



ECOLE POLYTECHNIQUE FÉDÉRALE DE LAUSANNE

$\mathbf{N} \mathbf{I} \mathbf{A} \mathbf{O} \mathbf{m} \mathbf{n} \mathbf{n}$	TITECTO
MASTER	I HESIS

Evaluation of Optical Aberrations using Phase Diversity

Author:
Jordan VOIRIN

Supervisors:
Dr. Laurent JOLISSAINT
Dr. Jean-Paul KNEIB

A thesis submitted in fulfillment of the requirements for the degree of Master in Applied Physics

in the

Astrophysics laboratory Basic Sciences

February 21, 2018

Declaration of Authorship

I, Jordan VOIRIN, declare that this thesis titled, "Evaluation of Optical Aberrations using Phase Diversity" and the work presented in it are my own. I confirm that:

- This work was done wholly or mainly while in candidature for a research degree at this University.
- Where any part of this thesis has previously been submitted for a degree or any other qualification at this University or any other institution, this has been clearly stated.
- Where I have consulted the published work of others, this is always clearly attributed.
- Where I have quoted from the work of others, the source is always given. With the exception of such quotations, this thesis is entirely my own work.
- I have acknowledged all main sources of help.
- Where the thesis is based on work done by myself jointly with others, I have made clear exactly what was done by others and what I have contributed myself.

Signed:			
Date:			

"For the sake of persons of ... different types, scientific truth should be presented in different forms, and should be regarded as equally scientific, whether it appears in the robust form and the vivid coloring of a physical illustration, or in the tenuity and paleness of a symbolic expression."

James Clerk Maxwell

Acknowledgements

The acknowledgments and the people to thank go here, don't forget to include your project advisor. . .

Contents

D	eclara	tion of Authorship	iii
A	cknov	vledgements	vii
In	trodu	ction	1
1	The	oretical background	3
	1.1	Scalar Diffraction Theory	3
		1.1.1 Scalar Field and Helmholtz equation	
		1.1.2 Rayleigh-Sommerfeld integral	3
		1.1.3 Fresnel approximation	4
		1.1.4 Fraunhofer approximation	5
		1.1.5 Converging lens introduction	5
	1.2	Imaging system	
		1.2.1 Impulse Response (IR)	
		1.2.2 Optical Transfer Function (OTF)	
		1.2.3 From Object to Image	
	1.3	Aberrations	9
		1.3.1 Sources of aberrations	9
		1.3.2 Zernike polynomials	9
2	Pha	se Diversity Experiment	11
	2.1	ONERA algorithm	11
	2.2	Experimental Setup	
		2.2.1 Light source	
		2.2.2 Entrance pupil	
		2.2.3 Pupil imaging system	
		2.2.4 Detectors	
	2.3	Data Acquisition	13
		2.3.1 Ximea Camera	13
		2.3.2 Shack-Hartman WFS	13
	2.4	Results	13
		2.4.1 Parallel plane plate	14
Δ	Ont	ical Component Datasheets	15
11	_	Pigtailed laser diode	16
	21.1	A.1.1 Power supply modification	
	A 2	Converging lens A220TM-A, f = 11 mm	
		Pinhole $10 \ \mu \text{m}$	
	A.4	Converging lens AL100200, $f = 200 \text{ mm}$	
	A.5	Converging lens AC254-100-A, f = 100 mm	
	A.6	Ximea Camera, MQ013MG-E2	
	A.7		
	1 1.7		

В	B Python Code				
	B.1	AlignementScriptXimeaCamera.py	23		
	B.2	AcquisAndSaveXimea.py	26		
Bi	bliog	raphy	29		

List of Figures

1.1	Diffraction Schemas	4
1.2	Schema of a imaging instrument, (Goodman, 1968, Chapter 6.1)	6
1.3	PSF of a perfect imaging system composed by a 3.6 mm pupil and a	
	focal length of 80 mm at a wavelength of 637.5 nm. The size, N, of the	
	PSF is 400 and the pixel size is $5.3 \mu m$.	7
1.4	OTF of a perfect imaging system composed by a 3.6 mm pupil and a	
	focal length of 80 mm at a wavelength of 637.5 nm	7
2.1	Experimental Setup Schema	11
2.3	Source schema and pinhole effect on the beam	12
2.2	Wavefront curvature	12
2.4	PSFs example of an alignment procedure	13

List of Tables

2.1 Optical Components		12
------------------------	--	----

Introduction

???

Chapter 1

Theoretical background

In this chapter, we will present the theory upon which this work is based. First, the light propagation formalism will be reminded through the scalar diffraction theory based on the Goodman (1968). Then we will describe in general an imaging system and its properties. And finally we will discuss the wavefront aberration theory.

1.1 Scalar Diffraction Theory

1.1.1 Scalar Field and Helmholtz equation

A monochromatic wave, at position P and time t, can be represented by a scalar field u(P,t) written as :

$$u(P,t) = A(P)exp\left[-j2\pi\nu t + j\phi(P)\right], \tag{1.1}$$

where A(P) and $\phi(P)$ are the amplitude and phase, respectively, of the wave at position P and ν is the wave frequency.

The spatial part of eqt. (1.1), also called phasor in the literature,

$$U(P) = A(P)e^{j\phi(P)}, \tag{1.2}$$

must verify the Helmotz equation:

$$(\nabla^2 + k^2)U = 0, (1.3)$$

where k is the wave number given by

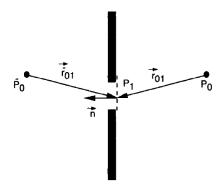
$$k = 2\pi n \frac{\nu}{c} = \frac{2\pi}{\lambda},\tag{1.4}$$

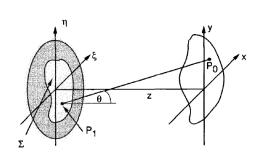
and λ is the wavelength in the dielectric medium.

