

The Seven Wonders of the World

Notes on 21st-century physics

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Version 0.1, 10 February 2024



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Cover image: the “Pale Blue Dot” image of the Earth taken by Voyager 1,
from right outside Pluto’s orbit

<https://science.nasa.gov/resource/voyager-1s-pale-blue-dot/>

Introduction

The loss implied in such an acquisition can be estimated only by those who have been compelled to unlearn a science that they might at length begin to learn it.

J. C. Maxwell 1878

Until some decades ago, the 18th-century physical notions typically taught in introductory Bachelor physics courses were enough to prepare an engineer for future specializations and jobs. Students who wanted to venture into modern theories, such as Relativity, were required to **re-learn** some of the most important physical notions – *Energy*, mass, time, entropy above all – which in these theories are quite different from the 18th-century ones. But at that time the modern theories still mostly had only theoretical, not practical, importance. So the re-learning efforts of the curious students could perhaps be justified.

That situation has changed today. Modern theories are an essential part of many everyday technologies, like nuclear reactors and the [Global Positioning System](#)¹ (see the entertaining discussion in [Taylor & Wheeler 2000 Project A](#), and also [Petit & Wolf 2005](#); [Fliegel & DiEsposti 1996](#); [Ashby 2002](#); [Müller et al. 2008](#)); and they are required for developing new technological possibilities, from [quantum computers](#)² to [solar sails](#)³ (see for instance the [CubeSail](#)⁴ mission). An engineering student (including communication and data engineering) may likely end up in a job that requires an understanding of modern physical notions. The diffusion of [large language models](#)⁵ will moreover require future engineers who actually **understand** those physical notions, not little monkeys who have been trained to manipulate some equations and to throw some technobabble around. Automated large language models are faster, cheaper, and more precise in doing the

"Numerous relativistic issues and effects play roles in the global positioning system, on which millions of drivers, hikers, sailors, and pilots depend to find out where they are. The GPS system is, in effect, a realization of Einstein's view of space and time. Indeed, the system cannot function properly without taking account of fundamental relativistic principles." [Ashby 2002](#)

latter kind of monkey activities. So why should one hire a human to do the same?

While moving from the older to the newer notions often requires re-learning efforts and conceptual re-orientations, the move in the opposite direction is less demanding. The modern physical notions are more encompassing than their 18th-century parents. Their understanding leads to an understanding of their older counterparts as approximate and special cases. Students, moreover, have often been hearing quite early from mass media about the new notions; for instance about the equivalence of mass and energy. Owing to this early exposure, students sometimes ask very intelligent and deep questions – such as “*should the mass of the body be included in its internal energy?*” – when exposed to the old notions.

It is therefore high time that introductory Bachelor physics courses be based on modern physical notions. Students should not be required to waste time and mental effort to learn something that they must unlearn and re-learn, only because of academia’s and teachers’ inertia.

Some teachers say “it would be too difficult for students to understand modern ideas, because they are too familiar with the old ones. This is why we need to teach the old and slightly incorrect ideas first, and correct them later”. I think that this kind of reasoning is scientifically unacceptable and leads to a vicious circle. Students are unfamiliar with new notions only because they were raised by a generation who was taught old ones. New notions become familiar only after a couple of generations learn them at an early stage. This is obvious if you consider that notions such as “energy”, “electromagnetic field”, “vector” are very familiar today, but were absolutely *unfamiliar* a couple centuries ago. If we had always taught what’s familiar, then we would still be teaching about the elements *air, earth, water, fire*, and that the Sun revolves around the Earth. Arguments in favour of teaching old notions are for the most part just pretexts for laziness.

The present lecture notes are an experiment and an attempt to introduce classical mechanics and thermodynamics from modern physical notions. The core equations remain the same, but the students should have a broader conception of their meaning and of the symbols that appear in them.

URLs for chapter *Introduction*

1. <https://www.gps.gov>
2. <https://www.ibm.com/topics/quantum-computing>
3. <https://www.planetary.org/sci-tech/lightsail>
4. <https://www.cubesail.us>
5. <https://www.ibm.com/topics/large-language-models>

Chapter 0

Guide to the student

Maths prerequisites

- Working familiarity with algebra, its operations, and their properties, including square roots; knowledge of exponentials and logarithms is useful.
- Working familiarity with solving equations and inequalities, linear and non-linear.
- Working familiarity with the study of functions of one real variable.
- Working familiarity with derivatives.
- Understanding of what an integral is, even if you won't be required to solve integrals.
- Working familiarity with vector calculus.
- Some familiarity with functions of many variables.
- Understanding of what partial derivatives are.

Physics prerequisites

Just some vague reminiscences of secondary/high-school physics should be enough.

It can be beneficial if you are familiar with basic physics notions such as *velocity*, *mass*, *force*, and similar ones; and if you vaguely remember about Newton's laws.

Structure of this text

👉 Graphical devices

The text includes the following graphical devices:

- Important notions and definitions are also given in **boldface**.
- Important-notion boxes:

Something

This is a definition or explanation of Something.

- Warnings and important points that require careful thinking:

❗ Something

Something you must be careful about.

- Exercises:

💡 Exercise 0.1

This isn't an exercise

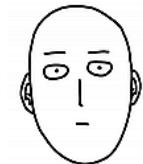
- Discussions and connections with more advanced physics:

👁️ How things really are in quantum physics

Just for your curiosity: you don't need to remember any of this.

👉 Side pictures and quotes

Pictures, graphs, or quotes related to the material are displayed on the right.



This is an image of Saitama, which actually has nothing to do with the text on the left.

👉 Hyperlinks and bibliography

Some words are hyperlinks, like this one about [One Punch Man¹](#); you also recognize them because they have a little footnote number. The links' URLs are listed at the end of each chapter, just in case you're reading a printed copy and wonder what the link was.

The text gives bibliographic references, like “Einstein 1905a”, to scientific literature. The references are listed in the final Bibliography on page 93.

These references are not part of the course and don't need to be “studied”. They are given for two reasons:

- For your own curiosity.
- To back up what's written in the text. In science you should not believe something just because you've read it somewhere. You should, as much as possible, go and check *for yourself how the logic behind the statement is proved and what is the experimental evidence behind the statement.*

👉 Notation and terminology

Mathematical notation, as well as notation for physical dimensions, strictly follows the standards given by the International System of Units (SI)², listed for example in ISO 2009 and ISO 2019.

“Believe nothing, O monks, merely because you have been told it, or because it is traditional, or because you yourselves have imagined it. Do not believe what your teacher tells you merely out of respect for the teacher.”

(attributed to Gautama Buddha)

URLs for chapter 0

1. <https://onepunchman.fandom.com>
2. <https://www.nist.gov/pml/special-publication-811>

Chapter 1

Physics?

If you think about it, many things we ordinarily do every day are some sort of magic. Think of how you can instantaneously see and speak with a person living on another continent, in real time, using just a small widget in the palm of your hand. Think of how you can instantaneously see where you are on the Earth, using the same widget. Think of how fast you can go to another country, by flying in a huge metal thing. Think of how you can command and interact with a purely fictitious animated world when you play on your computer. The list can go on forever. Other things are luckily less ordinary, but still inspire a lot of awe: think of the devastating power unleashed by something roughly as small as a tennis ball, in an atomic bomb.

We can do these astonishing things thanks to our understanding of how the world works. That's Physics.

Many things can be said and have been said about science and physics. Rather than repeating what's been already written in many excellent books, I invite you to take a break here and go read their introductions. Choose as you please; compare what they say; don't limit yourself to popular books.

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1.1 Several possible formalisms or “languages”

Physics can be expressed and written from wildly different points of view, using wildly different principles. Let's call these “different physics languages”; a more technical name is “physics formalisms”. One may approach a physics phenomenon or problem in terms of *Lagrangeans*, or

Hamiltonians, or *fibre bundles*, or *categories*, or *action principles*, or many other formalisms. These formalisms or languages are not completely separated; we know how to translate among them. In "doing" physics, one may jump among formalisms, because some ideas may be easier to express, or some results easier to find, in one formalism than another. No matter which physics formalism you choose, the results and the concrete applications are still the same. The choice is to a great extent subjective, based on your aesthetic tastes. You see that in "doing" physics you can express your personality and put your own artistic touch; this is why it's such a cool subject (and other scientific subjects are like this too).

In these notes I'm choosing one particular formalism: the one that for me is the most easily *visualizable*; because I believe that visualization can be beneficial in learning new things. Or maybe I'm choosing it just because I like it best. I encourage you to explore how the physics you've learned is expressed in other physics formalisms; maybe you'll like another physics formalism better.

The formalism we'll be using might be called "field theory". Roughly speaking it takes as starting point the ideas of space and time, or better spacetime, in which there are different kinds of "stuff". It expresses the regularity and patterns that we observe in physical phenomena as "budgets" about the different kinds of stuff, and of relations between these kinds. Please don't take the description just given too literally; it's just meant to give you a very vague idea of the field-theoretical viewpoint.

It goes without saying that all these "physics languages" are to a great extent mathematical.

One reason is that numbers allow us to convey information in a concise and precise way. Imagine you have to tell someone, who doesn't know Bergen, where in Bergen you are right now, to within 10 m. You can do that with a description, "... and there's a building called so-and-so which looks like so-and-so...", which would be lengthy and tricky. Or you can just give two numbers: latitude and longitude:

60.369 40, 5.3518 .

And in these two numbers all digits are important; for instance, the latitude is not 60.369 47.

But the most important reason is that mathematics allows us to describe and follow the patterns and variety of physical phenomena in a greatly concise and precise way. And to develop their relationships in a rigorous

$$\delta \int L dt = 0 \quad L = \frac{1}{2} mv^2$$

$$\mathbf{F} = \frac{d}{dt} m \mathbf{v} \quad \mathbf{F} = 0$$

Example of two different formalisms (red, blue) expressing the same physical phenomenon.

"this grand book, the universe, which stands continually open to our gaze. But the book cannot be understood unless one first learns to comprehend the language and read the letters in which it is composed. It is written in the language of mathematics"

Galilei 1623

way. All our present technology would have been impossible to discover, and would be impossible to realize, without the mathematical language of physics.

I invite you again to read what many good texts say about the relationship between physics and mathematics. No point repeating here what is said better elsewhere.

1.2 Quantities: primitive and derived

One topic must be briefly discussed because it's important for understanding the notes that follow. It's the distinction between *primitive* and *derived* quantities.

