

Introduction to Fourier transform and signal analysis

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Outline

1 Continuous Fourier transform

2 Discrete Fourier transform

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Orthogonal condition

- Any two vectors **a**, **b** satisfying following condition are mutually orthogonal.

$$\mathbf{a}^* \cdot \mathbf{b} = 0 \quad (1)$$

- Any two functions $a(x)$, $b(x)$ satisfying the following condition are mutually orthogonal.

$$\int a^*(x) \cdot b(x) dx = 0 \quad (2)$$

- * means complex conjugate.

Complete and orthogonal basis

- $\cos nx$ and $\sin mx$ are mutually orthogonal in which n and m are integers.

$$\begin{aligned}\int_{-\pi}^{\pi} \cos nx \cdot \sin mx dx &= 0 \\ \int_{-\pi}^{\pi} \cos nx \cdot \cos mx dx &= \pi \delta_{nm} \\ \int_{-\pi}^{\pi} \sin nx \cdot \sin mx dx &= \pi \delta_{nm}\end{aligned}\tag{3}$$

- δ_{nm} is Dirac-delta symbol. It means $\delta_{nn} = 1$ and $\delta_{nm} = 0$ when $n \neq m$.

Fourier series

Since $\cos nx$ and $\sin mx$ are mutually orthogonal, we can expand an arbitrary periodic function $f(x)$ by them. we shall have a series expansion of $f(x)$ which has 2π period.

$$\begin{aligned}
 f(x) &= a_0 + \sum_{k=1}^{\infty} (a_k \cos kx + b_k \sin kx) \\
 a_0 &= \frac{1}{2\pi} \int_{-\pi}^{\pi} f(x) dx \\
 a_k &= \frac{1}{\pi} \int_{-\pi}^{\pi} f(x) \cos kx dx \\
 b_k &= \frac{1}{\pi} \int_{-\pi}^{\pi} f(x) \sin kx dx
 \end{aligned} \tag{4}$$

Fourier series

If $f(x)$ has L period instead of 2π , x is replaced with $\pi x/L$.

$$\begin{aligned}
 f(x) &= a_0 + \sum_{k=1}^{\infty} \left(a_k \cos \frac{2k\pi x}{L} + b_k \sin \frac{2k\pi x}{L} \right) \\
 a_0 &= \frac{1}{L} \int_{-\frac{L}{2}}^{\frac{L}{2}} f(x) dx \\
 a_k &= \frac{2}{L} \int_{-\frac{L}{2}}^{\frac{L}{2}} f(x) \cos \frac{2k\pi x}{L} dx, k = 1, 2, \dots \\
 b_k &= \frac{2}{L} \int_{-\frac{L}{2}}^{\frac{L}{2}} f(x) \sin \frac{2k\pi x}{L} dx, k = 1, 2, \dots
 \end{aligned} \tag{5}$$

Fourier series of step function

$f(x)$ is a periodic function with 2π period and it's defined as follows.

$$\begin{aligned} f(x) &= 0, -\pi < x < 0 \\ f(x) &= h, 0 < x < \pi \end{aligned} \quad (6)$$

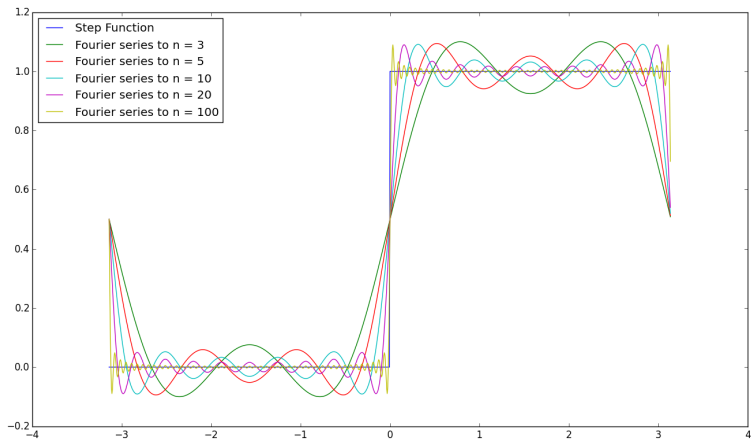
Fourier series expansion of $f(x)$ is

$$f(x) = \frac{h}{2} + \frac{2h}{\pi} \left(\frac{\sin x}{1} + \frac{\sin 3x}{3} + \frac{\sin 5x}{5} + \dots \right) \quad (7)$$

$f(x)$ is piecewise continuous within the periodic region. Fourier series of $f(x)$ converges at speed of $1/n$.