1.1.2 Rayleigh-Sommerfeld integral

Rayleigh and Sommerfeld developed a formalism using the Helmholtz equation and Green's Theorem to compute the induced diffraction by a plane screen. Let's suppose that we have a monochromatic source at \widetilde{P}_0 on the left of a plane screen with aperture Σ , the Rayleigh-Sommerfeld formula allows to compute the complex amplitude at P_0 on the right of the plane screen (see Figure 1.1a).

$$U(P_0) = \frac{1}{j\lambda} \iint_{\Sigma} U'(P_1) \frac{exp(jkr_{01})}{r_{01}} cos(\mathbf{n}, \mathbf{r_{01}}) d\mathbf{s}$$
 (1.5)





- (A) Rayleigh-Sommerfeld formulation of diffraction by a plane screen, (Goodman, 1968, Chapter 3.5).
- (B) Diffraction geometry, (Goodman, 1968, Chapter 4.1).

FIGURE 1.1: Diffraction Schemas

 $U'(P_1)$ is the complex amplitude on the screen, $cos(\mathbf{n}, \mathbf{r_{01}})$ is the cosine of the angle between the aperture plane normal toward the source and the vector \mathbf{r}_{01} = P_0P_1 given by

$$r_{01} = \sqrt{z^2 + (x - \xi)^2 + (y - \eta)^2}.$$
 (1.6)

We can rewrite eqt. (1.5) using $cos(\mathbf{n}, \mathbf{r}_{01}) = cos(\theta) = \frac{z}{r_{01}}$ and the coordinate systems (ξ, η) and (x, y), see Figure 1.1b,

$$U(x,y) = \frac{z}{j\lambda} \iint_{-\infty}^{\infty} U(\xi,\eta) \frac{exp(jkr_{01})}{r_{01}^2} d\xi d\eta.$$
 (1.7)

We can integrate from $-\infty$ to ∞ , using $U(\xi, \eta) = P(\xi, \eta)U'(\xi, \eta)$ where $P(\xi, \eta)$ is the pupil function. The latter equals to one in the pupil and zero outside.

Fresnel approximation

To reduce eqt. (1.7), also known as the Huygens-Fresnel principle, one can approximate the distance r_{01} using the taylor expansion of the square root :

$$r_{01} = z\sqrt{1 + \frac{x - \xi}{z} + \frac{y - \eta}{z}} \approx z \left[1 + \frac{1}{2} \left(\frac{x - \xi}{z} \right)^2 + \frac{1}{2} \left(\frac{y - \eta}{z} \right)^2 \right]$$
 (1.8)

To obtain the Fresnel approximation, one has to replace r_{01} by eqt. (1.8) in eqt. (1.7). At the denominator, only the first term z is kept, since the introduced error is small, but in the exponential everything is kept. Then the final expression is given by,

$$U(x,y) = \frac{e^{jkz}}{j\lambda z} \int_{-\infty}^{\infty} U(\xi,\eta) exp\left\{ j\frac{k}{2z} \left[(x-\xi)^2 + (y-\eta)^2 \right] \right\} d\xi d\eta. \tag{1.9}$$

In this form, the Fresnel approximation can be seen as a convolution between $U(\xi, \eta)$ and $h(x, y) = \frac{e^{jkz}}{j\lambda z} exp\left[\frac{jk}{2z}(x^2 + y^2)\right]$. Another form is found by developing $\left[(x - \xi)^2 + (y - \eta)^2\right]$,

$$U(x,y) = \frac{e^{jkz}}{j\lambda z} e^{j\frac{k}{2z}(x^2 + y^2)} \iint_{-\infty}^{\infty} \left\{ U(\xi, \eta) e^{j\frac{k}{2z}(\xi^2 + \eta^2)} \right\} e^{-j\frac{2\pi}{\lambda z}(x\xi + y\eta)} d\xi d\eta, \tag{1.10}$$

it is the Fourier transform of the complex field in the pupil multiplied by a quadratic phase exponential.

1.1.4 Fraunhofer approximation

In addition to the Fresnel approximation, we can introduce another approximation using the condition,

$$z >> \frac{k(\xi^2 + \eta^2)_{max}}{2}. (1.11)$$

If eqt. (1.11) is satisfied the Fresnel approximation simplifies, since the quadratic phase factor in (ξ, η) is approximately one on the entire pupil, as

$$U(x,y) = \frac{e^{jkz}}{j\lambda z} e^{j\frac{k}{2z}(x^2 + y^2)} \iint_{\infty} U(\xi,\eta) e^{-j\frac{2\pi}{\lambda z}(x\xi + y\eta)} d\xi d\eta.$$
 (1.12)

For instance, at a wavelength of 637.5 nm and a pupil diameter of 3.6 mm the Fraunhofer approximation constrain *z* to be greater than 63 meters to be valid.

1.1.5 Converging lens introduction

The Fraunhofer conditions are severe as shown above, but one can reduce the distance z by observing at the focal plane of a converging lens. Indeed, using the paraxial approximation, i.e. small angles with respect with the optical axis, the lens transmission function is given by,

$$t_{l}(\xi,\eta) = \exp\left[jkn\Delta_{0}\right] \exp\left[-jk\left(n-1\right) \frac{\xi^{2}+\eta^{2}}{2} \left(\frac{1}{R_{1}} - \frac{1}{R_{2}}\right)\right]$$
$$= \exp\left[-j\frac{k}{2f}(\xi^{2}+\eta^{2})\right], \tag{1.13}$$

where n is the refractive index of the lens material, R_1 and R_2 are the radii of curvature of the front and back surface of the lens, respectively and f is the focal length of the lens defined as,

$$\frac{1}{f} \equiv (n-1) \left(\frac{1}{R_1} - \frac{1}{R_2} \right) \tag{1.14}$$

