SRG operator evolution

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Abstract

Brief description of project.

CONTENTS

I. INTRODUCTION

Results on SRG-evolved operators from several NN potentials:

– How operators evolve from band- and block-diagonal SRG transformations.

– Operator evolution for different potentials (regulators, chiral order, etc.)

II. MATHEMATICAL/COMPUTATIONAL DETAILS

A. Building SRG unitary transformations

Brief description of how to make $U(s)$.

Diagonalize initial and evolved Hamiltonians which we will call $H(0)$ and $H(s)$, respectively. This gives $\psi_\alpha(0)$ and $\psi_\alpha(s)$ for each eigenvalue indexed by α. Then the SRG unitary transformation can be computed by taking a sum over outer products of the evolved and initial wave functions:

$$U(s) = \sum_{\alpha=1}^{N} |\psi_\alpha(s)\rangle \langle\psi_\alpha(0)| \,, \tag{1}$$

where N is the dimension of the Hamiltonian matrix. Here the weights are factored into the wave functions, thus $U(s)$ is unitless.

To evolve operators, we simply apply $U(s)$:

$$O(s) = U(s)O(0)U^\dagger(s), \tag{2}$$

where $O(0)$ is the bare operator.

B. Momentum projection operator: $a_q^\dagger a_q$

Applying $a_q^\dagger a_q$ to a wave function $\psi(k)$ returns $\psi(q)$. For the discrete case, $\psi(k_i)$ is an $N \times 1$ vector and $a_q^\dagger a_q(k_i, k_j)$ is an $N \times N$ matrix where i, $j = 1, \ldots, N$. Then $a_q^\dagger a_q$ acting on $\psi(k)$ is a matrix multiplication, implying a continuous integration over d^3k. Therefore, we include a factor of $1/(k^2 dk) \implies 1/(k_i k_j \sqrt{w_i w_j})$ in $a_q^\dagger a_q(k_i, k_j)$ where w represents the momentum weights. In matrix form,

$$a_q^\dagger a_q(k_i, k_j) = \frac{\delta_{k_i q} \delta_{k_j q}}{k_i k_j \sqrt{w_i w_j}}, \tag{3}$$

which has units fm^3. To evolve operators, we apply $U(s)$ at this point. For mesh-independent figures, we must divide by an additional factor of $k_i k_j \sqrt{w_i w_j}$.

C. Momentum distribution function: ϕ^2

We diagonalize the Hamiltonian for eigenvectors ϕ_α. In the ${}^3S_1$-${}^3D_1$ coupled channel, the S-component is given by $\phi_\alpha[: N]$ and the D-component by $\phi_\alpha[N :]$ where N is the length of the momentum mesh. Then the momentum distribution is given by,

$$|\phi(k)|^2 = |\phi_\alpha[: N]|^2 + |\phi_\alpha[N :]|^2. \tag{4}$$

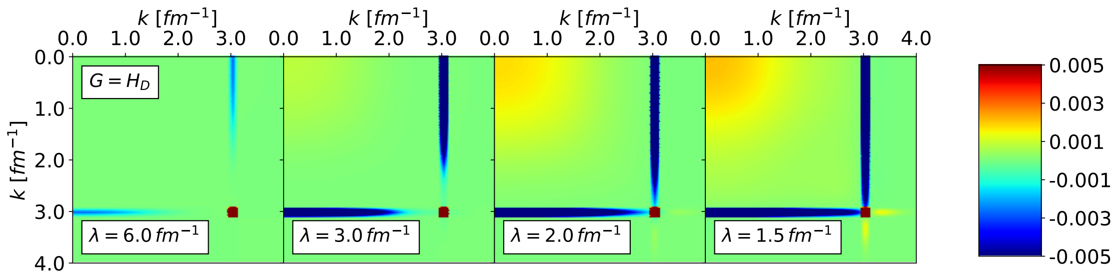
This satisfies the normalization condition $\sum_{i=1}^{N} |\phi(k_i)|^2 = 1$, implying that the factor $k^2 dk$ (or in the discrete case $k_i^2 w_i$) is factored into the wave function. For mesh-independent figures, divide by $k_i^2 w_i$.