I shall assume that you already know what a **physical quantity** is. Examples are: position, duration, velocity, pressure, energy, temperature.

A **derived quantity** is one that is defined in terms of other quantities. For example, velocity v (more precisely: average velocity) is defined as the ratio between a distance d (a vector) and a time duration t :

$$v := \frac{d}{t}$$

where the symbol “:=” means “is defined as” or “is defined by”. This means that in principle we could avoid using the word “velocity” and the symbol “ v ” altogether, and instead always speak about distance and duration, using their symbols. It would lead to very long sentences and formulae and would be extremely inconvenient, but it could be done. The definition of a derived quantity often tells us how that quantity can be measured.

A derived quantity is defined in terms of other quantities, and these may in turn be derived quantities, that is, defined in terms of still other quantities, and so on. But at some point this chain of definitions must come to an end, otherwise we would go around in circles.

A **primitive quantity** is one that we do not define in terms of other quantities. Primitive quantities are the building blocks from which we define all others. That they are not defined in terms of others doesn't mean that we cannot try to explain them. But such explanations must be taken as informal and heuristic. Primitive quantities are often explained through metaphors and by appealing to intuition. You must always be wary of such explanations, because they may fail you spectacularly in some situations.

Often we have a choice about which quantities should be primitive and which should be derived. For instance, *energy* can be defined, in a

“There is nothing that can be said by mathematical symbols and relations which cannot also be said by words. The converse, however, is false. Much that can be and is said by words cannot successfully be put into equations, because it is nonsense.”

Truesdell 1966

somewhat complicated way, in terms of quantities like *work* and *heat*, which would then need to be taken as primitive. Or we can take *energy* as primitive, and define *work* and *heat* in terms of it. This second choice can be more convenient to develop a physical theory. It often happens that a quantity is very convenient for building a theory, if used as primitive; but difficult to understand intuitively. Vice versa, a quantity can be very intuitive but lead to a complicated theory. Among quantities which we'll take as primitive are: *time*, *space* and *length*, *matter*, *energy*, *momentum*, *entropy*, *temperature*, and several others. All will be discussed soon.

1.3 Physical dimensions and units

Measurement is the process by which we determine the value of a physical quantity. Measurements can be extremely complex, and can extremely different even if they are about the same quantity. Consider the ways we can measure the mass of a football, compared to the ways we can measure the mass of the Sun.

To each quantity we associate a **physical dimension**. The term ‘dimension’ here has nothing to do with physical extension, as in “the dimensions of this box”; be careful not to confuse the two. Usually it’s clear which one is meant from the context. Physical dimensions help us avoid making operations that don’t make sense with some quantities. For example, it doesn’t make sense to sum up the volume of a body of water with its temperature; and indeed the volume has dimension *length*³, whereas temperature has dimension *temperature*, and quantities with different dimensions cannot be added up.

With each physical dimension we can associate a measurement **unit**, which expresses a basic standard for comparing the measurement results of similar quantities. For example, we could use the *minute* or the *second* as units to measure *time*.

One can choose a basic set of physical dimensions from which to define all others, and for these a set of standard units. Here we shall follow the [International System of Units \(SI\)](#)¹ (see also [NIST Special Publication 811](#)²).

The topics of measurement and physical dimensions, which are studied in *metrology* and in *dimensional analysis*, could occupy an entire course by themselves. I shall assume that you already know their basics notions and that you read about the SI.

The measurement of some physical quantities consists in just one number with associated physical dimension; we shall call such quantity a

scalar. The measurement of other physical quantities consists instead in a triplet of numbers with associated physical dimension; we shall call such quantity a **vector**.

! What's scalar or vector depends on the theory

Scalar and *vector* have very specific and slightly different meanings in different theories, so don't take the definitions used here as universal. For example, in these notes and in Newtonian mechanics we call *energy density* a scalar, but in general relativity it cannot be called a scalar.

1.4 Informal tips about units and maths

Importance of units

Units are very important and must always be written for several reasons.

First, a number without units doesn't tell us anything. If I tell you "the place is at a distance 100 from here", you have no idea how far the place is. "100" what? 100 metres? 100 kilometres? These are completely different distances.

Second, units give us useful information about mathematical objects and their physical relationships and measurement. If you see the expression "3 m/s", then there's a strong possibility that that's a velocity. If you see the expression "5 J/m²", then you have a hint that it could be measured by measuring an energy and then dividing by an area.

Third, because of the previous reason, keeping track of units often allows us to quickly catch errors in solving a physical problem.

Variables and units

When a physical quantity is denoted by a symbol or variable, keep in mind that a unit is "contained" in the symbol, so to speak. For example if the variable t denotes a time, then it includes some time unit, say seconds. This becomes apparent when we write the value of the symbol, for instance " $t = 120 \text{ s}$ ". The unit is not predetermined, but it must correspond to the dimension of that quantity. We could for instance write " $t = 2 \text{ min}$ " instead; the two expressions are completely equivalent.

This fact must be kept in mind when combining symbols. For example, if d is a distance and t is a time, then writing $v = d/t$ tells us that v is a

velocity, and it has appropriate units that come from d and t , for instance m/s.

Units otherwise behave just like *literal constants* for all mathematical purposes, just like the letter ‘ a ’ in the expression ‘ $a x$ ’. This is why they can, for instance, be simplified:

$$3 \text{ mol/s} \cdot 5 \text{ s} = 3 \frac{\text{mol}}{\text{s}} \cdot 5 \text{ s} = 15 \text{ mol} .$$

Mathematical functions and units

Particular care must be taken with trigonometric and exponential functions, such as $\sin()$, $\cos()$, $\tan()$, $\exp()$, $\log()$; **these functions only admit a dimensionless argument** (which for the trigonometric ones corresponds to *radians*). So there cannot be units like ‘s’ or ‘m’ within these functions: we must make sure that any units present within cancel out.

This makes sense, because we wouldn’t know how to interpret the argument otherwise. Suppose you read “ $\cos(60 \text{ s})$ ” somewhere: how much is that? If we say “just discard the unit”, we would have

$$\cos(60 \text{ s}) \stackrel{?}{=} \cos(60) \approx -0.95$$

but wait: $60 \text{ s} \equiv 1 \text{ min}$, so we could equivalently write “ $\cos(1 \text{ min})$ ”. Then, according to the hypothetical rule “just discard the unit”, we would have

$$\cos(60 \text{ s}) \equiv \cos(1 \text{ min}) \stackrel{?}{=} \cos(1) \approx +0.54$$

a completely different result!

For this reason an expression like ‘ $\cos(t)$ ’, with t denoting time, doesn’t make sense: there’s a time unit in the argument of $\cos()$. If we want to express an oscillation with time, we must write instead something like

$$\cos\left(\frac{t}{T}\right)$$

where T is the period of the oscillation, a symbol which also includes a time unit, which simplifies with the one in t . If the period of the oscillation is $T = 1 \text{ s}$ then we can also simply write

$$\cos(t/\text{s})$$

This expression is now unambiguous: suppose that $t = 60 \text{ s} \equiv 1 \text{ min}$, then

$$\begin{aligned} \cos(t/\text{s}) &= \cos(60 \text{ s}/\text{s}) = \cos(60) \approx -0.95 \\ &= \cos(1 \text{ min}/\text{s}) = \cos(1 \cdot 60 \text{ s}/\text{s}) = \cos(60) \approx -0.95 \end{aligned}$$

Also remember that **the results of trigonometric and exponential functions are dimensionless numbers**, so an expression like ‘ $3 \cos(\dots)$ ’ denotes a pure number, with no units. If you want to express that the result is a length, the appropriate units must appear. We can for instance write

$$L \cos(\dots)$$

where L is a length, and therefore includes some kind of length unit such as ‘m’. If this length is, say, $L = 2 \text{ m}$ we can also simply write

$$2 \cos(\dots) \text{ m}$$

Units and derivatives

When we follow the rules above, all other mathematical operations automatically take care of everything. The derivative, for instance, is calculated in the usual way, treating any visible units as literal constants. Let’s see a concrete example. This expression

$$x(t) = 2 \cos(t/\text{s}) \text{ m}$$

says that the position of some object oscillates with time, between the values -2 m and $+2 \text{ m}$. When $t = 0 \text{ s}$, the position is $x = +2 \text{ m}$. The position $x = -2 \text{ m}$ is reached when the argument of $\cos()$ is π , that is

$$t/\text{s} = \pi \quad \Rightarrow \quad t \approx 3.14 \text{ s}.$$

The velocity of the object (see § 2.4) is given by the derivative of this expression with respect to t . Let’s calculate it treating all unit symbols as literal constants:

$$\frac{dx(t)}{dt} = \frac{d}{dt} \left(2 \cos(t/\text{s}) \text{ m} \right) = 2 \underbrace{\left[-\sin(t/\text{s}) \cdot \frac{1}{\text{s}} \right]}_{\text{chain rule}} \text{ m} = -2 \sin(t/\text{s}) \text{ m/s}$$

and you see that the correct units for velocity have automatically appeared.

1.5 What is “fundamental” physics?

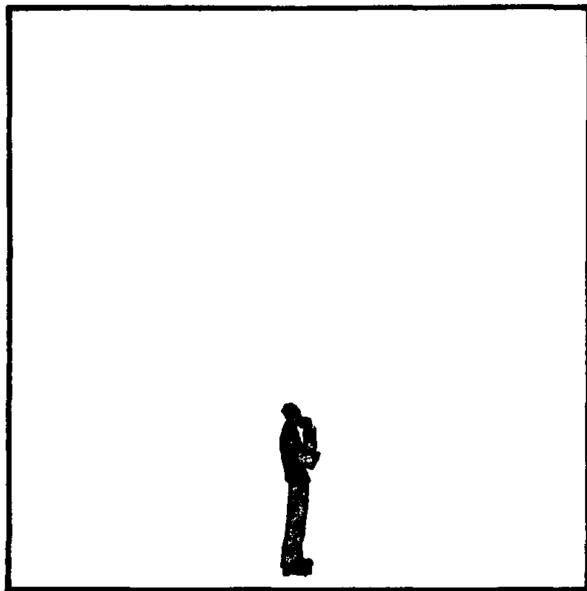
But what’s the “ultimate” goal of physics? What’s “fundamental” physics? The answer to this question is again subjective – also in this case physics lets you express your proclivities and personality. In the history of physics one can probably identify two main conceptions of “fundamental” physics.

