} \right\} \boldsymbol{\tau} \right\|^2 + \|\nabla \cdot \boldsymbol{\tau} - \boldsymbol{\beta} \cdot \nabla \boldsymbol{v}\|^2 \\ &+ \|\boldsymbol{\beta} \cdot \nabla \boldsymbol{v}\|^2 + \epsilon \|\nabla \boldsymbol{v}\|^2 \underbrace{+ \left(\frac{1}{|K|} \int_{K} \boldsymbol{v}\right)^2}_{\text{Zero mean term}} \end{aligned}$$

Stability Analysis

We follow Brezzi's theory for an abstract mixed problem:

$$\begin{cases} \mathbf{u} \in \mathbf{U}, p \in Q \\ \mathbf{a}(\mathbf{u}, \mathbf{w}) + \mathbf{c}(p, \mathbf{w}) &= l(\mathbf{w}) \quad \forall \mathbf{w} \in \mathbf{U} \\ \mathbf{c}(q, \mathbf{u}) &= g(q) \quad \forall q \in Q \end{cases}$$

where a, c, l, g denote the appropriate bilinear and linear forms. Note that $a(\mathbf{u}, \mathbf{w}) = b(\mathbf{u}, T\mathbf{w}) = (T\mathbf{u}, T\mathbf{w})_V$.

Let ψ denote the $\mathbf{H}(\text{div}, \Omega)$ extension of flux \hat{f} that realizes the minimum in the definition of the quotient (minimum energy extension) norm.

The norm for the Lagrange multipliers λ_K is implied by the quotient norm used for $H^{-1/2}(\Gamma_h)$ and continuity bound for form $c(p, \mathbf{w})$:

$$||\boldsymbol{\lambda}|| := \left(\sum_{\mathcal{K}} \mu(\mathcal{K}) \lambda_{\mathcal{K}}^2\right)^{1/2}$$

Inf Sup Condition

The inf-sup condition relating spaces \boldsymbol{U} and Q is

$$\sup_{\boldsymbol{w}\in\boldsymbol{U}}\frac{|c(p,\boldsymbol{w})|}{||\boldsymbol{w}||_{\boldsymbol{U}}}\geq \beta||p||_{Q}$$

Let

$$R: L^2(\Omega) \ni q \to \psi \in \textit{\textbf{H}}(\mathsf{div},\Omega) \cap \textit{\textbf{H}}^1(\Omega) = \textit{\textbf{H}}^1(\Omega)$$

be the continuous right inverse of the divergence operator constructed by Costabel and McIntosh $^{[3]}$. Let ψ_h denote the classical, lowest order Raviart-Thomas (RT) interpolant of function

$$\psi = R(\sum_K \lambda_K 1_K)$$

Note that $\operatorname{div}\psi_h = \operatorname{div}\psi = \lambda_K$ in element K.

Inf Sup Condition

Classical *h*-interpolation interpolation error estimate for the lowest error Raviart-Thomas elements and continuity of operator *R* imply the stability estimate:

$$\begin{split} ||\psi_h|| & \leq ||\psi_h - \psi|| + ||\psi|| \\ & \leq Ch||\psi||_{H^1} + ||\psi|| \\ & \leq C||\mathsf{div}\psi|| = C(\sum_K \mu(K)\lambda_K^2)^{1/2} \end{split}$$

Let \hat{f} be the trace of ψ_h , then

$$\begin{split} \sup_{\hat{f} \in H^{-1/2}(\Gamma_h)} \frac{|\sum_K \lambda_K \langle \hat{f}, \mathbf{1}_K \rangle_{\partial K}|}{||\hat{f}||_{H^{-1/2}(\Gamma_h)}} &\geq \frac{|\sum_K \lambda_K \int_K \operatorname{div} \psi_h \, \mathbf{1}_K|}{||\psi_h||_{H(\operatorname{div},\Omega)}} \\ &\geq \frac{1}{C} (\sum_K \mu(K) \lambda_K^2)^{1/2} \end{split}$$

Inf Sup in Kernel Condition

We characterize the "kernel" space:

$$\mathbf{U}_0 := \{ \mathbf{w} \in \mathbf{U} : c(q, \mathbf{w}) = 0 \quad \forall q \in Q \}$$

$$= \{ (u, \sigma, \hat{u}, \hat{t}) : \langle \hat{t}, 1_K \rangle = 0 \quad \forall K \}$$

