

SRG evolution of large-cutoff chiral effective field theory interactions

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APS DNP Meeting - Waikoloa Village, Hawaii

October 24, 2018



Motivation

- Revisit problem studied by Wendt et al.¹
 - SRG evolve non-local chiral potentials at leading-order and high cutoffs
- Since then there has been an explosion of chiral potentials
- At leading-order chiral effective field theory is cutoff dependent
- Does the SRG evolve these potentials to a decoupled, cutoff independent form?

¹K.A. Wendt, R.J. Furnstahl, R.J. Perry, Phys. Rev. C **83**, 034005 (2011)

Chiral effective field theory

- Regularization procedure is necessary in χ^{EFT} to tame divergences from high energy physics
- Several options used for regulator function (local, non-local; Gaussian, super-Gaussian; coordinate-space, momentum-space)
- Typical cutoffs: $\Lambda \sim 500 \text{ MeV} (\sim 2.5 \text{ fm}^{-1})$

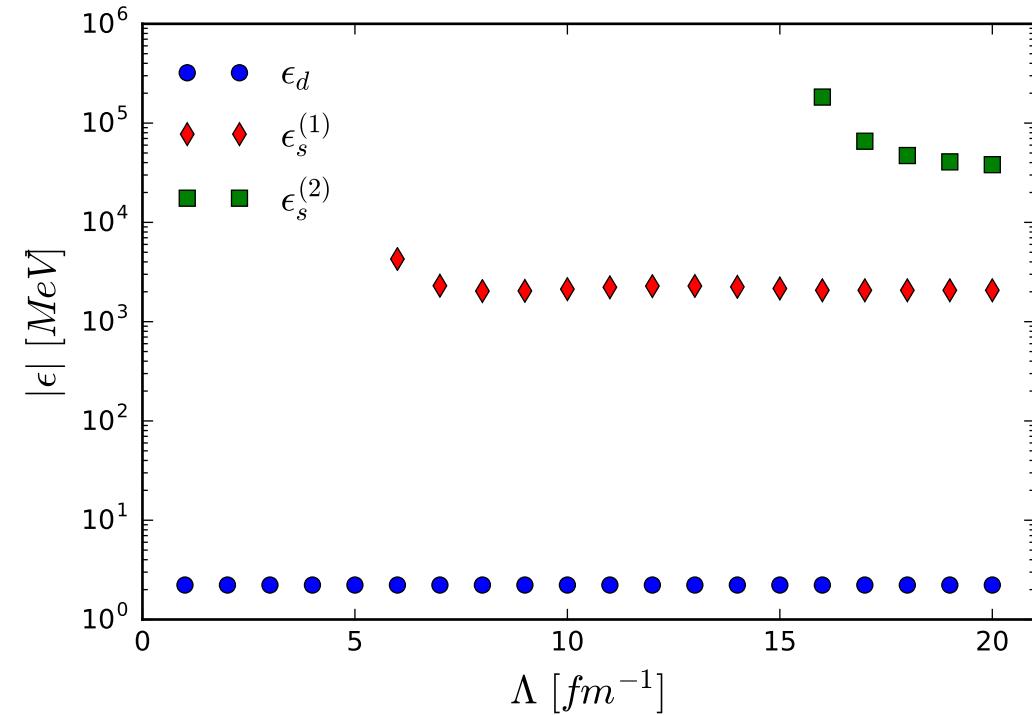
Chiral effective field theory

- Nogga et al.² took EFT cutoff Λ to much higher values to study cutoff dependence in **non-local** leading-order chiral nuclear interactions
- Tews et al.³ studied regulator and cutoff dependence for **local** potentials at LO
- **High cutoffs**
 - At short distances (high momenta), the Hamiltonian is highly singular featuring spurious deeply bound states (which are non-physical) from one-pion exchange attractive tensor forces

²A. Nogga, R.G.E. Timmermans, and U. van Kolck, Phys. Rev. C **72**, 054006 (2005), ³I. Tews, L. Huth, and A. Schwenk, Phys. Rev. C **98**, 024001 (2018)

Chiral effective field theory

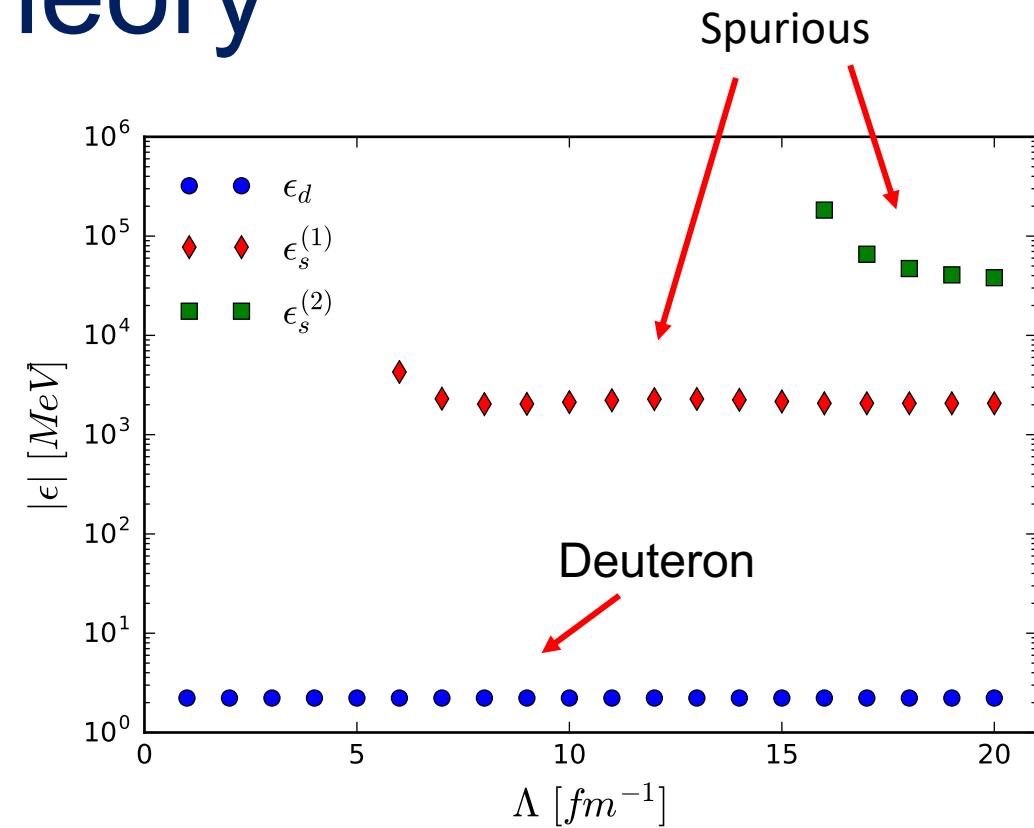
- Spurious deeply bound states
 - From EFT perspective, deeply bound states are outside of range of EFT
 - In practice, do they corrupt low-momentum physics?
- Here we consider non-local and semi-local chiral potentials at LO in the 3S_1 - 3D_1 channel



