Short-range correlation physics at low RG resolution

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ajt, S.K. Bogner, and R.J. Furnstahl, arXiv:2006.11186

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ajt, S.K. Bogner, and R.J. Furnstahl, in progress



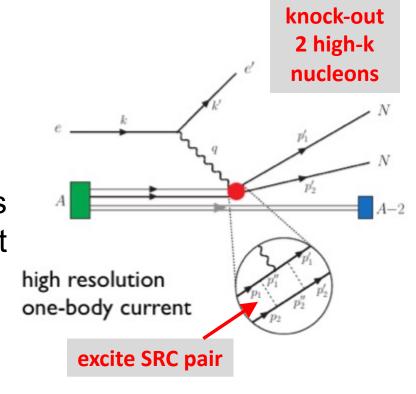


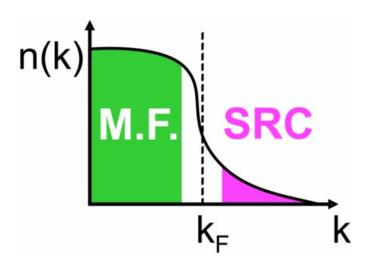




Motivation

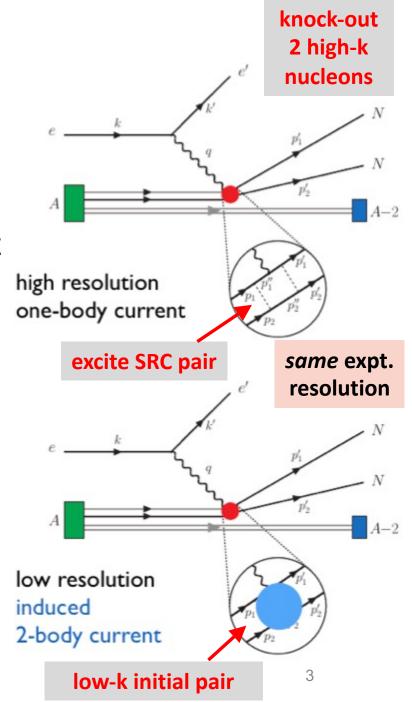
- Recent experiments have been able to isolate processes where short-range correlation (SRC) physics is dominant and well accounted for by SRC phenomenology
- SRC physics at high RG resolution
 - SRC pairs are components in the nuclear wave function with relative momenta above the Fermi momentum





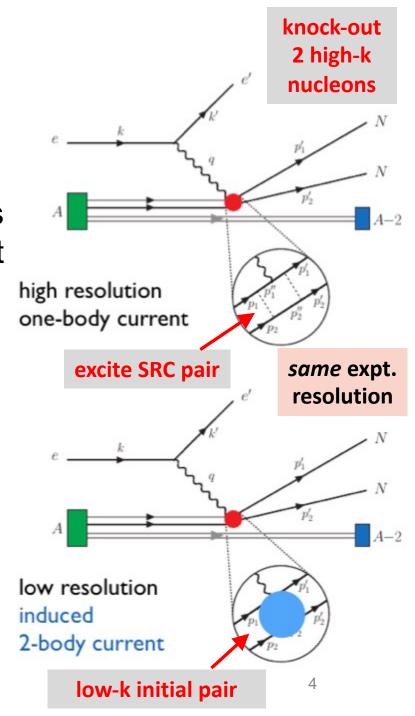
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 - The SRC *physics* is shifted into the reaction operators from the nuclear wave function (which becomes soft)
 - Operators do not become hard which simplifies calculations



Motivation

- Recent experiments have been able to isolate processes where short-range correlation (SRC) physics is dominant and well accounted for by SRC phenomenology
- SRC physics at low RG resolution
 - The SRC *physics* is shifted into the reaction operators from the nuclear wave function (which becomes soft)
 - Operators do not become hard which simplifies calculations
- Experimental resolution (set by momentum of probe) is the same in both pictures
- Same observables but different physical interpretation!



