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Pressure Control in Dissipative Particle Dynamics and its Application in Simulating Micro- and Nano-bubbles

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Outline



- Introduction
- Barostat in DPD & MDPD
 - Berendsen barostat in DPD
 - Partial Berendsen barostat in DPD & MDPD
- Bubble dynamics in DPD & MDPD
 - · Application in multi-component system
 - Bubble Collapse & Oscillation
- Summary

Supported by:



Introduction – Bubbles



• Cavitation Inception – Hydrodynamics of Propeller:

Strong Water, $\sigma_i = 0.93$

- Liquid to Vapor Phase Transition;
- Gaseous Nucleation.



Beginning sheet cavitation, Ut/c = -2.33



Developed sheet cavitation, Ut/c = -1.02

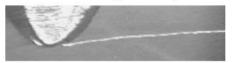
Arndt, Annu. Rev. Fluid Mech. 2002



Inception downstream of tip

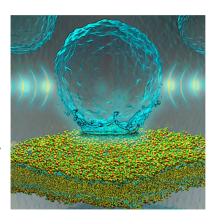


Ut/c = 0.22 (2.8 millisec)



Bubble reaches tip Ut/c = 0.73 (9.3 millisec)

- Sonoporation Drug Delivery:
 - Ultrasound Contrast Agent;
 - Drug Delivery & Noninvasive Therapy.

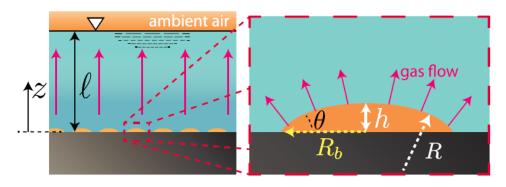


Fu *et al., J. Phys. Chem. Lett.* 2015

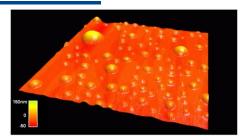
Introduction – Bubbles

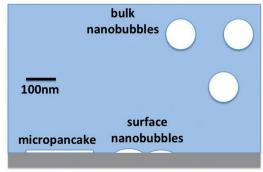


- Nanobubbles (surface or bulk):
 - Stability of their long life (days & weeks);
 - Current MD simulation is limited to several tens of nanometers, resulting in lifetimes of order 100 ns.



Weijs and Lohse, Phys. Rev. Lett., 2013





Seddon *et al., ChemPhysChem,* 2012 Lohse & Zhang, *Rev. Mod. Phys.,* 2015

Introduction – Dissipative Particle Dynamics



Macroscale

Mesoscale

Microscale

Continuum:

Finite difference/volume/element, volume of fluid, level set, immersed boundary methods ... Mesoscopic: DPD, LBM, MesoDyn ...

Microscopic: Density functional theory, molecular dynamics, Monte Carlo methods ...

 $1 \mu m$ 1 ms 10 nm 1ns

DPD

Colloidal suspensions [Koelman and Hoogerbrugge, 1993, Boek et al., 1997]

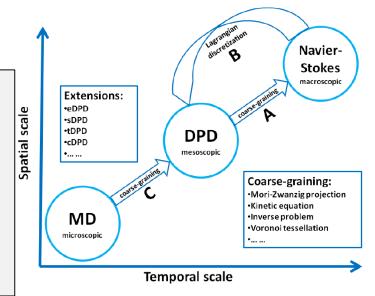
Ploymer solution [Kong et al., 1997, Spenley, 2000, Jiang et al., 2007]

Two-phase flow [Pan et al., 2016, Warren, 2003]

Surfactant [Groot, 2000, Groot, 2003, Rekvig et al., 2003]

Membranes [Yang and Ma, 2010, Dutt et al., 2011, Li et al., 2012]

Drug delivery [Tomasini and Tomassone, 2012, Patterson et al., 2011, Masoud and Alexeev, 2011, Delcea et al., 2011]



Standard DPD Method



Basic Theory

•
$$\frac{d\mathbf{r}_i}{dt} = \mathbf{v}_i$$
; $\frac{d\mathbf{v}_i}{dt} = \mathbf{f}_i + \mathbf{F}_e$.