Fourier series of step function

Notice that even we have $n = 100$, there are still overshoots phenomenon near 0. This is called Gibbs phenomenon.



Fourier series of triangular function

$f(x)$ is a periodic function with 2π period and it's defined as follows.

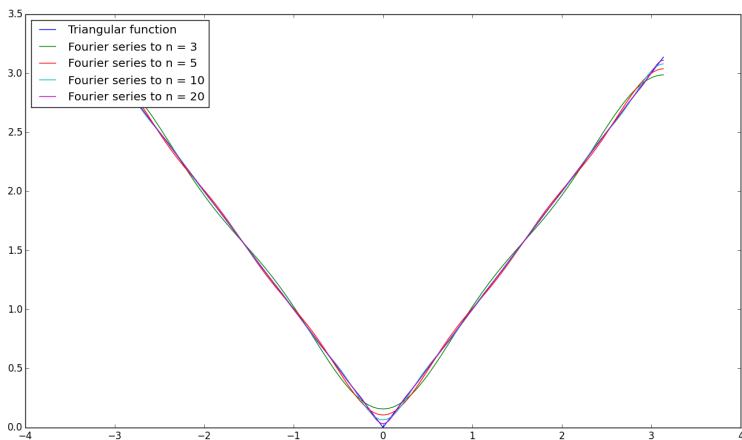
$$\begin{aligned} f(x) &= -x, -\pi < x < 0 \\ f(x) &= x, 0 < x < \pi \end{aligned} \quad (8)$$

Fourier series expansion of $f(x)$ is

$$f(x) = \frac{\pi}{2} - \frac{4}{\pi} \sum_{n=1,3,5,\dots} \left(\frac{\cos nx}{n^2} \right) \quad (9)$$

$f(x)$ is continuous and its derivative is piecewise continuous within the periodic region. Fourier series of $f(x)$ converges at speed of $1/n^2$.

Fourier series of triangular function



Complex Fourier series

Using Euler's formula, Eq. 4 becomes

$$f(x) = a_0 + \sum_{k=1}^{\infty} \left(\frac{a_k - ib_k}{2} e^{ikx} + \frac{a_k + ib_k}{2} e^{-ikx} \right)$$

Let $c_0 \equiv a_0$, $c_k \equiv \frac{a_k - ib_k}{2}$ and $c_{-k} \equiv \frac{a_k + ib_k}{2}$, we have

$$\begin{aligned} f(x) &= \sum_{m=-\infty}^{\infty} c_m e^{imx} \\ c_m &= \frac{1}{2\pi} \int_{-\pi}^{\pi} f(x) e^{-imx} dx \end{aligned} \quad (10)$$

e^{imx} and e^{inx} are also mutually orthogonal provided $n \neq m$ and it forms a complete set. Therefore, it can be used as orthogonal basis.

Complex Fourier series

If $f(x)$ has T period instead of 2π , x is replaced with $2\pi x/T$.

$$f(x) = \sum_{m=-\infty}^{\infty} c_m e^{i \frac{2\pi m x}{T}}$$

$$c_m = \frac{1}{T} \int_{-\frac{T}{2}}^{\frac{T}{2}} f(x) e^{-i \frac{2\pi m x}{T}} dx, m = 0, 1, 2, \dots \quad (11)$$

Fourier transform

from Eq. 11, we define variables $k \equiv \frac{2\pi m}{T}$, $\hat{f}(k) \equiv \frac{c_m T}{\sqrt{2\pi}}$ and

$$\Delta k \equiv \frac{2\pi(m+1)}{T} - \frac{2\pi m}{T} = \frac{2\pi}{T}.$$

We can have

$$f(x) = \frac{1}{\sqrt{2\pi}} \sum_{m=-\infty}^{\infty} \hat{f}(k) e^{ikx} \Delta k$$

$$\hat{f}(k) = \frac{1}{\sqrt{2\pi}} \int_{-\frac{T}{2}}^{\frac{T}{2}} f(x) e^{-ikx} dx$$

Fourier transform

Let $T \longrightarrow \infty$

$$f(x) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \hat{f}(k) e^{ikx} dk \quad (12)$$

$$\hat{f}(k) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} f(x) e^{-ikx} dx \quad (13)$$

Eq.13 is the *Fourier transform* of $f(x)$ and Eq.12 is the *inverse Fourier transform* of $\hat{f}(k)$.