Similarity Renormalization Group (SRG)

 Evolve operators to low RG resolution

$$O(s) = U(s)O(0)U^{\dagger}(s)$$

where $s = 0 \rightarrow \infty$ and U(s) is unitary

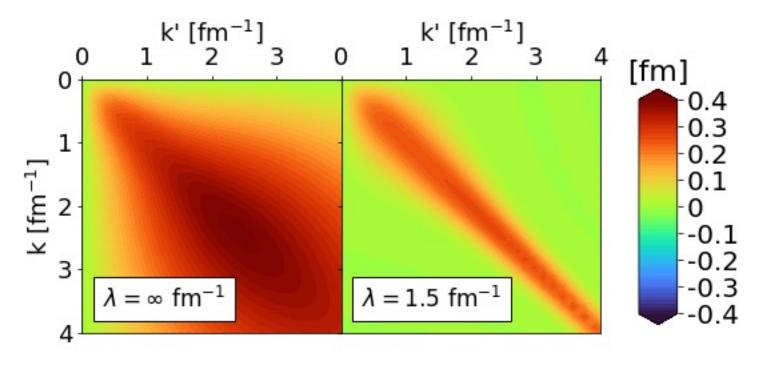


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• $\lambda = s^{-1/4}$ describes the decoupling scale of the RG evolved operator

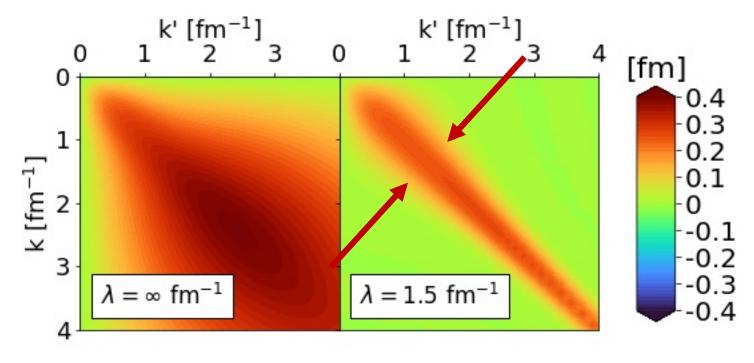


Fig. 1: Momentum space matrix elements of Argonne v18 (AV18) under SRG evolution in ¹P₁ channel.

Deuteron wave function at low RG resolution

- AV18 wave function has significant SRC
- What happens to the wave function at low RG resolution?

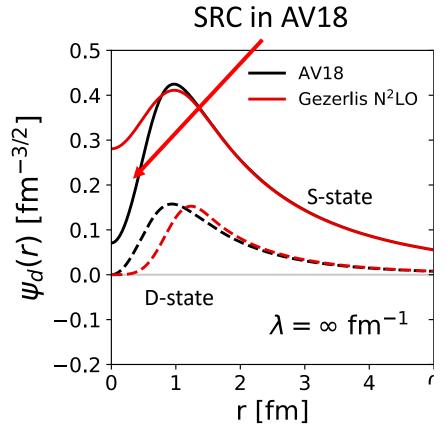


Fig. 2: SRG evolution of deuteron wave function in coordinate space for AV18 and Gezerlis N2LO¹.

Deuteron wave function at low RG resolution

- SRC physics in AV18 is gone from wave function at low RG resolution
- Deuteron wave functions become soft and D-state probability goes down
- Observables such as asymptotic D-S ratio are the same