- Bound states energies of non-local potential as a function of Λ

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SRG for LO chiral EFT with high cutoffs

- Wendt et al.¹ studied **band-diagonal** SRG evolution for **non-local** interaction at large-cutoffs
- SRG decouples low- and high-momenta by ‘rotating’ Hamiltonian:

$$\tilde{H}(s) = U(s) H U^\dagger(s)$$

where $s = 0 \rightarrow \infty$

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- In practice, solve differential flow equation

$$\frac{d\tilde{H}}{ds} = [\eta(s), \tilde{H}(s)]$$

$$\tilde{H}(s) = U(s) H U^\dagger(s)$$

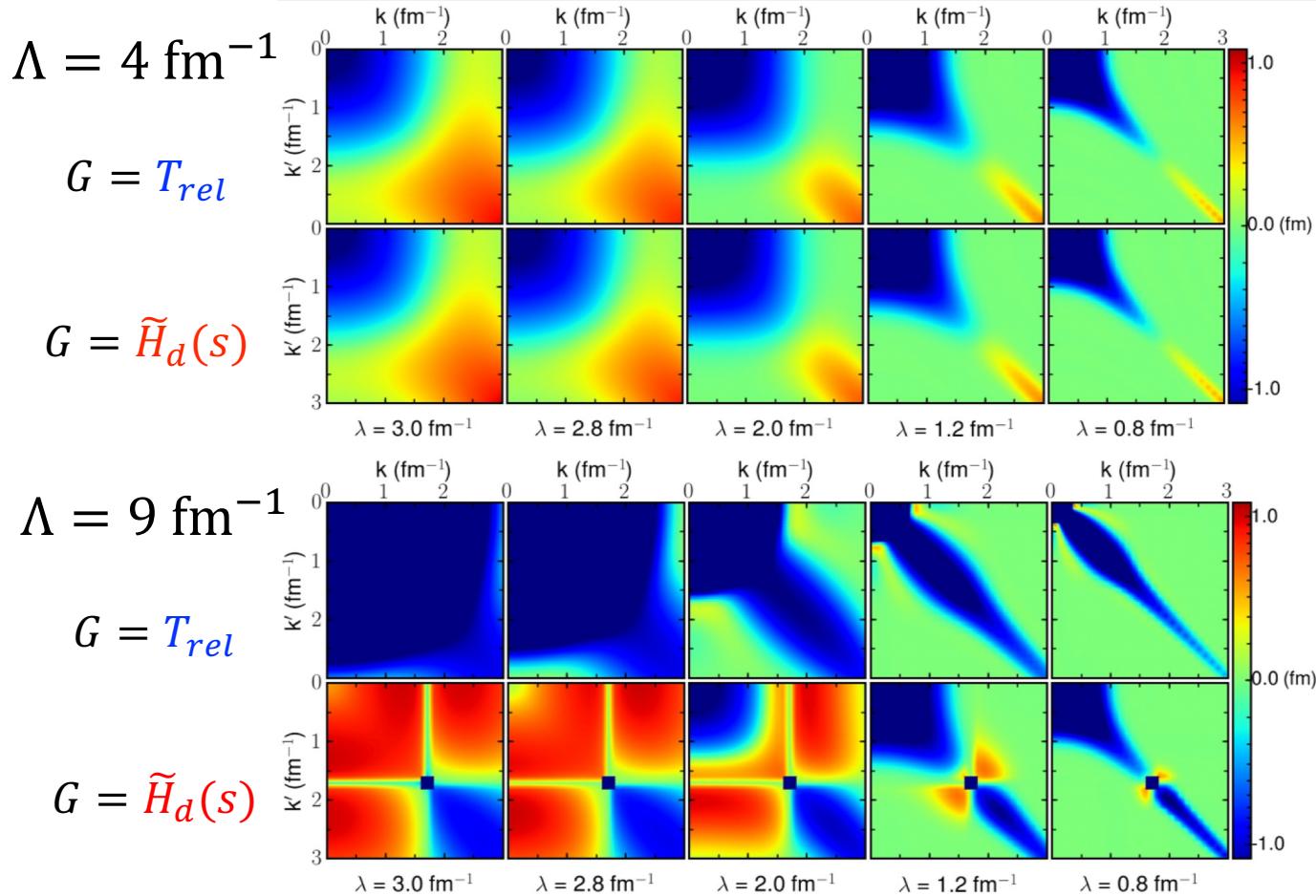
with SRG generator $\eta(s) \equiv \frac{dU}{ds} U^\dagger(s) = [G, \tilde{H}(s)]$

where $s = 0 \rightarrow \infty$

- $G = T_{rel}$ or $\tilde{H}_d(s)$ (Wegner) for band-diagonal and $G = \tilde{H}_{BD}(s)$ for block-diagonal decoupling

SRG for LO chiral EFT with high cutoffs

- For typical cutoffs, same $\tilde{H}(s)$ regardless of $G = T_{rel}$ or $\tilde{H}_d(s)$!
- For high cutoff $\Lambda = 9 \text{ fm}^{-1}$
 - **Wegner**: keeps spurious bound state safely outside low-momentum block
 - T_{rel} : spurious bound state distorts the low-momentum block of $\tilde{H}(s)$
 - Choice in G matters at high cutoffs!