III. RESULTS

Organize this according to the figures. Format should be description of the calculation (previous section), followed by the figure, followed by takeaways.

A. Entem-Machleidt N^{3}LO non-local potential

Add takeaways for these figures.



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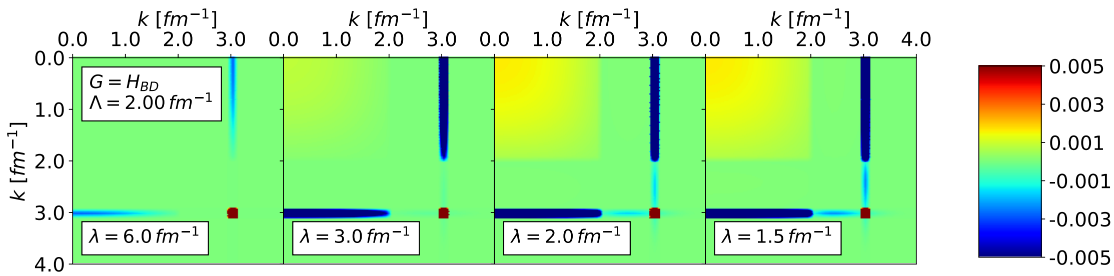

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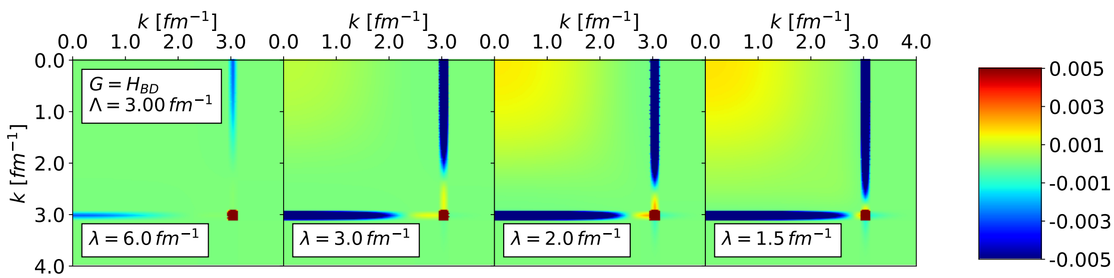


(c)

FIG. 1: Matrix elements of $\langle k|a_q^\dagger a_q|k'\rangle$ SRG-evolving in λ right to left under transformations from the Entem-Machleidt N^{3}LO non-local potential with the Wegner generator (a) and block-diagonal generators decoupling at $\Lambda = 2$ and 3 fm^{-1} (b and c). Here $q = 3$ fm^{-1}.

