

Phase Shift Calculation with Matlab

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Abstract

Adapted from handwritten notes.

I. PHASE SHIFT CALCULATION

- Let's recall the method from Landau's Quantum Mechanics II book, Section 18.3 [1]:
- This applies directly to uncoupled channels. We'll consider coupled channels separately.
- We solve for the R matrix (known as the K-matrix in other contexts), which has standing wave boundary conditions

$$R_l(k_0, k_0, E_{k_0}) = -\frac{\tan \delta_l}{\rho_T} \quad \text{where } \rho_T = 2\mu k_0, \mu = \frac{m_N}{2} \quad (1)$$

and R_l satisfies

$$R_l(k', k, E) = V_l(k', k) + \frac{2}{\pi} \mathbb{P} \int_0^\infty dp \frac{p^2 V_l(k', \rho) R_l(\rho, k, E)}{E - E_p} \quad (2)$$

- We work in units where $\frac{\hbar^2}{m} = 41.47105 \text{ MeV-fm}^2$ (for np only)
- Lab energy E_{lab} is related to COM momentum k by

$$E_{lab} = 2 \frac{\hbar^2}{m} k^2 \quad (3)$$

but the energy E_p is given by

$$E_p = \frac{\hbar^2}{2\mu} p^2 = \frac{\hbar^2}{m} p^2 \rightarrow p^2 \quad (4)$$

- So, in our standard units, R_l and V_l are in fm and ρ in fm^{-1} , so the dimensions in Eqn.(2) work out without any other factors.
- Then

$$\delta_l(E_{k_0}) = \tan^{-1}[-k_0 \cdot R_l(k_0, k_0, k_0^2)] \quad (5)$$

To Eqn.(2), we do the standard subtract and add trick to move the \mathbb{P} to an integral we can do

$$\begin{aligned} R_l(k', k) = & V_l(k', k) + \frac{2}{\pi} \int_0^\Lambda dp \frac{p^2 V_l(k', \rho) R_l(\rho, k) - k_0^2 V_l(k', k_0) R_l(k_0, k)}{k_0^2 - p^2} \quad (6) \\ & + \frac{2}{\pi} k_0^2 V_l(k', k_0) R_l(k_0, k) \underbrace{\mathbb{P} \int_0^\Lambda dp \frac{1}{k_0^2 - p^2}}_{\frac{1}{k_0} \tanh^{-1} \frac{k_0}{\Lambda}} \\ \int_0^\Lambda dp \frac{1}{k_0^2 - p^2} = & \int_0^\infty dp \frac{1}{k_0^2 - p^2} - \frac{1}{2k_0} \int_\Lambda^\infty dp \left(\frac{-1}{p - k_0} + \frac{1}{p + k_0} \right) \\ & - \frac{1}{2k_0} \log \frac{p + k_0}{p - k_0} \Big|_\Lambda^\infty = \frac{1}{2k_0} \log \frac{1 + \frac{k_0}{\Lambda}}{1 - \frac{k_0}{\Lambda}} = \frac{1}{2k_0} \log \frac{\Lambda + k_0}{\Lambda - k_0} \end{aligned}$$

(assuming $k_0 < \Lambda$)

- Now convert the integral to quadratic points $i = 1$ to N (and drop the l label):

$$R_l(k_i, k_0) = V(k_i, k_0) + \frac{2}{\pi} \sum_{j=1}^N \frac{k_j^2 V(k_i, k_j) R(k_j, k_0) w_j}{k_0^2 - k_j^2} \quad (7)$$

$$- \frac{2}{\pi} \left(\left[\sum_{m=1}^N \frac{w_m}{k_0^2 - k_m^2} \right] - \frac{1}{2k_0} \log \frac{\Lambda + k_0}{\Lambda - k_0} \right) k_0^2 V(k_i, k_0) R(k_0, k_0)$$

- We now add k_0 as the $(N+1)^{th}$ mesh point and define the vectors

$$\begin{aligned} \tilde{R}_i &= R(k_i, k_0) & i = 1, N+1 \text{ and } k_{N+1} \equiv k_0 \\ \tilde{V}_i &= V(k_i, k_0) \end{aligned} \quad (8)$$

Let's write down the standard operations by Landau.

He lets $i = N+1$ be the k_0 point

$$\Rightarrow R_i = R(k_i, k_0) \quad i = 1, N+1, \quad \text{also } V_i = V(k_i, k_0) \quad (9)$$

$N+1$ linear equations ($2\mu = 1$)

$$R_i = V_i - \frac{2}{\pi} \sum_{j=1}^N \frac{k_j^2 V_{ij} R_j w_j}{k_j^2 - k_0^2} \quad (10)$$

$$+ \frac{2}{\pi} \left(\sum_{m=1}^N \frac{w_m}{k_m^2 - k_0^2} + \frac{1}{2k_0} \log \frac{\Lambda + k_0}{\Lambda - k_0} \right) k_0^2 V_{ii} R_0$$

define

$$D_i = \begin{cases} \frac{2}{\pi} \frac{w_i k_i^2}{k_i^2 - k_0^2} & i = 1, N \\ -\frac{2}{\pi} k_0^2 \left(\sum_{j=1}^N \frac{w_j}{k_j^2 - k_0^2} + \frac{1}{2k_0} \log \frac{\Lambda + k_0}{\Lambda - k_0} \right) & i = N+1 \end{cases} \quad (11)$$

$$\Rightarrow R_i + \sum_{j=1}^{N+1} V_{ij} D_j R_j = V_i \quad (12)$$

or $F_{ij} \stackrel{def}{=} \delta_{ij} + D_j V_{ij}$ (no sum; F is an $(N+1) \times (N+1)$ matrix)

$$\Rightarrow F \vec{R} = \vec{V} \quad \leftarrow \text{or solve inversion problem}$$

$$\Rightarrow \vec{R} = F^{-1} \vec{V}$$

- OK, so for the starting attempt

V_vector

R_vector

D_vector

are formed. Actually, V_vector and D_vector are formed.

