

(3/12/21)

- Starting point is $\hat{n}_\lambda^z(\vec{q})$ after evaluation of contractions. (\vec{q} is not a relative momentum here!)

$$\begin{aligned}
 \hat{n}_\lambda^z(\vec{q}) = & \sqrt{\left\langle \sum_{\sigma} \Theta(k_f^z - q) \right\rangle} \quad \text{Term 1} \\
 & + \sum_{\sigma\sigma'z'} \int \frac{d^3 k}{(2\pi)^3} \left[\underbrace{\left\langle k_{\sigma z} \sigma z | \delta U | k_{\sigma' z'} \sigma' z' \right\rangle}_{\text{Term 2}} \right. \\
 & \left. + \left\langle k_{\sigma z} \sigma z | \delta U + k_{\sigma z} \sigma z \right\rangle \right] \Theta(k_f^z - q) \Theta(k_f^{z'} - |\vec{q} - \vec{k}|) \\
 & + \frac{1}{2} \sum_{\sigma\sigma''\sigma''' z'' z'''} \int \frac{d^3 k}{(2\pi)^3} \frac{d^3 K}{(2\pi)^3} \left\langle k_{\sigma z} \sigma z | \delta U | \vec{q} - \frac{1}{2} \vec{k} \sigma z'' z''' \right\rangle \times \\
 & \left. \left(\vec{q} - \frac{1}{2} \vec{k} \sigma z'' z''' | \delta U + \vec{k} \sigma z'' z''' \right) \Theta(k_f^{z'} - |\frac{1}{2} \vec{k} + \vec{k}|) \Theta(k_f^{z''} - |\frac{1}{2} \vec{k} - \vec{k}|) \right\rangle \quad (1)
 \end{aligned}$$

Term 1 is simple. Contribution is $2 \Theta(k_f^z - q)$.

No 2nd + 3rd term (multiply 2nd term by 2 for full 2+3 contribution).

$$\begin{aligned}
 & 2 \sum_{\sigma\sigma'z'} \int \frac{d^3 k}{(2\pi)^3} \left\langle k_{\sigma z} \sigma z | \delta U | k_{\sigma' z'} \sigma' z' \right\rangle \Theta(k_f^z - q) \Theta(k_f^{z'} - |\vec{q} - \vec{k}|) \\
 & = 2 \sum_{\sigma\sigma'z'} \sum_{SM_1} \sum_{SM_2} \sum_{TM_1} \sum_{TM_2} \sum_{LM_1} \sum_{LM_2} \sum_{JM_1} \sum_{JM_2} \int_0^\infty d\vec{h} h^2 d\vec{M} \vec{h}
 \end{aligned}$$

$$\langle \sigma\sigma' | S_{M_3} \rangle \langle \tau\tau' | T_{M_T} \rangle \langle \vec{k} | k L M_L \rangle \langle M_S M_L | J M_J \rangle \times \\ \langle k J L S T | S \bar{S} | k J' L' S' T' \rangle \langle J' M'_S | M'_S M'_L \rangle \langle k L' M'_L | \vec{k} \rangle \times \\ \langle T' M'_T | \tau\tau' \rangle \langle S' M'_S | \sigma\sigma' \rangle \Theta(k_F^z - q) \Theta(k_F^{z'} - |\vec{q} - 2\vec{k}|)$$

Apply $\sum_{\sigma\sigma'} \langle \sigma\sigma' | S_{M_3} \rangle \langle S' M'_S | \sigma\sigma' \rangle = \delta_{SS'} \delta_{M_3 M'_S}$ ~~$\sum_{S' M'_S}$~~

Apply $\int dS_{M_3} \langle \vec{k} | k L M_L \rangle \langle k L' M'_L | \vec{k} \rangle = \frac{2}{\pi} \delta_{LL'} \delta_{M_L M'_L}$ ~~$\sum_{L' M'_L}$~~