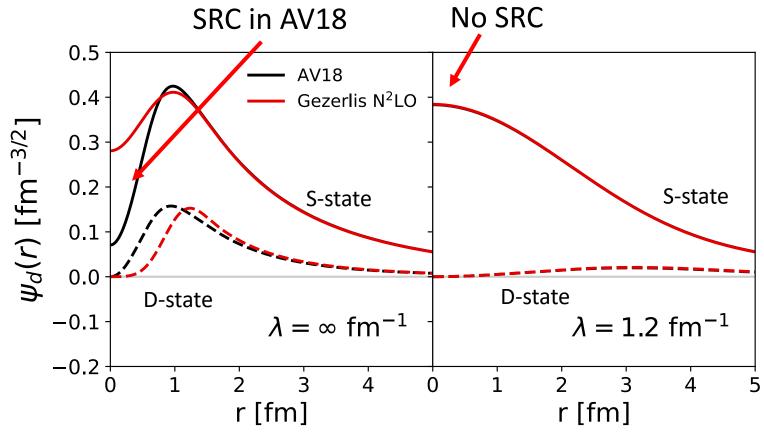


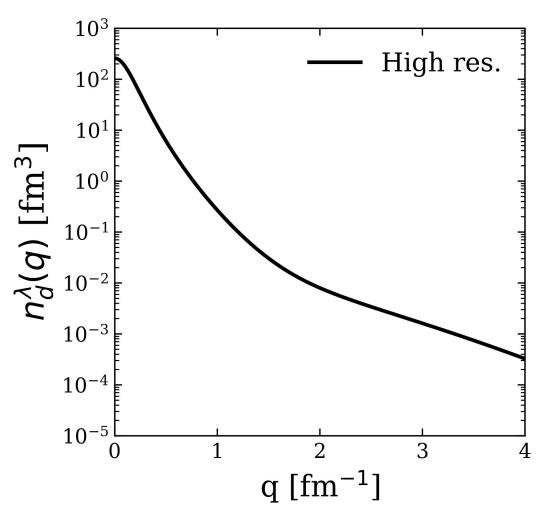
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- Soft wave functions at low RG resolution
 - Where does the SRC physics go?

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 - Where does the SRC physics go?
- SRC physics shifts to the operators $\langle \psi_f^{hi} | U_{\lambda}^{\dagger} U_{\lambda} O^{hi} U_{\lambda}^{\dagger} U_{\lambda} | \psi_i^{hi} \rangle$
- Apply SRG transformations to momentum distribution operator

$$n^{hi}(\boldsymbol{q}) = a_{\boldsymbol{q}}^{\dagger} a_{\boldsymbol{q}}$$

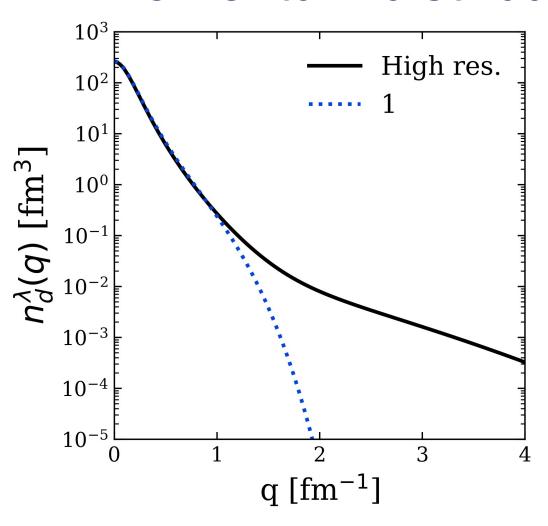
$$U_{\lambda} = 1 + \frac{1}{4} \sum_{\boldsymbol{K}, \boldsymbol{k}, \boldsymbol{k}'} \delta U_{\lambda}^{(2)}(\boldsymbol{k}, \boldsymbol{k}') a_{\boldsymbol{K}}^{\dagger} a_{\boldsymbol{K}}^{\dagger} a_{\boldsymbol{K}}^{\dagger} a_{\boldsymbol{K}}^{K} - \boldsymbol{k}' \frac{a_{\boldsymbol{K}}}{2} + \boldsymbol{k}'} + \cdots$$



$$n^{lo}(\boldsymbol{q}) = (1 + \delta U)a_{\boldsymbol{q}}^{\dagger}a_{\boldsymbol{q}}(1 + \delta U^{\dagger})$$

$$\langle \psi_d^{hi} | a_{\boldsymbol{q}}^{\dagger} a_{\boldsymbol{q}} | \psi_d^{hi} \rangle$$

Fig. 3: Contributions to deuteron momentum distribution with AV18 and $\lambda = 1.35$ fm⁻¹.