Properties of Fourier transform

$f(x)$, $g(x)$ and $h(x)$ are functions and their Fourier transforms are $\hat{f}(k)$, $\hat{g}(k)$ and $\hat{h}(k)$. a , b , x_0 and k_0 are real numbers.

- Linearity: If $h(x) = af(x) + bg(x)$, then Fourier transform of $h(x)$ equals to $\hat{h}(k) = a\hat{f}(k) + b\hat{g}(k)$.
- Translation: If $h(x) = f(x - x_0)$, then $\hat{h}(k) = \hat{f}(k)e^{-ikx_0}$
- Modulation: If $h(x) = e^{ik_0x}f(x)$, then $\hat{h}(k) = \hat{f}(k - k_0)$
- Scaling: If $h(x) = f(ax)$, then $\hat{h}(k) = \frac{1}{a}\hat{f}(\frac{k}{a})$
- Conjugation: If $h(x) = f^*(x)$, then $\hat{h}(k) = \hat{f}^*(-k)$. With this property, one can know that if $f(x)$ is real and then $\hat{f}^*(-k) = \hat{f}(k)$. One can also find that if $f(x)$ is real and then $|\hat{f}(k)| = |\hat{f}(-k)|$.

Properties of Fourier transform

- If $f(x)$ is even, then $\hat{f}(-k) = \hat{f}(k)$.
- If $f(x)$ is odd, then $\hat{f}(-k) = -\hat{f}(k)$.
- If $f(x)$ is real and even, then $\hat{f}(k)$ is real and even.
- If $f(x)$ is real and odd, then $\hat{f}(k)$ is imaginary and odd.
- If $f(x)$ is imaginary and even, then $\hat{f}(k)$ is imaginary and even.
- If $f(x)$ is imaginary and odd, then $\hat{f}(k)$ is real and odd.

Dirac delta function

Dirac delta function is a generalized function defined as the following equation.

$$\begin{aligned} f(0) &= \int_{-\infty}^{\infty} f(x)\delta(x)dx \\ \int_{-\infty}^{\infty} \delta(x)dx &= 1 \end{aligned} \tag{14}$$

The Dirac delta function can be loosely thought as a function which equals to infinite at $x = 0$ and to zero else where.

$$\delta(x) = \begin{cases} +\infty, & x = 0 \\ 0, & x \neq 0 \end{cases}$$

Dirac delta function

From Eq.13 and Eq.12

$$\begin{aligned}\hat{f}(k) &= \frac{1}{2\pi} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \hat{f}(k') e^{ik'x} dk' e^{-ikx} dx \\ &= \frac{1}{2\pi} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \hat{f}(k') e^{i(k'-k)x} dx dk'\end{aligned}$$

Comparing to "Dirac delta function", we have

$$\begin{aligned}\hat{f}(k) &= \int_{-\infty}^{\infty} \hat{f}(k') \delta(k' - k) dk' \\ \delta(k' - k) &= \frac{1}{2\pi} \int_{-\infty}^{\infty} e^{i(k'-k)x} dx\end{aligned}\tag{15}$$

Eq.15 doesn't converge by itself, it is only well defined as part of an integrand.