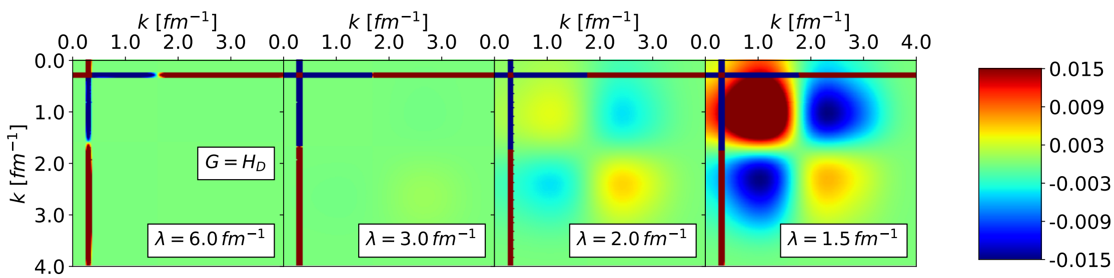
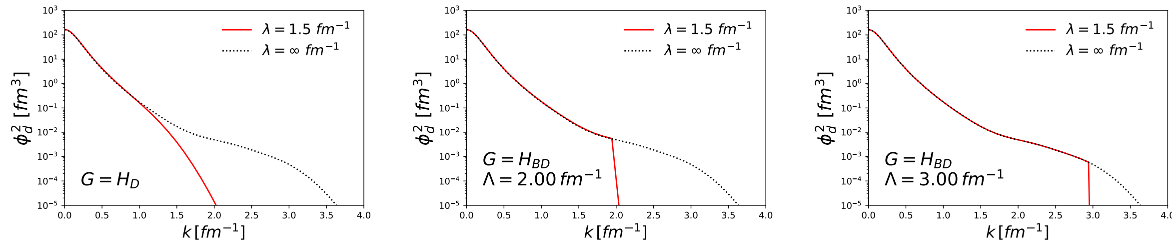
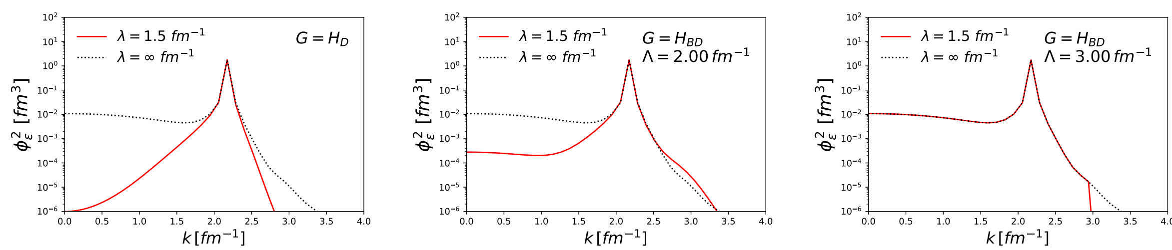


FIG. 2: Matrix elements of $\langle k|a_q^\dagger a_q|k'\rangle$ SRG-evolving in λ right to left under transformations from the Entem-Machleidt N^{3}LO non-local potential with the Wegner generator. Here $q = 0.3$ fm^{-1}.



(a) (b) (c)

FIG. 3: Momentum probability densities of the deuteron SRG-evolving the wave function to $\lambda = 1.5$ fm^{-1} from the Entem-Machleidt N^{3}LO non-local potential with the Wegner generator (a) and block-diagonal generators decoupling at $\Lambda = 2$ and 3 fm^{-1} (b and c). The black dotted line corresponds to the momentum probability density of the initial deuteron wave function.

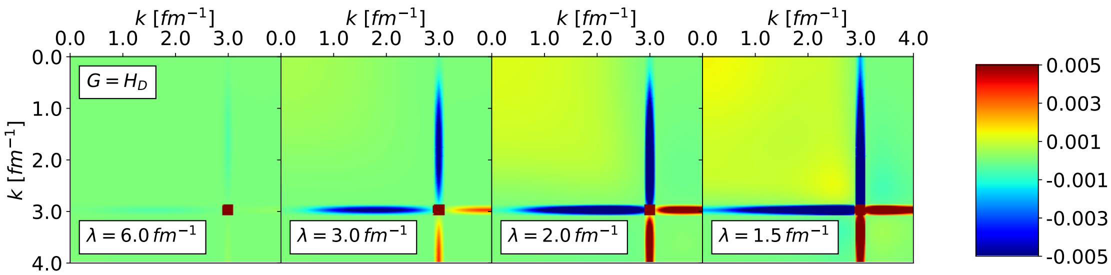


(a) (b) (c)

FIG. 4: Momentum probability densities of the continuum state at $\epsilon \approx 200$ MeV SRG-evolving the wave function to $\lambda = 1.5$ fm^{-1} from the Entem-Machleidt N^{3}LO non-local potential with the Wegner generator (a) and block-diagonal generators decoupling at $\Lambda = 2$ and 3 fm^{-1} (b and c). The black dotted line corresponds to the initial momentum probability density.

B. RKE N^{3}LO and N^{4}LO semi-local potentials

Add takeaways for these figures.