With $\boldsymbol{u} \in \boldsymbol{U}_0$, we have then:

$$\begin{split} \sup_{\boldsymbol{w} \in \boldsymbol{U}_0} \frac{|\boldsymbol{a}(\boldsymbol{u}, \boldsymbol{w})|}{||\boldsymbol{w}||_{\boldsymbol{U}}} &\geq \frac{|\boldsymbol{b}(\boldsymbol{u}, T\boldsymbol{u})|}{||\boldsymbol{u}||} = \frac{|\boldsymbol{b}(\boldsymbol{u}, T\boldsymbol{u})|}{||T\boldsymbol{u}||} \frac{||T\boldsymbol{u}||}{||\boldsymbol{u}||} \\ &= \sup_{(\boldsymbol{v}, \boldsymbol{\tau})} \frac{|\boldsymbol{b}((\boldsymbol{u}, \boldsymbol{\sigma}, \hat{\boldsymbol{u}}, \hat{\boldsymbol{t}}), (\boldsymbol{v}, \boldsymbol{\tau}))|}{||(\boldsymbol{v}, \boldsymbol{\tau})||} \frac{||T\boldsymbol{u}||}{||\boldsymbol{u}||} \geq \gamma^2 ||(\boldsymbol{u}, \boldsymbol{\sigma}, \hat{\boldsymbol{u}}, \hat{\boldsymbol{t}})|| \end{split}$$

where γ is the stability constant for the standard continuous DPG formulation.

The FE error is bounded by the best approximation error. Note that the exact Lagrange multipliers are zero, so the best approximation error involves only the solution $(u, \sigma, \hat{u}, \hat{t})$.

- We prove robustness of the restricted DPG method by switching to the energy norm in Brezzi's stability analysis.
- The inf-sup in kernel condition is simple. Upon replacing the original norm of solution ${\bf u}$ with the energy norm, γ and the continuity constant become one.
- In order to investigate the robustness of inf-sup constant β , we need to understand what the energy norm of flux variable \hat{f} is.
- For an element, K, we solve for the optimal test functions, $v_K \in H^1(K)$, and $\tau_K \in H(\text{div}, K)$ corresponding to flux \hat{f} :

$$((v_K, au_K), (\delta v, \delta au))_V = \left\langle \hat{f}, \delta v \right\rangle \quad \forall \delta v \in H^1(K), \ \delta au \in H(\operatorname{div}, K)$$

• The energy norm for \hat{f} is then

$$||\hat{f}||_E^2 = \sum_K ||(v_K, \tau_K)||_V^2$$

 Need to establish conditions under which the inf-sup constant is independent of viscosity.

$$\sup_{\hat{f}} \frac{|\sum_K \lambda_K \langle \hat{f}, \mathbf{1}_K \rangle|}{||\hat{f}||_E} \geq \beta \left(\sum_K \mu(K) \lambda_K^2\right)^{1/2}$$

- Select \hat{f} as the trace of the Raviart-Thomas interpolant ψ_h of $\psi = R(\sum_K \lambda_K 1_K)$.
- Proceed as in the previous analysis, but evaluation of the norm of \hat{f} requires a local solve.

$$((v, \tau), (\delta v, \delta \tau))_{V} = \left\langle \hat{f}, \delta v \right\rangle_{\partial K} = \int_{K} \operatorname{div} \psi_{h} \delta v = \int_{K} \operatorname{div} \psi \delta v$$
$$= \int_{K} \lambda_{K} \delta v = \lambda_{K} (1_{K}, \delta v)_{K} \quad \forall \delta v \in H^{1}(K) \, \forall \delta \tau \in \mathbf{H}(\operatorname{div}, K)$$

• We need an upper bound for the energy norm of (v_h, τ_h) :

$$\left(\sum_{K}||(v,\tau)||_{V}^{2}\right)^{1/2}$$

• Substituting (v, τ) for $(\delta v, \delta \tau)$ in the previous slide:

$$||(\mathbf{v}, \boldsymbol{\tau})||_V^2 = \lambda_K(\mathbf{1}_K, \mathbf{v}_k)$$

• With a robust stability estimate $(1_K, v_K) \leq C\mu(K)^{1/2}||(v, \tau)||_K$,

$$||(\mathbf{v}, \boldsymbol{\tau})||_{V} \leq C\mu(K)^{1/2}|\lambda_{K}|$$

$$\sum_{K} ||(v,\tau)||_{V}^{2} \leq C^{2} \sum_{K} \mu(K) \lambda_{K}^{2}$$

which leads to the robust estimate of inf-sup constant β .