For some physicists it is about finding the ultimate building blocks, so that one day we can say “... and these are the constituents, and they obey these equations”. The history of physics seems to show that this goal is overturned every few generations. And yet every generation says “Now we almost have the complete picture – it’s right behind the corner. It’s true that previous generations thought they almost had it, and turned out to be wrong. But *this time* is different, this time we have the real deal!”. The theoretical and particle physicist [Geffrey Chew³](#) depicted this situation as in fig. 1.1. For this reason some physicists are a little sceptical about this goal; maybe it’s a never-ending structure, with surprises at every deeper look.

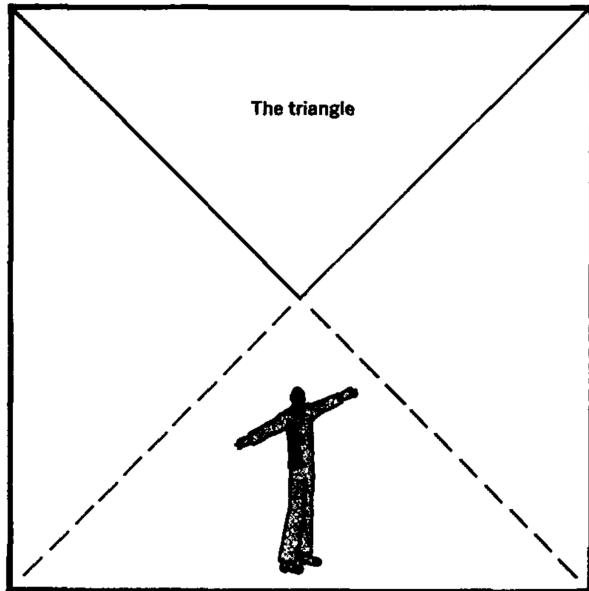
So for other physicists fundamental physics is about finding some point of view or mathematical structure that is rich enough to make useful predictions, and yet flexible enough to accommodate any new patterns or objects that we might discover. In a manner of speaking, it is about finding “patterns of patterns” or “laws about physical laws”.

The two conceptions above are not mutually exclusive, and both are always pursued, even if time-changing fashions may emphasize the one or the other.

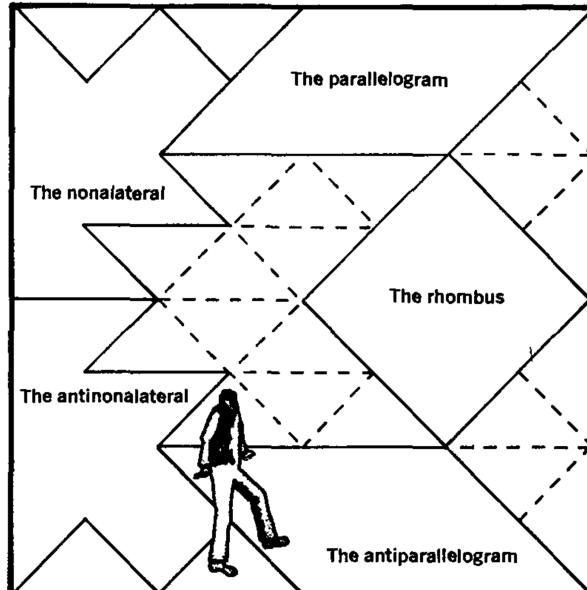
In these notes we take a point of view slightly closer to the second conception. This will also be reflected in the main division between physical laws that we’ll draw in ch. 5.



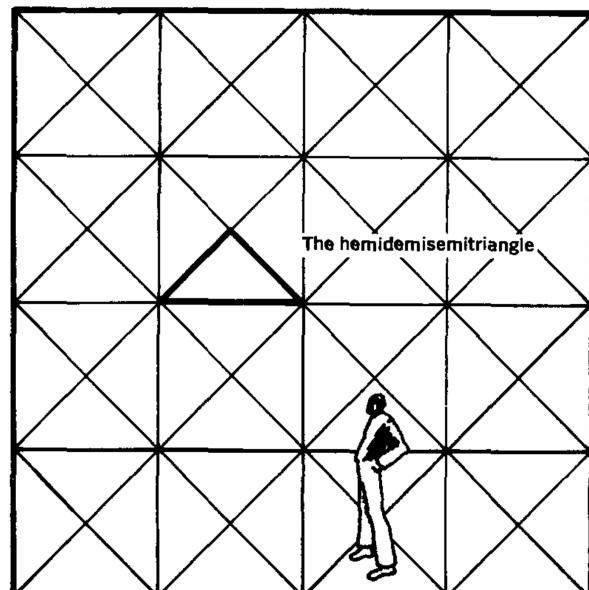
10 000 BC. The inhabitants of the paper square have no conception of the true nature of the universe they inhabit.



1900 AD. Physicists of the square discover a basic subdivision of their universe. They call it the "triangle" and consider it to be the fundamental building block of the universe.

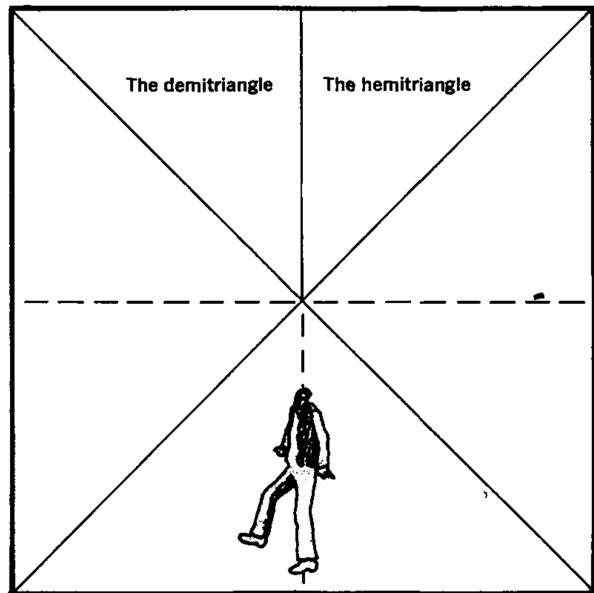


1960 AD. Physicists' conception of their universe is further clouded by new discoveries: the rhombus, the parallelogram, the antiparallelogram, the nonalateral and many others. It is unclear what these discoveries signify.

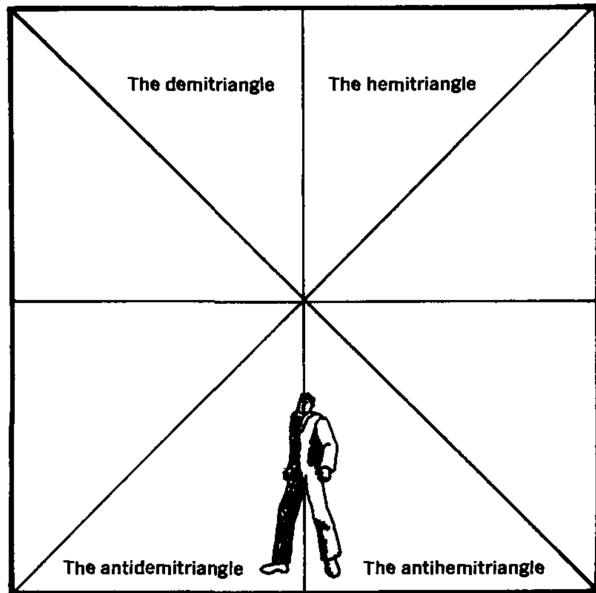


1970 AD. A new configuration, the "hemidemisemitriangle," is hypothesized, out of which all known configurations of the universe can be constructed. The hemidemisemitriangle is thought to be the fundamental building block of the universe.

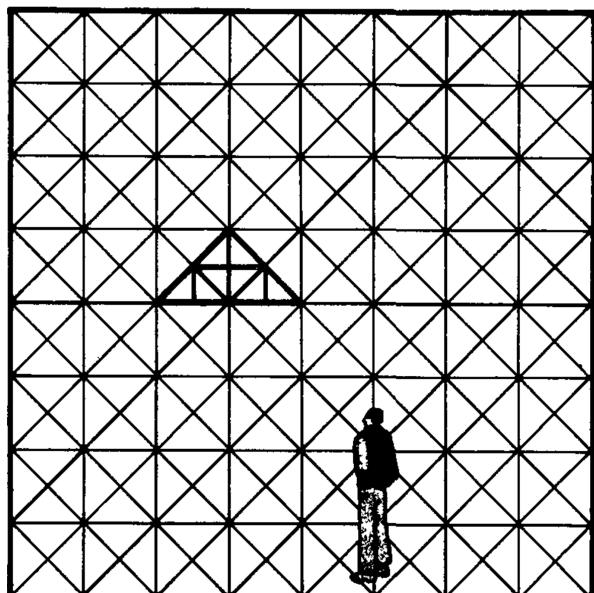
Figure 1.1 (Continues on p. 17) *The progress of "fundamental" physics*, from Chew 1970 as reproduced in Truesdell 1987



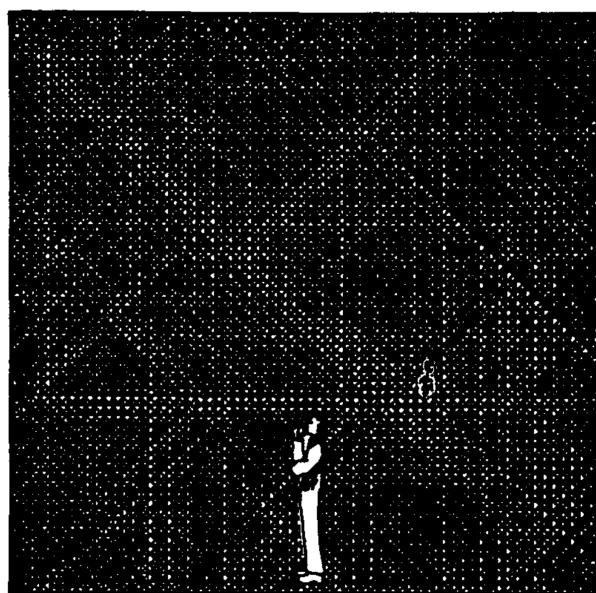
1930 AD. Physicists discover that the triangle can be split. Its parts are termed the "hemitriangle" and the "demitriangle." These are thought to be the fundamental building blocks of the universe.