SRG evolution of $V_\lambda(k, k')$ for several values of $\lambda = s^{-1/4}$

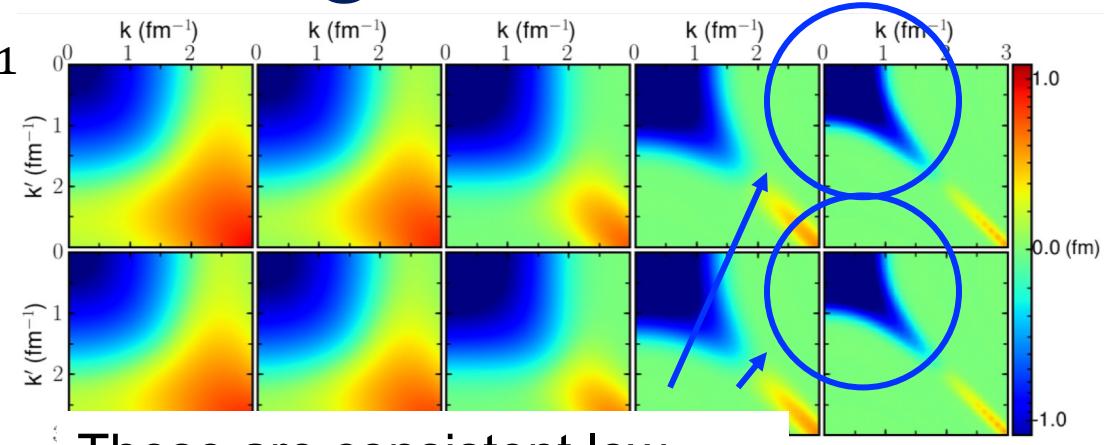
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$$\Lambda = 4 \text{ fm}^{-1}$$

$$G = T_{rel}$$

$$G = \tilde{H}_d(s)$$

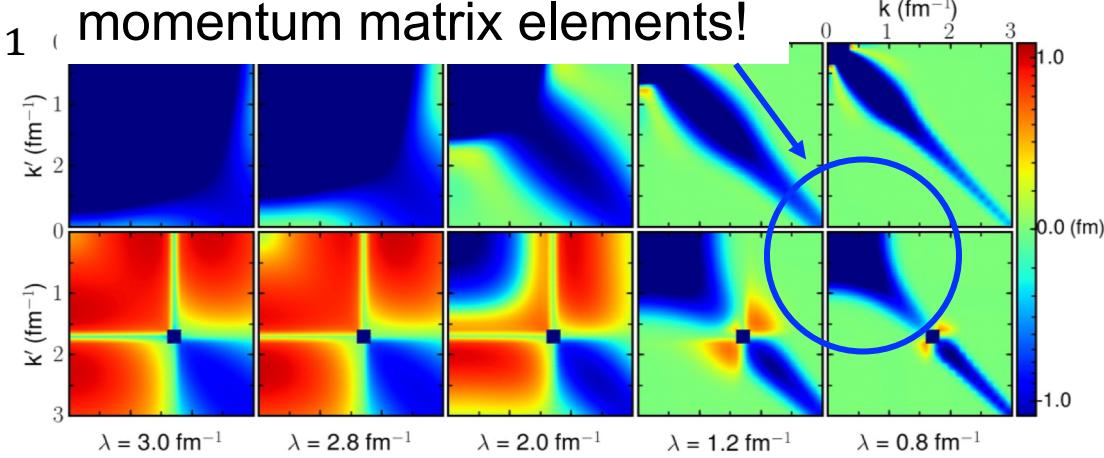


These are consistent low-momentum matrix elements!

$$\Lambda = 9 \text{ fm}^{-1}$$

$$G = T_{rel}$$

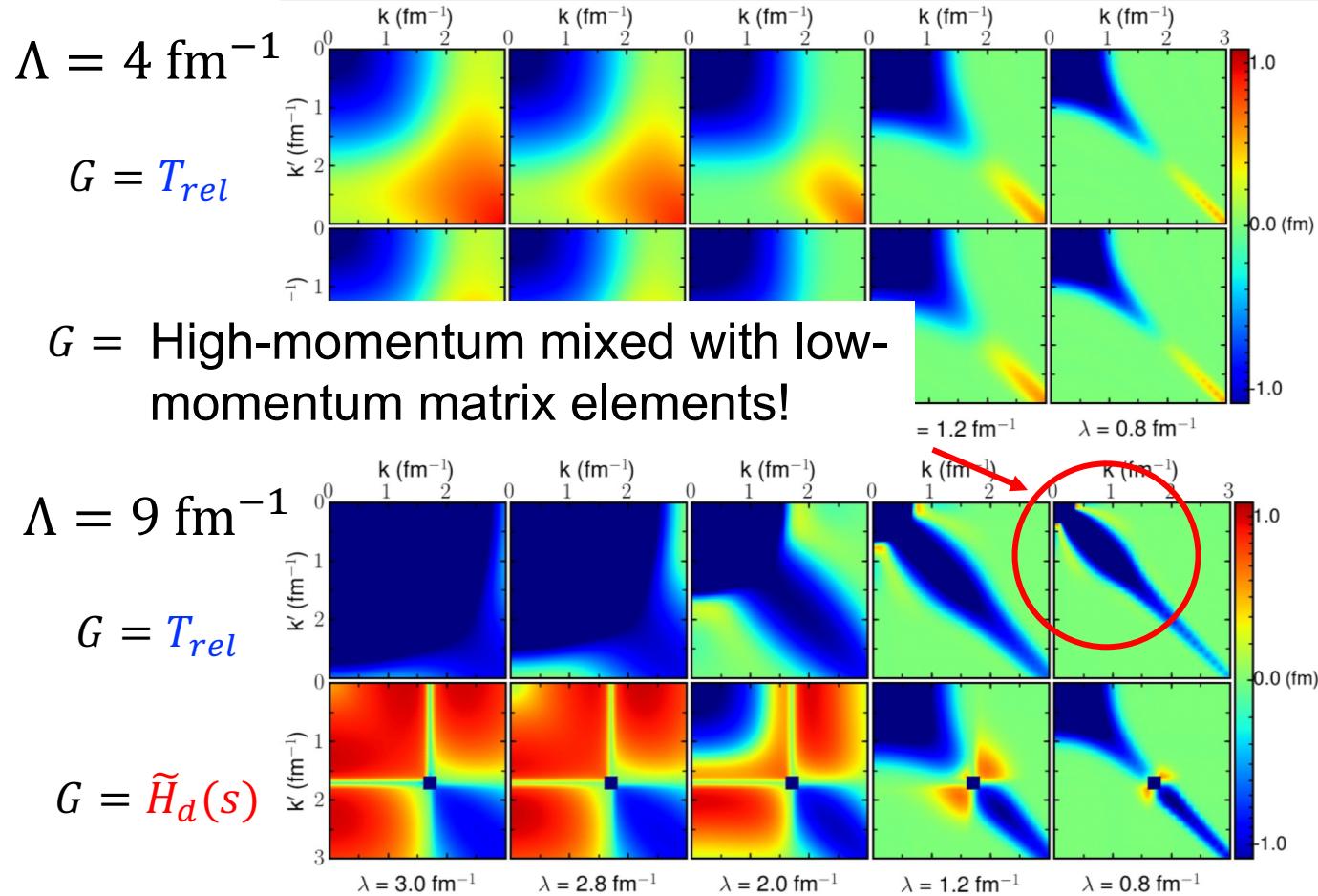
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Motivation

