

(3/12/21)

- Starting point is  $\hat{n}_\lambda^z(\vec{q})$  after evaluation of contractions. ( $\vec{q}$  is not a relative momentum here!)

$$\begin{aligned}
 \hat{n}_\lambda^z(\vec{q}) = & \sqrt{\left\langle \sum_{\sigma} \Theta(k_f^z - q) \right\rangle} \quad \text{Term 1} \\
 & + \sum_{\sigma\sigma'z'} \int \frac{d^3 k}{(2\pi)^3} \left[ \underbrace{\left\langle k_{\sigma z} \sigma z | \delta U | k_{\sigma' z'} \sigma' z' \right\rangle}_{\text{Term 2}} \right. \\
 & \left. + \left\langle k_{\sigma z} \sigma z | \delta U + k_{\sigma z} \sigma z \right\rangle \right] \Theta(k_f^z - q) \Theta(k_f^{z'} - |\vec{q} - \vec{k}|) \\
 & + \frac{1}{2} \sum_{\sigma\sigma''\sigma''' z'' z'''} \int \frac{d^3 k}{(2\pi)^3} \frac{d^3 K}{(2\pi)^3} \left\langle k_{\sigma z} \sigma z | \delta U | \vec{q} - \frac{1}{2} \vec{k} \sigma z'' z''' \right\rangle \times \\
 & \left. \left( \vec{q} - \frac{1}{2} \vec{k} \sigma z'' z''' | \delta U + \vec{k} \sigma z'' z''' \right) \Theta(k_f^{z'} - |\frac{1}{2} \vec{k} + \vec{k}|) \Theta(k_f^{z''} - |\frac{1}{2} \vec{k} - \vec{k}|) \right\rangle \quad (1)
 \end{aligned}$$

Term 1 is simple. Contribution is  $2 \Theta(k_f^z - q)$ .

No 2nd + 3rd term (multiply 2nd term by 2 for full 2+3 contribution).

$$\begin{aligned}
 & 2 \sum_{\sigma\sigma'z} \int \frac{d^3 k}{(2\pi)^3} \left\langle k_{\sigma z} \sigma z | \delta U | k_{\sigma' z'} \sigma' z' \right\rangle \Theta(k_f^z - q) \Theta(k_f^{z'} - |\vec{q} - \vec{k}|) \\
 & = 2 \sum_{\sigma\sigma'z} \sum_{SM_1} \sum_{SM_2} \sum_{TM_1} \sum_{TM_2} \sum_{LM_1} \sum_{LM_2} \sum_{JM_1} \sum_{JM_2} \int_0^\infty d\vec{h} h^2 d\vec{M} \vec{h}
 \end{aligned}$$

$$\langle \sigma\sigma' | S_{M_3} \rangle \langle \tau\tau' | T_{M_T} \rangle \langle \vec{h} | h L M_L \rangle \langle M_S M_L | J M_J \rangle \times \\ \langle k J L S T | S \vec{U} | k J' L' S' T' \rangle \langle J' M'_S | M'_S M'_L \rangle \langle k L' M'_L | \vec{h} \rangle \times \\ \langle T' M'_T | \tau\tau' \rangle \langle S M'_S | \sigma\sigma' \rangle \Theta(k_F^z - q) \Theta(k_F^{z'} - |\vec{q} - 2\vec{h}|)$$

Apply  $\sum_{\sigma\sigma'} \langle \sigma\sigma' | S_{M_3} \rangle \langle S' M'_S | \sigma\sigma' \rangle = \delta_{SS'} \delta_{M_3 M'_S}$   ~~$\sum_{S M'_S}$~~

Apply  $\int dS_L \langle \vec{h} | h L M_L \rangle \langle k L' M'_L | \vec{h} \rangle = \frac{2}{\pi} \delta_{LL'} \delta_{M_L M'_L}$   ~~$\sum_{L' M'_L}$~~

Apply  $\sum_{M_S M_L} \langle M_S M_L | J M_J \rangle \langle J' M'_S | M_S M_L \rangle = \delta_{JJ'} \delta_{M_S M'_S}$   ~~$\sum_{J' M'_S}$~~

$$= 2 \frac{2}{\pi} \sum_{\tau}, \sum_S \sum_{T_{M_T}} \sum_{T' M'_T} \sum_L \sum_{J M_J} \int_0^\infty dk h^2 \langle \tau\tau' | T_{M_T} \rangle \times$$

$$\langle k J L S T | S \vec{U} | k J' L' S' T' \rangle \langle T' M'_T | \tau\tau' \rangle \Theta(k_F^z - q) \Theta(k_F^{z'} - |\vec{q} - 2\vec{h}|)$$