We can define $U_l(\xi, \eta) = U(\xi, \eta)t_l(\xi, \eta)$, which represents the complex amplitude passing through a lens. Finally, replacing $U(\xi, \eta)$ by $U_l(\xi, \eta)$ in Fresnel approximation and setting the observing distance to the focal length of the converging lens, we recover the Fraunhofer approximation,

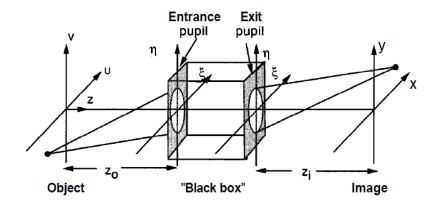


FIGURE 1.2: Schema of a imaging instrument, (Goodman, 1968, Chapter 6.1)

$$U(x,y) = \frac{e^{jkz}}{j\lambda z} e^{j\frac{k}{2z}(x^2+y^2)} \mathcal{F} \left\{ U(\xi,\eta) exp \left[j\frac{k}{2}(\xi^2+\eta^2)(\frac{1}{z}-\frac{1}{f}) \right] \right\}$$

$$\stackrel{z=f}{=} \frac{e^{jkz}}{j\lambda z} e^{j\frac{k}{2z}(x^2+y^2)} \mathcal{F} \left\{ U(\xi,\eta) \right\}. \tag{1.15}$$

1.2 Imaging system

An imaging system, such as a telescope, is used to acquire images of an object as perfectly as possible. An optical system, forming an instrument, is composed by lenses, mirrors, etc... and a detector (can be the human eye). A complex optical system can be reduced to a pupil, $P(\xi, \eta)$, and a focal length, f. The diffraction of the wave can be determined by the Fraunhofer approximation as long as the paraxial approximation is valid, see subsection 1.1.5. And the observed image of an incoherent object at the focal plane of the system is proportional to the square modulus of the complex amplitude U(x,y),

1.2.1 Impulse Response (IR)

The impulse response or point spread function (PSF), h(x, y; u, v), of an optical system is the field amplitude induced at coordinates (x, y) by a unit-amplitude point source at object coordinates (u, v). Using the linearity of the wave propagation, we can write the imaged amplitude as the superposition integral,

$$U(x,y) = \iint_{-\infty}^{\infty} h(x,y;u,v)U(u,v)dudv$$
 (1.16)

1.2.2 Optical Transfer Function (OTF)

The optical transfer function, OTF, is defined as the Fourier transform of the impulse response, see Figure 1.4. Using the Fourier transform properties it can also be given by the autocorrelation of the pupil function,

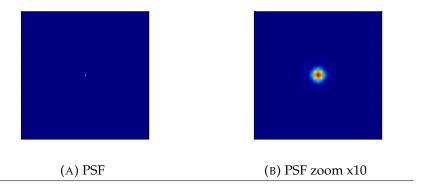


Figure 1.3: PSF of a perfect imaging system composed by a 3.6 mm pupil and a focal length of 80 mm at a wavelength of 637.5 nm. The size, N, of the PSF is 400 and the pixel size is $5.3~\mu m$.

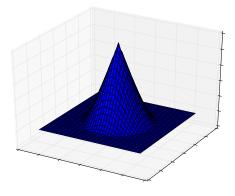


FIGURE 1.4: OTF of a perfect imaging system composed by a 3.6 mm pupil and a focal length of 80 mm at a wavelength of 637.5 nm.

$$\stackrel{\sim}{h_{optical}}(\xi,\eta) = \mathcal{F}\left\{h_{optical}(x,y)\right\} = (P \otimes P)(\xi,\eta),\tag{1.17}$$

where ξ and η are the conjugate variables of x and y with respect to the Fourier transform.

1.2.3 From Object to Image

A detector only senses the energy distribution produced by an electromagnetic wave. Therefore, an image is given by the square modulus of the complex amplitude at the focal plane,

$$i(x,y) = |U(x,y)|^2 = |\int_{-\infty}^{\infty} h(x,y;u,v)U(u,v)dudv|^2.$$
 (1.18)

This integral simplifies differently depending on the type of object we are observing. For a coherent object, Goodman (1968, Chapter 6.2) showed that the imaging is linear in complex amplitude, thus eqt. (1.18) becomes,

$$i(x,y) = \left| \iint_{-\infty}^{\infty} h(x - \widetilde{u}, y - \widetilde{v}) U(\widetilde{u}, \widetilde{v}) d\widetilde{u} d\widetilde{v} \right|^{2}, \tag{1.19}$$

where $(\tilde{u} = Mu, \tilde{v} = Mv)$ are the normalized object coordinates and M is the magnification of the imaging system. The image is given by the squared modulus of the convolution of the impulse response and the object complex amplitude.

For an incoherent object, Goodman (1968, Chapter 6.2) showed that the imaging is linear in intensity,

$$i(x,y) = \iint_{-\infty}^{\infty} |h(x - \widetilde{u}, y - \widetilde{v})|^2 o(\widetilde{u}, \widetilde{v}) d\widetilde{u} d\widetilde{v} = (h_{optical} \otimes o)(x, y), \tag{1.20}$$

where $o(x,y) = |U(x,y)|^2$. We can recognize the convolution of an object with an intensity impulse response, $h_{optical}(x,y) = |h(x,y)|^2$.

In this study, we will study stars radiation which are object with an incoherent emission. The stars are point source objects given there distances. The image that an optical system gives of a point source is called the point spread function, PSF, or impulse response, IR, of the system, see Figure 1.3. A point source is characterized by an infinite distance to the instrument and therefore the wave is planar, which means that the phasor is reduced to $U(\xi, \eta) = P(\xi, \eta)$. The PSF or IR is given by,

$$h_{optical}(x, y) = | [\mathcal{F} \{ P(\xi, \eta) \}] (x, y) |^2$$
 (1.21)

The domain where the impulse response is invariant under translation is called the **isoplanatic domain**.

