# Short-range correlation physics at low RG resolution

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ajt, S.K. Bogner, and R.J. Furnstahl, arXiv:2006.11186

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ajt, S.K. Bogner, and R.J. Furnstahl, in progress



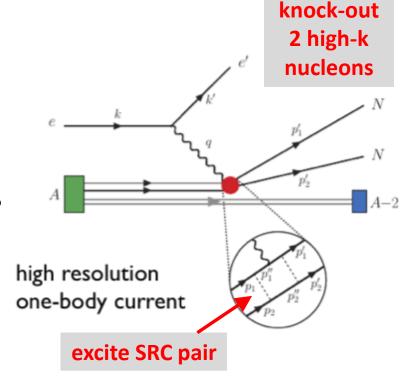


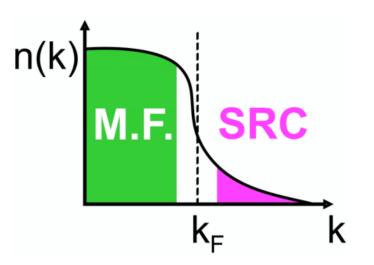




## Motivation

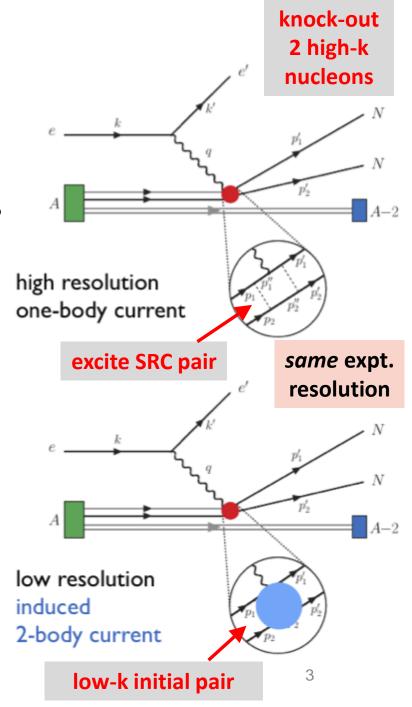
- Recent experiments have been able to isolate processes where short-range correlation (SRC) physics is dominant and well accounted for by SRC phenomenology
- High RG resolution description of SRC physics
  - SRC pairs are components in the nuclear wave function with relative momenta above the Fermi momentum





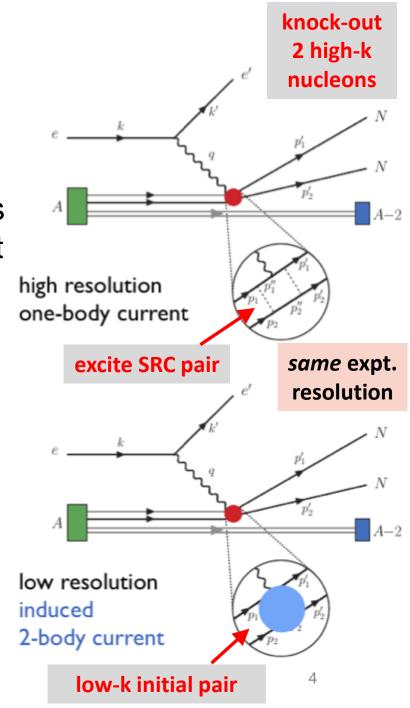
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- Low RG resolution description of SRC physics
  - Using renormalization group (RG) methods we can tune the scale to low RG resolution
  - The SRC *physics* is shifted into the reaction operators from the nuclear wave function (which becomes soft)



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  - Using renormalization group (RG) methods we can tune the scale to low RG resolution
  - The SRC *physics* is shifted into the reaction operators from the nuclear wave function (which becomes soft)
- Experimental resolution (set by momentum of probe) is the same in both pictures
- Same observables but different physical interpretation!