•
$$\mathbf{f}_i = \sum_{j \neq i} (\mathbf{F}_{ij}^C + \mathbf{F}_{ij}^D + \mathbf{F}_{ij}^R).$$

•
$$\mathbf{F}_{ij}^{C} = \begin{cases} a_{ij} (1 - r_{ij}) \hat{\mathbf{r}}_{ij}, & r_{ij} < 1 \\ 0, & r_{ij} \ge 1 \end{cases}$$
;

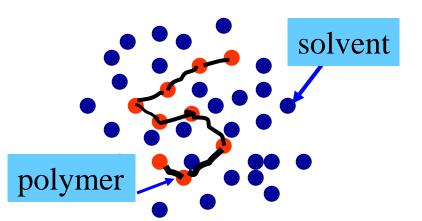
•
$$\mathbf{F}_{ij}^D = -\gamma w^D (r_{ij}) \mathbf{v}_{ij} \cdot \mathbf{r}_{ij} \mathbf{r}_{ij}$$
;

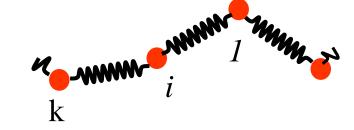
•
$$\mathbf{F}_{ij}^R = \sigma w^R (r_{ij}) \theta_{ij} \hat{\mathbf{r}}_{ij}$$
.

• Fluctuation-Dissipation Theorem:

•
$$\gamma = \frac{\sigma^2}{2k_BT}$$

•
$$w^{D}(r) = [w^{R}(r)]^{2} = \begin{cases} (1 - r/r_{C})^{s}, r_{ij} < r_{C} \\ 0, r_{ij} \ge r_{C} \end{cases}$$





Many-body DPD



Basic Theory

•
$$\frac{d\mathbf{r}_i}{dt} = \mathbf{v}_i$$
; $\frac{d\mathbf{v}_i}{dt} = \mathbf{f}_i + \mathbf{F}_e$.

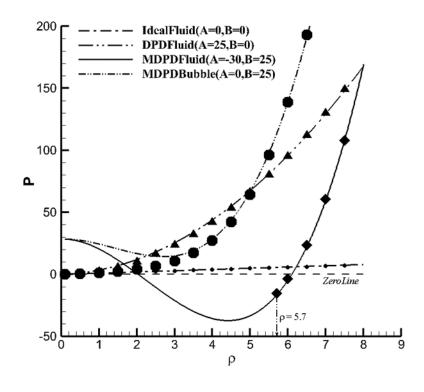
•
$$\mathbf{f}_i = \sum_{j \neq i} (\mathbf{F}_{ij}^C + \mathbf{F}_{ij}^D + \mathbf{F}_{ij}^R).$$

•
$$\mathbf{F}_{ij}^C = [A_{ij}w^C(r_{ij}) + B_{ij}(\bar{\rho}_i +$$

• Equations of State (EoS):

• **DPD**:
$$P = \rho k_B T + \alpha A \rho^2$$

• **MDPD**:
$$P = \rho k_B T + \alpha A \rho^2 + 2\alpha B r_d^4 (\rho^3 - c\rho^2 + d)$$



Introduction – Pressure control

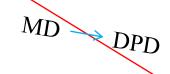


• The boundary condition(BC) in Particle methods

Cube system

- Periodic BC
- No-slip wall BC
- Pressure BC?

• ...





DPD

- Berendsen barostat [Atashafrooz and Mehdipour, 2016, Seaton et al., 2013]
- SSA: Shardlow-splitting algorithm [Lisal et al., 2011]
- Langevin piston approach [Jakobsen, 2005]

Pressure control (barostat)

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Berendsen Barostat theory

• The Berendsen barostat's scale factor,

$$\mu = \left[1 - \frac{\Delta t}{\tau_p} (P - P_0)\right]^{1/3}$$

• Consider a cubic system which contains N molecules and its volume is $V = L^3$.