- High cutoff potentials were used as a test case for the Magnus expansion (analogous problem with IMSRG intruder states)
 - Magnus agreed with SRG result but struggled to evolve very high cutoffs ($\Lambda \sim 20 \text{ fm}^{-1}$)
- In the process, this led to several questions

Motivation

- How do different regulator functions (non-local, semi-local, local) affect SRG evolution?
- Do matrix elements of potential approach universal form at lower λ ?
- How does the spurious bound state decouple for Wegner generators?
- What happens to spurious bound state(s) with block-diagonal SRG evolution?
- Can we gain insight into renormalizability for chiral EFT formulations?

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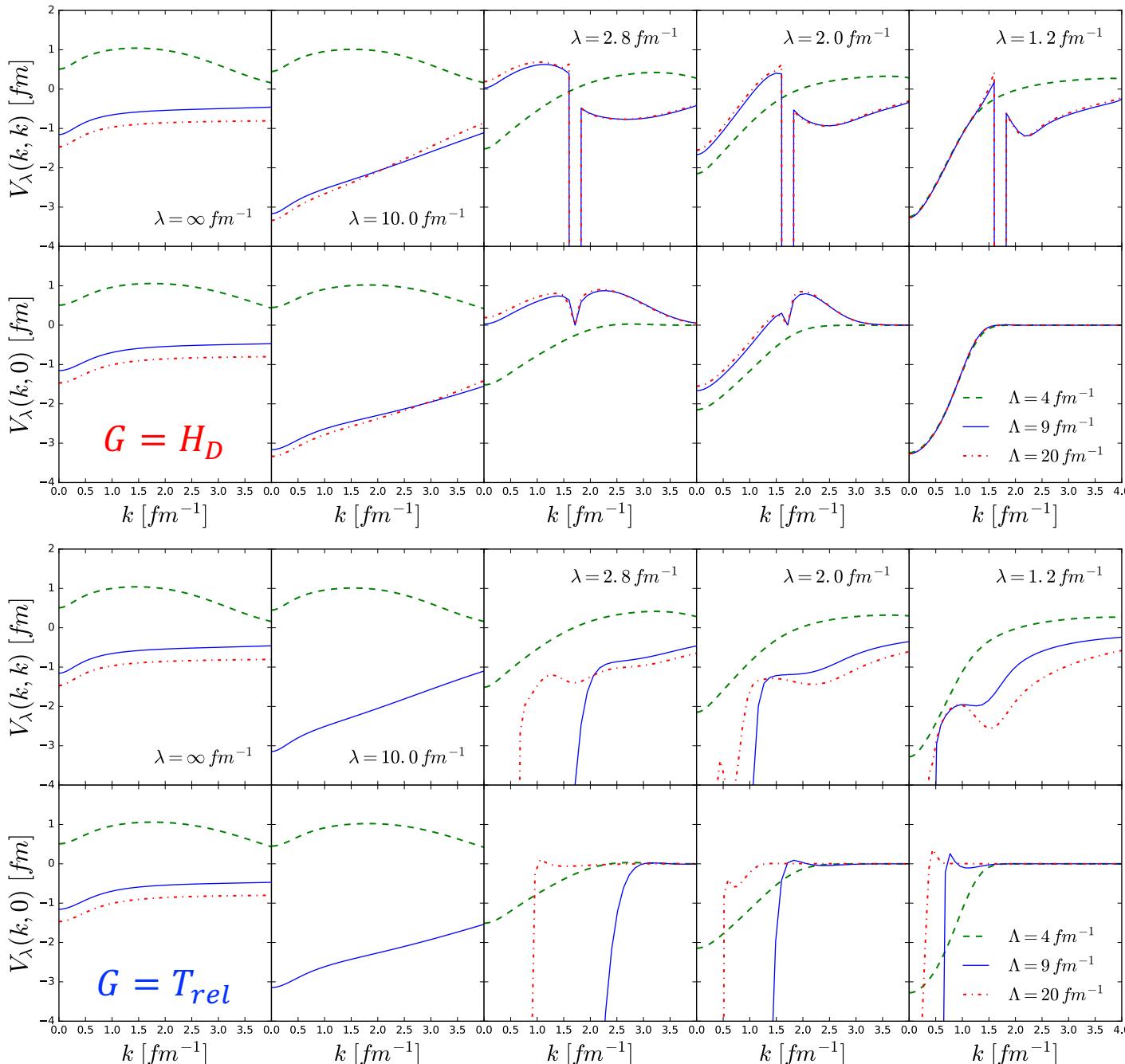
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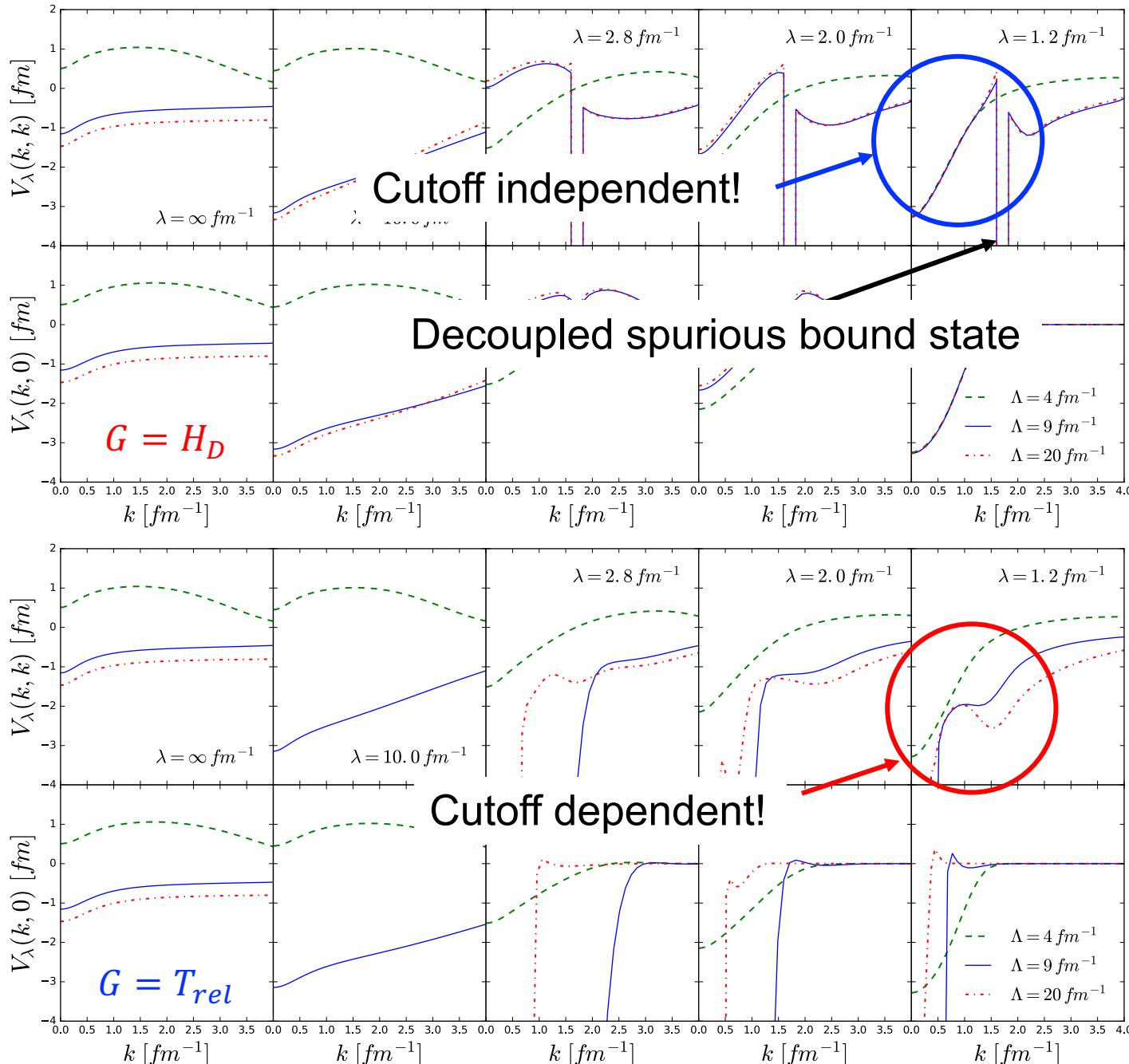
Band-diagonal evolution