$G$  depends on  $|\vec{q} - 2\vec{h}| = \sqrt{q^2 + 4h^2 - 4\vec{q} \cdot \vec{h}}$

Average over angles  $\int_{-1}^1 d(\vec{q} \cdot \vec{h}) [\dots] / 2$  diagonal!

$$= \frac{2}{\pi} (2S+1) \underbrace{\sum_{\tau}, \sum_{T_{M_T}} \sum_{T' M'_T} \sum_{S L S'}}_{\text{SL}} \int_0^\infty dk h^2 \int_{-1}^1 d(\vec{q} \cdot \vec{h}) \times$$

$$\langle \tau\tau' | T_{M_T} \rangle \langle k J L S T | S \vec{U} | k J' L' S' T' \rangle \langle T' M'_T | \tau\tau' \rangle \times \\ \Theta(k_F^z - q) \Theta(k_F^{z'} - |\vec{q} - 2\vec{h}|) \quad (2)$$

$$\tau = +\frac{1}{2}$$

$$\tau' = +\frac{1}{2} \rightarrow T=1 M_T=1 CGS=1$$

$$\Rightarrow ^1S_0 \rightarrow S=0, L=0, J=0$$

$$\tau' = -\frac{1}{2} \rightarrow T=1 M_T=0 CGS=\frac{1}{\sqrt{2}} \quad \langle \frac{1}{2} - \frac{1}{2} | 110 \rangle = \frac{1}{\sqrt{2}}$$

$$\Rightarrow ^1S_0 \rightarrow S=0, L=0, J=0$$

$$\rightarrow T=0 M_T=0 CGS=\frac{1}{\sqrt{2}}$$

$$\Rightarrow ^3S_1 \rightarrow S=1 L=0 J=1$$

Do we want  $L > 0$ ?

$$\hat{n}_\lambda^p(q) \approx \frac{3}{\pi} \int_0^\infty dk k^2 \int_{-1}^1 d(\vec{q} \cdot \vec{k}) \left\{ \delta U_{1S_0}(k, k) \Theta(k_F^p - |\vec{q} - 2\vec{k}|) \right. \\ \left. + \frac{1}{2} \Theta(k_F^p - |\vec{q} - 2\vec{k}|) \right] + \frac{3}{2} \delta U_{3S_1-3S_1}(k, k) \Theta(k_F^p - |\vec{q} - 2\vec{k}|) \\ \times \Theta(k_F^p - q) \quad (S\text{-waves only}) \quad (3)$$

Likewise

$$\hat{n}_\lambda^n(q) \approx \frac{3}{\pi} \int_0^\infty dk k^2 \int_{-1}^1 d(\vec{q} \cdot \vec{k}) \left\{ \delta U_{1S_0}(k, k) \left[ G(k_F^n - |\vec{q} - 2\vec{k}|) \right. \right. \\ \left. \left. + \frac{1}{2} G(k_F^n - |\vec{q} - 2\vec{k}|) \right] + \frac{3}{2} \delta U_{3S_1-3S_1}(k, k) \Theta(k_F^n - |\vec{q} - 2\vec{k}|) \right\} \times \\ \Theta(k_F^n - q) \quad (n) \quad (4)$$

$$\text{To do WPA } 4\pi \int_0^\infty dr r^2 \hat{n}_\lambda^p(q; k_F^p, k_F^n) / Z$$

$\hookrightarrow$  Determined by  $\delta$  functions!

Term 4 :

$$\frac{1}{2} \sum_{\sigma\sigma''\sigma''' \tau\tau''\tau'''} \sum_{\substack{\sigma\sigma''\sigma''' \\ \tau\tau''\tau'''}} \int \frac{d^3k}{(2\pi)^3} \frac{d^3K}{(2\pi)^3} \langle \hat{k} | \sigma'' \tau'' \sigma''' \tau''' | \delta U | \vec{q} - \frac{1}{2}\vec{K} \sigma'' \tau''' \rangle \times$$

$$\langle \vec{q} - \frac{1}{2}\vec{K} \sigma'' \tau''' | \delta J^+ | \hat{k} | \sigma'' \tau'' \sigma''' \rangle G(k_F^2 - |\frac{1}{2}\vec{K} + \vec{k}|) \times \\ G(k_F^2 - |\frac{1}{2}\vec{K} - \vec{k}|)$$