1950 AD. Mirror images of the hemitriangle and the demitriangle are discovered. These are termed "antihemimtangle" and "antidemimtangle."



1975 AD. The hemidemisemimtangle is discovered. The following year the hemidemisemimtangle is split.



2000 AD. The inhabitants of this paper square have no conception of the true nature of the universe they inhabit.

URLs for chapter 1

1. <https://www.nist.gov/pml/owm/metric-si/si-units>
2. <https://www.nist.gov/pml/special-publication-811>
3. <https://www.physics.lbl.gov/rememberinggeoffreychew/>

Chapter 2

Building blocks: Time and space

If we want to describe the motion of a material point, we give the values of its coordinates as a function of time. However, we should keep in mind that for such a mathematical description to have physical meaning, we first have to clarify what is to be understood here by “time”. We have to bear in mind that all our propositions involving time are always propositions about simultaneous events. If, for example, I say that “the train arrives here at 7 o’clock”, that means, more or less, “the pointing of the small hand of my clock to 7 and the arrival of the train are simultaneous events”.

A. Einstein 1905a

2.1 Time

Time is a primitive quantity. We understand the notion of time intuitively, even if it’s difficult to explain (that’s why it’s taken as primitive). In 1905, with the theory of relativity, part of our everyday intuition about this notion was seriously shaken. For many years afterwards our old intuition could still be used in practice and in applications. But the new, correct intuition is becoming more and more important in everyday life and technologies. For example, GPS navigation – which we use everyday from leisure activities such as hiking or sightseeing to more critical ones such as aeroplane landing – critically depends on the correct notion and intuition of time.

Let’s see how our traditional intuition goes astray with a concrete experiment. Here’s Alice, Bob, and Charlie. They have extremely precise

clocks built in exactly the same way. They stay very close to one another and synchronize their clocks. Still keeping close, they go around, maybe on an aeroplane or space ship, and all the time they check their clocks. They notice that their clocks stay perfectly synchronized all the time, no matter where they go and what they do.

At some point they separate, each one going around independently. One of them might stay in place, another might take a helicopter, and another might go for a trip on Mars and back.

Alice and Bob at some point meet again, and compare their clocks. They see that their clocks aren't synchronized anymore; the difference could be as small as microseconds, or as large as years. In fact, if this time discrepancy is large, they would notice that they themselves have aged differently; so time discrepancy doesn't affect the clocks only. Let's say for concreteness that Alice's clock is ahead of Bob's, or equivalently that Bob's is behind Alice's. Note the following aspects:

First, neither Alice or Bob can say "my clock was wrong": neither has noticed anything strange about the "passage of time".

Second, they might wonder what's the time on Charlie's clock. But Charlie is at some distance away. They could decide to contact Charlie via radio, say, and ask "what shows your clock right *now*?". But they would notice that there's a delay, even if extremely small, in the radio transmission; so it's unclear to what time would Charlie's answer apply. If we say "let's account for the radio-signal speed", we see that there's a logical problem: speed is distance divided by time, and here we have a problem in exactly determining what's the "correct" time! So we would be reasoning in circles. Besides, even neglecting these difficulties, Charlie's answer could reveal a time that completely different from Alice's and from Bob's – it could be years ahead or behind both of theirs!

Third, if they now stay together, they will see that their clocks remain exactly synchronized, besides the discrepancy they noticed when they met. This discrepancy doesn't increase or decrease. They may even retrace together Alice's and Bob's previous trips; their clocks still remain synchronized.

The experience just described will occur again any time two or more of them meet. There could be a hundred observers like Alice, Bob, Charlie, initially at the same place and synchronized. Whenever two or more of them meet after having been separated, they will notice discrepancies in their clocks. But their clocks will have exactly the same time lapses as long as they stay together. See the illustration in fig. 2.1

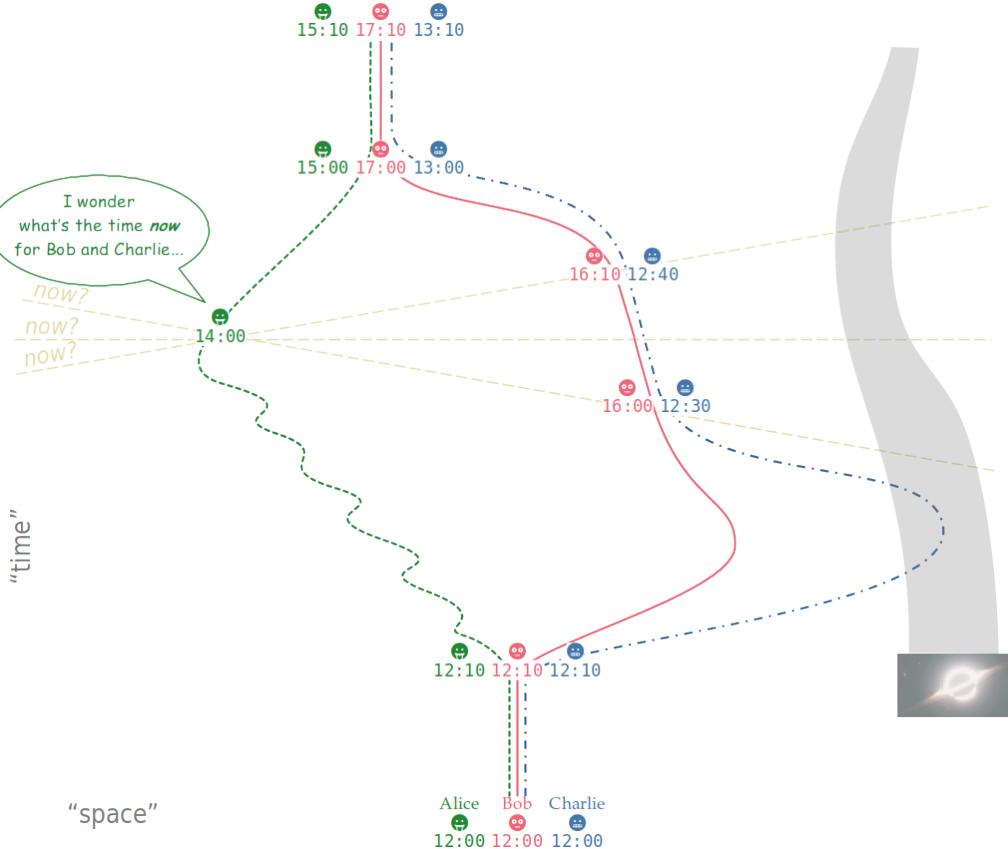


Figure 2.1 Illustration of the experiences of Alice (dashed 😎), Bob (solid 😊), Charlie (dot-dashed 😌) with time. The page represents a two-dimensional spacetime, and is followed from bottom to top.

Bottom: Alice, Bob, Charlie stay close and observe their clocks are perfectly synchronized from 12:00 to 12:10, then they separate.

Right: Charlie visits a region near a strong mass-energy source. Upon meeting again with Bob, the two notice their clocks differ: 16:00 for Bob, 12:30 for Charlie. Yet this clock difference stays the same while they travel together for 10 min.

Left: Alice wanders around travelling at high speed with respect to the fixed stars. At some point she wonders what's the time "now" for Bob and Charlie. Obviously this question doesn't make sense, because (1) when Bob and Charlie are together their clocks differ – impossible to say what's "the" time at their position; (2) not clear which instant in Bob & Charlie's trajectory should be considered as "now" (yellow dashed lines).

Top: When all three meet again, their clocks have completely different readings; and they themselves have aged differently. But their clocks run again at the same rate as long as they stay close.

Consider for a moment an imaginary world in which these experiments had given a different kind of result. According to Newtonian mechanics, whenever two or more initially synchronized observers like Alice, Bob, Charlie had met, their clocks would have always shown exactly the same time. If one year, 23 days, 8 hours, 9 minutes, and 3.045 399 283 240 992 663 02 seconds have passed for you since you last met Alice, you'd see that exactly the same amount of time has passed for her when you two meet again. If you think about it, in this case it would have been somewhat natural to think "right now, the clocks of far-away Alice, Bob, Charlie must show the same time as mine" (even though you have no real experimental way of confirming that).

But that's an imaginary world. In our world is the more complicated situation described initially that holds. Only one conclusion can be drawn from these experimental results: **Time is not some sort of universal quantity. It is, so to speak, "local" to a person or clock, or to a group of persons or clocks that stick together.** This also means that *it doesn't make sense to ask questions like "what can be the time for far-away Charlie, right now?"*

The time measured by a specific observer is called the **proper time** of that observer. Luckily we know more about how the proper times of separated observers can differ when they meet again. It turns out – according to our current understanding – that the time differences depend, roughly speaking, on how fast the observers are moving with respect to one another and to matter around the universe, and on how much energy is contained in the regions they travel. The general theory of relativity gives us the equations determining any such proper-time differences.

The situation depicted in the experiments above is real. It can be measured, for example, comparing initially synchronized clocks that have been put in aeroplanes flying in different directions. Most importantly, it affects everyday relevant technologies such as the Global Positioning System. Formulae from general relativity appear in your phone's GPS software; see for instance § 20.3.3.3 of the Interface Control Document IS-GPS-200 at <https://www.gps.gov/technical/icwg/>. It must also be taken into account in the establishment and synchronization of time in our everyday equipments:

International Atomic Time (TAI) is based on more than 250 atomic clocks distributed worldwide that provide its stability, whereas a small number of primary frequency standards provide its accuracy. Universal Coordinated Time, which is the basis of all legal time scales, is derived from TAI. To

"In 1976, the International Astronomical Union introduced relativistic concepts of time and the transformations between various time scales and reference systems. [...] Now [...] it is necessary to base all astrometry, reference systems, ephemerides, and observational reduction procedures on consistent relativistic grounds. This means that relativity must be accepted in its entirety, and that concepts, as well as practical problems, must be approached from a relativistic point of view."

Kovalevsky & Seidelmann 2004

allow the construction of TAI and the general dissemination of time, clocks separated by thousands of kilometres must be compared and synchronized. [...] The achieved performances of atomic clocks and time transfer techniques imply that the definition of time scales and the clock comparison procedures must be considered within the framework of general relativity.