In presence of aberrations, which will be discussed in the section 1.3, the wavefront is deformed with respect to the perfect planar or spheric form. The phasor at the exit pupil characterizing a point source is modified as following,

$$U(\xi,\eta) = P(\xi,\eta)e^{j\phi_{Ab}(\xi,\eta)},\tag{1.22}$$

where $\phi_{Ab}(\xi, \eta)$ is the dephasing caused by the aberration present between the object plane and the image plane.

1.3. Aberrations 9

1.3 Aberrations

- 1.3.1 Sources of aberrations
- 1.3.2 Zernike polynomials

Chapter 2

Phase Diversity Experiment

Here, we will describe the experiment put in place in the optical laboratory at the HEIG-VD to reconstruct wavefronts with unknown static aberrations introduced using phase screens. At first, we study the behavior of the phase diversity algorithm put in place by Mugnier, Blanc, and Idier (2006) at ONERA with respect to number of averaging images, in other words noise level, and number of Zernike coefficients retrieved. Then we test the algorithm using a known aberration introduced by a parallel plane plate in the beam comparing the result to Zemax simulation. And finally, we introduce the phase screen to have random aberrations in the pupil and try to compare the phase diversity results with the Shack Hartman wavefront sensor results.

2.1 ONERA algorithm

2.2 Experimental Setup

The design of the experiment was already done by Bouxin (2017). The system is built according to her plans and specifications.

The experiment is mounted on a pressurized legs optical table. The assembly contains six main components: a light source, an entrance pupil, an imaging system, a converging lens to focus the beam on the camera, a camera and a wavefront sensor.

2.2.1 Light source

The final application of the phase diversity will be to characterize the optical aberrations induced by the imperfect optical path to a scientific detector of a telescope. For this reason, the light source has to simulate a distant star aberration-free wavefront. A dis-

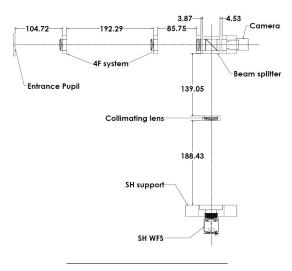


FIGURE 2.1: Experimental setup schema with the relevant distances, (Bouxin, 2017).

tant star wavefront is considered planar since the object distance, z, is far greater than the telescope size, r, see Fig. 2.2. The source of our experiment must then be characterized by a planar wavefront.

#	Components	Model	Reference
1	Pigtailed laser diode	Thorlabs, LPS-635-FC	A.1
2	Converging lens, f = 11 mm	Thorlabs, A220TM-A	A.2
3	Pinhole, 10 μm	Thorlabs, P10S	A.3
4	Converging lens, f = 200 mm	Thorlabs, AL100200	A.4
5	3.2 mm Hole milled in metal sheet		
6	Converging lens, f = 100 mm	Thorlabs, AC254-100-A	A.5
7	Converging lens, f = 80 mm		
8	Camera CMOS	Ximea, MQ013MG-E2	A.6
9	Converging lens, f = 100 mm		
10	Shack-Hartman WFS	Thorlabs, WFS150-5C	A.7

TABLE 2.1: Optical Components

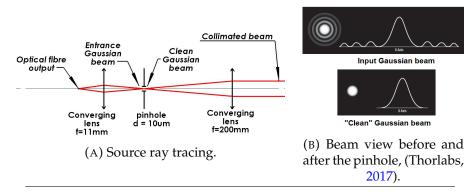


FIGURE 2.3: Source schema and pinhole effect on the beam.

In order to obtain such a planar wavefront at the entrance pupil, the light source consist of a "pigtailed laser diode", a f=11mm converging lens, a pinhole and a f=200 mm converging lens, see Table 2.1. The pigtailed laser diode emits a Gaussian beam centred at 637.5 nm slightly diverging. The converging lens concentrates the beam at the center of the 10μ m pinhole to filter the noise. The second converging lens collimates the beam, obtaining a collimated beam with a planar wavefront, see Fig. 2.3a and 2.3b.

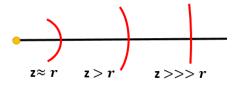


FIGURE 2.2: Wavefront curvature for different point source's distances, *z*. *r* represents the characteristic size of the arc of interest.

2.2.2 Entrance pupil

The entrance pupil of our optical system is a circular aperture of 3.2 mm diameter placed after the collimating lens of the light source. It is milled in a metal plate and centred in his support, to avoid positioning with a XY table. The diameter is chosen in available material to fit in the different detector's surfaces.

2.2.3 Pupil imaging system

The phase diversity technique requires PSFs images as input, which means that the beam as to be focused onto the detector surface. To analyse the aberration in the pupil plane, one needs to focus an image of the beam passing through the entrance



FIGURE 2.4: PSFs example of an alignment procedure

pupil. The simplest assembly to achieve this goal is the 4F system, which consist of two converging lenses of focal 100 mm. The two lenses are separated by 200 mm, see Fig. 2.1. This places the image of the entrance pupil 100 mm after the second converging lens.

2.2.4 Detectors

The image of the entrance pupil, obtained with the 4F system, is focused onto a CMOS Ximea camera by a f=80 mm converging lens to acquire the PSFs for the phase diversity wavefront retrieval. The camera has a surface composed by 1280×1024 pixels of 5.3 μ m, see Appendix A.6. It is mounted on sliding support in order to be able to acquire in/out-of-focus images. A beam splitter is placed in the converging beam to separate it in two. The second beam is collimated and a Shack-Hartman WFS is placed on the entrance pupil image plane, to check the results of the phase diversity wavefront retrieval. The Shack-Hartman WFS has a 39 X 31 lenslets grid and a CCD with a resolution of 1280×1024 pixels of $4.65 \ \mu$ m, see Appendix A.7.