# Similarity Renormalization Group (SRG)

 Evolve operators to low RG resolution

$$O(s) = U(s)O(0)U^{\dagger}(s)$$

where  $s = 0 \rightarrow \infty$  and U(s) is unitary

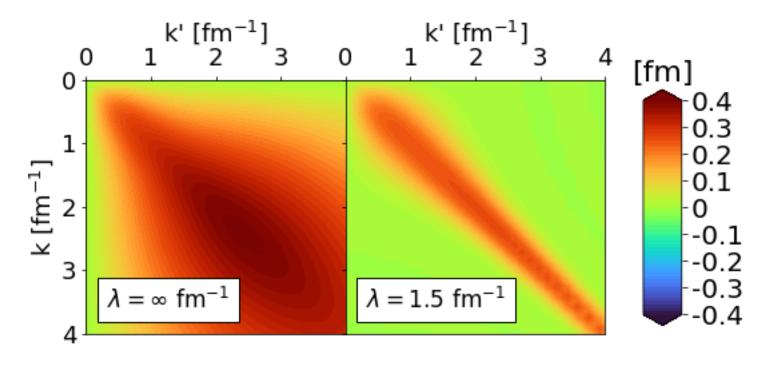


Fig. 1: Momentum space matrix elements of Argonne v18 (AV18) under SRG evolution in <sup>1</sup>P<sub>1</sub> channel.

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•  $\lambda = s^{-1/4}$  describes the decoupling scale of the RG evolved operator

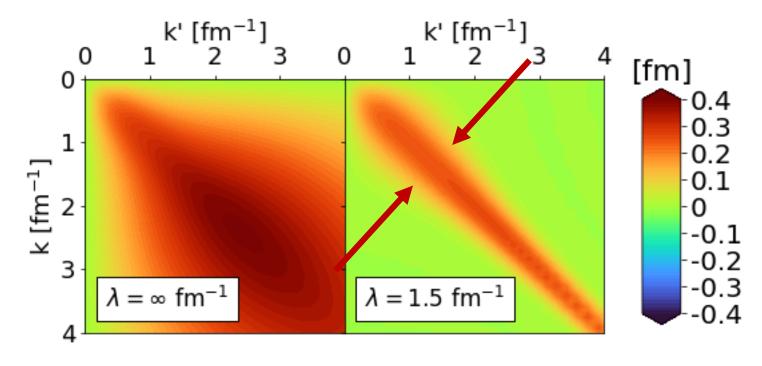


Fig. 1: Momentum space matrix elements of Argonne v18 (AV18) under SRG evolution in <sup>1</sup>P₁ channel.

#### Deuteron wave function at low RG resolution

- AV18 wave function has significant SRC
- What happens to the wave function at low RG resolution?

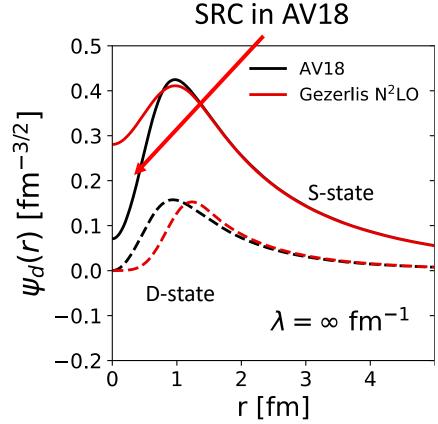


Fig. 2: SRG evolution of deuteron wave function in coordinate space for AV18 and Gezerlis N2LO<sup>1</sup>.

#### Deuteron wave function at low RG resolution

- SRC physics in AV18 is gone from wave function at low RG resolution
- Deuteron wave functions become soft and D-state probability goes down
- Observables such as asymptotic D-S ratio are the same