$$= \frac{1}{2} \sum_{\sigma\sigma''\sigma''' \tau\tau''\tau'''} \sum_{\substack{\sigma\sigma''\sigma''' \\ \tau\tau''\tau'''}} \sum_{S M_S \dots S'' M_S'' T M_T \dots T'' M_T'' L M_L \dots L'' M_L'' I M_I \dots I'' M_I''} \times$$

$$\int_0^\infty dk k^2 \int_0^\infty dK K^2 \int dR_K \int dR_{\vec{k}} \langle \sigma'' \sigma''' | S M_S \rangle \langle \tau'' \tau''' | T M_T \rangle \times$$

$$\langle \hat{k} | k L M_L \rangle \langle M_S M_L | I M_T \rangle \langle k J L S T | \delta U | |\vec{q} - \frac{1}{2}\vec{K}| J' L' S' T' \rangle \times \\ \langle J' M_S' | M_S' M_L' \rangle \langle |\vec{q} - \frac{1}{2}\vec{K}| L' M_L' | \vec{q} - \frac{1}{2}\vec{K} \rangle \langle T' M_T' | \tau'' \tau''' \rangle \langle S' M_S' | \sigma'' \sigma''' \rangle \times \\ \langle \sigma'' \sigma''' | S'' M_S'' \rangle \langle \tau'' \tau''' | T' M_T' \rangle \langle \vec{q} - \frac{1}{2}\vec{K} | |\vec{q} - \frac{1}{2}\vec{K}| L'' M_L'' \rangle \langle M_S'' M_L'' | I'' M_I'' \rangle \times \\ \langle |\vec{q} - \frac{1}{2}\vec{K}| J'' L'' S'' T'' | \delta U^+ | \hat{k} | J'' L'' S'' T'' \rangle \langle I'' M_I'' | M_S'' M_L'' \rangle \times \\ \langle k L'' M_L'' | \hat{k} \rangle \langle T' M_T' | \tau'' \tau''' \rangle \langle S'' M_S'' | \sigma'' \sigma''' \rangle G(k_F^2 - |\frac{1}{2}\vec{K} + \vec{k}|) \times \\ G(k_F^2 - |\frac{1}{2}\vec{K} - \vec{k}|) \quad (5)$$

- Apply  $\sum_{\sigma\sigma''} \langle \sigma'' \sigma''' | S M_S \rangle \langle S'' M_S'' | \sigma'' \sigma''' \rangle = \delta_{S''S} \delta_{M_S M_S''}$

$$\sum_{\sigma\sigma''} \langle \sigma'' \sigma''' | S' M_S' \rangle \langle S'' M_S'' | \sigma'' \sigma''' \rangle = \delta_{S'S} \delta_{M_S' M_S''}$$

- Apply  $\int dR_{\vec{k}} \langle \hat{k} | k L M_L \rangle \langle k L'' M_L'' | \hat{k} \rangle = \frac{2}{\pi} \delta_{L''L} \delta_{M_L M_L''}$

- Approximate  $|q - \frac{1}{2}\vec{K}| \approx q^2 + \frac{K^2}{4}$

- Then  $\int \frac{dR_{\vec{q} \cdot \vec{K}}}{4\pi} \langle |q - \frac{1}{2}\vec{K}| L'M'_L | \vec{q} - \vec{K} \rangle \langle \vec{q} - \frac{1}{2}\vec{K} | L''M''_L \rangle$   
 $= \frac{1}{4\pi} \frac{2}{\pi} \delta_{L'L''} \delta_{M'_L M''_L}$

- Apply  $\sum_{M_S M_L} \langle M_S M_L | JM_J \rangle \langle J''M''_J | M_S M_L \rangle = \delta_{J''J} \delta_{M_S M''_J}$

$$\sum_{M'_S M'_L} \langle J'M'_J | M'_S M'_L \rangle \langle M'_S M'_L | J''M''_J \rangle = \delta_{J''J} \delta_{M'_S M''_J}$$

$$= \frac{1}{2} \frac{1}{4\pi} \left(\frac{3}{\pi}\right)^2 \sum_{T''T'''} \sum_{TM_T \dots T''M''_T} \sum_{SS'} \sum_{LL'} \sum_{JM_J} \sum_{J''M''_J} \int_0^\infty dk k^2 \int_0^\infty dR_{\vec{K}}$$