2.3 Data Acquisition

2.3.1 Ximea Camera

The ONERA algorithm takes at least one focused and one defocused PSFs, as described in section 2.1. The PSFs are acquired using a python script which uses an open-source library to control the ximea camera, pyXimea¹, available on GitHub. The acquisition is done following these steps:

- 1. The first step in order to acquire PSFs is to determine the position of the camera's focus point using the python script AlignementScriptXimeaCamera.py, see Appendix B.1. This script let's you acquire consecutively PSFs at different camera's positions and computes their FWHM. It finally returns the minimum FWHM and the camera's position, see Figure 2.4.
- 2. Once you have the focus point position

2.3.2 Shack-Hartman WFS

2.4 Results

This section presents the results of the phase diversity experiment, with the introduction of different sources of aberration.

https://github.com/pupil-labs/pyximea

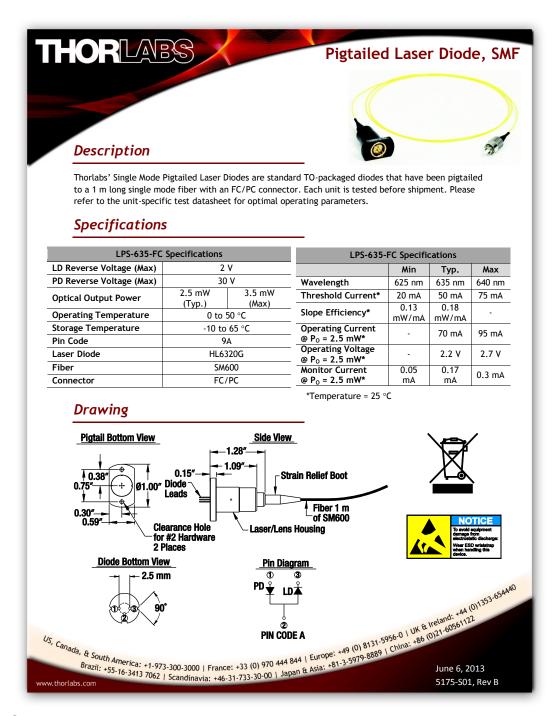
2.4.1 Parallel plane plate

The first source of aberration studied in this work is a tilted parallel plane plate which is used as a calibrated source of astigmatism.

Appendix A

Optical Component Datasheets

A.1 Pigtailed laser diode

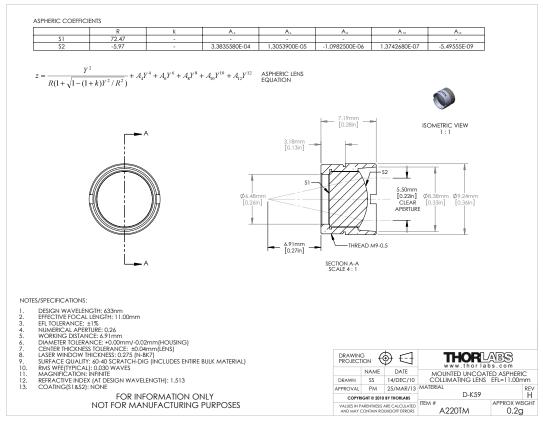


A.1.1 Power supply modification

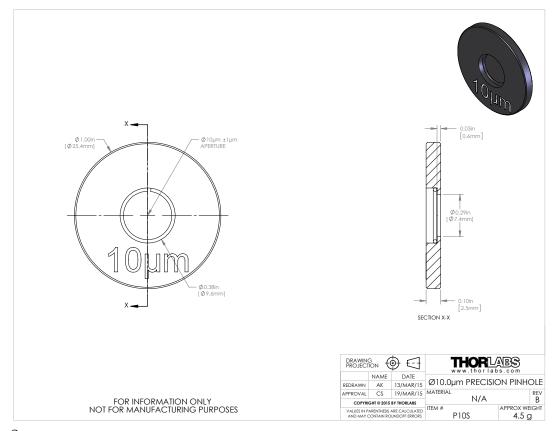
The former student, ??? its name ???, mounted a diode driver card to power it. Unfortunately, for the phase diversity experiment, chapter, the power was too high and the Ximea camera was always saturated. So I modified the driver circuit and added two resistances to lower the current so that the detector do not reach the saturation.

!!! mettre la photo du driver et de la modif !!!

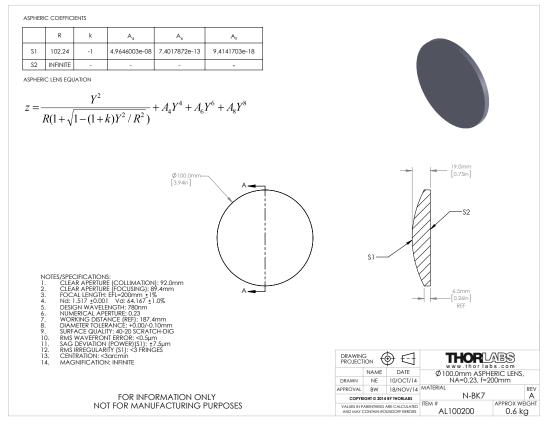
A.2 Converging lens A220TM-A, f = 11 mm



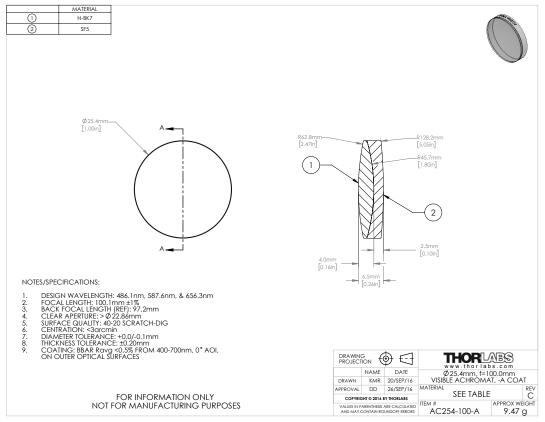
A.3 Pinhole 10 μ m



A.4 Converging lens AL100200, f = 200 mm



A.5 Converging lens AC254-100-A, f = 100 mm



Ximea Camera, MQ013MG-E2



On-chip binning: 1x1, 2x2 Image data interface: USB 3.0 GPIO IN, OUT Data I/O: Power requirements: 0.9 Watt Lens mount: C or CS Mount Weight: 26 grams **Dimensions WxHxD:** 26 x 26 x 26 mm Operating 50 °C environment:

Customs tariff code: 8525.80 30 (EU) / 8525.80 40 (USA)

ECCN: EAR99

Source: www.ximea.com/en/products/usb3-vision-cameras-xiq-line/mq013mg-e2

Shack-Hartmann wavefront sensor, WFS150-5C

8 Appendix

8.1 Technical Data

8.1.1 WFS150/300

Item #	WF\$150-5C	WFS150-7AR	WF\$300-14AR		
Microlenses					
Microlens Array	MLA150M-5C	MLA150M-7AR	MLA300M-14AR		
Substrate Material		Fused Silica (Quartz)			
Number of Active Lenslets	Software Selectable				
Max. Number of Lenslets	39 x 31		19 x 15		
Camera	Camera				
Sensor Type		CCD			
Resolution	max. 1280 x 1024 pixels, Software Selectable				
Aperture Size	5.95 mm x 4.76 mm				
Pixel Size	4.65 μm x 4.65 μm				
Shutter	Global				
Exposure Range	79 µs - 65 ms				
Frame Rate	max. 15 Hz				
Image Digitization	8 bit				
Wavefront Measurement					
Wavefront Accuracy 1)	λ/15 rms @ 633 nm		λ/50 rms @ 633 nm		
Wavefront Sensitivity 2)	λ/50 rms @ 633 nm		λ/150 rms @ 633 nm		
Wavefront Dynamic Range 3)	> 100 λ @	0 633 nm	> 50 λ @ 633 nm		
Local Wavefront Curvature 4)	> 7.4 mm	> 10.0 mm	> 40.0 mm		
External Trigger Input					
Save Static Voltage level	0 to 30 V DC				
LOW Level	0.0 V to 2.0 V				
HIGH Level	5.0 V to 24 V				
Input current	> 10 mA				
Min Pulse Width	100 μs				
Min. Slew Rate	35 V / msec				
Common Specifications					
Optical Input	C-Mount				
Power Supply	<1.5 W, via USB				
Operating Temperature Range ⁵)	+5 to +35 °C				
Storage Temperature Range	-40 to 70 °C				
Warm-Up Time for Rated Accuracy	15 min				
Dimensions (W x H x D)	32.0 mm x 40.4 mm x 45.5 mm				
Weight	0.1 kg				

¹⁾ Absolute accuracy using internal reference. Measured for spherical wavefronts of known RoC.

All technical data are valid at 23 \pm 5°C and 45 \pm 15% rel. humidity (non condensing)

@ 2007 - 2015 Thorlabs GmbH

99

Source: WFS Series Operation Manual, www.thorlabs.com

²⁾ Typical relative accuracy. Achievable after, and with respect to a user calibration, 10 image averages

³⁾ Over entire aperture of wavefront sensor
4) Radius of wavefront curvature over single lenslet aperture