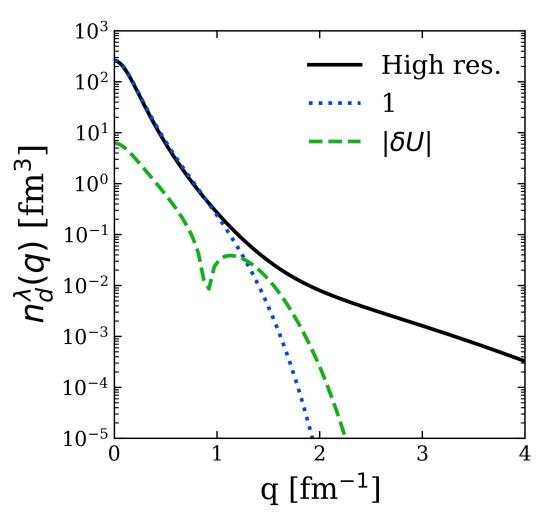


$$n^{lo}(\boldsymbol{q}) = (1 + \delta U)a_{\boldsymbol{q}}^{\dagger}a_{\boldsymbol{q}}(1 + \delta U^{\dagger})$$

$$\langle \psi_d^{hi} | a_q^{\dagger} a_q | \psi_d^{hi} \rangle$$

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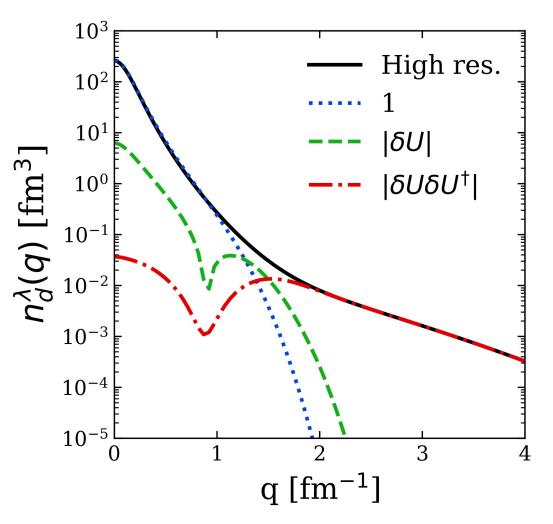
$$n^{lo}(\boldsymbol{q}) = (1 + \delta U)a_{\boldsymbol{q}}^{\dagger}a_{\boldsymbol{q}}(1 + \delta U^{\dagger})$$

$$\langle \psi_{d}^{hi} | a_{q}^{\dagger} a_{q} | \psi_{d}^{hi} \rangle$$

$$\langle \psi_{d}^{lo} | a_{q}^{\dagger} a_{q} | \psi_{d}^{lo} \rangle$$

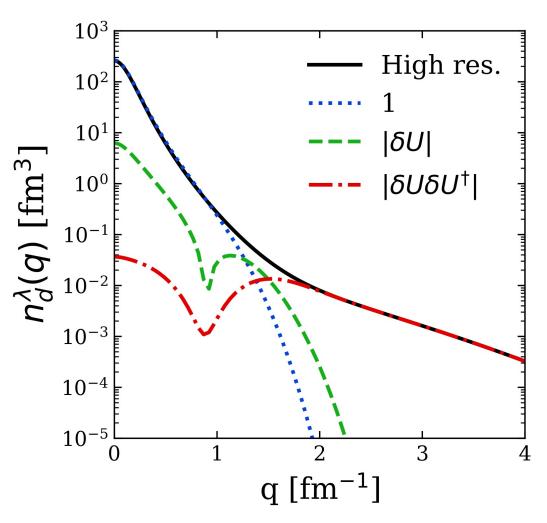
$$\langle \psi_{d}^{lo} | \delta U a_{q}^{\dagger} a_{q} + a_{q}^{\dagger} a_{q} \delta U^{\dagger} | \psi_{d}^{lo} \rangle$$

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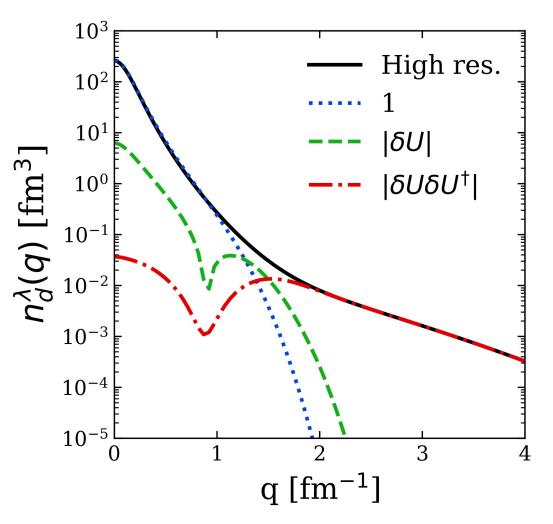
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• For high-q, the $\delta U_{\lambda} \delta U_{\lambda}^{\dagger}$ 2-body term dominates