Convolution theory

Considering two functions $f(x)$ and $g(x)$ with their Fourier transform $F(k)$ and $G(k)$. We define an operation

$$f * g = \int_{-\infty}^{\infty} g(y)f(x-y)dy \quad (16)$$

as the convolution of the two functions $f(x)$ and $g(x)$ over the interval $\{-\infty \sim \infty\}$. It satisfies the following relation:

$$f * g = \int_{-\infty}^{\infty} F(k)G(k)e^{ikx} dt \quad (17)$$

Let $h(x)$ be $f * g$ and $\hat{h}(k)$ be the Fourier transform of $h(x)$, we have

$$\hat{h}(k) = \sqrt{2\pi}F(k)G(k) \quad (18)$$

Parseval relation

$$\begin{aligned}
 \int_{-\infty}^{\infty} f(x)g(x)^* dx &= \int_{-\infty}^{\infty} \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} F(k)e^{ikx} dk \frac{1}{\sqrt{2\pi}} \\
 &\quad \int_{-\infty}^{\infty} G^*(k')e^{-ik'x} dk' dx \\
 &= \int_{-\infty}^{\infty} \frac{1}{2\pi} \int_{-\infty}^{\infty} F(k)G^*(k')e^{i(k-k')x} dk dk'
 \end{aligned}$$

By using Eq. 15, we have the Parseval's relation.

$$\int_{-\infty}^{\infty} f(x)g^*(x)dx = \int_{-\infty}^{\infty} F(k)G^*(k)dk \quad (19)$$

Calculating inner product of two functions gets same result as the inner product of their Fourier transform.

Cross-correlation

Considering two functions $f(x)$ and $g(x)$ with their Fourier transform $F(k)$ and $G(k)$. We define cross-correlation as

$$(f \star g)(x) = \int_{-\infty}^{\infty} f^*(x+y)g(x)dy \quad (20)$$

as the cross-correlation of the two functions $f(x)$ and $g(x)$ over the interval $\{-\infty \sim \infty\}$. It satisfies the following relation: Let $h(x)$ be $f \star g$ and $\hat{h}(k)$ be the Fourier transform of $h(x)$, we have

$$\hat{h}(k) = \sqrt{2\pi}F^*(k)G(k) \quad (21)$$

Autocorrelation is the cross-correlation of the signal with itself.

$$(f \star f)(x) = \int_{-\infty}^{\infty} f^*(x+y)f(x)dy \quad (22)$$

Uncertainty principle

One important properties of Fourier transform is the uncertainty principle. It states that the more concentrated $f(x)$ is, the more spread its Fourier transform $\hat{f}(k)$ is.

Without loss of generality, we consider $f(x)$ as a normalized function which means $\int_{-\infty}^{\infty} |f(x)|^2 dx = 1$, we have uncertainty relation:

$$\left(\int_{-\infty}^{\infty} (x - x_0)^2 |f(x)|^2 dx \right) \left(\int_{-\infty}^{\infty} (k - k_0)^2 |\hat{f}(k)|^2 dk \right) \geq \frac{1}{16\pi^2} \quad (23)$$

for any x_0 and $k_0 \in \mathbf{R}$. [3]

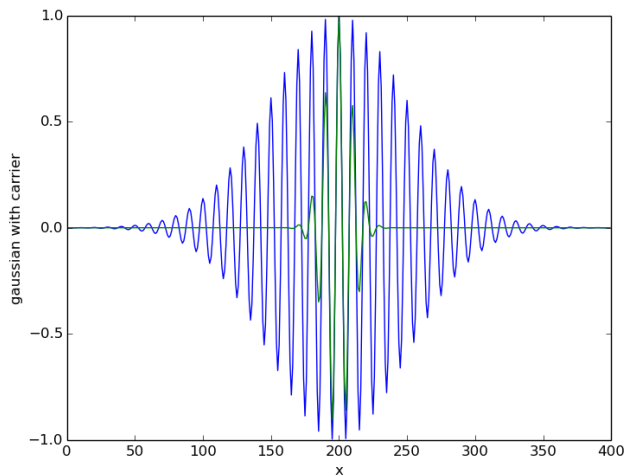
Fourier transform of a Gaussian pulse

$$\begin{aligned}
 f(x) &= f_0 e^{\frac{-x^2}{2\sigma^2}} e^{ik_0 x} \\
 \hat{f}(k) &= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} f(x) e^{-ikx} dx \\
 &= \frac{f_0}{1/\sigma^2} e^{\frac{-(k_0 - k)^2}{2/\sigma^2}} \\
 |\hat{f}(k)|^2 &\propto e^{\frac{-(k_0 - k)^2}{1/\sigma^2}}
 \end{aligned}$$

Wider the $f(x)$ spread, the more concentrated $\hat{f}(k)$ is.