Finally, we need a continuity estimate on

$$\sum_{K} \lambda_{K} \langle \hat{\mathbf{f}}, \mathbf{1}_{K} \rangle$$

• Testing with $(1_K, \mathbf{0})$ in the local problem, we obtain

$$((\mathbf{v},\boldsymbol{\tau}),(\mathbf{1}_K,\mathbf{0}))_V=\langle \hat{\mathbf{f}},\mathbf{1}_K\rangle_{\partial K}$$

• With a robust estimate $|((v, \tau), (1_K, \mathbf{0}))_V| \leq C\mu(K)^{1/2}||(v, \tau)||_V$,

$$egin{aligned} |\sum_{\mathcal{K}} \lambda_{\mathcal{K}} \langle \hat{f}, 1_{\mathcal{K}} \rangle| &\leq C \left(\sum_{\mathcal{K}} \mu(\mathcal{K}) \lambda_{\mathcal{K}}^2 \right)^{1/2} \left(\sum_{\mathcal{K}} ||(v, au)||_V^2 \right)^{1/2} \\ &= C \left(\sum_{\mathcal{K}} \mu(\mathcal{K}) \lambda_{\mathcal{K}}^2 \right)^{1/2} ||\hat{f}||_E \end{aligned}$$

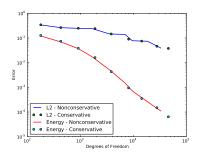
Erickson-Johnson Problem

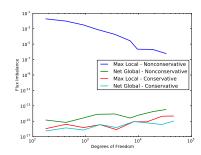
On domain $\Omega = [0, 1]^2$, with $\beta = (1, 0)^T$, f = 0 and boundary conditions

$$\hat{f} = u_0, \quad \beta_n \leq 0, \qquad \hat{u} = 0, \quad \beta_n > 0$$

Separation of variabes gives an analytic solution

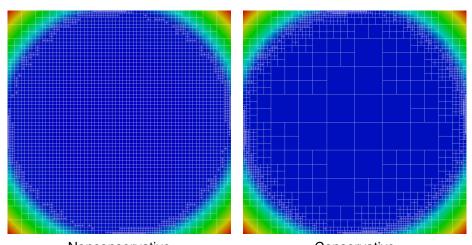
$$u(x,y) = C_0 + \sum_{n=1}^{\infty} C_n \frac{\exp(r_2(x-1)) - \exp(r_1(x-1))}{r_1 \exp(-r_2) - r_2 \exp(-r_1)} \cos(n\pi y)$$





Vortex Problem

After 6 refinements, $\epsilon = 10^{-4}$, $\beta = (-y, x)^T$

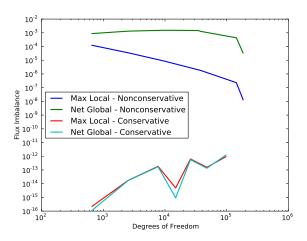


Nonconservative

Conservative

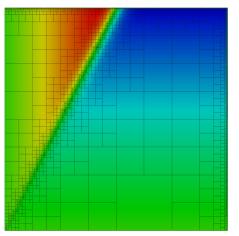
Vortex Problem

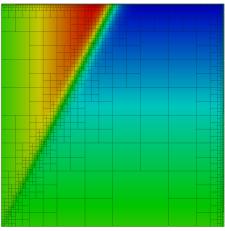
After 6 refinements, $\epsilon = 10^{-4}$, $\beta = (-y, x)^T$



Discontinuous Source Problem

After 8 refinements,
$$\epsilon = 10^{-3}$$
, $\beta = (0.5, 1)^T / \sqrt{1.25}$, $\hat{g} = \begin{cases} 1, & y >= 2x \\ 0, & y < 2x \end{cases}$





Nonconservative

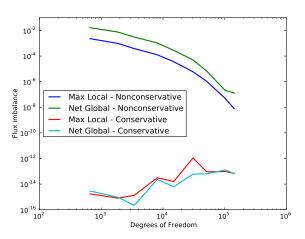
Conservative

Truman, E. Ellis Locally Conservative DPG

July 25, 2013

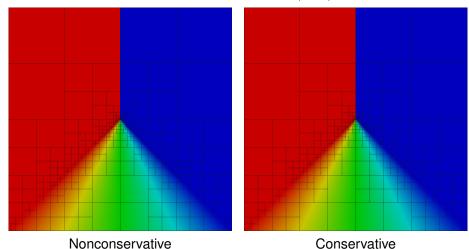
Discontinuous Source Problem

After 8 refinements,
$$\epsilon = 10^{-3}$$
, $\beta = (0.5, 1)^T / \sqrt{1.25}$, $\hat{g} = \begin{cases} 1, & y >= 2x \\ 0, & y < 2x \end{cases}$