Apply $\sum_{M_S M_L} \langle M_S M_L | J M_J \rangle \langle J' M'_S | M_S M_L \rangle = \delta_{JJ'} \delta_{M_S M'_S}$ ~~$\sum_{J' M'_S}$~~

$$= 2 \frac{2}{\pi} \sum_{\tau}, \sum_S \sum_{T_{M_T}} \sum_{T' M'_T} \sum_L \sum_{J M_J} \int_0^\infty dk k^2 \langle \tau\tau' | T_{M_T} \rangle \times$$

$$\langle k J L S T | S \bar{S} | k J' L' S' T' \rangle \langle T' M'_T | \tau\tau' \rangle \Theta(k_F^z - q) \Theta(k_F^{z'} - |\vec{q} - 2\vec{k}|)$$

G depends on $|\vec{q} - 2\vec{k}| = \sqrt{q^2 + 4k^2 - 4\vec{q} \cdot \vec{k}}$

Average over angles $\int_{-1}^1 d(\vec{q} \cdot \vec{k}) [\dots] / 2$ diagonal!

$$= \frac{2}{\pi} (2S+1) \sum_{\tau}, \sum_{T_{M_T}} \sum_{T' M'_T} \sum_{S L S} \int_0^\infty dk k^2 \int_{-1}^1 d(\vec{q} \cdot \vec{k}) \times$$

$$\langle \tau\tau' | T_{M_T} \rangle \langle k J L S T | S \bar{S} | k J' L' S' T' \rangle \langle T' M'_T | \tau\tau' \rangle \times \\ \Theta(k_F^z - q) \Theta(k_F^{z'} - |\vec{q} - 2\vec{k}|) \quad (2)$$

$$\tau = +\frac{1}{2}$$

$$\tau' = +\frac{1}{2} \rightarrow T=1 M_T=1 CGS=1$$

$$\Rightarrow ^1S_0 \rightarrow S=0, L=0, J=0$$

$$\tau' = -\frac{1}{2} \rightarrow T=1 M_T=0 CGS=\frac{1}{\sqrt{2}} \quad \langle \frac{1}{2} - \frac{1}{2} | 110 \rangle = \frac{1}{\sqrt{2}}$$

$$\Rightarrow ^1S_0 \rightarrow S=0, L=0, J=0$$

$$\rightarrow T=0 M_T=0 CGS=\frac{1}{\sqrt{2}}$$

$$\Rightarrow ^3S_1 \rightarrow S=1 L=0 J=1$$

Do we want $L > 0$?

$$\hat{n}_\lambda^p(q) \approx \frac{3}{\pi} \int_0^\infty dk k^2 \int_{-1}^1 d(\vec{q} \cdot \vec{k}) \left\{ \delta U_{^1S_0}(k, k) \Theta(k_F^p - |\vec{q} - 2\vec{k}|) \right. \\ \left. + \frac{1}{2} \Theta(k_F^p - |\vec{q} - 2\vec{k}|) \right] + \frac{1}{2} \delta U_{^3S_1 - ^3S_1}(k, k) \Theta(k_F^p - |\vec{q} - 2\vec{k}|) \\ \times \Theta(k_F^p - q) \quad (S\text{-waves only}) \quad (3)$$

Likewise

$$\hat{n}_\lambda^n(q) \approx \frac{3}{\pi} \int_0^\infty dk k^2 \int_{-1}^1 d(\vec{q} \cdot \vec{k}) \left\{ \delta U_{^1S_0}(k, k) \left[G(k_F^n - |\vec{q} - 2\vec{k}|) \right. \right. \\ \left. \left. + \frac{1}{2} G(k_F^n - |\vec{q} - 2\vec{k}|) \right] + \frac{1}{2} \delta U_{^3S_1 - ^3S_1}(k, k) \Theta(k_F^n - |\vec{q} - 2\vec{k}|) \right\} \times \\ \Theta(k_F^n - q) \quad (4)$$

$$\text{To do WPA } 4\pi \int_0^\infty dr r^2 \hat{n}_\lambda^p(q; k_F^p, k_F^n) / Z$$

\hookrightarrow Determined by δ functions!