$$r_i \rightarrow \mu r_i (i = 1, 2, 3, ..., N)$$

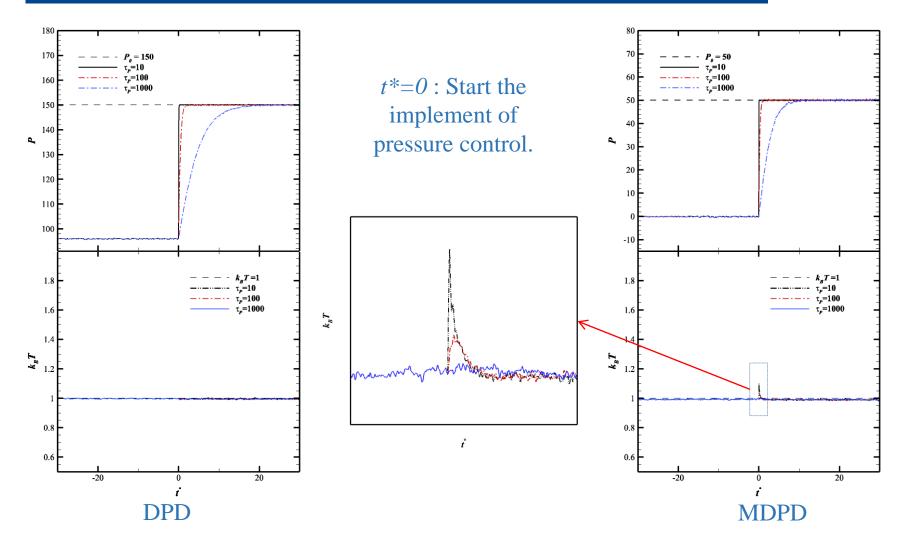
 $L \rightarrow \mu L (L = L_x = L_y = L_z)$

$$\Leftrightarrow \begin{cases} \left(r_{x}, r_{y}, r_{z}\right)_{i} \rightarrow \left(\mu r_{x}, \mu r_{y}, \mu r_{z}\right)_{i} \\ \left(L_{x}, L_{y}, L_{z}\right) \rightarrow \left(\mu L_{x}, \mu L_{y}, \mu L_{z}\right) \end{cases}$$

Here, t is the time step, τ_p is the "rise time" of the barostat, and P_0 is the desired pressure.

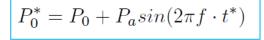
Berendsen barostat applys in single component system



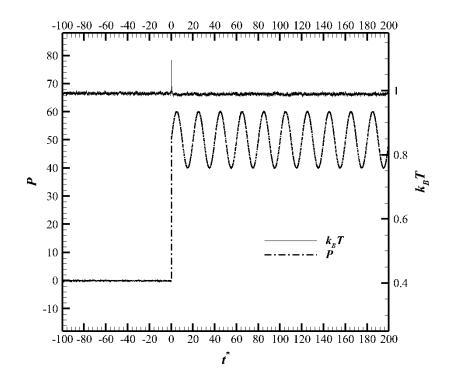


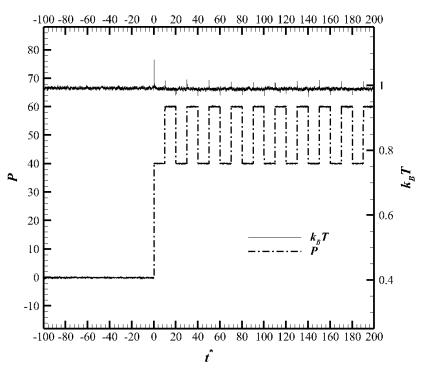


Berendsen barostat - Nonequilibrium dynamics



$$P_0^* = P_0 - P_a(-1)^{[2f \cdot t^*]}$$





Here, f is the frequency, P_a is the amplitude. And P_0 is the constant part of desired pressure P_0 *

Berendsen barostat limitation



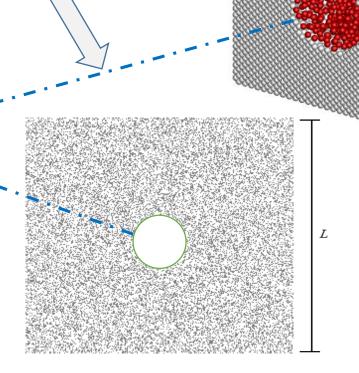
Particle displacement rescaling

$$r_i \rightarrow \mu r_i \left(i = 1, 2, 3, ..., N \right)$$

Region size rescaling

$$L \to \mu L \left(L = L_x = L_y = L_z \right)$$

Bring an artificial effect on the multi-component system interface.