Appendix B

Python Code

B.1 AlignementScriptXimeaCamera.py

```
1 ##Script to compute the FWHM of the beam on the camera averaging
2 #over "nbrImgAveraging" images and see which position minimizes it.
4 from ximea import xiapi
5 import numpy as np
6 from matplotlib import pyplot as plt
7 import scipy.optimize as opt
8 import datetime
9 import functionsXimea as fX
10 import seaborn as sns
11 import os
12 sns.set()
13 #%% instanciation
15 dataFolderPath =
      'C:/Users/Jojo/Desktop/PdM-HEIG/Science/data/PD/phaseScreen/alignement/'
      'C:/Users/Jojo/Desktop/PdM-HEIG/Science/fig/PD/phaseScreen/alignement/'
17 #create the matrix grid of the detector CCD
x = np.linspace(0,1280,1280)
y = np.linspace(0,1024,1024)
y = \text{np.meshgrid}(x, y)
22 #initial guess for the fit depending on the position of the beam in the CCD
23 initial_guess = [250, 481, 706, 3, 3]
25 #number of image to average
26 nbrImgAveraging = 10
_{\rm 28} #%%data acquisition and treatment
30 #create instance for first connected camera
31 cam = xiapi.Camera()
32 #start communication
33 print('Opening camera...')
34 cam.open_device()
35 #settings
36 cam.set_imgdataformat('XI_MONO8') #XIMEA format 8 bits per pixel
37 cam.set_gain(0)
38 #create instance of Image to store image data and metadata
39 img = xiapi.Image()
40 #start data acquisition
```

```
41 print('Starting data acquisition...')
42 if cam.get_acquisition_status() == 'XI_OFF':
          cam.start_acquisition()
43
44
45 cam.set_exposure(fX.determineUnsaturatedExposureTime(cam,img,[60,10000],1))
47 #instanciation for the while loop
48 answer ='y'
49 i=0
50 relativePos = []
51 data = []
52 data_fitted = []
53 \text{ FWHMx} = []
54 \text{ FWHMy} = []
55 \times 0 = []
56 \text{ y0} = []
57 sigmaX0 = []
58 sigmaY0 = []
60 while answer == 'y':
61
      try:
62
          relativePos.append(float(raw_input('What is the position on the
              screw [mm] ? ')))
      except ValueError:
64
          print('Not a float number')
65
      [tmpdata,stdData] = fX.acquireImg(cam,img,nbrImgAveraging)
      data.append(tmpdata)
68
      #Fit the img data on the 2D Gaussian to compute the FWHM
      print('Fitting 2D Gaussian...')
70
      popt, pcov = opt.curve_fit(fX.TwoDGaussian, (x,y), data[i].ravel(), p0 =
          initial_guess)
      print('Fitting done')
72
73
      FWHMx.append(2*np.sqrt(2*np.log(2))*popt[3])
74
      FWHMy.append(2*np.sqrt(2*np.log(2))*popt[4])
75
      x0.append(popt[2])
76
      sigmaX0.append(popt[4])
      y0.append(popt[1])
      sigmaY0.append(popt[3])
80
      print 'Fig %d : (x,y) = (\%3.2f,\%3.2f), FWHM x = \%3.2f, FWHM y = \%3.2f'
          %(i,x0[i],y0[i],FWHMx[i],FWHMy[i])
82
      data_fitted.append(fX.TwoDGaussian((x, y),
83
          popt[0],popt[1],popt[2],popt[3],popt[4]).reshape(1024, 1280))
84
      #plot the beamspot
85
      fig, ax = plt.subplots(1, 1)
86
      ax.imshow(data[i], cmap=plt.cm.jet,origin='bottom',
          extent=(x.min(), x.max(), y.min(), y.max()))
88
      ax.contour(x, y, data_fitted[i], 5, colors='w',linewidths=0.8)
89
      plt.xlim( (popt[2]-4*popt[4], popt[2]+4*popt[4]) )
      plt.ylim((popt[1]-4*popt[3], popt[1]+4*popt[3]))
      plt.show()
92
93
```

```
#ask if the person wants to acquire a new image to improve the alignement
95
      pressedkey = raw_input('Do you want to acquire an other image [y (yes)
          or n (no)]: ')
      if (pressedkey =='n'):
96
          answer = pressedkey
97
      #increase i
98
      i+=1
99
101 #stop data acquisition
  print('Stopping acquisition...')
103 cam.stop_acquisition()
105 #stop communication
106 cam.close_device()
108 #convert list to np.array
109 relativePos = np.array(relativePos)
110 data = np.array(data)
111 FWHMx = np.array(FWHMx)
112 FWHMy = np.array(FWHMy)
113 \times 0 = np.array(x0)
114 y0 = np.array(y0)
sigmaX0 = np.array(sigmaX0)
116 sigmaY0 = np.array(sigmaY0)
118 #plot the FWHM vs. relPos
fig, ax = plt.subplots(1,1)
120 ind = np.argsort(relativePos)
ax.plot(relativePos[ind],(np.sqrt(FWHMx**2+FWHMy**2))[ind])
122 ax.set_xlabel('Position [mm]')
123 ax.set_ylabel('FWHM [px]')
124 ax.grid()
125 date = datetime.datetime.today()
126 if not os.path.isdir(plotFolderPath):
      os.makedirs(plotFolderPath)
128 plt.savefig(plotFolderPath+date.strftime('%Y%m%d%H%M%S')+'FWHM_pos.pdf')
129 plt.savefig(plotFolderPath+date.strftime('%Y%m%d%H%M%S')+'FWHM_pos.png')
130
indOfMinFWHM = np.argmin(np.sqrt(FWHMx**2+FWHMy**2))
134 fig, axarr = plt.subplots(1,np.size(data,0))
135 #plot all the images besides each other
136 for iImg in ind:
      axarr[iImg].imshow(data[iImg], cmap=plt.cm.jet,origin='bottom',
137
          extent=(x.min(), x.max(), y.min(), y.max()))
138
139 #
       axarr[iImg].contour(x, y, data_fitted[iImg], 5,
      colors='w',linewidths=0.8)
      axarr[iImg].set_xlim( (x0[iImg]-12, x0[iImg]+12) )
140
      axarr[iImg].set_ylim( (y0[iImg]-12, y0[iImg]+12) )
141
      axarr[iImg].set_yticklabels('', visible=False)
      axarr[iImg].set_xticklabels('', visible=False)
143
      axarr[iImg].set_title('%5.3f mm'%relativePos[iImg],fontsize=8)
144
      if iImg == indOfMinFWHM:
145
          axarr[iImg].set_frame_on(True)
          for pos in ['top', 'bottom', 'right', 'left']:
147
              axarr[iImg].spines[pos].set_edgecolor('r')
148
```

```
axarr[iImg].spines[pos].set_linewidth(2)

else:

axarr[iImg].set_frame_on(False)

plt.show()

adate = datetime.datetime.today()

plt.savefig(plotFolderPath+date.strftime('%Y%m%d%H%M%S')+'ImgPSF.pdf')

plt.savefig(plotFolderPath+date.strftime('%Y%m%d%H%M%S')+'ImgPSF.png')

#save data

if not os.path.isdir(dataFolderPath):

os.makedirs(dataFolderPath)

date = datetime.datetime.today()

np.save(dataFolderPath+date.strftime('%Y%m%d%H%M%S')+'data.npy',data)

np.save(dataFolderPath+date.strftime('%Y%m%d%H%M%S')+'relativePos.npy',relativePos)
```