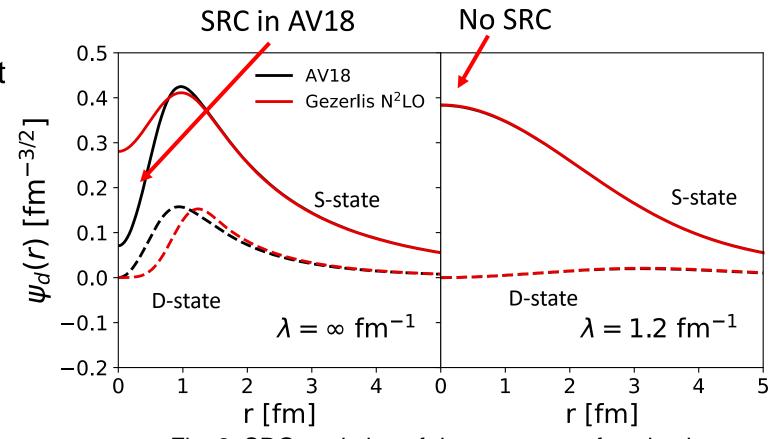


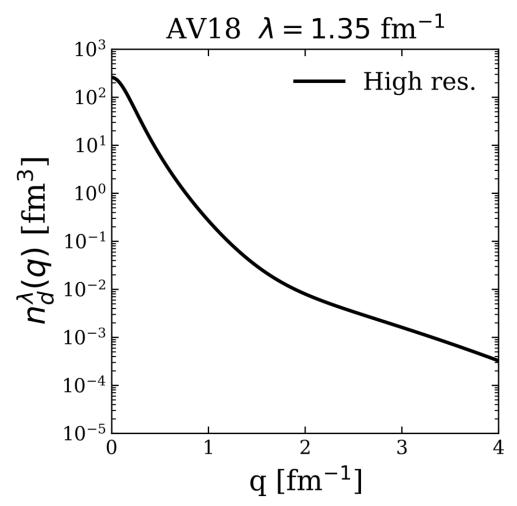
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- Soft wave functions at low RG resolution
  - Where does the SRC physics go?

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  - Where does the SRC physics go?
- SRC physics shifts to the operators  $\langle \psi_A^{hi} | U_\lambda^{\dagger} U_\lambda O^{hi} U_\lambda^{\dagger} U_\lambda | \psi_A^{hi} \rangle$
- Apply SRG transformations to momentum distribution operator

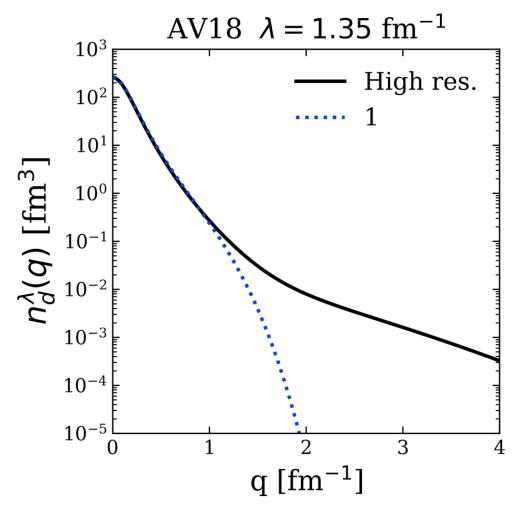
$$n^{hi}(\boldsymbol{q}) = a_{\boldsymbol{q}}^{\dagger} a_{\boldsymbol{q}}$$

$$U_{\lambda} = 1 + \frac{1}{4} \sum_{\boldsymbol{K}, \boldsymbol{k}, \boldsymbol{k}'} \delta U_{\lambda}^{(2)}(\boldsymbol{k}, \boldsymbol{k}') a_{\boldsymbol{K} + \boldsymbol{k}}^{\dagger} a_{\boldsymbol{K} - \boldsymbol{k}'}^{\dagger} a_{\boldsymbol{K} - \boldsymbol{k}'} a_{\boldsymbol{K} + \boldsymbol{k}'}^{\boldsymbol{K} + \boldsymbol{k}'} + \cdots$$



$$n^{lo}(\boldsymbol{q}) = (1 + \delta U)a_{\boldsymbol{q}}^{\dagger}a_{\boldsymbol{q}}(1 + \delta U^{\dagger})$$

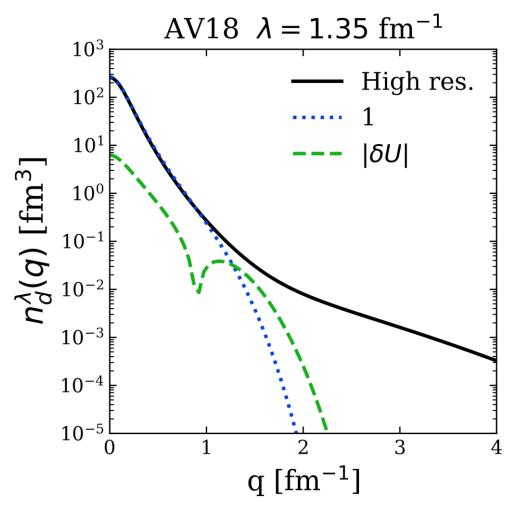
$$\langle \psi_d^{hi} | a_{\boldsymbol{q}}^{\dagger} a_{\boldsymbol{q}} | \psi_d^{hi} \rangle$$