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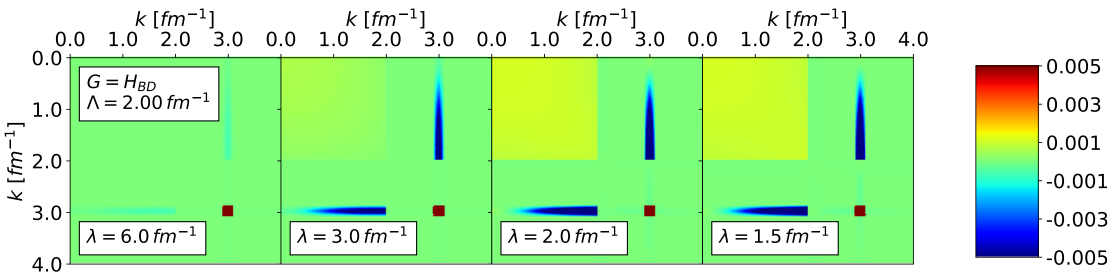
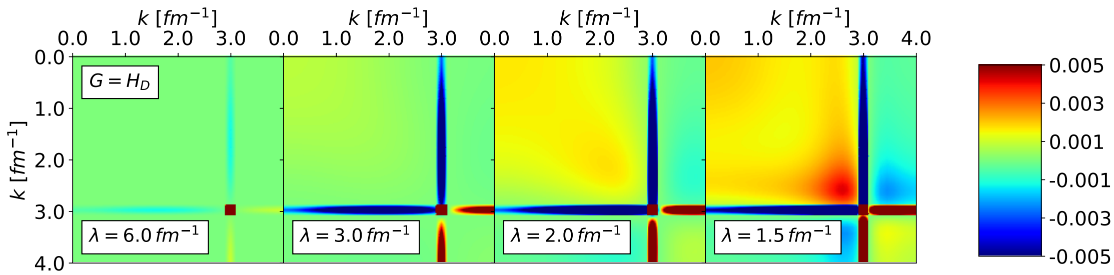


(b)

FIG. 5: Matrix elements of $\langle k|a_q^\dagger a_q|k'\rangle$ SRG-evolving in λ right to left under transformations from the RKE N^{3}LO semi-local potential with the Wegner generator (a) and block-diagonal generator decoupling at $\Lambda = 2$ fm^{-1} (b). Here $q = 3$ fm^{-1} and the EFT cutoff is 450 MeV.



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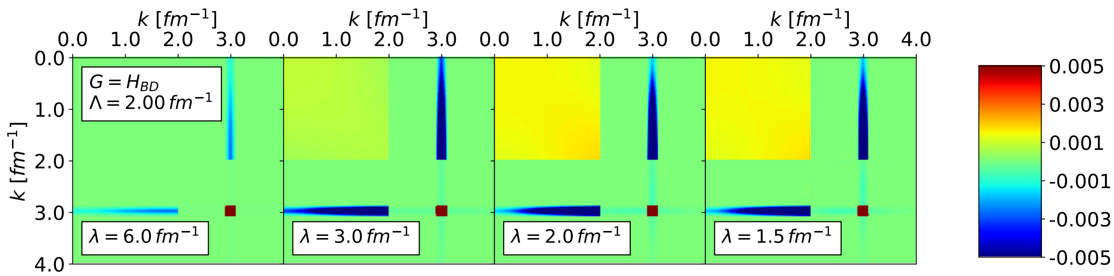
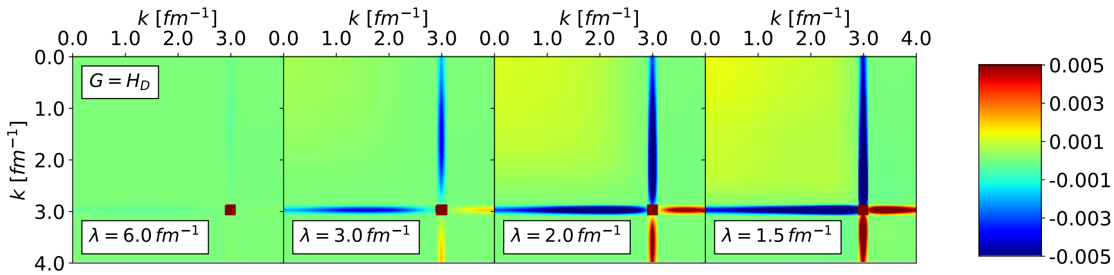


(b)

FIG. 6: Matrix elements of $\langle k|a_q^\dagger a_q|k'\rangle$ SRG-evolving in λ right to left under transformations from the RKE N^{3}LO semi-local potential with the Wegner generator (a) and block-diagonal generator decoupling at $\Lambda = 2$ fm^{-1} (b). Here $q = 3$ fm^{-1} and the EFT cutoff is 500 MeV.