- Non-local regulator
 - Matrix elements of $V_\lambda(k, k')$ are independent of Λ at low-momentum for Wegner evolution
 - Wegner generator decouples the spurious bound state at $k \sim 1.83 \text{ fm}^{-1}$
- But is mesh dependent!



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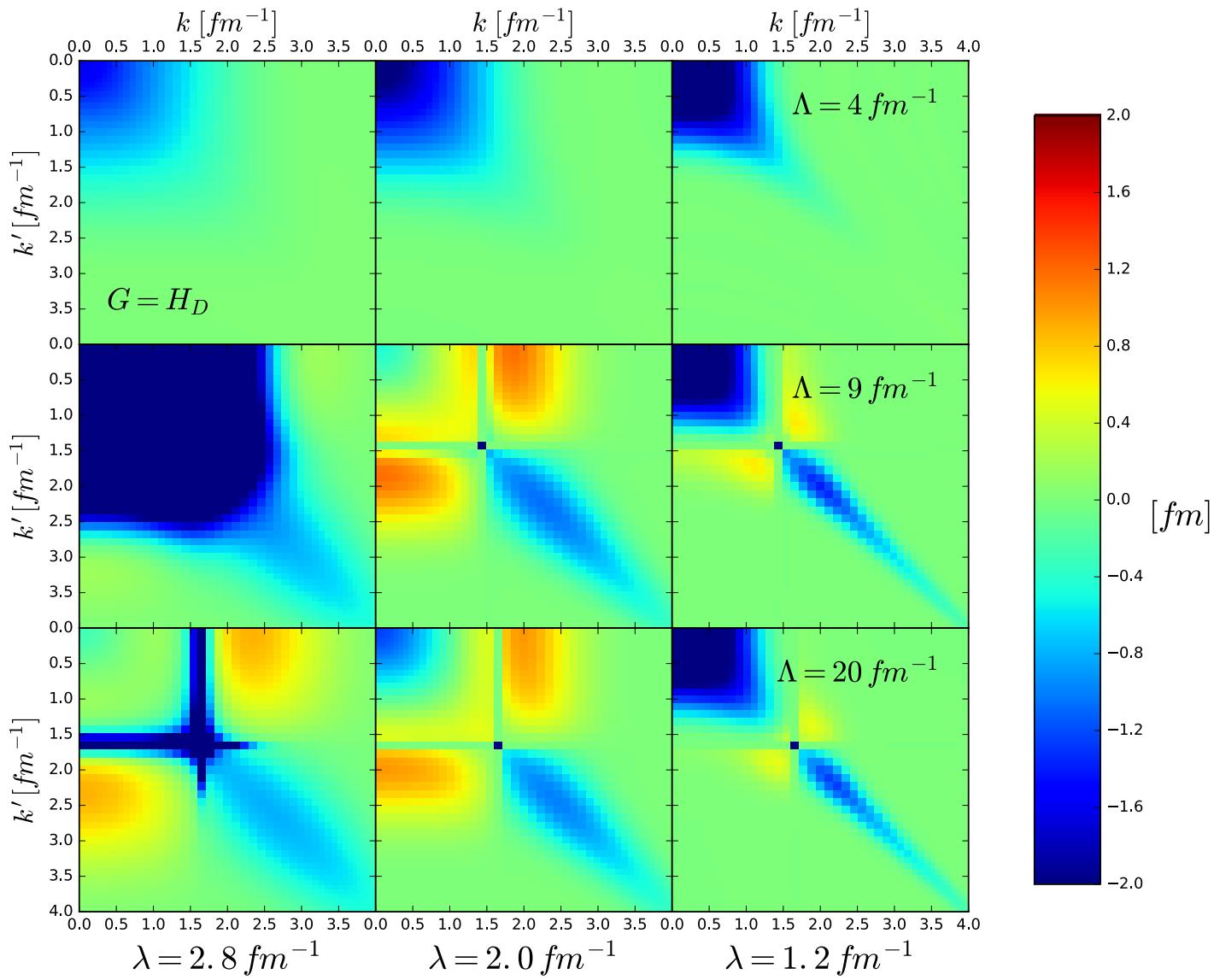
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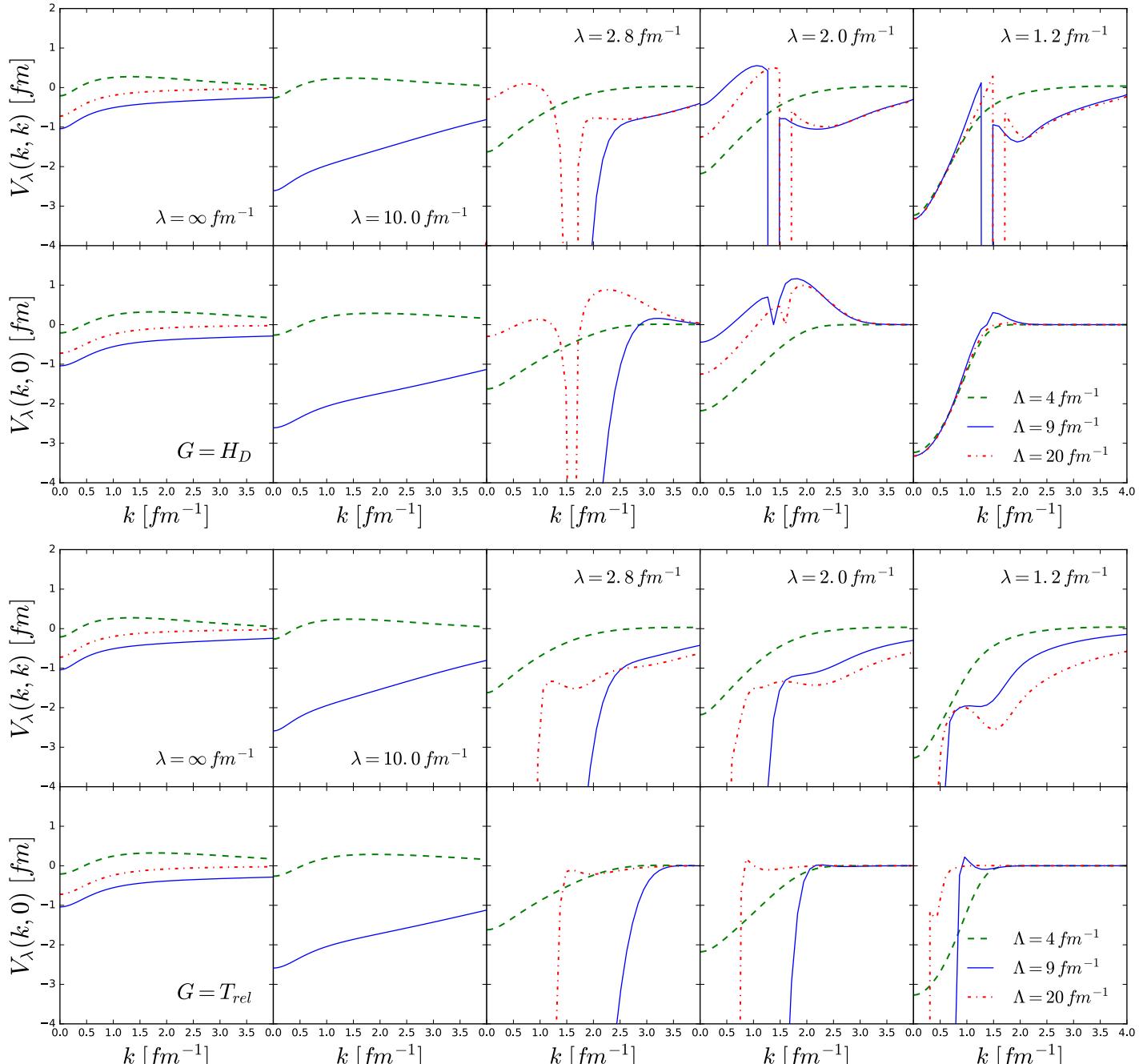
- **Semi-local regulator**

- Local one-pion exchange and non-local contact interaction
- Similar to non-local case:
Wegner generator decouples the spurious bound state at $k \sim 1.49$ and 1.71 fm^{-1} for $\Lambda = 9$ and 20 fm^{-1} respectively