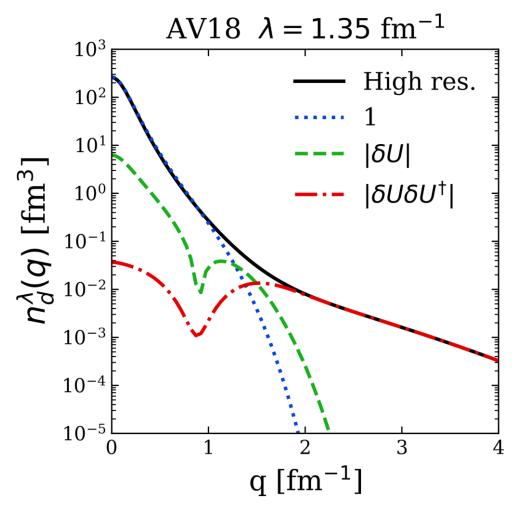
$$n^{lo}(\boldsymbol{q}) = (1 + \delta U)a_{\boldsymbol{q}}^{\dagger}a_{\boldsymbol{q}}(1 + \delta U^{\dagger})$$

$$\langle \psi_d^{hi} | a_q^{\dagger} a_q | \psi_d^{hi} \rangle$$

$$\langle \psi_d^{lo} | a_q^{\dagger} a_q | \psi_d^{lo} \rangle$$



$$n^{lo}(\boldsymbol{q}) = (1 + \delta U)a_{\boldsymbol{q}}^{\dagger}a_{\boldsymbol{q}}(1 + \delta U^{\dagger})$$



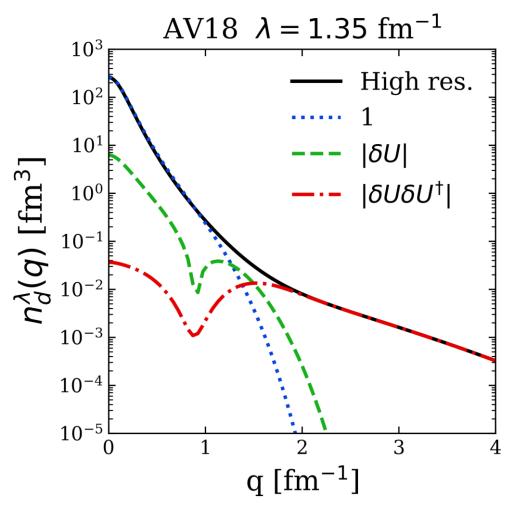
$$n^{lo}(\boldsymbol{q}) = (1 + \delta U)a_{\boldsymbol{q}}^{\dagger}a_{\boldsymbol{q}}(1 + \delta U^{\dagger})$$

$$\langle \psi_{d}^{hi} | a_{q}^{\dagger} a_{q} | \psi_{d}^{hi} \rangle$$

$$\langle \psi_{d}^{lo} | a_{q}^{\dagger} a_{q} | \psi_{d}^{lo} \rangle$$

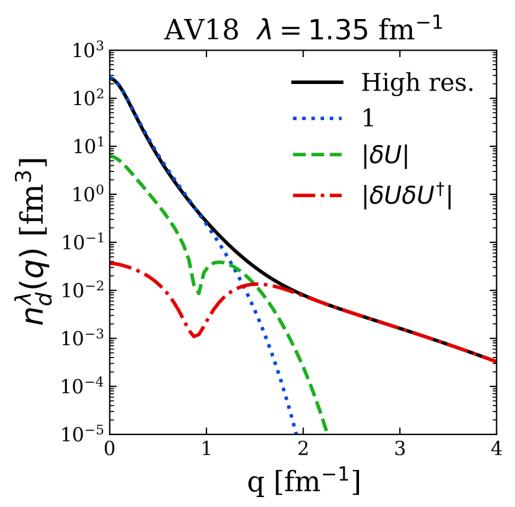
$$\langle \psi_{d}^{lo} | \delta U a_{q}^{\dagger} a_{q} + a_{q}^{\dagger} a_{q} \delta U^{\dagger} | \psi_{d}^{lo} \rangle$$

$$\langle \psi_{d}^{lo} | \delta U a_{q}^{\dagger} a_{q} \delta U^{\dagger} | \psi_{d}^{lo} \rangle$$