Term 4 :

$$\frac{1}{2} \sum_{\sigma\sigma''\sigma''' \tau\tau''\tau'''} \sum_{SM_S S'M'_S Tm_T T'm'_T Lm_L L'm'_L L''m''_L L'''m'''_L} \int \frac{d^3k}{(2\pi)^3} \frac{d^3K}{(2\pi)^3} \langle k \sigma \tau \sigma'' \tau'' | \delta U | \vec{q} - \frac{1}{2}\vec{K} \sigma \tau \sigma'' \tau'' \rangle \times \\ \langle \vec{q} - \frac{1}{2}\vec{K} \sigma \tau \sigma''' \tau''' | \delta U^+ | k \sigma' \tau' \sigma'' \tau'' \rangle G(k_F^\tau - |\frac{1}{2}\vec{R} + \vec{k}|) \times \\ G(k_F^{\tau''} - |\frac{1}{2}\vec{R} - \vec{k}|)$$

δU diagonal in J, S, T, M_J, M_S, M_T

$$= \frac{1}{2} \sum_{\sigma\sigma''\sigma''' \tau\tau''\tau'''} \sum_{SM_S S'M'_S Tm_T T'm'_T Lm_L L'm'_L L''m''_L L'''m'''_L} \sum_{JM_J} \sum_{JM_S} \sum_{JM_T} \sum_{J'm'_J}$$

$$\int_0^\infty dk k^2 \int_0^\infty dK K^2 \int d\Omega_K \int d\Omega_{k_F} \langle \sigma \sigma'' | S M_S \rangle \langle \tau \tau'' | T M_T \rangle \times$$

$$\langle k | h L M_L \rangle \langle M_S M_L | J M_J \rangle \langle k J L S T | \delta U | |\vec{q} - \frac{1}{2}\vec{K}| J L' S T \rangle \times \\ \langle J M_S | M_S M_L \rangle \langle |\vec{q} - \frac{1}{2}\vec{K}| L' M'_L | \vec{q} - \frac{1}{2}\vec{K} \rangle \langle T M_T | \tau \tau''' \rangle \langle S M_S | \sigma \sigma''' \rangle \times \\ \langle \sigma \sigma''' | S' M'_S \rangle \langle \tau \tau''' | T' M'_T \rangle \langle \vec{q} - \frac{1}{2}\vec{K} | |\vec{q} - \frac{1}{2}\vec{K}| L'' m''_L \rangle \langle M'_S M''_L | J' M'_J \rangle \times \\ \langle |\vec{q} - \frac{1}{2}\vec{K}| J' L'' S' T' | \delta U^+ | k J' L''' S' T' \rangle \langle J' M'_S | M'_S M'''_L \rangle \times \\ \langle k L''' M'''_L | k \rangle \langle T M_T | \tau' \tau'' \rangle \langle S' M'_S | \sigma' \sigma'' \rangle G(k_F^\tau - |\frac{1}{2}\vec{R} + \vec{k}|) \times \\ G(k_F^{\tau''} - |\frac{1}{2}\vec{R} - \vec{k}|) \quad (5)$$

- Apply $\sum_{\sigma\sigma''} \langle \sigma \sigma'' | S M_S \rangle \langle S' M'_S | \sigma' \sigma'' \rangle = \delta_{SS'} \delta_{MM'} \cancel{\sum_{S \neq S'}}$

$$\sum_{\sigma\sigma''} \langle \sigma \sigma''' | S M_S \rangle \langle S M_S | \sigma \sigma''' \rangle = 1$$

$$- \text{Apply} \quad \int d\Omega_{\vec{h}} \langle \vec{h} | h_L M_L \rangle \langle h_L'' M_L'' | \vec{h} \rangle = \frac{2}{\pi} \delta_{LL''} \delta_{MM''}$$

$$\sum_{L'' M_L''}$$

$$= \frac{1}{2} \frac{3}{\pi} \sum_{\tau'' \tau''' T M_T} \sum_{S M_S} \sum_{L M_L} \sum_{L' M_L'} \sum_{L'' M_L''} \sum_{J M_J} \sum_{J' M_J'} \int_0^\infty dk k^2 \int_0^\infty dK K^2 \int d\Omega_{\vec{k}} \times$$