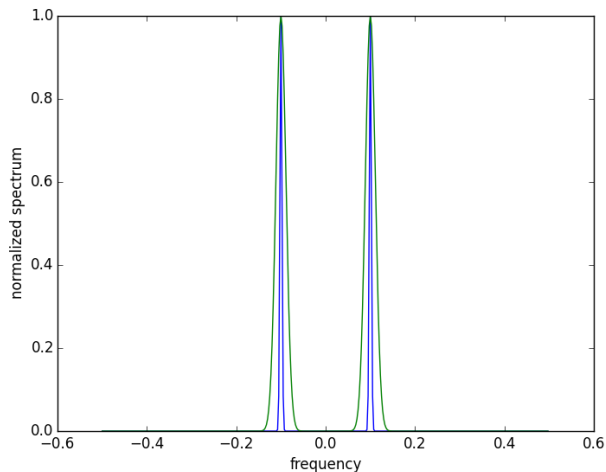
Fourier transform of a Gaussian pulse

Signals with different width.



Fourier transform of a Gaussian pulse

The bandwidth of the signals are different as well.



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- 1 Continuous Fourier transform
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- 3 References

Nyquist critical frequency

Critical sampling of a sine wave is two sample points per cycle. This leads to *Nyquist critical frequency* f_c .

$$f_c = \frac{1}{2\Delta} \quad (24)$$

In above equation, Δ is the sampling interval.

Sampling theorem: If a continuous signal $h(t)$ sampled with interval Δ happens to be bandwidth limited to frequencies smaller than f_c . $h(t)$ is completely determined by its samples h_n . In fact, $h(t)$ is given by

$$h(t) = \Delta \sum_{n=-\infty}^{\infty} h_n \frac{\sin[2\pi f_c(t - n\Delta)]}{\pi(t - n\Delta)} \quad (25)$$

It's known as Whittaker - Shannon interpolation formula.

Discrete Fourier transform

Signal $h(t)$ is sampled with N consecutive values and sampling interval Δ . We have $h_k \equiv h(t_k)$ and $t_k \equiv k * \Delta$, $k = 0, 1, 2, \dots, N - 1$.

With N discrete input, we evidently can only output independent values no more than N . Therefore, we seek for frequencies with values

$$f_n \equiv \frac{n}{N\Delta}, n = -\frac{N}{2}, \dots, \frac{N}{2} \quad (26)$$

Discrete Fourier transform

Fourier transform of Signal $h(t)$ is $H(f)$. We have discrete Fourier transform H_n .

$$\begin{aligned}
 H(f_n) &= \int_{-\infty}^{\infty} h(t) e^{-i2\pi f_n t} dt \approx \Delta \sum_{k=0}^{N-1} h_k e^{-i2\pi f_n t_k} \\
 &= \Delta \sum_{k=0}^{N-1} h_k e^{-i2\pi kn/N} \\
 H_n &\equiv \sum_{k=0}^{N-1} h_k e^{-i2\pi kn/N}
 \end{aligned} \tag{27}$$

Inverse Fourier transform is

$$h_k \equiv \frac{1}{N} \sum_{n=0}^{N-1} H_n e^{i2\pi kn/N} \tag{28}$$

Periodicity of discrete Fourier transform

From Eq.27, if we substitute n with $n + N$, we have $H_n = H_{n+N}$. Therefore, discrete Fourier transform has periodicity of N .

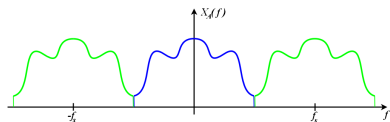
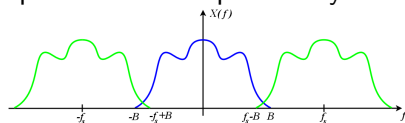
$$\begin{aligned}
 H_{n+N} &= \sum_{k=0}^{N-1} h_k e^{-i2\pi k(n+N)/N} \\
 &= \sum_{k=0}^{N-1} h_k e^{-i2\pi k(n)/N} e^{-i2\pi kN/N} \\
 &= H_n
 \end{aligned} \tag{29}$$

Critical frequency f_c corresponds to $\frac{1}{2\Delta}$.