Inviscid Burgers' Equation

$$\frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} = 0 \quad \Leftrightarrow \quad \nabla_{x,t} \cdot \left(\begin{array}{c} \frac{u^2}{2} \\ u \end{array} \right) = 0$$

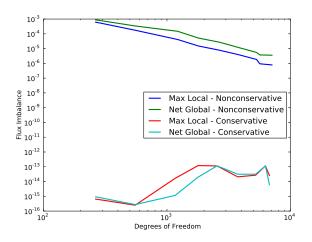


Truman. E. Ellis Locally Conservative DPG

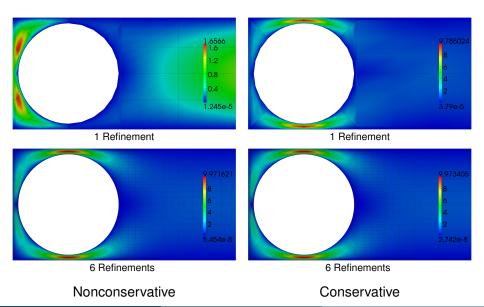
July 25, 2013

17/20

Inviscid Burgers' Equation



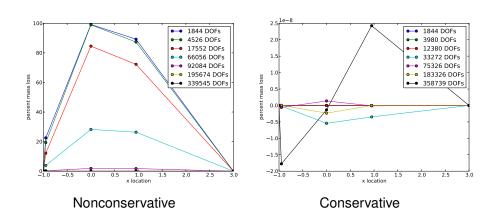
Stokes Flow Around a Cylinder



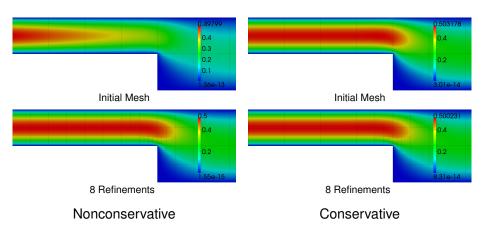
Truman. E. Ellis Locally Conservative DPG July 25, 2013

18 / 20

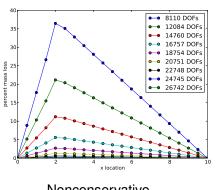
Stokes Flow Around a Cylinder



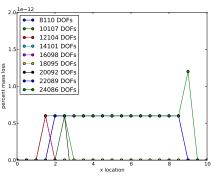
Stokes Flow Backward Facing Step



Stokes Flow Backward Facing Step



Nonconservative



Conservative

Summary

What have we done?

- We've turned our minimization problem into a saddlepoint problem.
- The change is computationally feasible.
- · Mathematically, it gets rid of troublesome term.

Does it make a difference?

- Enforcement changes refinement strategy.
- Some improvement on condition number for local solves.
- Standard DPG is nearly conservative for many of the problems considered, but seems to suffer from mass loss similar to other LSFEM methods.
- For some problems, local conservation allows us to converge to a reasonable solution on a much coarser mesh.



J. Chan, N. Heuer, T Bui-Thanh, and L. Demkowicz.

Robust DPG method for convection-dominated diffusion problems II: A natural inflow condition. Technical Report 21, ICES, 2012.



C. L. Chang and John J. Nelson.

Least-squares finite element method for the Stokes problem with zero residual of mass conservation.

SIAM J. Num. Anal., 34:480-489, 1997.



M. Costabel and A. McIntosh.

On Bogovskiĭ and regularized Poincaré integral operators for de Rham complexes on Lipschitz domains.

Mathematische Zeitschrift, 265(2):297-320, 2010.



T.E. Ellis, L. Demkowicz, and J. Chan.

Locally conservative discontinuous Petrov-Galerkin finite elements for fluid problems. Technical Report xx, ICES, 2013.



D. Moro, N.C. Nguyen, and J. Peraire.

A hybridized discontinuous Petrov-Galerkin scheme for scalar conservation laws. *Int.J. Num. Meth. Eng.*, 2011.