$$\langle \tau'' \tau''' | T M_T \rangle \langle M_S M_L | J M_J \rangle \langle h J L S T | S \tilde{U} | \vec{q} - \frac{1}{2} \vec{k} | J L' S T \rangle \langle J M_J | M_S M_L' \rangle \times$$

$$\langle \vec{q} - \frac{1}{2} \vec{k} | L' M_L' | \vec{q} - \frac{1}{2} \vec{k} \rangle \langle T M_T | \tau \tau''' \rangle \langle \tau \tau''' | T' M_T' \rangle \langle \vec{q} - \frac{1}{2} \vec{k} | \vec{q} - \frac{1}{2} \vec{k} | L'' M_L'' \rangle \times$$

$$\langle M_S M_L'' | J' M_J' \rangle \langle \vec{q} - \frac{1}{2} \vec{k} | J' L'' S T' | S \tilde{U}' | \vec{k} J' L' S T' \rangle \times$$

$$\langle J' M_J' | M_S M_L \rangle \langle T' M_T' | \tau \tau''' \rangle G(k_F' - |\frac{1}{2} \vec{k} + \vec{r}|) G(k_F'' - |\frac{1}{2} \vec{k} - \vec{r}|)$$

- No angular dependence on \vec{k} $\rightarrow 4\pi$

- Average over $d(\vec{q} \cdot \vec{k})$ and $d(\vec{k} \cdot \vec{k})$

$$\text{where } |\vec{q} - \frac{1}{2} \vec{k}| = \sqrt{q^2 + \frac{K^2}{4} - \vec{q} \cdot \vec{k}},$$

$$|\frac{1}{2} \vec{k} \pm \vec{k}| = \sqrt{\frac{K^2}{4} + k^2 \pm \vec{k} \cdot \vec{k}}, \quad \langle \vec{q} - \frac{1}{2} \vec{k} | L' M_L' | \vec{q} - \frac{1}{2} \vec{k} \rangle$$

$$= \sqrt{\frac{2}{\pi}} Y_{L' M_L'}^*(R_{\vec{q} \cdot \vec{k}})$$

$$= \frac{1}{2} \left(\frac{2}{\pi} \right)^2 4\pi \sum_{\tau'' \tau''' T M_T} \sum_{S M_S} \sum_{L M_L} \sum_{L' M_L'} \sum_{L'' M_L''} \sum_{J M_J} \int_0^\infty dk k^2 \int_0^\infty dK K^2 \int \frac{d\Omega_{\vec{q} \cdot \vec{k}}}{4\pi} \times$$

$$\int_{-1}^1 d(\frac{\vec{k} \cdot \vec{k}}{2}) \langle \tau'' \tau''' | T M_T \rangle \langle M_S M_L | J M_J \rangle \langle h J L S T | S \tilde{U} | \vec{q} - \frac{1}{2} \vec{k} | J L' S T \rangle \times$$

$$\langle J M_J | M_S M_L' \rangle Y_{L' M_L'}^*(R_{\vec{q} \cdot \vec{k}}) \langle T M_T | \tau \tau''' \rangle \langle \tau \tau''' | T' M_T' \rangle Y_{L'' M_L''}^*(R_{\vec{q} \cdot \vec{k}}) \times$$

$$\langle M_s M_L'' | J'M_J' \rangle \langle 1g - \frac{1}{2}\vec{k} | J'L''S\tau' | \delta\vec{v}^+ | \vec{k} J'L'S\tau' \rangle \langle J'm_s | m_{sL} \rangle \times \\ \langle \tau'm_\tau | \tau'\tau'' \rangle \Theta(k_R^{x'} - 1\frac{1}{2}\vec{k} + \vec{a}_1) \Theta(k_F^{x''} - 1\frac{1}{2}\vec{k} - \vec{a}_1) \quad (6)$$