• For high-q, the  $\delta U_{\lambda} \delta U_{\lambda}^{\dagger}$  term dominates

$$\approx \sum_{\boldsymbol{K},\boldsymbol{k},\boldsymbol{k}'} \delta U_{\lambda}(\boldsymbol{k},\boldsymbol{q}) \delta U_{\lambda}^{\dagger}(\boldsymbol{q},\boldsymbol{k}') a_{\underline{K}+\boldsymbol{k}}^{\dagger} a_{\underline{K}-\boldsymbol{k}'}^{\dagger} a_{\underline{K}-\boldsymbol{k}'}^{K} a_{\underline{K}+\boldsymbol{k}'}^{K}$$

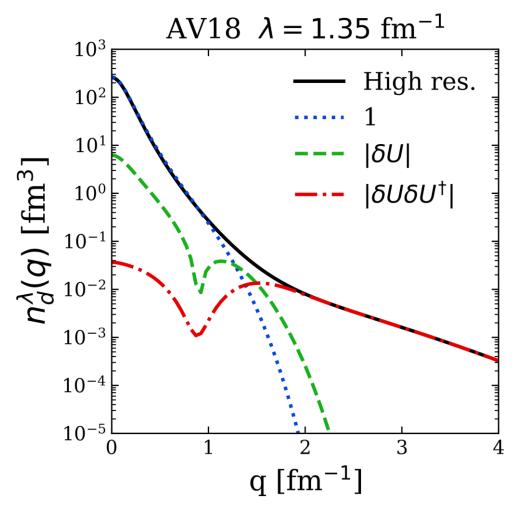


• For high-q, the  $\delta U_{\lambda} \delta U_{\lambda}^{\dagger}$  term dominates

$$\approx \sum_{K,k,k'} \delta U_{\lambda}(k,q) \delta U_{\lambda}^{\dagger}(q,k') a_{K}^{\dagger} a_{K}^{\dagger} a_{K}^{\dagger} a_{K}^{\dagger} a_{K}^{\dagger} a_{K}^{\dagger} a_{K}^{\dagger} a_{K}^{\dagger} a_{K}^{\dagger}$$

Factorization:  $\delta U_{\lambda}(\mathbf{k}, \mathbf{q}) \approx F_{\lambda}^{lo}(\mathbf{k}) F_{\lambda}^{hi}(\mathbf{q})$ 

$$\approx \left| F_{\lambda}^{hi}(\boldsymbol{q}) \right|^2 \sum_{\boldsymbol{K},\boldsymbol{k},\boldsymbol{k}'}^{\lambda} F_{\lambda}^{lo}(\boldsymbol{k}) F_{\lambda}^{lo}(\boldsymbol{k}') a_{\boldsymbol{K}+\boldsymbol{k}}^{\dagger} a_{\boldsymbol{K}-\boldsymbol{k}'}^{\dagger} a_{\boldsymbol{K}-\boldsymbol{k}'}^{K} a_{\boldsymbol{K}+\boldsymbol{k}'}^{K}$$