We can see that discrete Fourier transform has f_s period where $f_s = 1/\Delta = 2 * f_c$ is the sampling frequency.

Aliasing

If we have a signal with its bandlimit larger than f_c , we have following spectrum due to periodicity of DFT.

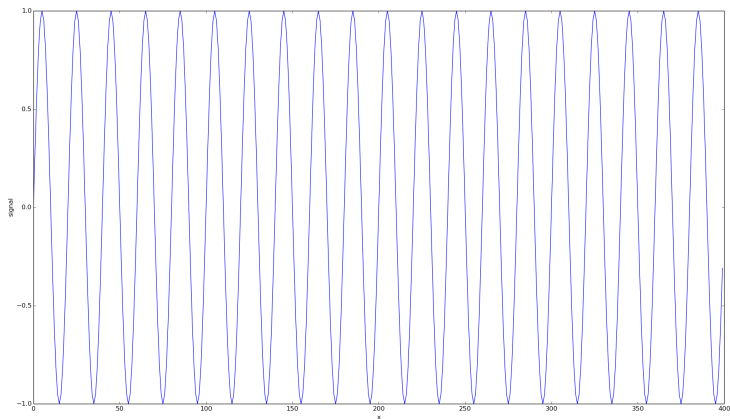


[5]

Aliased frequency is $f - N * f_s$ where N is an integer.

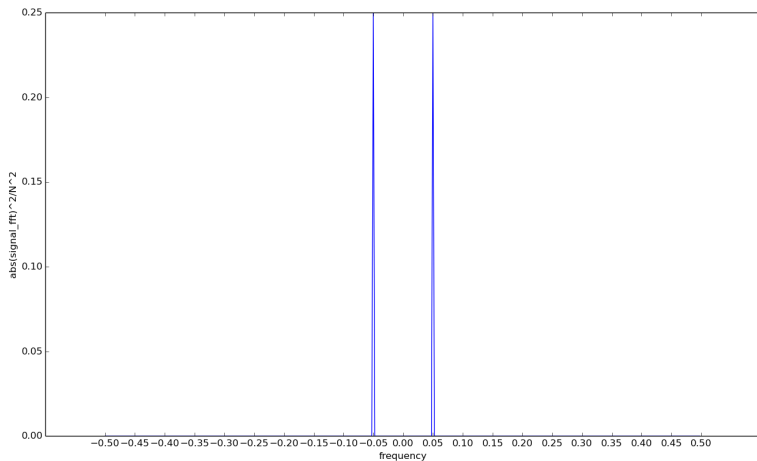
Aliasing example: original signal

Let's say we have a sinusoidal signal of frequency 0.05. The sampling interval is 1. We have the signal



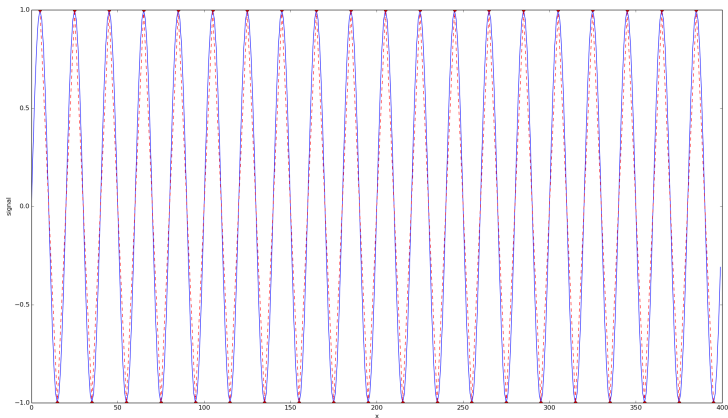
Aliasing example: spectrum of original signal

and we have its spectrum



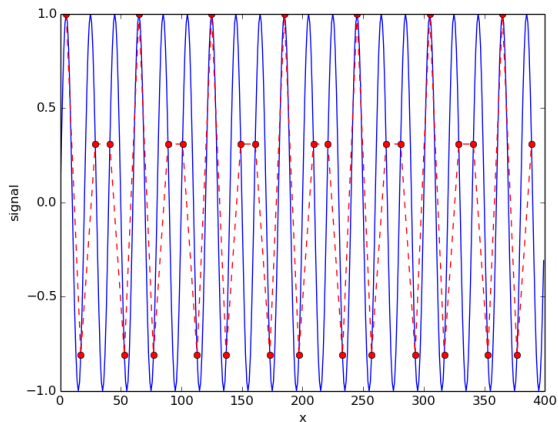
Aliasing example: critical sampling of original signal

The critical sampling interval of the original signal is 10 which is half of the signal period.