$$\begin{aligned} & \langle \varepsilon' \varepsilon'' | TM_T \rangle \langle k J L S T | S \bar{O} | q - \frac{1}{2}\vec{K} | J' L' S' T' \rangle \langle T' M' | \varepsilon \varepsilon'' \rangle \\ & \times \langle \varepsilon \varepsilon''' | T'' M''_T \rangle \langle q - \frac{1}{2}\vec{K} | J' L' S' T'' | S \bar{O}^+ | k J L S T'' \rangle \langle T'' M''_T | \varepsilon \varepsilon'' \rangle \\ & \times \Theta(k_p^{T'} - |\frac{1}{2}\vec{K} + \vec{\omega}|) \Theta(k_F^{T''} - |\frac{1}{2}\vec{K} - \vec{\omega}|) \end{aligned}$$

- Now apply  $S\bar{O}$ ,  $S\bar{O}^+$  bring diagonal in  
 $J M_J \quad S M_S \quad T M_T$

- Average  $\int_{-1}^1 \frac{d(\hat{K} \cdot \vec{k})}{2}$

-  $\int dR_{\vec{K}} = 4\pi \quad - \sum_{M_J} \rightarrow (2J+1)$

$$= \frac{1}{2} \frac{1}{4\pi} \cancel{4\pi} \left(\frac{2}{\pi}\right)^2 (2J+1) \sum_{\tau'' \tau''' \tau''''} \sum_{TM_T T'M'_T} \sum_S \sum_{LL'} \sum_J \times$$

$$\int_0^\infty dk k^2 \int_0^\infty dK K^2 \int_{-1}^1 \frac{d(\vec{k} \cdot \vec{k}')}{2} \langle \tau \tau'' | TM_T \rangle \times$$

$$\langle k_J L S T | \delta \vec{r} | 1 \vec{r} - \frac{1}{2} \vec{k} | J L' S T' \rangle \langle T M_T | \tau \tau'' \rangle \langle \tau \tau''' | T' M'_T \rangle \times$$

$$\langle 1 \vec{r} - \frac{1}{2} \vec{k} | J L' S T' | \delta \vec{r}' | k_J L S T' \rangle \langle T M'_T | \tau \tau'' \rangle \times$$

$$\Theta(k_F^{z'} - (\frac{1}{2} \vec{k} + \vec{k}') \cdot \vec{k}) \Theta(k_F^{z''} - \frac{1}{2} \vec{k} - \vec{k}') \quad (6)$$

Step through cases:

$$S = 0 \Rightarrow T = 1 \text{ and } T' = 1$$

$$S = 1 \Rightarrow T = 0 \text{ and } T' = 0$$

} So no point  
in summing over  
 $T M'_T$

$$\text{Fix } \tau = +\frac{1}{2}$$

$$\tau''' = +\frac{1}{2} \Rightarrow T = 1 \quad M_T = 1 \Rightarrow {}^1S_0, CGS = 1$$

$$\rightarrow \tau' = +\frac{1}{2} \quad \tau'' = +\frac{1}{2} \quad CGS = 1$$

$$\tau''' = -\frac{1}{2} \Rightarrow T = 1 \quad M_T = 0 \Rightarrow {}^1S_0, CGS = \frac{1}{\sqrt{2}}$$

$$\tau' = \pm \frac{1}{2} \quad \tau'' = \mp \frac{1}{2} \Rightarrow CGS = \pm \frac{1}{\sqrt{2}} \quad \Rightarrow T = 0 \quad M_T = 0 \Rightarrow {}^3S_1, CGS = \frac{1}{\sqrt{2}}$$

$$\tau' = \pm \frac{1}{2} \quad \tau'' = \mp \frac{1}{2} \Rightarrow CGS = \frac{1}{\sqrt{2}}$$

