

General Relativity Crib Sheet

Section 2.5 Spacetime

$$2.66 \quad g_{\mu\nu} \approx \eta_{\mu\nu} + \frac{1}{2} (\partial_\rho \partial_\sigma g_{\mu\nu})_P x^\rho x^\sigma \quad \text{1st order approximation of } g_{\mu\nu}$$

Rules: Special Relativity → General Relativity:

Tensor: No change

Partial differentiation (comma) → Covariant differentiation (semicolon) (eq 2.53)

Total derivative (d/du) → Absolute derivative (D/du)

$\eta_{\mu\nu} \rightarrow g_{\mu\nu}$

$$2.69 \quad c^2 (d\tau)^2 \equiv g_{\mu\nu} dx^\mu dx^\nu \quad \text{Revised definition of proper time}$$

$$2.70 \quad f^\mu = \frac{D p^\mu}{d\tau} \quad \text{Newton's 2nd law of motion (formerly, } F = ma)$$

$$2.71 \quad \frac{d^2 x^\mu}{d\tau^2} + \Gamma_{\nu\sigma}^\mu \frac{dx^\sigma}{d\tau} \frac{dx^\nu}{d\tau} = 0 \quad \text{Equation of motion of a free particle}$$

$$2.72 \quad \frac{d^2 x^\mu}{du^2} + \Gamma_{\nu\sigma}^\mu \frac{dx^\sigma}{du} \frac{dx^\nu}{du} \stackrel{(2.12)}{=} 0 \quad \text{Photon eq of motion, } u \text{ affine parameter}$$

$$2.73 \quad g_{\mu\nu} \frac{dx^\mu}{du} \frac{dx^\nu}{du} = 0 \quad \Leftrightarrow \quad \text{Speed of photon} = c$$

Section 2.7 Gravitational potential and the geodesic

$$2.74 \quad g_{\mu\nu} \equiv \eta_{\mu\nu} + h_{\mu\nu}$$

To approximate flat, Special Relativity spacetime, assume:

$$2.7-2 \quad h_{\mu\nu} \approx 0 \text{ to the 1st order in 2.74 (Compare with 2.66)}$$

$$2.73 \quad g^{\mu\nu} = \eta^{\mu\nu} - h^{\mu\nu} \text{ to first order}$$

$$2.78 \quad \Gamma_{\nu\sigma}^\mu \approx \frac{1}{2} \eta^{\mu\rho} (\partial_\nu h_{\rho\sigma} + \partial_\sigma h_{\nu\rho} - \partial_\rho h_{\nu\sigma}) \approx 0 \text{ to the 1st order}$$

$$2.7-5 \quad \left| \frac{dx^i}{dt} \right| \ll c \quad \text{slow-moving particles}$$

$$2.7-6 \quad |\partial_0 h_{\mu i}| \ll |\partial_i h_{\mu 0}| \text{ for all } i \text{ and } \mu \quad \text{quasi-static}$$

$$2.83 \quad g_{00} \approx 1 + \frac{2V}{c^2} \quad \text{Relation between } g_{00} \text{ and acceleration potential } V$$

Section 2.8 Newton's law of universal gravitation

Schwarzschild solution:

$$\begin{aligned}
 2.8-1 \quad c^2 d\tau^2 &= \left(1 - \frac{2GM}{rc^2}\right) c^2 dt^2 - \left(1 - \frac{2GM}{rc^2}\right)^{-1} dr^2 - r^2 d\theta^2 - r^2 \sin^2 \theta d\phi^2 \\
 \Rightarrow \mathbf{F} = m\mathbf{a} &= -\frac{GM}{r^2} \hat{\mathbf{r}} \text{ for small values of } \frac{GM}{rc^2} \\
 \equiv \text{Newton's eq for universal gravitation}
 \end{aligned}$$

Section 2.9 A rotating reference system

$$\begin{aligned}
 \text{Euler Force} &= -m \frac{d\boldsymbol{\omega}}{dt} \times \mathbf{r}' && (\text{Merry-go-round}) \\
 \text{Coriolis Force} &= -2m (\boldsymbol{\omega} \times \mathbf{v}') && (\text{Coordinate center on equator}) \\
 \text{Centrifugal Force} &= -m \boldsymbol{\omega} \times (\boldsymbol{\omega} \times \mathbf{r}') && (\text{Outward force of ball on a string})
 \end{aligned}$$

where the primed reference system is a rotating coordinate system, and

$\boldsymbol{\omega}$ is the angular velocity of the rotating frame relative to the inertial system,