(Petit & Wolf 2005)

In most everyday situations for us, who live on or nearby Earth and move at speeds much smaller than c with respect to one another, the discrepancies between our proper times are so small that cannot be measured with ordinary clocks or with our internal clocks. Consider a person walking 10 m away from you and then immediately walking back to you, at 1 m/s. The time elapsed for you will be 20 s, but for that person will be 19.999 999 999 999 999 889 s, a difference of 10^{-16} s, which is the error of an atomic clock. If human beings still exist in some decades or centuries, with space travel they will probably have to deal more and more with proper-time discrepancies also in everyday life.

For the most part of the rest of these notes, we won't need to deal with differences in proper time. But I recommend that you keep present how time really works, and that these small time discrepancies exist and occur all the time along your *worldline*.

Time has physical dimension of time and we shall for the most part measure it using the unit *second*, symbol 's'.

2.2 Space

Together with the notion of time, also the notion of space loses some of its traditional intuition. Several observers in motion with respect to one another will generally disagree on the dimensions of an approximately rigid object in their vicinity. For objects that are far away from an observer, the very notion of "distance" becomes tricky has different and non-equivalent definitions; one must be very careful on which definition is being used.

We shall not delve further into these peculiarities of time and space. Keep simply in mind that phenomena happen in **spacetime**, and that there's no way to attribute a universal time, nor a universal position in space, to a physical event. There is one absolute: whoever locally measures the speed of light, will find the value

$$c := 299\,792\,458 \text{ m/s} \quad (2.1)$$

"The plot for Cesium [...] characterizes the best orbiting clocks in the GPS system. What this means is that after initializing a Cesium clock, and leaving it alone for a day, it should be correct to within [...] 4 nanoseconds. Relativistic effects are huge compared to this."

Ashby 2003

This value is exact by definition, and serves as a way to define a local notion of space and distance.

Space has physical dimension of length and we shall for the most part measure it using the unit *metre*, symbol ‘m’.

2.3 Coordinate systems and events

It is necessary to have a way for distinguishing physical events and phenomena and locating them in spacetime. This is achieved through a **coordinate system**.

A coordinate system assigns four numerical labels to every point in spacetime. Often these labels have some kind of physical meaning – such as the proper time elapsed for a specific clock, or the distance from some event as measured by a specific observer – but they don’t need to.

A coordinate system also solves the problems coming from proper-time and space discrepancies among different observers. We can assign to every physical event a *coordinate time* and a *coordinate spatial position*, which are the same for all observers, because decided by agreement. Coordinate time doesn’t have a strict physical meaning, and will generally be different from the proper times registered by different observers. It can nevertheless be used for “doing physics”, and it is the time we shall most often use in our equations. A coordinate time commonly used for Earth-physics purposes is [Universal Coordinated Time \(UTC\)](#)¹. The clock on your phone, and on devices that get synchronized via internet, shows UTC, not your proper time. An observer on Earth at 0 m over sea level, and not moving, measures a proper time exactly equal to UTC (besides small variation coming from the movements of Solar System bodies). Observers at other altitudes or moving with respect to Earth’s surface notice that their proper times are slightly different from UTC.

Up to now we have often used the word ‘event’, informally taking its meaning for granted. Let’s be more precise: we call **event** or **spacetime point** an extremely small region of space – a point – that only lasts for a very small lapse of time – an instant. The name ‘event’ is used because typically we approximately identify such a point and instant by means of a physical phenomenon of limited spatial extension and short duration, such as the collision of two subatomic particles. The sudden burst of a very small soap bubble can be considered as an event in some circumstances; but something like “a tennis ball” cannot be considered as an event, mainly because a tennis ball exists for quite a long time, not just for a short instant.



From a four-dimensional spacetime point of view, a tennis ball could be characterized as a line: a **worldline**.

We shall often denote the four coordinates of a coordinate system by the letters

$$(t, x, y, z)$$

where t is a coordinate time, usually UTC, and (x, y, z) determine a spatial position. The triplet of spatial coordinates is often denoted by the vector \mathbf{r} :

$$\mathbf{r} := (x, y, z).$$

It is always important to specify how the coordinate system you're using is defined. The definition of the spatial coordinates (x, y, z) is typically different from problem to problem. We shall typically use coordinates that form $\frac{\pi}{2}$ rad $\equiv 90^\circ$ angles with one another; but their directions and their origin – that is, where they have value $x = y = z = 0$ m – always depend on the problem, so make sure you always specify them.

Whenever we speak of a “region of space” or of a “surface in space”, we mean a 3D or 2D region at some specific coordinate time t .

Some physical phenomena happen along a line, in one dimension. In this case we can omit two of the spatial coordinates, assuming they have some constant, unimportant values. In these cases we can simply write, for instance, (t, x) as our coordinates.

2.4 Velocity and acceleration

In some situations the spatial coordinates $\mathbf{r} = (x, y, z)$ may turn out to be functions of the time coordinate t ; the typical example is when we describe how the spatial position of a small body changes with coordinate time. We can write this functional dependence in different ways, for instance

$$\mathbf{r}(t) \quad \text{or} \quad (x(t), y(t), z(t)).$$

So \mathbf{r} is a vector function of time, which simply means that we have a collection of three functions of time.

If we take the derivative of each coordinate with respect to the time t , we obtain the **coordinate velocity**

$$\mathbf{v}(t) := \frac{d}{dt} \mathbf{r}(t) = \left(\frac{d}{dt} x(t), \frac{d}{dt} y(t), \frac{d}{dt} z(t) \right)$$

which is also a vector.

Dot-notation for time derivative

The derivative of some quantity with respect to coordinate time is often denoted by a **dot** over the quantity. So we can also write

$$\mathbf{v}(t) = \dot{\mathbf{r}}(t) = (\dot{x}(t), \dot{y}(t), \dot{z}(t))$$

The coordinate velocity is usually different from the *physical velocity*, which an observer would measure using proper time and space, for instance using bouncing light rays. In many everyday situations the difference between coordinate and physical velocity is so small that it can be neglected. But in situations involving subatomic particles at high speed, for example, one must take into account that the two velocities are different.

Taking the time derivative once more we obtain the coordinate acceleration, also a vector:

$$\mathbf{a}(t) := \frac{d}{dt} \mathbf{v}(t) = \frac{d^2}{dt^2} \mathbf{r}(t) = \left(\frac{d^2}{dt^2} x(t), \frac{d^2}{dt^2} y(t), \frac{d^2}{dt^2} z(t) \right)$$

which is also a vector.

Acceleration in relativity theory

In relativity theory, acceleration acquires a special physical significance, because it includes the effect of gravity, and its calculation does not involve just a time derivative. For instance, let's say that you are standing still on the ground, and let's use a coordinate system where x points in front of view, y to your left, and z points upwards. Then your coordinate velocity is $\mathbf{v} = (0, 0, 0)$ m/s also according to relativity theory. But the spatial part of your acceleration is approximately $(0, 0, 9.8)$ m/s², not zero!

The definitions and values of acceleration according to relativity theory and according to Newtonian mechanics are therefore quite different even in everyday situations. In these notes we'll mean the Newtonian definition of acceleration, unless stated otherwise.

URLs for chapter 2

1. [https://www.nist.gov/pml/time-and-frequency-division/time-realization/utc
nist-time-scale-0/](https://www.nist.gov/pml/time-and-frequency-division/time-realization/utc-nist-time-scale-0/)

Chapter 3

Building blocks: “Stuff”

if the skill of the mathematician has enabled the experimentalist to see that the quantities which he has measured are connected by necessary relations, the discoveries of physics have revealed to the mathematician new forms of quantities which he could never have imagined for himself.

J. Clerk Maxwell 1870

3.1 Seven primitive quantities with three basic properties

Besides time and space, our physics formalism includes around seven more quantities that we take as primitive:

Seven primitive quantities

matter
electric charge
magnetic flux
energy
momentum
angular momentum
entropy

Technically they are called *fields*, for reasons we shall shortly see.

Recall that primitive quantities cannot be defined: we can only try to understand them intuitively, for example through their properties. Some of the seven primitive quantities are easier to grasp intuitively than others.

But all seven primitive quantities have three basic properties in common. First of all, of each we can ask, or rather, measure:

📘 The two basic measurements that can be made on the seven quantities

1. How much of this quantity is in a particular three-dimensional region of space at a particular time instant?
2. How much of this quantity flows through a particular two-dimensional surface during a particular time lapse?

We can ask these questions of any region of space and any time lapse, and the surface in the second question can be moving and deforming. The results of the two measurements above are numbers, which in general can be positive or negative, for scalar quantities; or vectors for vector quantities.

Often we consider the second question, about the flow through a surface, in the case of a very short lapse of time, and divide the total flow by that time lapse. So we have an alternative form of the second measurement:

📘 Flux of a substance through a surface

- 2b. How much of this quantity is flowing through a particular two-dimensional surface per unit time, at a particular time instant?

This is called the **flux** of the quantity through that surface.

The third property common to the seven quantities regards the two measurements above.

📘 Extensivity of the seven quantities

If we consider two or more separate volumes, the amount of quantity in the total volume is equal to the sum of the amounts in the separate volumes.

Analogously for the flux through separate surfaces.

We say that each of the seven quantities is an **extensive** quantity

The basic measurements above can't in general be made, and don't even make sense, for some other quantities. For instance, we cannot ask "what's the total amount of *temperature* in this region?", or "how much velocity is flowing through this surface?".

Thanks to the three properties above, each of the seven quantities can be intuitively visualized as some kind of “stuff” that can be present at each spacetime point, fills regions of space, and flows through surfaces. This visualization is useful, but also comes with some warnings which we’ll discuss later.

What’s remarkable about matter, electric charge, magnetic flux, energy, momentum, angular momentum, and entropy, is that *they are common to all our main physical theories*, approximate or not: from Newtonian mechanics to general relativity and quantum theory; from subatomic scales to cosmological scales. The physical meaning and mathematical characterization of these quantities can be slightly different depending on the physical theory and scale. For example, in quantum theory they are mathematically represented by so-called operators rather than functions; and at molecular scales entropy has a meaning connected with probability theory. Yet, these seven quantities are really universal to our present way of doing physics and of describing physical phenomena around and within us.

Let us make a first acquaintance with these seven quantities. The discussion that follows is meant as an introduction; we shall repeat and say more about each quantity in later chapters.