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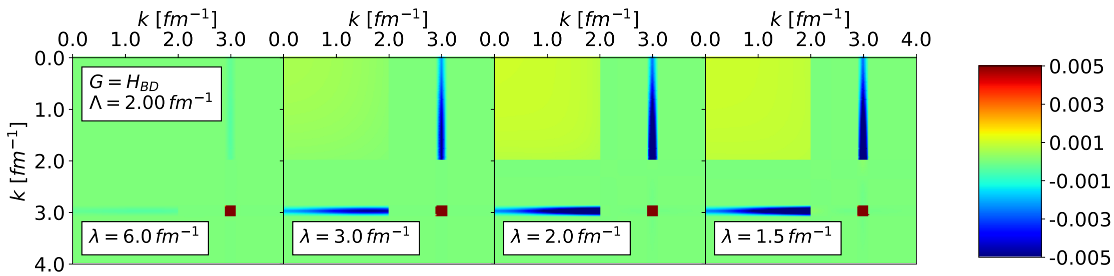
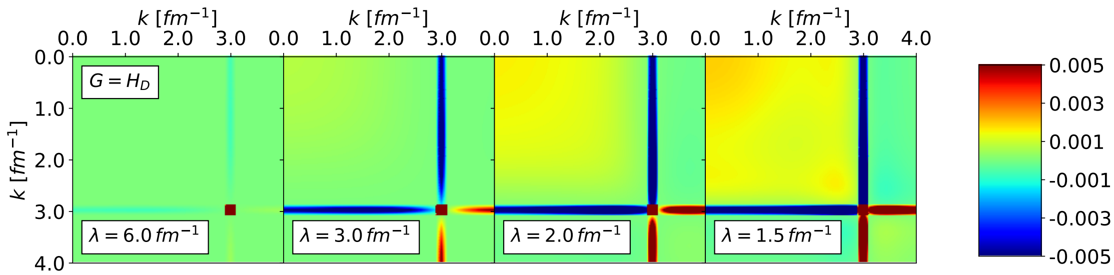


(b)

FIG. 7: Matrix elements of $\langle k|a_q^\dagger a_q|k'\rangle$ SRG-evolving in λ right to left under transformations from the RKE N^{4}LO semi-local potential with the Wegner generator (a) and block-diagonal generator decoupling at $\Lambda = 2$ fm^{-1} (b). Here $q = 3$ fm^{-1} and the EFT cutoff is 450 MeV.



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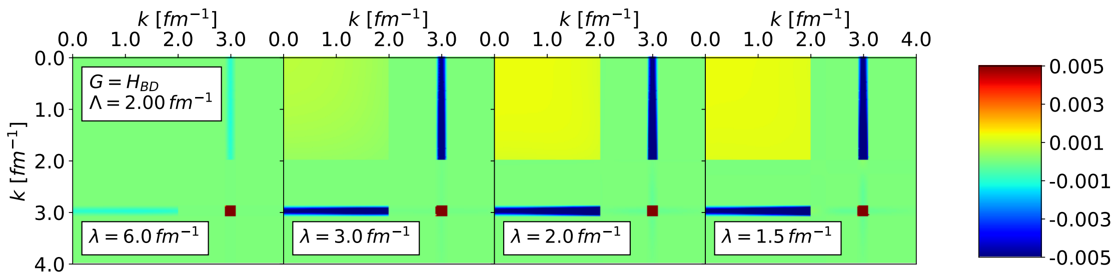