B.2 AcquisAndSaveXimea.py

```
1 #%% Script to acquire images average over nbrImgAveraging images and save
      them into fits file
3 from ximea import xiapi
4 import datetime
5 import functionsXimea as fX
6 import winsound
7 import numpy as np
9 #%%instanciation
10 #number of image to average
11 nbrImgAveraging = 5000
12 numberOfFinalImages = 1
14 #Cropping information
15 \text{ sizeImg} = 256
17 #Parameter of camera and saving
18 folderPathCropped =
      '.../.../data/PD/astigmatism/angle_study_3/wth/cropped/20/'
19 darkFolderPathCropped =
      '../../data/dark/astigmatism/angle_study_3/wth/cropped/20/'
20 folderPathFull = '.../.../data/PD/astigmatism/angle_study_3/wth/full/20/'
21 darkFolderPathFull =
      '../../data/dark/astigmatism/angle_study_3/wth/full/20/'
22 nameCamera = 'Ximea'
23 focusPos = 11.63
25 #Sound
26 duration = 1000 # millisecond
27 \text{ freq} = 2000 \# \text{Hz}
29 #initial guess for the fit depending on the position of the beam in the CCD
30 initial_guess = [250, 468, 954, 3, 3]
```

```
32 #-----
33 #%% data acquisition
      -----
34
35 #Opening the connection to the camera
36 cam = xiapi.Camera()
37 cam.open_device()
38 cam.set_imgdataformat('XI_MONO8') #XIMEA format 8 bits per pixel
39 cam.set_gain(0)
41 img = xiapi.Image()
42 if cam.get_acquisition_status() == 'XI_OFF':
     cam.start_acquisition()
44 #%% exposition
45 \text{ cond} = 1
46 while bool(cond):
     source = ''
     winsound.Beep(freq, duration)
     source = int(raw_input('Is the source turned on and at focus point
49
         (usually %5.3f mm) (yes = 1) ? '%focusPos))
     if source == 1:
         cond = 0
51
     else:
52
         print 'Please turn on the source and place the camera on the focus
             point (%5.3f mm)', focusPos
54
55 if bool(source):
     #Set exposure time
     cam.set_exposure(fX.determineUnsaturatedExposureTime(cam,img,[1,10000],1))
57
     #get centroid
     centroid = fX.acquirePSFCentroid(cam,img,initial_guess)
     print 'centroid at (%d, %d)' %(centroid[0],centroid[1])
62 #%%Acquire images at different camera position
64 acquire = 1
65 while bool(acquire):
     cond = 1
     while bool(cond):
         dark = ''
         winsound.Beep(freq, duration)
69
         dark = int(raw_input('Is the source turned off (yes = 1) ? '))
70
         if dark == 1:
             cond = 0
72
         else:
73
             print 'Please shut down the source.'
     winsound.Beep(freq, duration)
76
     pos = float(raw_input('What is the position of the camera in mm focused
77
         (\%5.3f \text{ mm}) \text{ dephase } 2Pi \text{ (pos+ = } \%5.3f \text{ mm, pos- = } \%5.3f) ?
         '%(focusPos,focusPos+3.19,focusPos-3.19)))
78
     if bool(dark):
79
         print 'Acquiring dark image...'
         # Acquire dark images
81
82
         [darkData,stdDarkData] = fX.acquireImg(cam,img,nbrImgAveraging)
         print 'Cropping'
83
```

```
[darkdataCropped,stddarkDataCropped] =
84
              fX.cropAroundPSF(darkData,stdDarkData,centroid,sizeImg,sizeImg)
          print 'saving'
85
          fX.saveImg2Fits(datetime.datetime.today(),darkFolderPathCropped,nameCamera,darkdataCro
86
          fX.saveImg2Fits(datetime.datetime.today(),darkFolderPathFull,nameCamera,darkData,stdDa
      #Acquire images -----
      cond = 1
      while bool(cond):
91
          source = ''
92
          winsound.Beep(freq, duration)
93
          source = int(raw_input('Is the source turned on (yes = 1) ? '))
          if source == 1:
             cond = 0
          else:
             print 'Please place turn on the camera'
      if bool(source):
100
          print 'Acquiring images...'
101
          # Acquire focused images
102
          for iImg in range(numberOfFinalImages):
103
              imgNumber = iImg+1
104
             print 'Acquiring Image %d'%imgNumber
              [data,stdData] = fX.acquireImg(cam,img,nbrImgAveraging)
             print 'Cropping'
              [dataCropped,stdDataCropped] =
108
                 fX.cropAndCenterPSF(data-darkData,stdData+stdDarkData,sizeImg,initial_guess)
             print 'Saving'
             fX.saveImg2Fits(datetime.datetime.today(),folderPathCropped,nameCamera,dataCropped
             fX.saveImg2Fits(datetime.datetime.today(),folderPathFull,nameCamera,data-darkData,s
      cond = 1
      while bool(cond):
114
          acquire = ''
          winsound.Beep(freq, duration)
          acquire = int(raw_input('Do you want to acquire at an other camera
117
              position (yes = 1, no = 0) ? '))
          if acquire == 1:
              cond = 0
          elif acquire == 0:
120
              cond = 0
          else:
              print 'please answer with 0 or 1 for no or yes, respectively'
124
126 ##Stop the acquisition
127 cam.stop_acquisition()
128 cam.close_device()
130 print 'Acquisition finished'
```

Bibliography

- Bouxin, A. (2017). "Phasor diversity to measure the static aberrations of an optical system". MA thesis. HEIG-VD.
- Goodman, Joseph W. (1968). *Introduction to Fourier Optics*. Ed. by L. Cox and John M. Morris. 2nd. McGraw-Hill Companies, Inc.
- Mugnier, Laurent M., Amandine Blanc, and Jérôme Idier (2006). "Phase Diversity: A Technique for Wave-Front Sensing and for Diffraction-Limited Imaging". In: ed. by Peter Hawkes. Vol. 141. Advances in Imaging and Electron Physics. Elsevier, pp. 1 –76. DOI: https://doi.org/10.1016/S1076-5670(05)41001-0. URL: http://www.sciencedirect.com/science/article/pii/S1076567005410010.
- Thorlabs (2017). *Principles of Spatial Filters*. Thorlabs. URL: https://www.thorlabs.com/newgrouppage9.cfm?objectgroup_id=1400.