$$\approx \sum_{\mathbf{K},\mathbf{k},\mathbf{k}'} \delta U_{\lambda}(\mathbf{k},\mathbf{q}) \delta U_{\lambda}^{\dagger}(\mathbf{q},\mathbf{k}') a_{\mathbf{K}+\mathbf{k}}^{\dagger} a_{\mathbf{K}-\mathbf{k}}^{\dagger} a_{\mathbf{K}-\mathbf{k}'}^{\dagger} a_{\mathbf{K}+\mathbf{k}'}^{\dagger}$$

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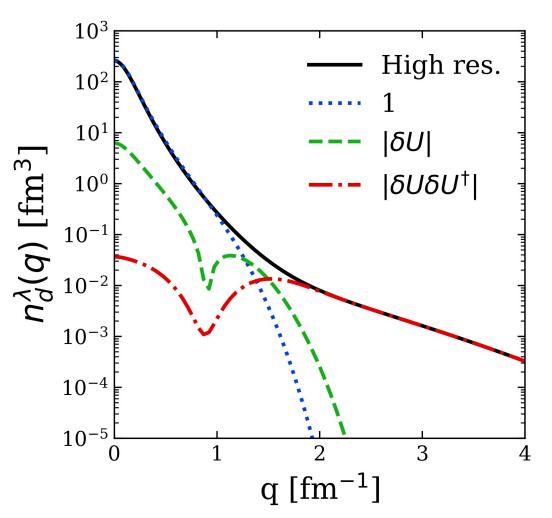
• For high-q, the $\delta U_{\lambda} \delta U_{\lambda}^{\dagger}$ 2-body term dominates

$$\approx \sum_{K,k,k'} \delta U_{\lambda}(\mathbf{k},\mathbf{q}) \delta U_{\lambda}^{\dagger}(\mathbf{q},\mathbf{k}') a_{K}^{\dagger} a_{K}^{\dagger} a_{K}^{\dagger} a_{K}^{\dagger} a_{K}^{\dagger} a_{K}^{\dagger} a_{K}^{\dagger} a_{K}^{\dagger}$$

Factorization: $\delta U_{\lambda}(\mathbf{k}, \mathbf{q}) \approx F_{\lambda}^{lo}(\mathbf{k}) F_{\lambda}^{hi}(\mathbf{q})$

$$\approx \left| F_{\lambda}^{hi}(\boldsymbol{q}) \right|^2 \sum_{\boldsymbol{K},\boldsymbol{k},\boldsymbol{k}'}^{\lambda} F_{\lambda}^{lo}(\boldsymbol{k}) F_{\lambda}^{lo}(\boldsymbol{k}') a_{\boldsymbol{K}+\boldsymbol{k}}^{\dagger} a_{\boldsymbol{K}-\boldsymbol{k}'}^{\dagger} a_{\boldsymbol{K}-\boldsymbol{k}'}^{K} a_{\boldsymbol{K}+\boldsymbol{k}'}^{K}$$

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$$\approx \sum_{\lambda} \delta U_{\lambda}(\mathbf{k}, \mathbf{q}) \delta U_{\lambda}^{\dagger}(\mathbf{q}, \mathbf{k}') a_{\underline{K} + k}^{\dagger} a_{\underline{K} - k}^{\dagger} a_{\underline{K} - k'}^{K} a_{\underline{K} + k'}^{K}$$

Apply this strategy to nuclear momentum distributions using $S_{\lambda}^{hi}(q)$ local density approximation (LDA)!