\mathbf{v}' is the velocity of an object in the primed system,

\mathbf{r}' is the position vector of an object in the primed system,

m is the mass of the object.

$$2.84 \quad c^2 d\tau^2 = c^2 dT^2 - dX^2 - dY^2 - dZ^2 \quad \text{Line element for non-rotating frame}$$

$$\begin{aligned}
 2.85 \quad T &= t \\
 X &= x \cos \omega t - y \sin \omega t \\
 Y &= x \sin \omega t + y \cos \omega t && \text{Coordinates for Rotating inertial frame } K' \\
 Z &= z
 \end{aligned}$$

$$\begin{aligned}
 2.87 \quad \ddot{t} &= 0 \\
 \ddot{x} - \omega^2 x t^2 - 2\omega \dot{y} t &= 0 \\
 \ddot{y} - \omega^2 y t^2 + 2\omega \dot{x} t &= 0 && \text{Equations of motion in } K' \text{ for free particle} \\
 \ddot{z} &= 0
 \end{aligned}$$

$$2.89 \quad m \frac{d^2 \mathbf{r}}{dt^2} = -m\boldsymbol{\omega} \times (\boldsymbol{\omega} \times \mathbf{r}) - 2m\boldsymbol{\omega} \times \frac{d\mathbf{r}}{dt} \quad \text{Breakout of 2 of the forces}$$

Section 3.1 The stress tensor and fluid motion

Flat spacetime

$$3.2 \quad \boxed{T^{\mu\nu} \equiv \left(\rho + \frac{P}{c^2}\right) u^\mu u^\nu - P \eta^{\mu\nu}} \quad \text{Energy-momentum-stress tensor}$$

$$3.1-10 \quad T^{\mu\nu} u_\nu = c^2 \rho u^\mu \quad T^{\mu\nu} \text{ is symmetric and } \rho, P, \text{ and } u^\mu \text{ are the fluid fields}$$

$$3.5 \quad (\rho u^\mu)_{;\mu} + \frac{P}{c^2} u^\mu_{;\mu} = 0 \quad \text{Relativistic continuity eq for a perfect fluid}$$

$$3.6 \quad \left(\rho + \frac{P}{c^2}\right) u^\nu_{;\mu} u^\mu = \left(\eta^{\mu\nu} - \frac{1}{c^2} u^\mu u^\nu\right) P_{;\mu} \quad \text{Eq of motion of a perfect fluid}$$

$$3.7 \quad \partial \rho / \partial t + \nabla \cdot (\rho \mathbf{v}) = 0 \quad \text{Classical continuity eq for a perfect fluid}$$

$$3.9 \quad \rho \left(\frac{\partial}{\partial t} + \mathbf{v} \cdot \nabla \right) \mathbf{v} = -\nabla P \quad \text{Classical eq of motion of a perfect fluid}$$

$$3.1-11 \quad \text{div } T^{\mu\nu} \equiv \nabla \cdot T^{\mu\nu} \stackrel{(2.61D)}{=} T^{\mu\nu}_{;\mu} \quad \text{Divergence}$$

Th 3.1.1 The divergence $T^{\mu\nu}_{;\mu} = 0$ iff equations (3.5) and (3.6) hold