Then F_matrix is formed.

Then $R_vector = inv(F_matrix) * V_vector$

Most everything is given.

We step through the k_0 points we want, probably by stepping through E_{lab}

$\Rightarrow E_{lab}$ in MeV

$$k_0 = \sqrt{\frac{E_{lab}}{2(\hbar^2/m)}}$$

Given k_0 , D_vector is built. We'll use $i=N+1$ as the k_0 point.

$V_i = V(k_i, k_0)$ so we interpolate all of these, and its a double interpolation to get V_0

Then filling the matrix F_{ij} we have all that we need.

- Extending to coupled channels

Once we get r_{11} , r_{12} , $r_{22} \Leftarrow$ from extremes of R_matrix

Then (using radians)

$$\begin{aligned} \varepsilon &= \frac{1}{2} \text{atan} \left(\frac{2.0 r_{12}}{r_{11} - r_{22}} \right) \\ r_\varepsilon &= \frac{r_{11} - r_{22}}{\cos(2.0 \varepsilon)} \\ \delta_a &= -\text{atan} \left(\frac{1}{2} k_0 (r_{11} + r_{22} + r_\varepsilon) \right) \\ \delta_b &= -\text{atan} \left(\frac{1}{2} k_0 (r_{11} + r_{22} - r_\varepsilon) \right) \\ \bar{\varepsilon} &= \frac{1}{2} \text{asin} (\sin(2\varepsilon) \sin(\delta_a - \delta_b)) \\ \bar{\delta}_a &= \frac{1}{2} \text{atan} \left(\delta_a + \delta_b + \text{asin} \left(\frac{\tan(2\bar{\varepsilon})}{\tan(2\varepsilon)} \right) \right) \\ \bar{\delta}_b &= \frac{1}{2} \text{atan} \left(\delta_a + \delta_b - \text{asin} \left(\frac{\tan(2\bar{\varepsilon})}{\tan(2\varepsilon)} \right) \right) \end{aligned} \tag{13}$$

- Interpolate v_{11} , v_{12} , v_{21} , v_{22} onto next higher size

Then form big $\begin{pmatrix} v_{11} & v_{12} \\ v_{21} & v_{22} \end{pmatrix}$ matrix

F is a big matrix. Is it better to recast as vector?

What is V? No, a matrix is fine.

- Matrix form of phase shift code

If we take the discretized form without restricting to k_0 column

$$i, j = 1, N + 1 \rightarrow R(k_i, k_j) = V($$

$$m = max = N$$

R_{ij} corresponding to $k_0 \dots$

$$R_{ij} = V_{ij} + \frac{2}{\pi} \sum_{l=1}^N \frac{k_l^2 V_{il} R_{lj} w_l}{k_0^2 - k_l^2} - \frac{2}{\pi} \left(\sum_{l=1}^N \frac{w_l}{k_0^2 - k_l^2} - \frac{1}{2k_0} \log \frac{\Lambda + k_0}{\Lambda - k_0} \right) k_0^2 V_{i,N+1} R_{N+1,j} \quad (14)$$

$$R_{ij} + \frac{2}{\pi} \sum_{l=1}^N \frac{k_l^2 V_{il} R_{lj} w_l}{k_l^2 - k_0^2} - \frac{2}{\pi} \left(\sum_{l=1}^N \frac{w_l}{k_l^2 - k_0^2} + \frac{1}{2k_0} \log \frac{\Lambda + k_0}{\Lambda - k_0} \right) k_0^2 V_{i,N+1} R_{N+1,j} = V_{ij} \quad (15)$$

$$D_i = \begin{cases} \frac{2}{\pi} \frac{w_i k_i^2}{k_i^2 - k_0^2} & i = 1, N \\ -\frac{2}{\pi} k_0^2 \left(\sum_{l=1}^N \frac{w_l}{k_l^2 - k_0^2} + \frac{1}{2k_0} \log \frac{\Lambda + k_0}{\Lambda - k_0} \right) & i = N + 1 \end{cases} \Rightarrow \text{same as before} \quad (16)$$

$$R_{ij} + \sum_{l=1}^{N+1} V_{il} D_l R_{lj} = V_{ij} \quad (17)$$

$$\sum_l F_{il} R_{lj} = V_{ij} \quad (18)$$

$$\Rightarrow \vec{R} = F^{-1} \vec{V}$$

Now simply as matrices

Then still the $N+1, N+1$ element of R that we want

OK that works immediately

[1] R. H. Landau, *Quantum Mechanics II* (John Wiley & Sons, Inc., New York, 1996), 2nd ed.