3.2 Matter

Matter is a *scalar* quantity, with SI dimension [amount of substance](#)¹, and measured in units of [moles \(mol\)](#)². Its flux is therefore measured in units of *moles per second* (mol/s). In statistical mechanics and particle physics, matter is often simply counted and so measured in dimensionless units, rather than moles.

Matter is probably the easiest quantity to grasp intuitively; what we ordinarily call “stuff” is matter. It is usually classified into several kinds; the classification depends on the physical phenomena and physical theory one works with, and is related to whether the kinds of matter can be considered separately conserved, as we’ll discuss in § 6.1.

In everyday phenomena not involving [radioactivity](#)³ or [nuclear energy](#)⁴, the different kinds of matter approximately correspond to the non-radioactive [chemical elements](#)⁵: [hydrogen](#)⁶, [helium](#)⁷, [lithium](#)⁸, and so on. Note that some common everyday devices, such as smoke detectors, do involve radioactivity.

In phenomena involving radioactivity or nuclear energy, the different kinds of matter correspond approximately to *baryons*⁹, such as protons and neutrons; and *leptons*¹⁰, such as electrons. In particle physics, even more subtle classifications of matter are made, into kinds that seem to be conserved, such as electronic-leptons, muonic-leptons, and others. This kind of research is still open, but it seems that the total amount of baryonic and leptonic matter, independently of the kinds into which it can be classified, is always conserved. Note that we're using the term 'matter' in a sense that includes anti-matter, such as positrons.

The total amount of matter in a region *can be negative*. "Negative matter" is what's usually called *anti-matter*¹¹. Anti-matter appears in small amounts in everyday life, for example in connection with common radioactivity processes. It is also created and used in medicine, in *positron-emission tomography (PET)*¹² scans.

In these notes we shall usually not consider distinctions between different kinds of matter, making some exceptions for discussions about chemical reactions and nuclear phenomena.

Matter: notation

The amount of matter in a region is usually denoted with N , and flux of matter with J . In chemistry we usually specify what kind of matter we are speaking about, writing for instance $N_{\text{Ca}} = 5.3 \text{ mol}$, to indicate an amount of 5.3 mol of *calcium*¹³ atoms.

Matter is different from mass or energy

It is important to clearly distinguish matter from *mass* or *energy*. Mass can be considered a property of matter, but the two are different. In nuclear reactions, for instance, the mass of some amount of matter may change, even if the amount of matter stays the same.

As far as we know, the total amount of energy associated with an amount of matter is always positive, whether the amount of matter is positive or negative (antimatter). This is the reason why antimatter "falls" just like positive matter, a fact that has been experimentally confirmed: see Anderson et al. 2023.

Exercise 3.1

According to statements on [symmetrymagazine.org](https://www.symmetrymagazine.org/)¹⁴ and [quantum-diaries.org](https://www.quantum-diaries.org/)¹⁵,

The average banana (rich in potassium) produces a positron roughly once every 75 minutes.

Unfortunately the original site where this statement was discussed, and the corresponding calculation made, seems not to exist anymore.

1. Do a little research and find out whether this statement is true.
2. From your research, approximately quantify the flux of positrons around an ordinary banana, expressing it in particles/s.



How many positrons do bananas produce?

3.3 Electric charge and magnetic flux

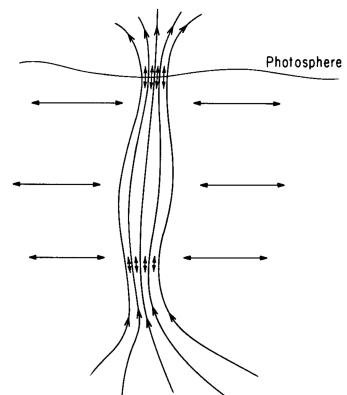
Electric charge is a *scalar* quantity, with SI dimension [electric charge](#)¹⁶ (equivalent to electric current×time), and measured in units of [coulombs \(C\)](#)¹⁷. Flux of electric charge is called *electric current*, and measured in units of [amperes \(A = C/s\)](#)¹⁸.

Electric charge is a quantity that is easily grasped in our everyday experience, and doesn't need much comments.

Magnetic flux is a *vector* quantity, with SI dimension [magnetic flux](#)¹⁹, and measured in units of [webers \(Wb\)](#)²⁰.

The electromagnetic field is most commonly represented by vectors associated to each point in space. But it can also be interpreted and visualized as a collection of moving, oriented tubes or lines, either closed or extending indefinitely; somewhat analogously to how we visualize matter and charge, as moving blobs or points, but with one more dimension. This interpretation goes back to Faraday (1846), Maxwell (1855), and later Dirac (1955) among others, and today is conveniently used in some fields such as [solar physics](#)²¹, for example to study [sunspots](#)²² (see Ryutova 2018). From this point of view, the *voltage* turns out to be the flux of the magnetic flux; it is indeed measured in *volts* (V), equivalent to webers/second.

Electromagnetism and this particular visualization of it are very fascinating topics, but we shall not discuss them in these notes.



"sketch of the magnetic lines of force in a magnetic filament extending up through the photosphere." Parker 1974a

3.4 Energy

Energy is a *scalar* quantity, with SI dimension energy, and measured in units of *joules* (J); its flux is measured in units of *watts* (W = J/s).

Equivalently we can speak of mass, with SI dimension mass, and measured in units of *kilograms* (kg)²³; its flux is measured in *kilograms per second* (kg/s).

The notion of energy is easily grasped today, and often we speak about it as something that “flows”, is “transported”, “converted”, “stored”, and similar visualizations. This intuition will be enough in these notes. The notion of *mass* is also very intuitive in our everyday life; we associate it with the difficulty in setting objects into motion, or with the weight of objects.

From Relativity Theory and experimentally we know that *energy and mass are the same quantity*, and in these notes we shall emphasize this experimental fact.

Why we cannot differentiate between energy and mass

Let's see some examples of why it is impossible to make a real distinction between energy and mass. The following examples have been simplified in some of their aspects, but their main point is valid.

"we are led to the more general conclusion: The mass of a body is a measure of its energy content; if the energy changes by L, the mass changes in the same sense by $L/9 \cdot 10^{20}$, if the energy is measured in ergs and the mass in grams." Einstein 1905b

Heated gas Imagine we have a box with a given amount of gas, say 1 mol of oxygen molecules. Using an extremely precise weighing scale, we observe that the mass of the gas is, say, exactly

$$0.031\,999\,540\,000\,000\,000 \text{ kg} .$$

Now we heat the gas, providing 60 J, while making sure that not a single molecule of oxygen gets in or out of the box. The temperature of the gas increases by around 3 K. We observe that the weight measured by the scale increases while we heat the gas, reaching the new value

$$0.031\,999\,540\,000\,000\,668 \text{ kg} .$$

Clearly the mass has increased, but no molecules were added! The additional mass is the 60 J that we provided to the gas. Energy has weight, energy is mass.

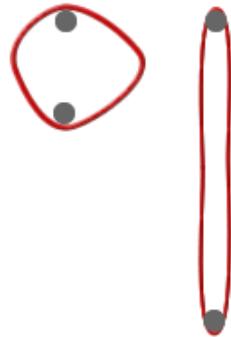
Stretched rubber band Take a common rubber band, and imagine again that we have an extremely precise weighing scale. The rubber band, unstretched, has a mass of exactly

0.000 500 000 000 000 000 000 kg .

Now we stretch the band a little. According to Newtonian mechanics it acquires ‘elastic energy’, which increases its initial internal energy by, say, 0.3 J. Now we weigh the rubber band again, while stretched, and observe a mass of approximately

0.000 500 000 000 000 003 338 kg .

The extremely small difference of around 3×10^{-18} kg from the initial mass is exactly the elastic energy that we provided by stretching. Energy has weight; energy is mass.



When we stretch a rubber band, its mass increases slightly – even if the amount of rubber remains exactly the same.

Now set the unstretched band in motion, so that according to Newtonian mechanics it acquires a kinetic energy of, say, 0.2 J. If we could weigh the band while in motion (but without moving the weighing scale), we would measure a mass of approximately

0.000 500 000 000 002 225 kg .

The small difference from the initial mass is the new kinetic energy of the band.

This case is actually connected with the example of the gas above. If we observed the gas at a molecular level, we would interpret the energy of 60 J provided to it as additional kinetic energy of its molecules. The increase in weight was exactly this additional kinetic energy.

Fission and atomic bombs The atomic bomb²⁵ is a dark example of the fact that mass is energy. In the case of nuclear fission, if we weigh the amount of nuclear material, say within a box, before and after fission, we observe that its mass has decreased. But we also observe that a great amount of energy has been released out of the box.



Hydrogen Bomb Test, 1954 (National Museum of Nuclear Science & History²⁴)

Electric heater As a final example consider a 1000 W electric heater, which is radiating 1000 J in one second. The heater is also losing around 0.000 000 000 01 kg of mass every second owing to this heat radiation – although it's also acquiring the same amount of mass as electromagnetic energy.

The equivalence between energy and mass is given by the famous formula $E = mc^2$, where c is the speed of light, eq. (2.1). In their respective units this gives

$$\begin{aligned} 1 \text{ kg} &= 89\,875\,517\,873\,681\,764 \text{ J} \quad (\text{exactly}) \\ 1 \text{ J} &\approx 0.000\,000\,000\,000\,000\,011\,126\,5 \text{ kg} \end{aligned}$$

To grasp these numbers, consider that the mass of the rubber band in the example above, 0.5 g, is comparable to the energy released by the [atomic bomb over Hiroshima](#)²⁶.

The practical distinction between mass and energy

From the examples above it is clear that the energy changes that we deal with and use in everyday situations are extremely small changes in mass – so small that we can't measure them within the precision of ordinary weighing scales. Vice versa, the masses that we ordinarily deal with are huge amounts of energy. If we spoke only of 'energy' or only of 'mass' in all situations, and used only one unit, either joules or kilograms, we would have to work with very impractical numbers. 'Mass' can be considered a convenient term for 'energy' when huge amounts of it are involved, concentrated in small regions of space, and ordinary small energy changes can be neglected. In these notes we shall often use the expressions 'energy-mass' and 'mass-energy' to remember that these two words denote the same thing.

Different 'forms' of energy

We often speak of different *forms* of energy-mass. The most important for us will be **internal energy**, **kinetic energy**, **gravitational potential energy**, to be discussed later; another important one is *electromagnetic energy*. We also speak of different forms of energy flux; the most important for us will be **heat** and **work**.