Curved spacetime

$$3.1-17 \quad (\rho u^\mu)_{;\mu} + \frac{P}{c^2} u^\mu_{;\mu} = 0 \quad \text{Continuity equation}$$

$$3.1-18 \quad \left(\rho + \frac{P}{c^2}\right) u^\nu_{;\mu} u^\mu = \left(g^{\mu\nu} - \frac{1}{c^2} u^\mu u^\nu\right) P_{;\mu} \quad \text{Eq of motion of a perfect fluid}$$

$$3.10 \quad T^{\mu\nu} \equiv \left(\rho + \frac{P}{c^2}\right) u^\mu u^\nu - P g^{\mu\nu} \quad \text{Stress tensor}$$

$$3.1.19 \quad T_{\mu\nu} = \left(\rho + \frac{P}{c^2}\right) u_\mu u_\nu - P g_{\mu\nu} \quad \text{and}$$

$$T_\mu{}^\nu = \left(\rho + \frac{P}{c^2}\right) u_\mu u^\nu - P \delta_\mu^\nu \quad \text{Associated Stress Tensors}$$

$$3.11 \quad T^{\mu\nu}_{;\mu} = 0 \quad \text{Vanishing divergence}$$

Section 3.2 The curvature tensor and related tensors

$$3.2-1 \quad \lambda_{a;b} \stackrel{(2.56)}{=} \partial_b \lambda_a - \Gamma_{ab}^d \lambda_d \quad \text{Covariant derivative of a vector}$$

$$3.2-2 \quad \lambda_{a;bc} = \partial_c (\lambda_{a;b}) - \Gamma_{ac}^e \lambda_{e;b} - \Gamma_{bc}^e \lambda_{a;e} \quad \text{2nd covariant derivative}$$

$$3.13 \quad R^d_{abc} \equiv \partial_b \Gamma_{ac}^d - \partial_c \Gamma_{ab}^d + \Gamma_{ac}^e \Gamma_{eb}^d - \Gamma_{ab}^e \Gamma_{ec}^d \quad \text{Curvature tensor}$$

$$3.2-4 \quad R_{abcd} \equiv g_{ae} R^e_{bcd} \quad \text{and} \quad R^a{}_{bcd} \equiv g^{be} R^a_{ecd} \quad \text{Assoc curvature tensors}$$

Th 3.2.3 Can interchange order of covariant diff of a tensor iff $R^d_{abc} = 0$

$$3.14 \quad R^a_{bcd} + R^a_{cdb} + R^a_{dbc} = 0 \quad \text{Cyclic identity}$$

$$3.15 \quad R_{abcd} = \frac{1}{2} [\partial_d \partial_a g_{bc} - \partial_d \partial_b g_{ac} + \partial_c \partial_b g_{ad} - \partial_c \partial_a g_{bd}]$$

$$-g^{ef} [\Gamma_{eac} \Gamma_{fbd} - \Gamma_{ead} \Gamma_{fbc}]$$

3.16 $R_{abcd} = -R_{bacd}$

3.17 $R_{abcd} = -R_{abdc}$

3.2-5 $R_{abcd} = R_{badc}$

3.18 $R_{abcd} = R_{cdab}$

3.19 $R^a_{acd} = 0$

3.20 $R^a_{bcd;e} + R^a_{bde;c} + R^a_{bec;d} = 0$ **Bianchi identity**

3.21 $R_{ab} \equiv R^c_{abc}$ **Ricci tensor**

3.2-9 $R_{ab} = R_{ba}$ Symmetric in a and b

3.2-12 $R^a_b = R^c{}^a_{bc}$

3.22 $R \equiv R^a_a = g^{ab} R_{ba}$ **Curvature scalar**

3.23 $G_{ab} \equiv R_{ab} - \frac{1}{2} R g_{ab}$ **Einstein tensor**

3.2-13 $G^a_b = R^a_b - \frac{1}{2} R \delta^a_b$ **Associated Einstein tensor**

3.2-14 $G^{ab} = G^a_c g^{cb} = R^{ab} - \frac{1}{2} R g^{ab}$ **Associated Einstein tensor**