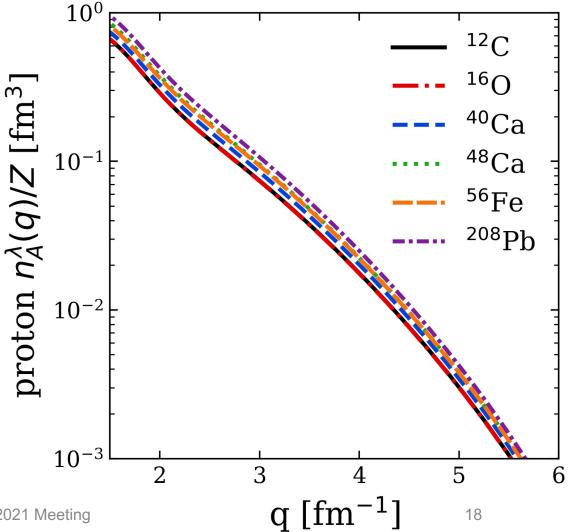
$$\approx \left| F_{\lambda}^{lo}(\boldsymbol{q}) \right|^2 \sum_{\boldsymbol{K},\boldsymbol{k},\boldsymbol{k}'} F_{\lambda}^{lo}(\boldsymbol{k}) F_{\lambda}^{lo}(\boldsymbol{k}') a_{\boldsymbol{K}}^{\dagger} a_{\boldsymbol{K}}^{\dagger} a_{\boldsymbol{K}}^{\dagger} a_{\boldsymbol{K}}^{K} a_{\boldsymbol{K}'}^{K} a_{\boldsymbol{K}+\boldsymbol{k}'}^{K}$$

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Universality

• High-q tail collapses to universal function $\approx \left| F_{\lambda}^{hi}(q) \right|^2$ fixed by 2-body

Fig. 4: Proton momentum distribution under LDA with AV18, $\lambda = 1.35$ fm⁻¹, and densities from Skyrme potential SLy4 using the HFBRAD code¹.



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- Low RG resolution calculations reproduce momentum distributions of AV18 data¹ (high RG resolution calculation)
 - Low RG works well with simple approximations and is systematically improvable
 - Absolute normalization still a work in progress (scaled up by one overall factor)

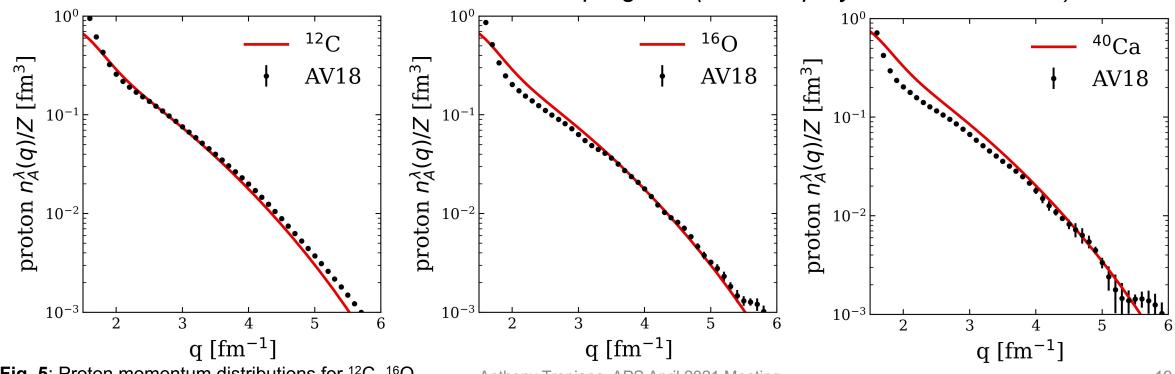


Fig. 5: Proton momentum distributions for 12 C, 16 O, and 40 Ca under LDA with AV18 and $\lambda = 1.35$ fm⁻¹.