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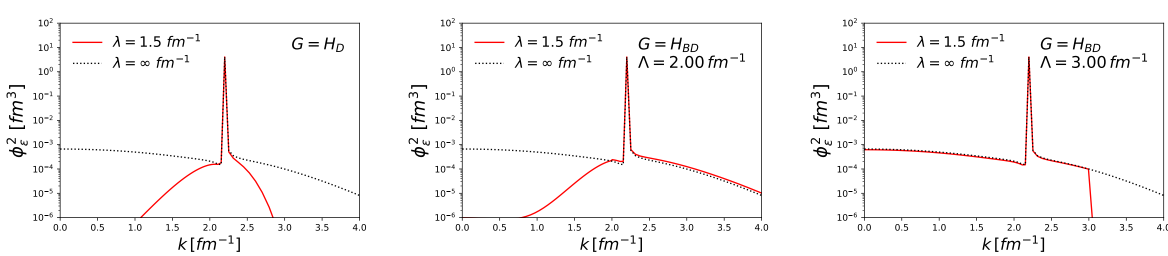
FIG. 8: Matrix elements of $\langle k|a_q^\dagger a_q|k'\rangle$ SRG-evolving in λ right to left under transformations from the RKE N^{4}LO semi-local potential with the Wegner generator (a) and block-diagonal generator decoupling at $\Lambda = 2$ fm^{-1} (b). Here $q = 3$ fm^{-1} and the EFT cutoff is 500 MeV.

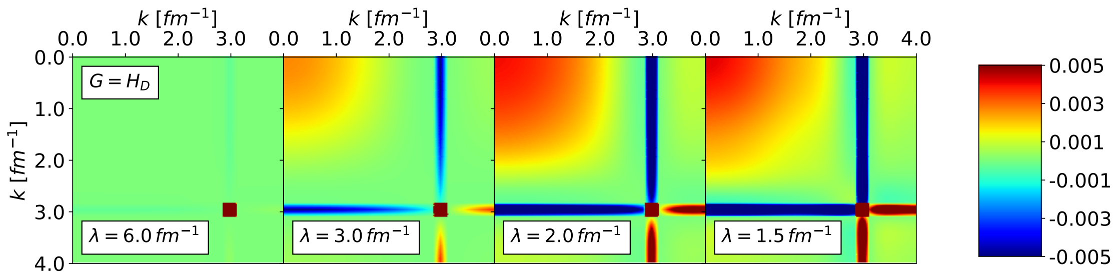


(a) (b) (c)

FIG. 9: Momentum probability densities of the continuum state at $\epsilon \approx 200$ MeV SRG-evolving the wave function to $\lambda = 1.5$ fm^{-1} from the RKE N^{4}LO semi-local potential with the Wegner generator (a) and block-diagonal generators decoupling at $\Lambda = 2$ and 3 fm^{-1} (b and c). The black dotted line corresponds to the initial momentum probability density. Here the EFT cutoff is 450 MeV.

C. Gezerlis N^{2}LO local potentials

Add takeaways for these figures.



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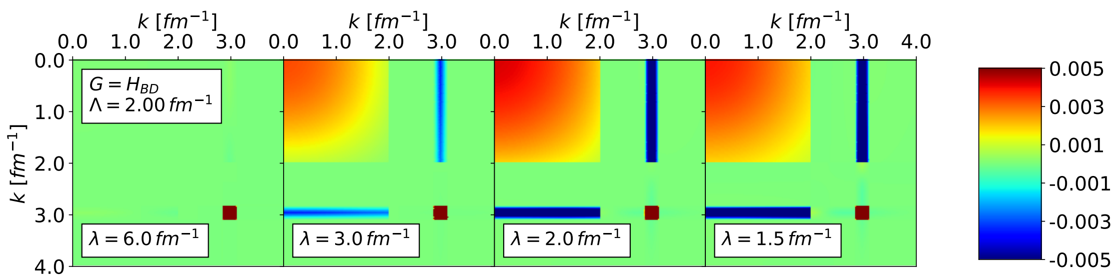