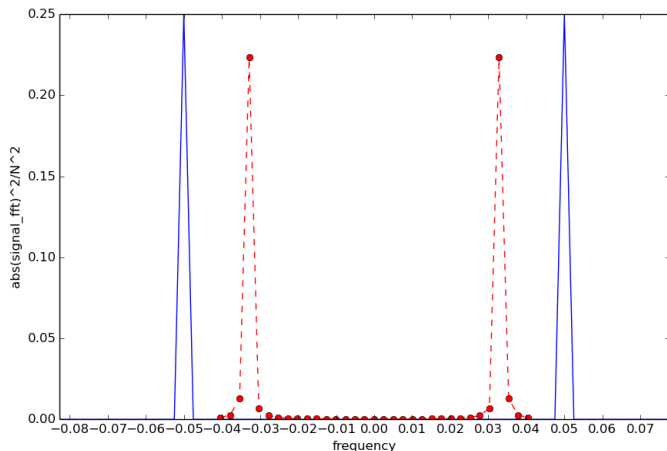
Aliasing example: under sampling of original signal

If we sampled the original sinusoidal signal with period 12, aliasing happens.



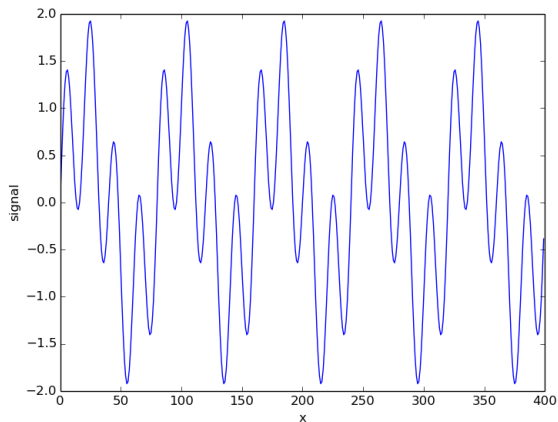
Aliasing example: DFT of under sampled signal

f_c of downsampled signal is $\frac{1}{2 \times 12}$, aliased frequency is $f - 2 * f_c = -0.03333$ and it has symmetric spectrum due to real signal.



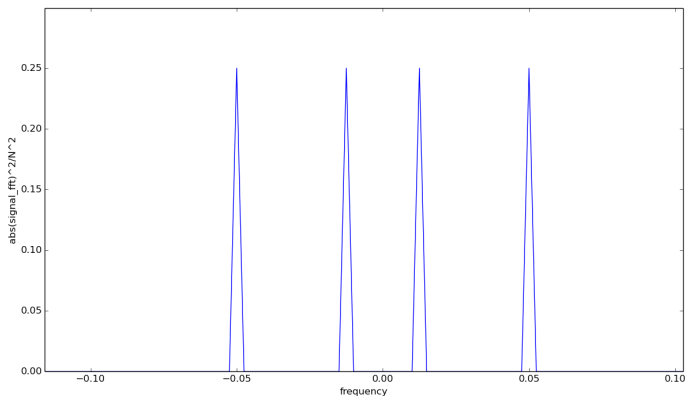
Aliasing example: two frequency signal

Let's say we have a signal containing two sinusoidal signal of frequency 0.05 and 0.0125. The sampling interval is 1. We have the signal