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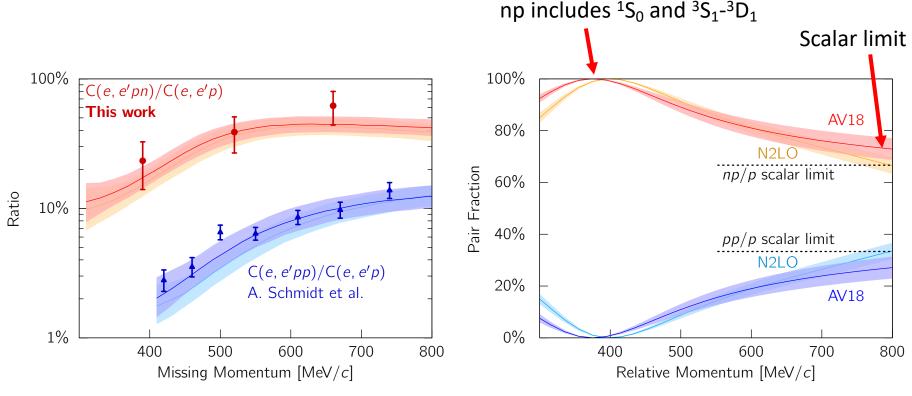
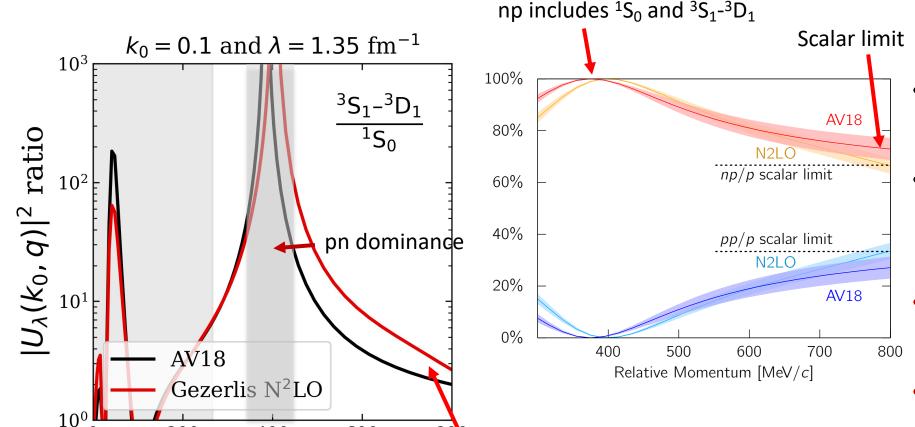


Fig. 6: (a) Ratio of two-nucleon to single-nucleon electron-scattering cross sections for carbon as a function of missing momentum. (b) Fraction of np to p and pp to p pairs versus the relative momentum. Figure from CLAS collaboration publication¹.

- At high RG resolution, the tensor force and the repulsive core of the NN interaction kicks nucleon pairs into SRCs
- np dominates because the tensor force requires spin triplet pairs (pp are spin singlets)
- Do we describe this physics at low RG resolution?



800 Scalar limit

- At low RG resolution, SRCs are suppressed in the wave function
- Consider the ratio of 3S_1 - 3D_1 to 1S_0 evolved momentum projection operator $a_q^{\dagger}a_q$
- This physics is established in the 2body system!
- Can apply to any nucleus!

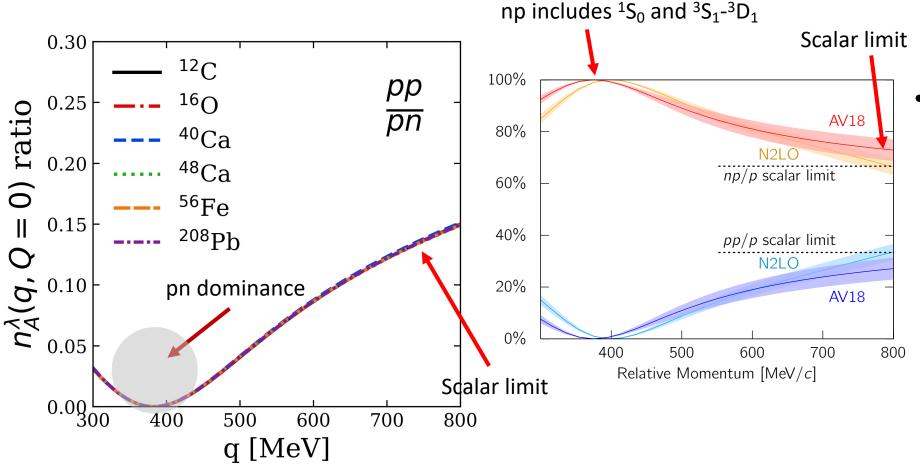
Fig. 7: ${}^{3}S_{1}$ - ${}^{3}D_{1}$ to ${}^{1}S_{0}$ ratio of SRG-evolved momentum projection operators $a_{a}^{\dagger}a_{a}$.