THE UNIVERSE

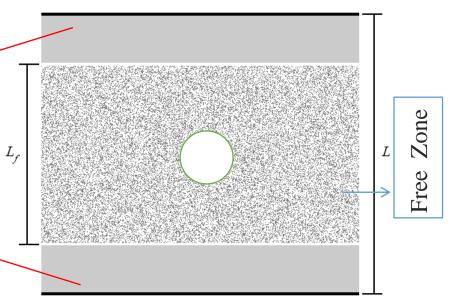
Partial Berendsen Barostat

• Consider a cubic system which contains N molecules and its volume is $V = L^3$.

$$\begin{array}{l}
r_{i} \to \mu r_{i} (i = 1, 2, 3, ..., N) \\
L \to \mu L (L = L_{x} = L_{y} = L_{z})
\end{array}$$

$$\Leftrightarrow \begin{cases}
(r_{x}, r_{y}, r_{z})_{i} \to (\mu r_{x}, \mu r_{y}, \mu r_{z})_{i} \\
(L_{x}, L_{y}, L_{z}) \to (\mu L_{x}, \mu L_{y}, \mu L_{z})
\end{cases}$$

Control Zone (Far-field Zone)



The Berendsen barostat's scale factor,

$$\mu = \left[1 - \frac{\Delta t}{\tau_p} (P - P_0)\right]^{1/3}$$

Here, t is the time step, τ_p is the "rise time" of the barostat, and P_0 is the desired pressure.

Define a new variable: $\delta = \frac{L_f}{L}$

Outline

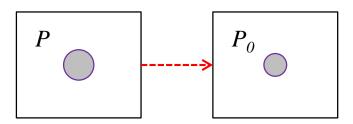


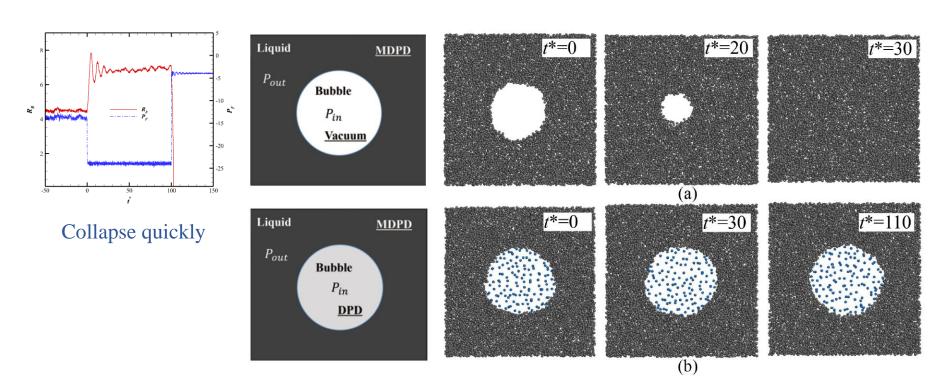
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Vacuum and Gaseous bubble



• Control the surrounding fluid pressure become the constatnt value P_o .



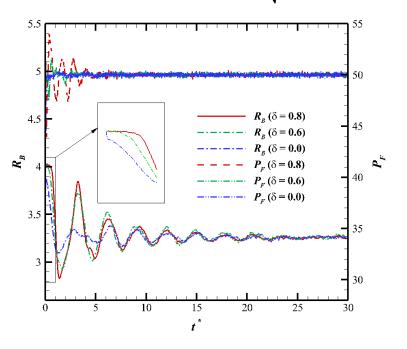


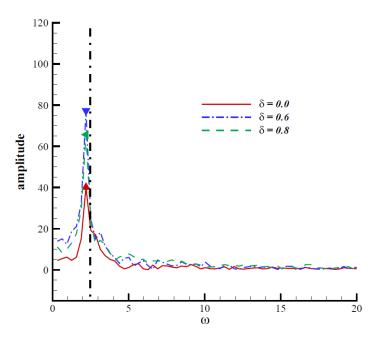
Microbubble oscillation



- Bubble Dynamics
 - Natural frequency from **Rayleigh Plesset** Equation (C. E. Brennen, *Cavitation and Bubble Dynamics*, 2013)