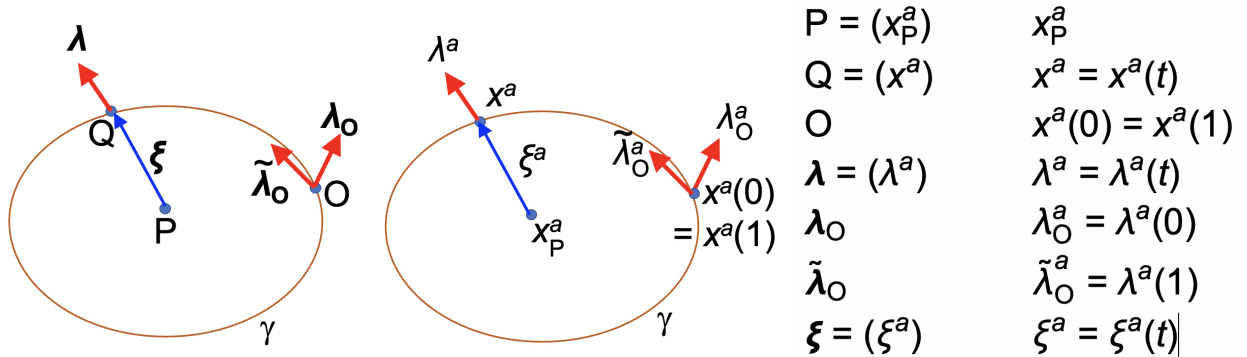
Th 3.2.5 $G^{ab}_{;a} = 0$ **Vanishing divergence**

3.2-20 $G^b_{c;b} = 0$ **Vanishing divergence**

3.2-21 $G^{ba}_{;b} = 0$ **Vanishing divergence**

Section 3.3 Curvature and parallel transport

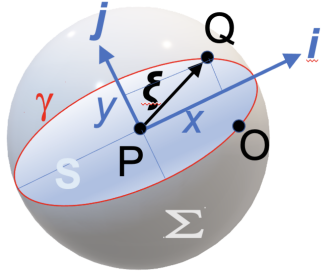
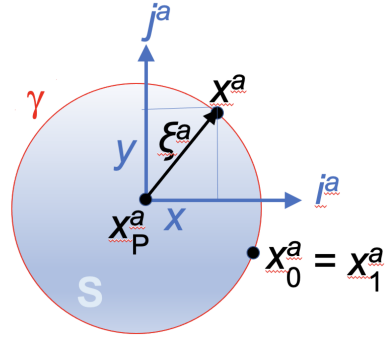
Parallel transport enables computation of the curvature components R^a_{bcd}



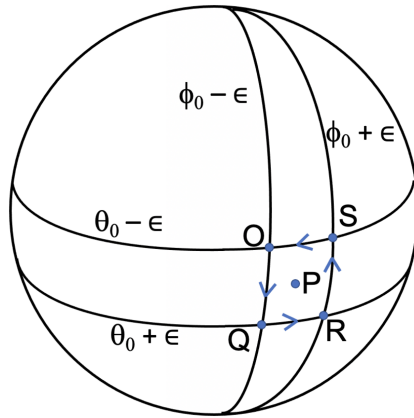
3.24 $\frac{d\lambda^a}{dt} = -\Gamma^a_{bc} \lambda^b \frac{dx^c}{dt}$ Defn of parallel transport of λ^a

3.29 $f^{cd} \equiv \oint \xi^c d\xi^d = \frac{1}{2} \oint [\xi^c d\xi^d - \xi^d d\xi^c]$ where $\xi^a \equiv x^a - x^a_P$

3.30 $\Delta\lambda^a \equiv \tilde{\lambda}^a_O - \lambda^a_O = -\frac{1}{2} (R^a_{bcd})_P \lambda^b_O f^{cd}$ 2nd order approx of $\Delta\lambda$

Small sphere Σ of radius a Cross section S with area A

$$3.32 \quad \frac{\Delta \lambda^a}{A} = - (R^a_{bcd})_P \lambda^b_O i^c j^d \quad \text{Geometric version of (3.30)}$$



$$3.3-14 \quad \Delta \lambda^a \approx \frac{4 \epsilon^2}{a} \delta^a_2$$

$$3.3-15 \quad A \approx 4 \epsilon^2 a^2 \sin \theta_0,$$

$$3.3-18 \quad - (R^a_{bcd})_P \lambda^b_O i^c j^d = - (R^a_{112})_P \frac{1}{a^3 \sin \theta_0}$$

$$3.3-19 \quad (R^a_{112})_P = - \delta^a_2$$

$$3.3-21 \quad (R_{1b12})_P = g_{2b}$$

$$\Rightarrow (R_{1212})_P = (R_{2121})_P = -(R_{2112})_P = -(R_{1221})_P = g_{22} \quad \text{and the rest are zero}$$