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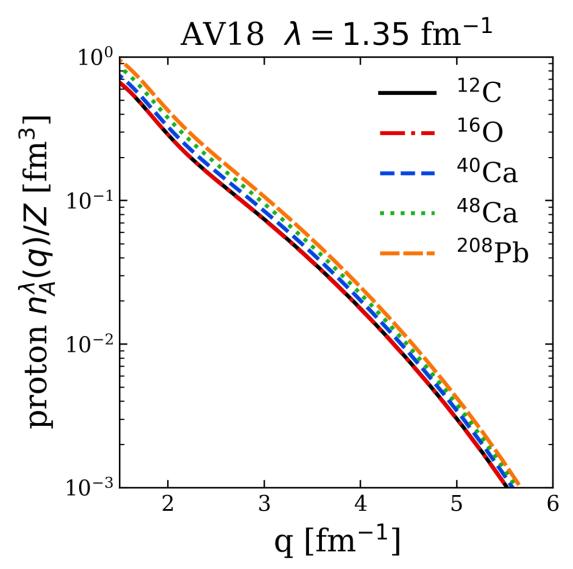
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Apply this strategy to nuclear momentum distributions using local density approximation (LDA)!

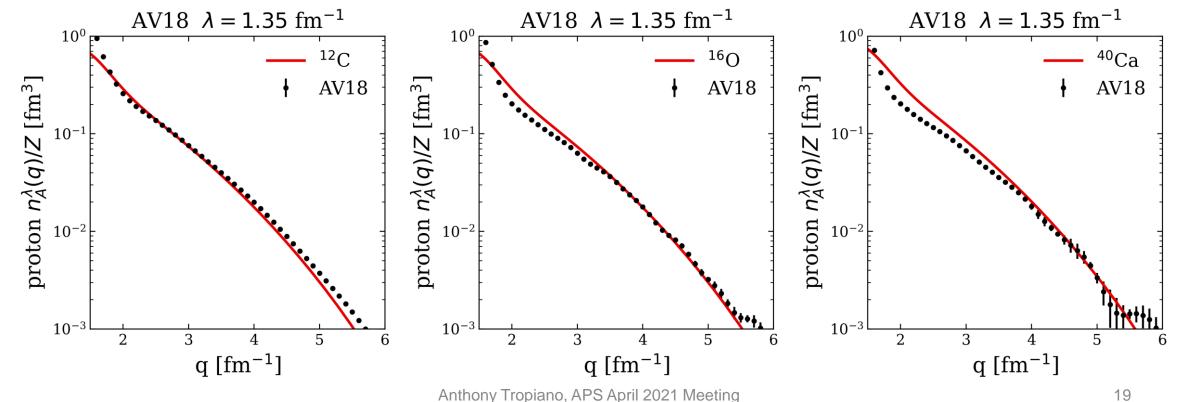
$$\approx \left| F_{\lambda}^{lo}(\boldsymbol{q}) \right|^2 \sum_{\boldsymbol{K},\boldsymbol{k},\boldsymbol{k}'} F_{\lambda}^{lo}(\boldsymbol{k}) F_{\lambda}^{lo}(\boldsymbol{k}') a_{\boldsymbol{K}+\boldsymbol{k}}^{\dagger} a_{\boldsymbol{K}-\boldsymbol{k}'}^{\dagger} a_{\boldsymbol{K}-\boldsymbol{k}'}^{\dagger} a_{\boldsymbol{K}+\boldsymbol{k}'}^{\dagger}$$

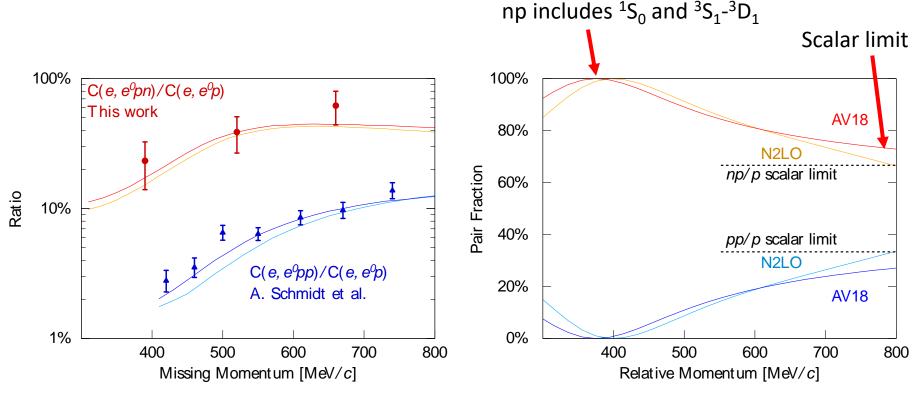
 $F_{a}^{hi}(q)$ 

- Universality
  - High-q tail collapses to universal function  $\approx \left| F_{\lambda}^{lo}(\mathbf{q}) \right|^2$  fixed by 2-body