400

q [MeV]

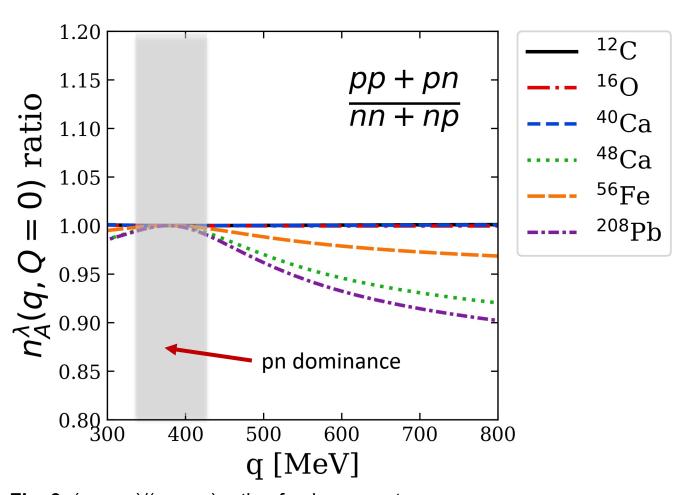
600

200



 Reproduces the characteristics of cross section ratios using low RG resolution operator with simple approximations

Fig. 8: pp/pn ratio of pair momentum distributions under LDA with AV18 and $\lambda = 1.35$ fm⁻¹.



- Ratio ~1 independent of N/Z in pn dominant region
- Ratio < 1 for nuclei where N > Z and outside pn dominant region

Fig. 9: (pp+pn)/(nn+np) ratio of pair momentum distributions under LDA with AV18 and $\lambda = 1.35$ fm⁻¹. Anthony Tropiano, APS April 2021 Meeting

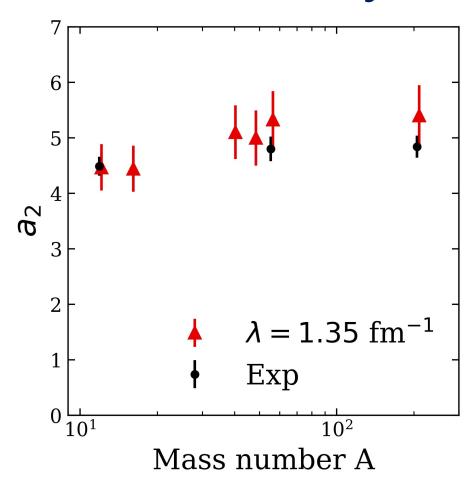


Fig. 10: a_2 scale factors using single-nucleon momentum distributions under LDA with AV18 and $\lambda = 1.35 \text{ fm}^{-1}$ compared to experimental values¹.

SRC scale factors

$$a_2 = \lim_{q \to \infty} \frac{P^A(q)}{P^d(q)} \approx \frac{\int_{\Delta p^{high}} dq P^A(q)}{\int_{\Delta p^{high}} dq P^d(q)}$$

where $P^A(q)$ is the single-nucleon probability distribution in nucleus A with error bars from varying Δp^{high}

 Decent agreement with experiment¹ and LCA calculations² but need to further test systematics

¹B. Schmookler et al. (CLAS), Nature **566**, 354 (2019)

²J. Ryckebusch et al., Phys. Rev. C **100**, 054620 (2019)

Summary and outlook

- Simple approximations work and are systematically improvable at low RG resolution
- Results suggest that we can analyze high-energy nuclear reactions using low RG resolution structure (e.g., shell model) and consistently evolved operators
 - Matching resolution scale between structure and reactions is crucial!
- Ongoing work:
 - Extend to cross sections and test scale/scheme dependence of extracted properties
 - Further investigate how final state interactions and physical interpretations depend on the RG scale
 - Apply to more complicated knock-out reactions (SRG with optical potentials) see Mostofa Hisham's talk (X13.00004)