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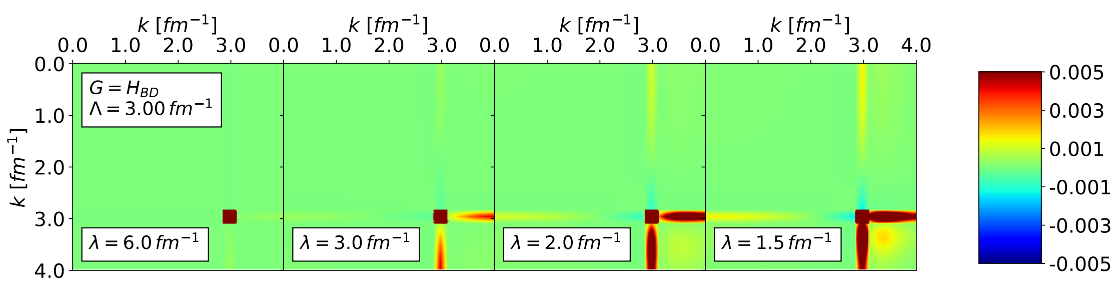
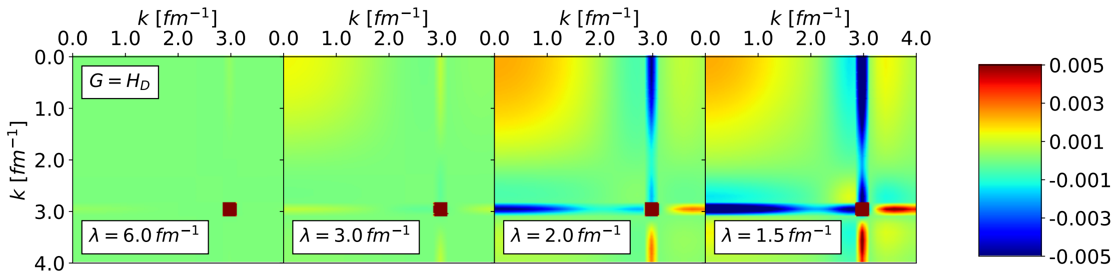


(c)

FIG. 10: Matrix elements of $\langle k|a_q^\dagger a_q|k'\rangle$ SRG-evolving in λ right to left under transformations from the Gezerlis et al. N^{2}LO local potential with the Wegner generator (a) and block-diagonal generators decoupling at $\Lambda = 2$ and 3 fm^{-1} (b and c). Here $q = 3$ fm^{-1} and the EFT cutoff is 1 fm.



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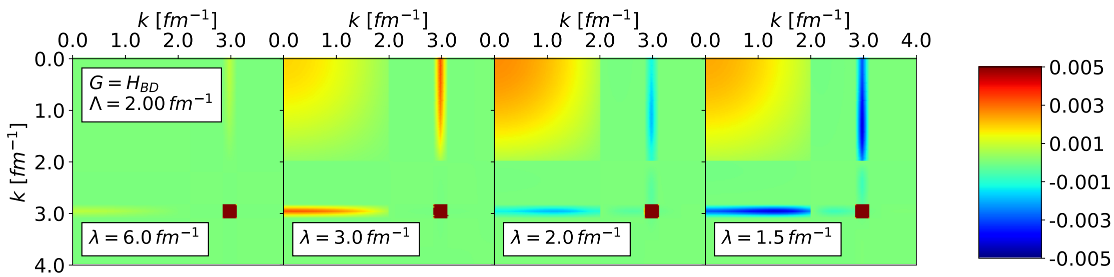

(b)

FIG. 11: Matrix elements of $\langle k|a_q^\dagger a_q|k'\rangle$ SRG-evolving in λ right to left under transformations from the Gezerlis et al. N^{2}LO local potential with the Wegner generator (a) and block-diagonal generator decoupling at $\Lambda = 2$ fm^{-1} (b). Here $q = 3$ fm^{-1} and the EFT cutoff is 1.2 fm.

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[2] P. Reinert, H. Krebs, and E. Epelbaum, Eur. Phys. J. A **54**, 86 (2018), arXiv:1711.08821 [nucl-th].

[3] A. Gezerlis, I. Tews, E. Epelbaum, M. Freunek, S. Gandolfi, K. Hebeler, A. Nogga, and A. Schwenk, Phys. Rev. C **90**, 054323 (2014), arXiv:1406.0454 [nucl-th].