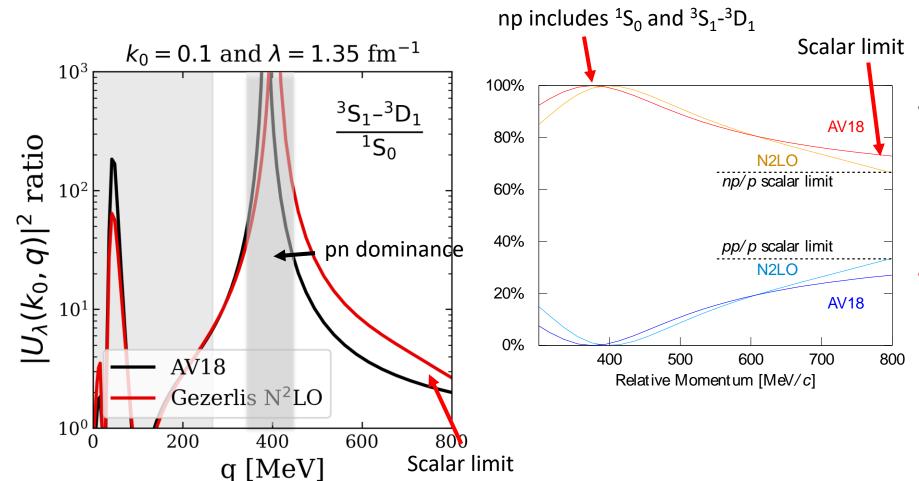
- Low RG resolution calculations reproduce momentum distributions of AV18 data (high RG resolution calculation)
  - Absolute normalization still a work in progress (scaled up by one overall factor)



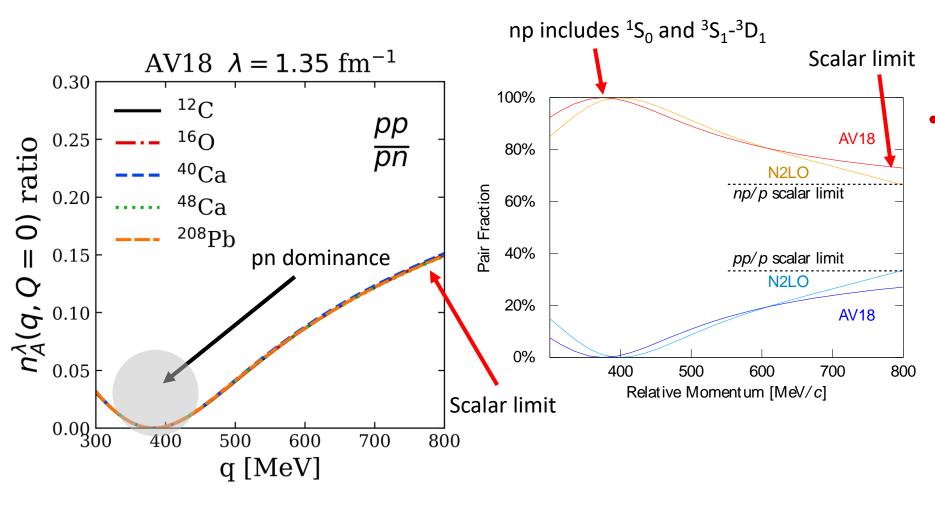


- pp pairs are spinsinglets whereas the tensor force requires spin-triplets
- pn pairs dominate in region of tensor force