! 'Forms' of energy-mass are coordinate- and observer-dependent

Although the distinction of energy forms is often useful, it's important to keep in mind that it is not absolute: *it depends on the coordinate system, on the scale of observation, and on the theory used*.

For example, in one system of spatial coordinates we say that a volume has 11 J of internal energy and 2 J of kinetic energy, but in another system of spatial coordinates this corresponds to 13 J of internal energy and 0 J of kinetic energy. The same happens if we keep the initial system of coordinates but change our observation scale, for instance modelling a gas no longer as a uniform fluid but as a collection of molecules. The same is true of energy flux: what we call 'heat' on one observation scale is 'work' on a finer scale. In phenomena that involve matter and electromagnetic fields, any separation between 'energy of matter' and 'energy of electromagnetic field' is often arbitrary.

So whenever you hear about a distinction between forms of energy or energy fluxes, pay attention which coordinate system and which observation scale are being used.

Energy-mass is present wherever there's matter, charge, or magnetic flux, and vice versa. *Its specific relation with these quantities depends on the coordinate system* (§ 2.3) that we choose. It is extremely important to keep in mind this dependence. For instance, in the examples above we said that the unstretched rubber band in motion had a mass-energy of 0.000 500 000 000 000 002 kg; but if we choose a coordinate system in which the band is not moving, its mass-energy would be 0.000 500 000 000 000 000 kg. Owing to this coordinate dependence, and depending on whether matter or magnetic flux is present in a region, we often conveniently partition an amount of energy into different "forms" such as rest energy, internal energy, kinetic energy, gravitational potential energy, electromagnetic energy. This is only a conceptual division.

Also energy flux can be partitioned depending on the coordinate system and on whether there's also a flux of matter. This leads to the important distinction between two forms of energy flux: **heat** and **work**.

We shall see later how to define and calculate these different forms of energy and energy flux.

Energy-mass: notation

The amount of energy in a region is usually denoted with E , or with M if we describe it as mass. Internal energy is denoted U , kinetic energy E_k , potential energy E_p .

The total flux of energy will be denoted by Φ ; the flux in the form of heat, by Q ; and in the form of work, by W .

Exercise 3.2

In an hour, 14 people pass through a door. Taking the average human weight to be 62 kg (Walpole et al. 2012), what's the average *energy* flux, in J/s, through that door?

3.5 Momentum

Momentum, also called *linear momentum* to distinguish it from angular momentum, is a *vector* quantity. Its SI dimension and units can be written in several equivalent ways; we shall keep in mind especially these three:

$$\begin{array}{ccc} \text{force} \times \text{time} & \equiv & \text{mass} \times \text{length/time} \\ \text{N} \cdot \text{s} & \equiv & \text{kg} \cdot \text{m/s} \end{array} \quad \equiv \quad \begin{array}{c} \text{energy} \times \text{time/length} \\ \text{J} \cdot \text{s/m} \end{array}$$

Since it is a vector quantity, it is usually expressed with three numbers, typically the x -, y -, and z -components. In simplified problems where only one or two dimensions are relevant, only the relevant components are reported.

Momentum is a subtle quantity. Textbooks that focus on Newtonian mechanics *define* it as the product of the mass of a body times the body's velocity, usually written " $\mathbf{p} = m\mathbf{v}$ ". This relation, however, is only valid in some circumstances; it cannot be used as a general definition. First, it is true that where there's matter in motion there's also momentum; but this momentum is given by 'mass \times velocity' only as an approximation. Second, there can be momentum also where there is no matter; for instance where there is magnetic flux. If you prepare a sealed glass box which is completely empty of matter within (you create a vacuum), there still is momentum within the box if light or electromagnetic waves are present. We shall discuss later some technological uses of this fact.

It is therefore convenient to separate your idea of momentum from the idea of "bodies moving", and keep in mind that the latter idea is just a particular case of momentum. Yet, momentum is somehow associated with motion; motion of matter or motion of electromagnetic fields.

Momentum is energy flux

According to Relativity Theory, momentum *is* energy flux (such as heat), and energy flux *is* momentum. If we represent momentum with the symbol \mathbf{P} , and energy flux with the symbol \mathbf{Q} , the equivalence between

them is given by

$$\mathbf{Q} = \mathbf{P}c^2 \quad (3.1)$$

Compare this formula with $E = mc^2$. From this point of view, you can think of momentum as “energy in motion”. This is consistent with our discussion about mass-energy: since mass is energy, the Newtonian expression “ $m\mathbf{v}$ ” indicates energy in motion, or a flux of energy. On a sunny day, if you close your eyes and feel the Sun’s heat on your face, what you are feeling is actually a flow of momentum. And when you kick a ball, you’re setting a huge bundle of energy in motion, and that’s why the ball has acquired momentum.

Flux of momentum is what we call **force**. Force is a notion that we understand intuitively and that we can “feel” with our own bodies. In Newtonian mechanics it’s often taken as primitive, and is represented by a vector. But it can also be defined and interpreted as a *flux of momentum*. This point of view can be very useful and even more intuitive in some situations. We shall discuss it more in depth in § ∞ .

Momentum: notation

The amount of momentum in a region is usually denoted with \mathbf{P} . The flux of momentum is also called force and denoted with \mathbf{F} .

3.6 Angular momentum

Angular momentum, also called *moment of momentum* or *rotational momentum*, is a *vector* quantity. Its SI dimension and units can be written in several equivalent ways; we shall keep in mind especially these three:

$$\begin{array}{ccc} \text{force} \times \text{length} \times \text{time} & \equiv & \text{mass} \times \text{length}^2/\text{time} \\ \text{N} \cdot \text{m} \cdot \text{s} & & \text{kg} \cdot \text{m}^2/\text{s} \end{array} \equiv \text{energy} \times \text{time} \quad \text{J} \cdot \text{s}$$

It is usually expressed with three numbers, typically the x -, y -, and z -components.

Angular momentum is perhaps an even subtler quantity than momentum. Luckily there’s a way of calculating it in terms of momentum, so even if we don’t fully grasp it intuitively, we can still calculate it. If the momentum *in a small volume* is denoted by the vector $\mathbf{P} = (P_x, P_y, P_z)$

and the position vector by $\mathbf{r} = (x, y, z)$, then the angular momentum $\mathbf{L} = (L_x, L_y, L_z)$ with respect to the origin of coordinates is given by

$$\mathbf{L} = \mathbf{r} \times \mathbf{P}$$

or equivalently

$$\begin{cases} L_x = y P_z - z P_y \\ L_y = z P_x - x P_z \\ L_z = x P_y - y P_x \end{cases} \quad (3.2a)$$

Instead of calling the components " (L_x, L_y, L_z) ", we can also call them " (L_{yz}, L_{zx}, L_{xy}) ", as some books do. The latter names make the formulae above easier to remember:

$$\begin{cases} L_{yz} = y P_z - z P_y \\ L_{zx} = z P_x - x P_z \\ L_{xy} = x P_y - y P_x \end{cases} \quad (3.2b)$$

Choose whichever you prefer.

The formula above also suggests that we can somehow think of angular momentum as "energy in rotation". Just like momentum, angular momentum is something that is associated not only with ordinary bodies (matter), but also with charge and magnetic flux, and therefore electromagnetic radiation.

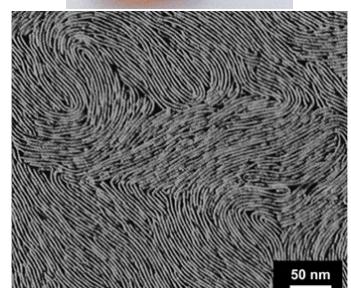
The flux of angular momentum is also called the *torque*, and bears a relation to the flux of momentum similar to the formula above. If the force – flux of momentum – is denoted by the vector $\mathbf{F} = (F_x, F_y, F_z)$ and the position vector by $\mathbf{r} = (x, y, z)$, then the flux of angular momentum, or torque, $\boldsymbol{\tau} = (\tau_x, \tau_y, \tau_z)$ with respect to the origin of coordinates is given by

$$\boldsymbol{\tau} = \mathbf{r} \times \mathbf{F}$$

or equivalently

$$\begin{cases} \tau_x = y F_z - z F_y \\ \tau_y = z F_x - x F_z \\ \tau_z = x F_y - y F_x \end{cases} \quad \text{or} \quad \begin{cases} \tau_{yz} = y F_z - z F_y \\ \tau_{zx} = z F_x - x F_z \\ \tau_{xy} = x F_y - y F_x \end{cases} \quad (3.3)$$

You may wonder: "Do we really need angular momentum? after all it just looks like something constructed from momentum". The answer is yes, we really need it, for two reasons. First, angular momentum obeys an important universal law which is independent from those obeyed by energy and by momentum (Truesdell 1968a tells some of the story of how this was



Some liquid polymers (**top**: Liquid Diethoxymethane Polysulfide) need to be described with a special kind of angular momentum.

discovered). Second, for some physical phenomena, for example involving liquid **polymers**²⁷, elementary particles, or electromagnetic radiation, the angular momentum includes an additional part, called *spin* or *intrinsic angular momentum*, that is not related to linear momentum. In the present notes we shall not use this more general kind of angular momentum.

📘 Angular momentum: notation

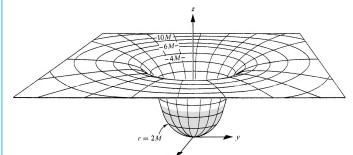
The amount of angular momentum in a region is usually denoted with **L** . The torque or flux of angular momentum is denoted with **τ** .

⌚ What are energy, momentum, angular momentum?

From the discussions and formulae above, it seems that energy-mass, momentum, angular momentum are quite closely related to one another.

Relativity Theory indeed shows that they are different aspects of one single geometric object, called *energy-momentum tensor*. They are like its “shadows”, that we can observe by looking at it from different points of view in time and space (almost literally: they are projections of that object). This is also why their values get intermixed if we change our system of coordinates.

General Relativity gives a new meaning to these quantities: they are *particular curvatures of spacetime*. They express how spacetime is curved in different directions. So whenever we measure, say, the energy or the momentum of some object or of some electromagnetic radiation, we are actually measuring how much that object or radiation is curving spacetime in a particular way. We shall discuss further connections between these three quantities and spacetime in §☒.