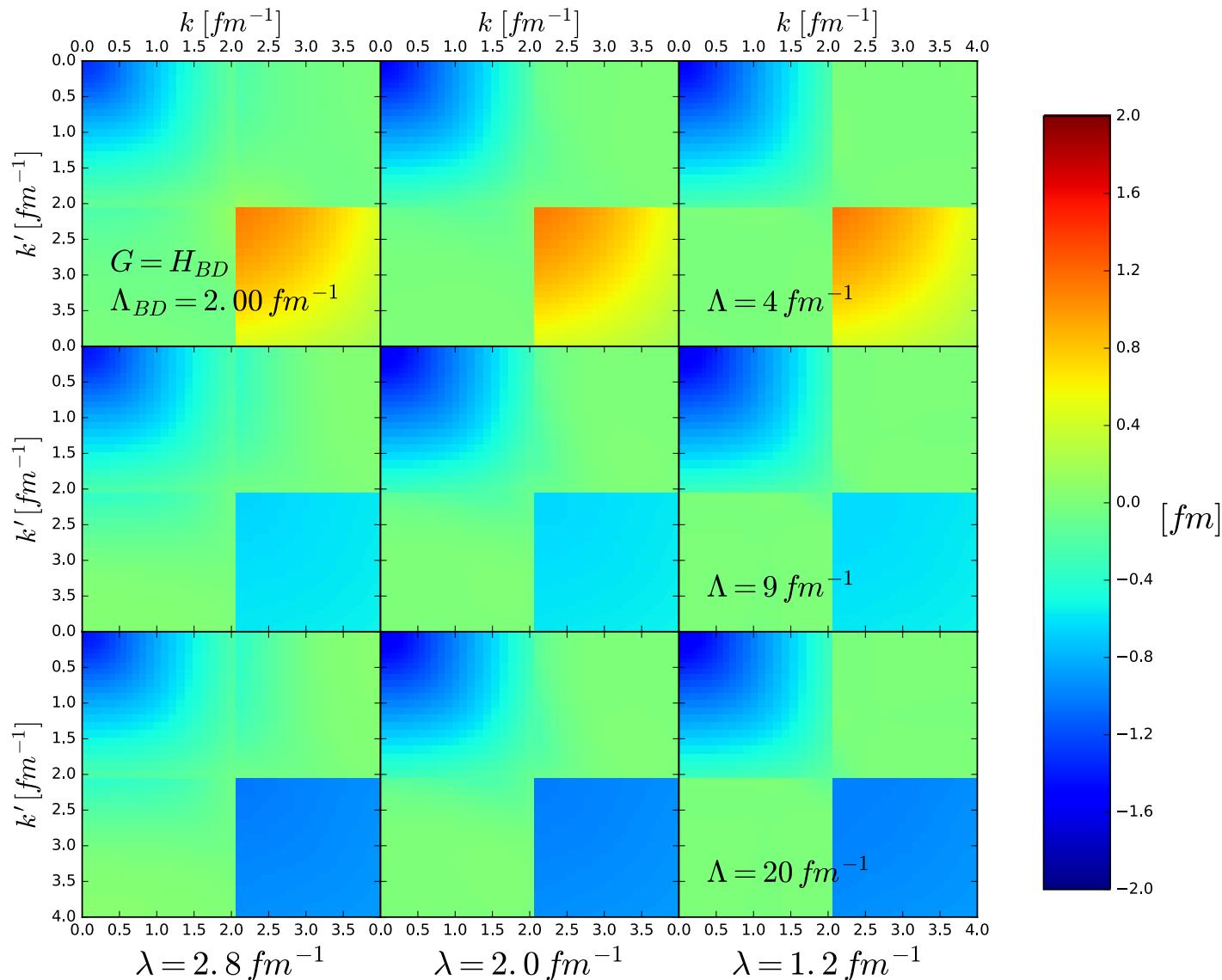
Band-diagonal evolution

- Semi-local regulator
 - Worse matrix element equivalence than non-local case (ε_s decoupled at lower k) for Wegner
 - T_{rel} behavior is the same as before
 - For band-diagonal generators, decoupling of ε_s depends on momentum mesh and regulator
 - Is decoupling cleaner for block-diagonal evolution?