$$\omega_N = \frac{1}{R_0} \sqrt{\frac{1}{\rho_l} \left(3k P_{g0} - \frac{2S}{R_0} \right)}, \ P_g(R) = P_{g0} \left(\frac{R_0}{R} \right)^{3k}$$





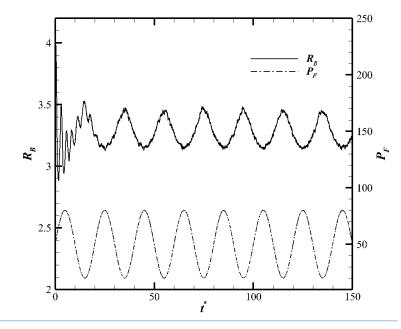
Bubble oscillation and collapse



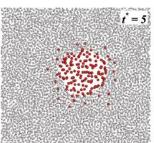
• Control the surrounding fluid pressure become the fluctuating value P_0^* .

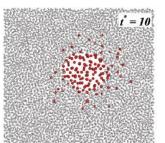
$$P_0^* = P_0 + P_a sin(2\pi f \cdot t^*)$$

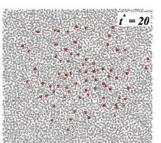
 Partial Berendsen barostat could be good barostat in nonequilibrium dynamics.











Summary



- The original **Berendsen barostat** works well in the (M)DPD simulation of the single-component system under **constant pressure** condition and **nonequilibrium dynamic** process;
- A partial Berendsen barostat is proposed to study the **multi-component** system in (M)DPD simulation;
- The partial Berendsen barostat could be a good candidate for the study on single or few droplets/bubbles under certain pressure control in **nonequilibrium dynamics**.

Thanks.



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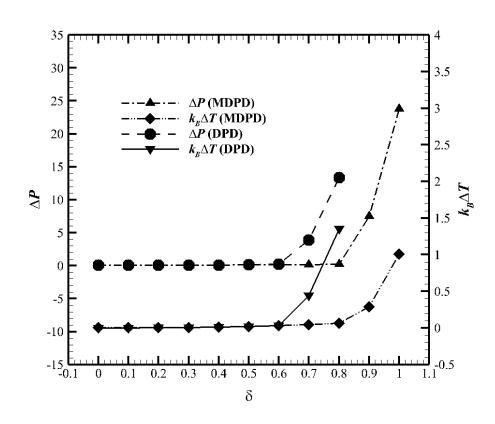
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Partial Berendsen barostat

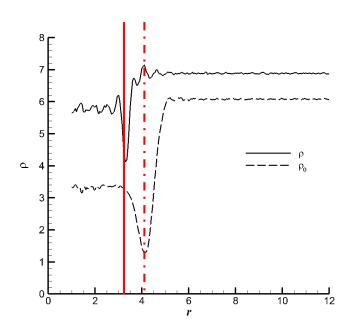


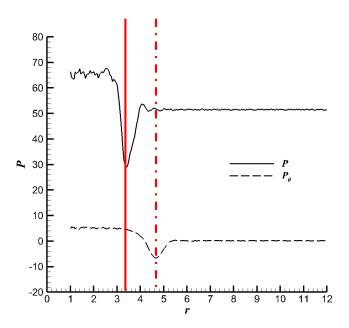
- Applys in (M)DPD Single-component system
 - A critical value of δ exist, when it smaller than the value, the barostat will work well.
 - Too great partial degree(greater than the critical value of δ) will lead to invalid the thermostat.



Gaseous bubble







- Improving the surrounding fluid pressure by the partial Berendsen barostat.
- Bubble shrink and become smaller, but the inside pressure improve with the outside fluid pressure.
- The surface tension increased because that the size of bubble dereased.