No schematic of terms for  $\tau = +\frac{1}{2}, L = 0$

$$\begin{aligned}
& \langle +\frac{1}{2} +\frac{1}{2} | 11 \rangle \delta \tilde{U}_{1S_0} \langle 11 | +\frac{1}{2} +\frac{1}{2} \rangle \langle +\frac{1}{2} +\frac{1}{2} | 11 \rangle \delta \tilde{U}_{1S_0}^+ \langle 10 | +\frac{1}{2} +\frac{1}{2} \rangle \\
& + \langle +\frac{1}{2} -\frac{1}{2} | 10 \rangle \delta \tilde{U}_{1S_0} \langle 10 | +\frac{1}{2} -\frac{1}{2} \rangle \langle +\frac{1}{2} -\frac{1}{2} | 10 \rangle \delta \tilde{U}_{1S_0}^+ \langle 10 | +\frac{1}{2} -\frac{1}{2} \rangle \\
& + \langle -\frac{1}{2} +\frac{1}{2} | 10 \rangle \delta \tilde{U}_{1S_0} \langle 10 | +\frac{1}{2} -\frac{1}{2} \rangle \langle +\frac{1}{2} -\frac{1}{2} | 10 \rangle \delta \tilde{U}_{1S_0}^+ \langle 00 | -\frac{1}{2} +\frac{1}{2} \rangle \\
& + \langle +\frac{1}{2} -\frac{1}{2} | 00 \rangle \delta \tilde{U}_{3S_1-3S_1} \langle 00 | +\frac{1}{2} -\frac{1}{2} \rangle \langle +\frac{1}{2} -\frac{1}{2} | 00 \rangle \delta \tilde{U}_{3S_1-3S_1}^+ \langle 00 | +\frac{1}{2} -\frac{1}{2} \rangle \\
& + \langle -\frac{1}{2} +\frac{1}{2} | 00 \rangle \delta \tilde{U}_{3S_1-3S_1} \langle 00 | +\frac{1}{2} -\frac{1}{2} \rangle \langle +\frac{1}{2} -\frac{1}{2} | 00 \rangle \delta \tilde{U}_{3S_1-3S_1}^+ \langle 00 | +\frac{1}{2} -\frac{1}{2} \rangle \\
& + \langle +\frac{1}{2} -\frac{1}{2} | 00 \rangle \delta \tilde{U}_{3D_1-3P_1} \langle 00 | +\frac{1}{2} -\frac{1}{2} \rangle \langle +\frac{1}{2} -\frac{1}{2} | 00 \rangle \delta \tilde{U}_{3D_1-3P_1}^+ \langle 00 | +\frac{1}{2} -\frac{1}{2} \rangle \\
& + \langle +\frac{1}{2} +\frac{1}{2} | 00 \rangle \dots \langle 00 | +\frac{1}{2} +\frac{1}{2} \rangle \\
= & \delta \tilde{U}_{1S_0} \delta \tilde{U}_{1S_0}^+ \Big|_{\substack{\tau'=\rho \\ \tau''=\rho}} + \frac{1}{4} \delta \tilde{U}_{1S_0} \delta \tilde{U}_{1S_0}^+ \Big|_{\substack{\tau'=\rho \\ \tau''=\eta}} \\
& + \frac{1}{4\rho} \delta \tilde{U}_{1S_0} \delta \tilde{U}_{1S_0}^+ \Big|_{\substack{\tau'=\eta \\ \tau''=\rho}} + \frac{3}{4\rho} \left\{ \delta \tilde{U}_{3S_1-3S_1} \delta \tilde{U}_{3S_1-3S_1}^+ + \right. \\
& \left. \delta \tilde{U}_{3S_1-3P_1} \delta \tilde{U}_{3P_1-3S_1}^+ \right\} \Big|_{\substack{\tau'=\rho \\ \tau''=\eta}} + \frac{3}{4\eta} \left[ \delta \tilde{U}_{3S_1-3S_1} \delta \tilde{U}_{3S_1-3S_1}^+ + \right. \\
& \left. \delta \tilde{U}_{3S_1-3P_1} \delta \tilde{U}_{3P_1-3S_1}^+ \right] \Big|_{\substack{\tau'=\eta \\ \tau''=\rho}}
\end{aligned}$$

Then we finally have

$$\begin{aligned}
& = \frac{1}{2} \left( \frac{2}{\pi} \right)^2 \sum_{x=3S_1, 3P_1} \int_0^\infty dk k^2 \int_0^\infty dK K^2 \int_{-1}^1 \frac{J(K \cdot k)}{\lambda} \times \\
& \left[ \delta U_{1S_0}(k, |\vec{k} - \frac{1}{2}\vec{K}|) \delta U_{1S_0}^+(|\vec{k} - \frac{1}{2}\vec{K}|, k) \Theta(k_p - |\frac{1}{2}\vec{K} + \vec{k}|) \times \right.
\end{aligned}$$