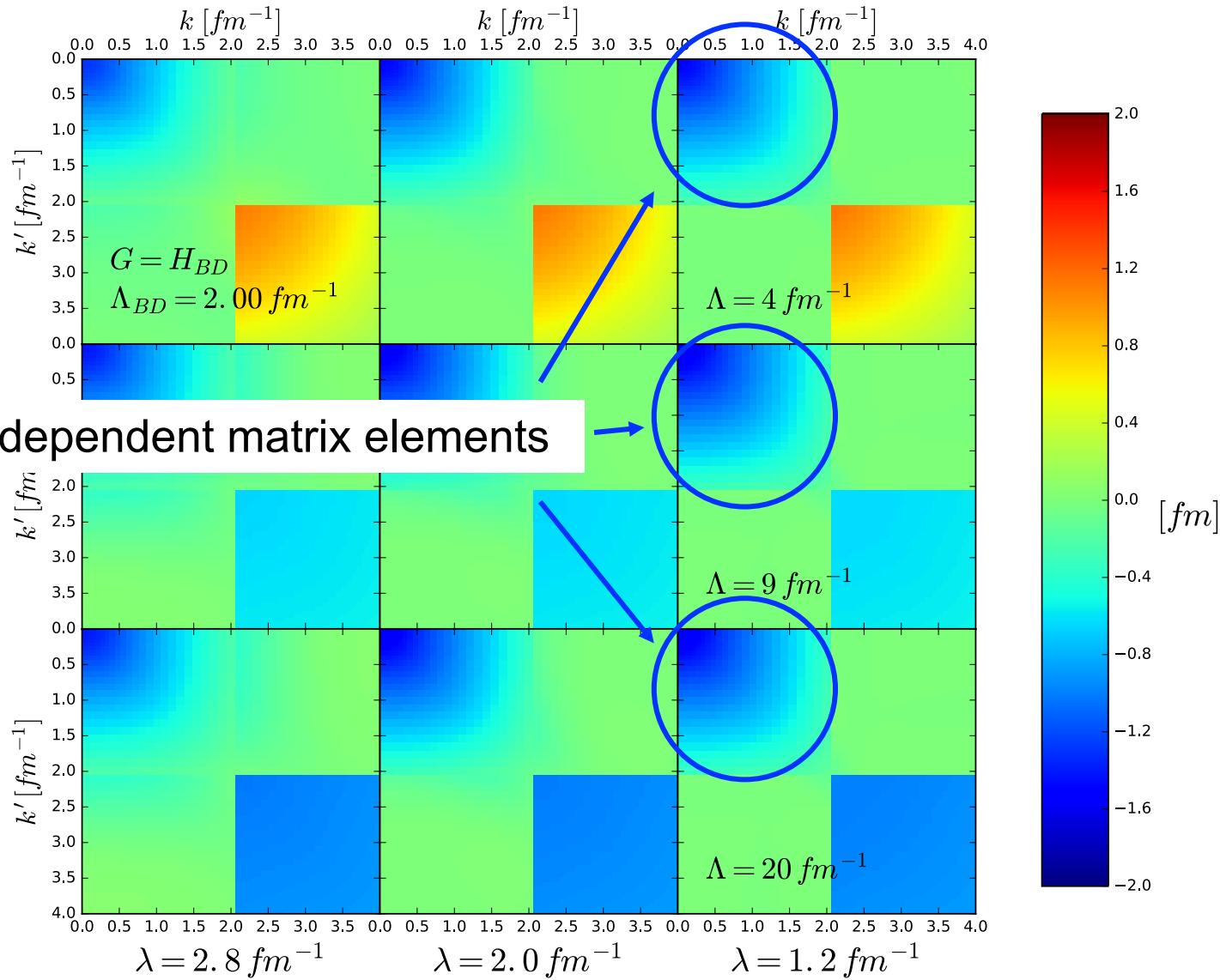
Block-diagonal evolution

- $G = \begin{pmatrix} PH(s)P & 0 \\ 0 & QH(s)Q \end{pmatrix}$ with projection operator P and $Q = 1 - P$
- Decouples Hamiltonian into sub-blocks at momentum Λ_{BD}
- Diagonalize sub-blocks to see where ε_s decouples
 - ε_s decouples at $k \sim 4.47 \text{ fm}^{-1}$ much higher than $G = \tilde{H}_d$ where $k \sim 1.83 \text{ fm}^{-1}$



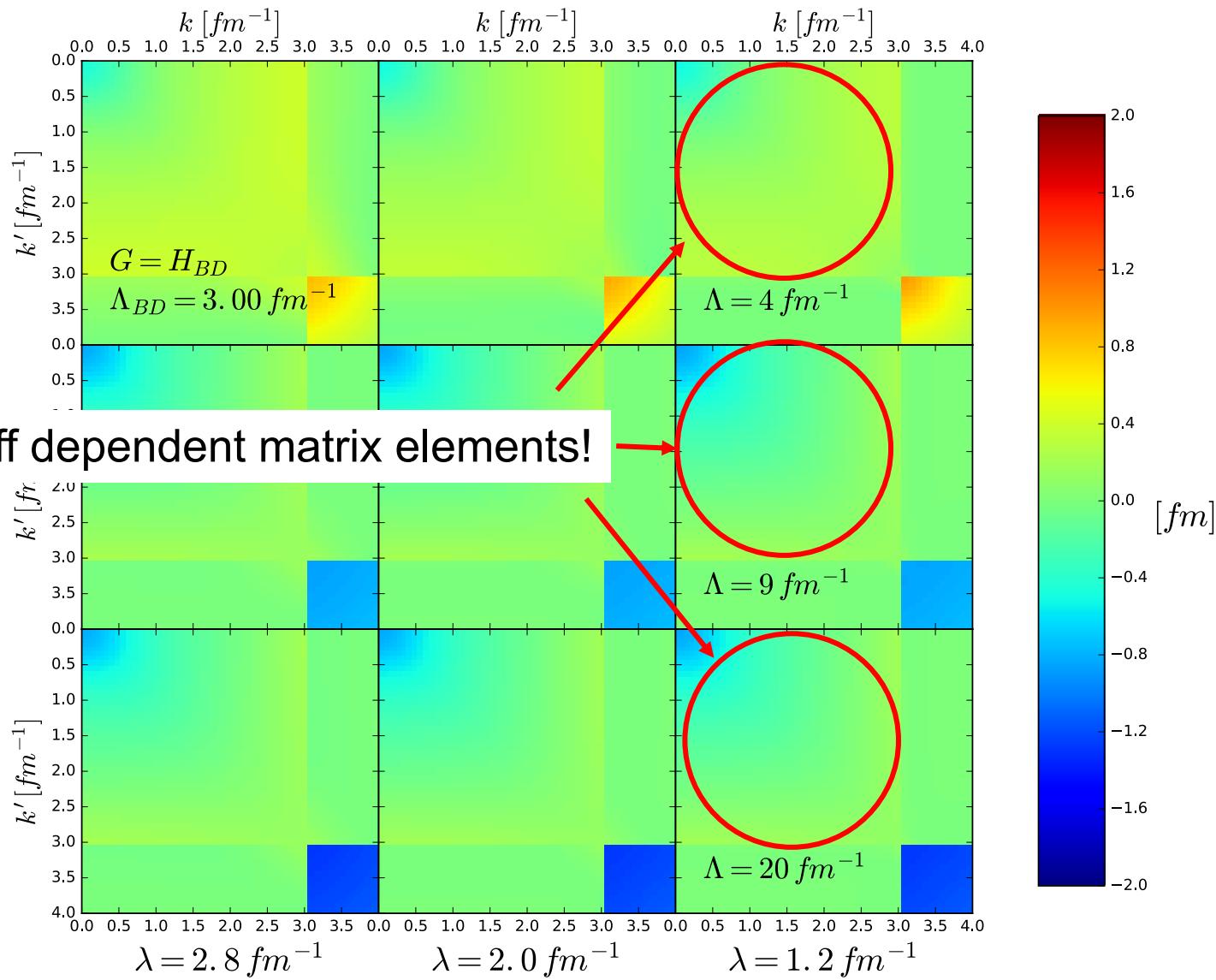
Block-diagonal evolution

- For $\Lambda_{BD} = 2 \text{ fm}^{-1}$ low-momentum matrix elements of evolved potential are independent of EFT Λ



Block-diagonal evolution

- For $\Lambda_{BD} = 3 \text{ fm}^{-1}$ low-momentum matrix elements are not equivalent
 - Low-momenta couples to modes affected by spurious bound state!



Conclusion

- **Summary**
 - SRG-evolved potentials flow to universal low-momentum matrix elements when ε_s decouples from low-momenta (band-diagonal Wegner or block-diagonal with low Λ_{BD})
 - Non-local and semi-local potentials have similar SRG evolution but ε_s decouples differently
 - Block-diagonal generators decouples ε_s at higher momentum than band-diagonal generators

Conclusion

- Future work
 - SRG evolve local chiral potential and other channels
 - Determine what controls where spurious bound state decouples
 - Use smooth cutoff for block-diagonal generator (cf. sharp cutoff)
 - Explore issues of renormalizability in χ^{EFT}