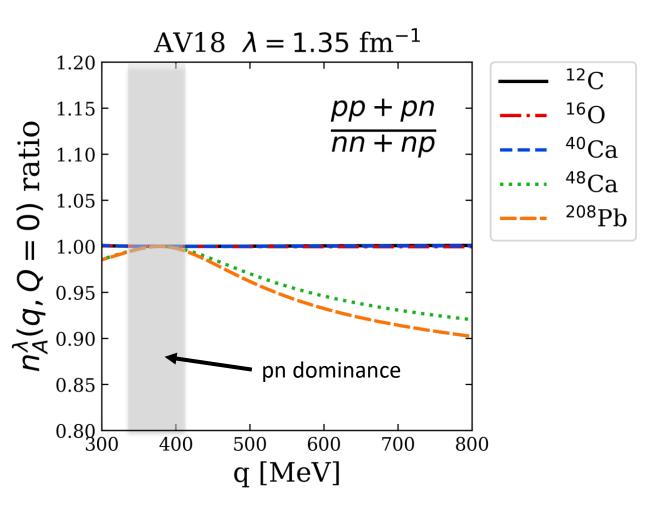
Fig.: (a) Ratio of two-nucleon to single-nucleon electron-scattering cross sections for carbon as a function of missing momentum. (b) Fraction of np to p and pp to p pairs versus the relative momentum. Figure from CLAS collaboration publication<sup>1</sup>. (add ref)



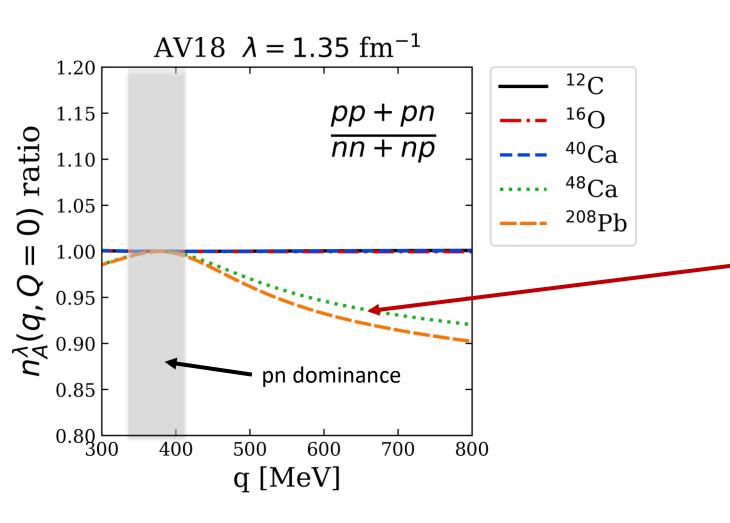
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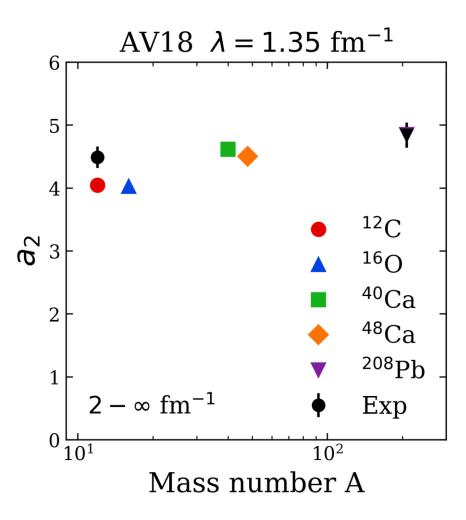
Low RG resolution picture reproduces characteristics of SRC experiments



 Ratio should be ~1 independent of N/Z in pn dominant region



- Ratio should be ~1 independent of N/Z in pn dominant region
- Outside pn dominant
   region ratio < 1 for nuclei where N > Z



SRC scale factors

$$a_2 = \lim_{q \to \infty} \frac{P^A(q)}{P^d(q)} \approx \frac{\int_2^{\infty} dq P^A(q)}{\int_2^{\infty} dq P^d(q)}$$

where  $P^A(q)$  is the single-nucleon probability distribution in nucleus A

 Decent agreement with experiment and LCA calculations (add ref.) but need to further test systematics

## Summary and outlook

- Results suggest that we can analyze high-energy nuclear reactions using low RG resolution structure (e.g., shell model) and consistently evolved operators
  - Matching resolution scale between structure and reactions is crucial!
- Ongoing work:
  - Extend to cross sections and test scale and scheme dependence of extracted properties
  - Further investigate how final state interactions (FSI's) and physical interpretations depend on the RG scale
  - Apply to knock-out reactions (optical potentials) see Mostofa Hisham's talk (add time/session)

## **Extras**