$$\begin{aligned}
& G(k_F^P - |\frac{1}{2}\vec{k} - \vec{\kappa}|) + \frac{1}{4} \delta U_{1S_0}(k, |\vec{q} - \frac{1}{2}\vec{k}|) \delta U_{1S_0}(|\vec{q} - \frac{1}{2}\vec{k}|) \times \\
& \left( G(k_F^P - |\frac{1}{2}\vec{k} + \vec{\kappa}|) G(k_F^P - |\frac{1}{2}\vec{k} - \vec{\kappa}|) + G(k_F^P - |\frac{1}{2}\vec{k} + \vec{\kappa}|) \right. \\
& \left. G(k_F^P - |\frac{1}{2}\vec{k} - \vec{\kappa}|) \right) \\
& + \frac{3}{4} \delta U_{3S_1}^+ \left( k, |\vec{q} - \frac{1}{2}\vec{k}| \right) \delta U_{3S_1}^+ (|\vec{q} - \frac{1}{2}\vec{k}|, k) \times \\
& \left. \left( G(k_F^P - |\frac{1}{2}\vec{k} + \vec{\kappa}|) G(k_F^P - |\frac{1}{2}\vec{k} - \vec{\kappa}|) + G(k_F^P - |\frac{1}{2}\vec{k} + \vec{\kappa}|) G(k_F^P - |\frac{1}{2}\vec{k} - \vec{\kappa}|) \right) \right] \\
& (7)
\end{aligned}$$

Taking only S-wave contributions :

$$\begin{aligned}
& \hat{A}_\lambda^P(q) = V \left[ 2G(k_F^P - q) + \frac{2}{\pi} \int_0^\infty dk k^2 \int_{-1}^1 d(\vec{q} \cdot \vec{k}) \left\{ \delta U_{1S_0}(k, k) \left[ G(k_F^P - |\vec{q} - 2\vec{k}|) \right. \right. \right. \\
& \left. \left. \left. + \frac{1}{2} G(k_F^P - |\vec{q} - 2\vec{k}|) \right] + \frac{3}{2} \delta U_{3S_1}^+(k, k) G(k_F^P - |\vec{q} - 2\vec{k}|) \right] G(k_F^P - q) \right. + \\
& \left. \frac{1}{4} \left( \frac{2}{\pi} \right)^2 \sum_{X=3S_1, 3D_1} \int_0^\infty dk k^2 \int_0^\infty dK K^2 \int_{-1}^1 d(\vec{k} \cdot \vec{k}) \left\{ \delta U_{1S_0}(k, |\vec{q} - \frac{1}{2}\vec{k}|) \right. \right. \\
& \left. \left. \delta U_{1S_0}^+(|\vec{q} - \frac{1}{2}\vec{k}|, k) G(k_F^P - |\frac{1}{2}\vec{k} + \vec{\kappa}|) G(k_F^P - |\frac{1}{2}\vec{k} - \vec{\kappa}|) \right. \right. \\
& \left. \left. + \frac{1}{4} \delta U_{1S_0}(k, |\vec{q} - \frac{1}{2}\vec{k}|) \delta U_{1S_0}^+(|\vec{q} - \frac{1}{2}\vec{k}|, k) \left( G(k_F^P - |\frac{1}{2}\vec{k} + \vec{\kappa}|) G(k_F^P - |\frac{1}{2}\vec{k} - \vec{\kappa}|) \right. \right. \right. \\
& \left. \left. \left. + G(k_F^P - |\frac{1}{2}\vec{k} + \vec{\kappa}|) G(k_F^P - |\frac{1}{2}\vec{k} - \vec{\kappa}|) \right) \right) \right]
\end{aligned}$$

$$\begin{aligned}
 & \frac{3}{q} \delta U_{3s_1-x} \left( k, \left| \vec{q} - \frac{1}{2} \vec{K} \right| \right) \delta U_{x=3s_1}^+ \left( \left| \vec{q} - \frac{1}{2} \vec{K} \right|, k \right) \times \\
 & \left. \left( G(k_F^F - \left| \frac{1}{2} \vec{K} + \vec{k} \right|) G(k_F^F - \left| \frac{1}{2} \vec{K} - \vec{k} \right|) + G(k_F^F - \left| \frac{1}{2} \vec{K} + \vec{k} \right|) \right. \right. \times \\
 & \left. \left. G(k_F^F - \left| \frac{1}{2} \vec{K} - \vec{k} \right|) \right) \right\} \quad ] \quad (8)
 \end{aligned}$$

\* For  $\hat{n}_\lambda(q)$  swap every instance of  $p$  with  $\lambda$  and vice versa.