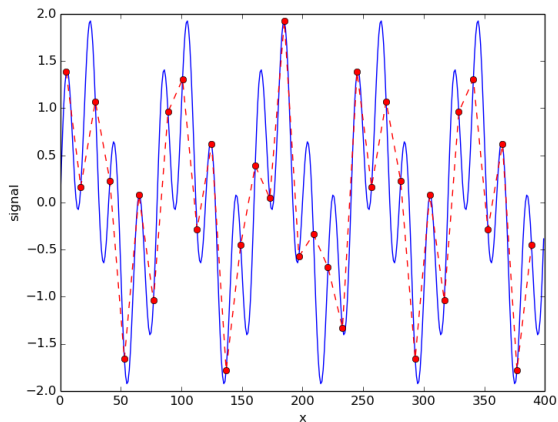
Aliasing example: spectrum of two frequency signal

and we have its spectrum



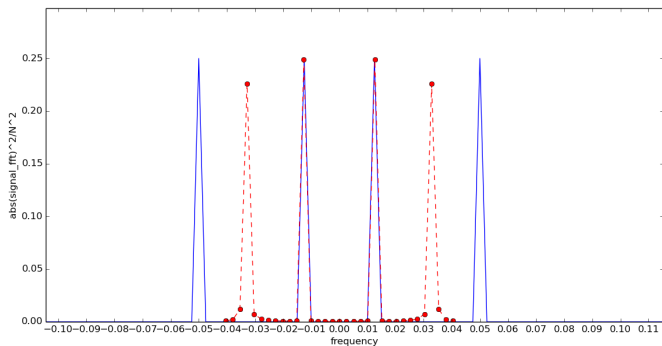
Aliasing example: downsampled two frequency signal

Doing same undersampling with interval 12.



Aliasing example: DFT of downsampled signal

We have the spectrum of downsampled signal.



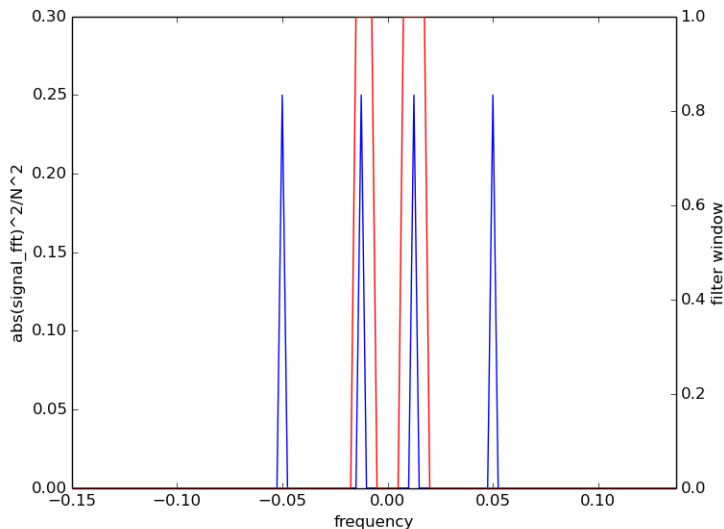
Filtering

We now want to get one of the two frequency out of the signal. We will adapt a proper rectangular window to the spectrum.

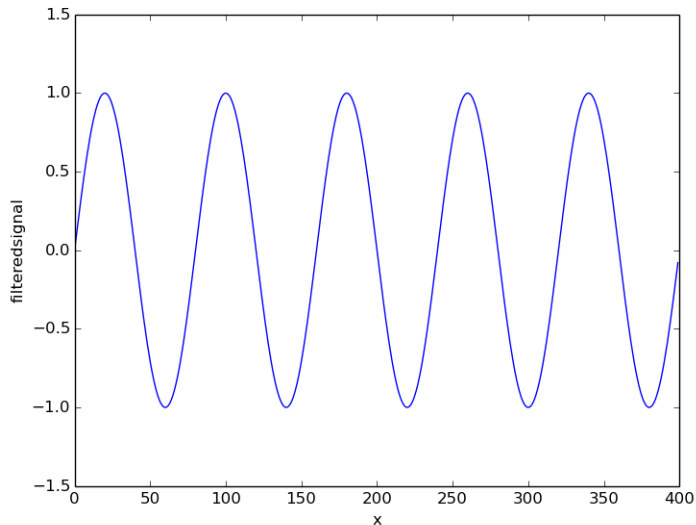
Assuming we have a filter function $w(f)$ and a multi-frequency signal $f(t)$, we simply do following steps to get the frequency band we want.

$$\mathcal{F}^{-1}\{w(f)\mathcal{F}\{f(t)\}\} \quad (30)$$

Filtering example: filtering window and signal spectrum










Filtering example: filtered signal



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