Energy, momentum, angular momentum are particular curvatures of spacetime.

3.7 Entropy

Entropy is a *scalar* quantity, with SI dimension energy/temperature, and measured in units of *joules/kelvins* (J/K)²⁸. From this definition it would seem that entropy is derived from temperature. However, although temperature is taken as primitive by the SI, the **definition of temperature**²⁹ actually depends on a fixed value of *Boltzmann's constant*³⁰, which has the dimension of entropy.

Entropy is probably the most difficult quantity to grasp intuitively. Many seemingly intuitive descriptions given in some textbooks are, unfortunately, unhelpful and even misleading. One particularly misleading intuition is that entropy would be a “measure of disorder”. Besides the fact that “disorder” is very vague and subjective, it turns out that some physical phenomena, for example with liquid crystals³¹, can be considered more “disordered”, and yet have *lower* entropy, than others. See also the example in the side figure. We shall discuss more about such phenomena later on.

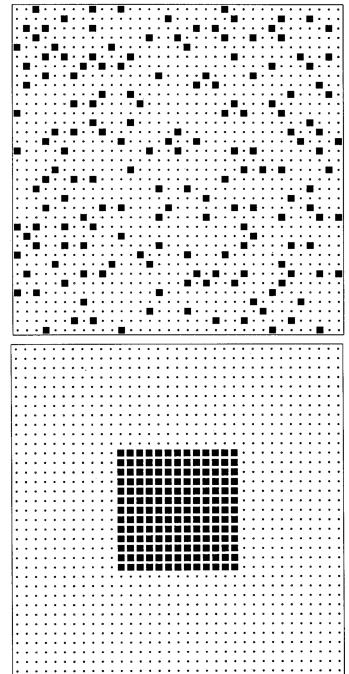
In these notes we shall rely on the idea that *entropy expresses a limit on the flux of energy into matter*. Said in simpler but more imprecise words, entropy is a bound on how fast we can heat up a body. We shall develop this idea further later.

One reason why entropy is difficult to grasp intuitively is that it has very different physical and mathematical aspects depending on the spatial scales and physical theory that we use to describe physical phenomena.

In many “continuum” phenomena, that is, phenomena where the molecular constitution of matter is not visible or not taken into account, entropy is treated as a “stuff-like” quantity similar to energy or electric charge. But there are difficulties also in this case. For some phenomena, for example involving non-elastic materials such as a simple paper clip, it is possible to introduce several entropies having different values – and not just because of a change in measuring scale – all of which can serve their purpose perfectly fine.

In molecular phenomena involving statistical mechanics, on the other hand, entropy is no longer a physical notion, but a *probabilistic* and *statistical* one, related to guesses and inferences that we make about the physical phenomenon. Yet from many points of view it has roles similar to those of the entropy used in continuum phenomena.

We shall see later that the physical laws for entropy have also a different status with respect to the laws for the other six main quantities: they are, so to speak, “laws about laws”.



Microscopic configurations of a lattice gas. **Top:** a configuration coming from a *low-entropy* state. **Bottom:** a configuration coming from a *high-entropy* state (Styer 2000).

Entropy: notation

The amount of entropy in a region is usually denoted with S . We shall see that the flux of entropy is tightly related to heat, and we won't need a special symbol for it.

3.8 Auxiliary quantities

Besides the seven principal quantities, other auxiliary quantities appear in some physical theories. Important examples are **temperature** and **metric**. Most auxiliary quantities don’t have the “stuff-like” properties of the seven principal quantities. For instance, we cannot ask “what’s the total amount of temperature in this region?”. We shall later discuss and use some auxiliary quantities, especially temperature.

The dimensions, units, and scalar or vector character of these quantities are summarized in table 3.1.

Quantity	SI Dimension	Unit
Time	time	<i>second</i> s
Length	length	<i>metre</i> m
Temperature	temperature	<ikelvin< i=""> K</ikelvin<>
Matter	amount of substance	<i>mole</i> mol
Electric charge	electric charge	<i>coulomb</i> C
Magnetic flux	magnetic flux	<i>weber</i> Wb
Energy	energy, mass	<i>joule</i> J, <i>kilogram</i> kg
Momentum	force · time, mass · length/time, energy · time/length	N · s, kg · m/s, J · s/m
Angular momentum	force · length · time, mass · length ² /time, energy · time	N · m · s, kg · m ² /s, J · s
Entropy	energy/temperature	J/K

Table 3.1 Dimensions and units of the main physical quantities used in these notes.
 Quantities in **boldface** are vectors, the others are scalars

URLs for chapter 3

1. <https://doi.org/10.1351/goldbook.A00297>
2. <https://doi.org/10.1351/goldbook.M03980>
3. <https://www.iaea.org/newscenter/news/what-are-radioactive-sources>
4. <https://www.iaea.org/newscenter/news/what-is-nuclear-energy-the-science-of-nuclear-power>
5. <https://www.ciaaw.org/radioactive-elements.htm>
6. <https://pubchem.ncbi.nlm.nih.gov/element/Hydrogen>
7. <https://pubchem.ncbi.nlm.nih.gov/element/Helium>
8. <https://pubchem.ncbi.nlm.nih.gov/element/Lithium>
9. <http://hyperphysics.phy-astr.gsu.edu/hbase/Particles/hadron.html#c6>
10. <http://hyperphysics.phy-astr.gsu.edu/hbase/Particles/lepton.html#c1>
11. <https://www.britannica.com/science/antimatter>
12. <https://www.britannica.com/topic/positron-emission-tomography>
13. <https://pubchem.ncbi.nlm.nih.gov/element/Calcium>
14. <https://www.symmetrymagazine.org/2009/07/23/antimatter-from-bananas>
15. <https://www.quantumdiaries.org/2009/07/21/positrons-from-bananas/>
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17. <https://doi.org/10.1351/goldbook.C01365>
18. <https://doi.org/10.1351/goldbook.A00300>
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23. <https://doi.org/10.1351/goldbook.K03391>
24. <https://nuclearmuseum.pastperfectonline.com/Archive/716477C1-5E7A-485C-8BE1-857919471563>
25. <https://www.britannica.com/science/nuclear-fission>
26. <https://www.britannica.com/story/atomic-bombing-of-hiroshima>
27. <https://www.britannica.com/science/polymer>
28. <https://doi.org/10.1351/goldbook.C01365>
29. <https://doi.org/10.1351/goldbook.K03374>
30. <https://doi.org/10.1351/goldbook.B00695>
31. <https://www.britannica.com/science/liquid-crystal>

Chapter 4

Fluxes and volume integrals

For the sake of persons of these different types, scientific truth should be presented in different forms, and should be regarded as equally scientific, whether it appears in the robust form and the vivid colouring of a physical illustration, or in the tenuity and paleness of a symbolical expression.

J. Clerk Maxwell 1870

In ch. 3 we said that the main seven quantities – matter, electric charge, magnetic flux, energy, momentum, angular momentum, and entropy – have three common properties related to their measurement:

- (1) we can measure the amount of quantity within a three-dimensional region, at a specific time instant
- (2a) we can measure the amount of quantity flowing through a two-dimensional surface during a time lapse...
- (2b) ... or alternatively we can measure the amount of quantity flowing through a two-dimensional surface *per unit time*, at a particular time instant
- (3) the amount in a volume consisting of separate volumes is equal to the total of the separate amounts; similarly for the flux through separate surfaces.

volume integral and flux

We can call measurement (1) the **volume integral** of the quantity, but we won't need to use this term very often.

Most important, we call measurement (2b) the **flux** of the quantity. Remember that in order to calculate the flux through a surface we need to specify how that surface is moving.

These two notions and measurements are very intuitive; that's also why it's convenient to base our physics upon them. In this chapter we straighten some details about their definition and also about our intuition.

4.1 Control volumes and control surfaces

A first question that comes to mind is: how are the volumes and surfaces chosen?

The choice of volume – and therefore of the surface bounding it – for a volume integral, and the choice of surface for a flux *are completely arbitrary*, and they *can be completely imaginary*. Since they are under our control, they are called **control volumes** and **control surfaces**.

For example, consider a classroom and the people in it. In your imagination you can divide the classroom into two halves, say the front and the rear half. You can then ask or measure simply by counting: 1. how many people are, right now, in the rear half; 2. how many people are crossing the imaginary division between the front and rear half during one minute, starting from now.

A control volume and a control surface don't need to be static: they can move and deform.

We can consider a control volume that exists for just one time instant, because the volume integral in a control volume *does not depend on the instantaneous motion of the volume*.

For a control surface the situation is slightly different. The flux through a control surface *depends also on the motion of the surface*. As a trivial example, consider a glass surface, and a person on one side of it, moving with a high velocity directed towards the surface. Will the person crash on the glass? We can't say for sure. The glass surface could be a glass wall in a building, which is not moving; in which case the person will likely crash on it. Or it could be the windscreens of a car, which is moving together with the person, who's the driver; in which case the person won't crash on it.

So we can't just imagine a surface that exists for one time instant: we need to imagine it for a very short time lapse, to be able to say how it's moving. If someone asks you what's the flux through a control surface at



The golf ball has a velocity directed towards the metal surface on the left. Will it hit the surface? We don't know unless we're told how the surface is moving.

a given instant, but they don't tell you what's the motion of the surface, then the flux is unknown.

The fact that we can choose control volumes and control surfaces arbitrarily gives us a lot of power in solving physics problems and in making predictions. Typically they are chosen so as to simplify the equations that describe the physical situation, and to focus on details of interest.

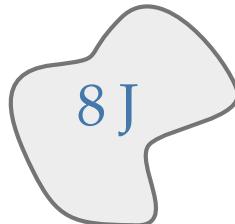
When we study some solid object, like a football, a rocket, or a planet, we typically choose a control volume that tightly encloses the object. When we study something flowing or moving, such as a fluid material or an electromagnetic field, we typically divide the space of interest into small control volumes and surfaces, constructing a *mesh*; this mesh can even be refined in regions that are of special interest.

4.2 Intuition and visualization for scalar quantities

Let's first consider a scalar quantity: matter, electric charge, energy, or entropy.

Volume integrals

A volume integral for, say, energy could be represented like this:

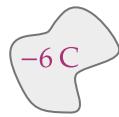


we have eliminated one spatial dimension for simplicity, considering the analogous two-dimensional idea. The volume is in light grey, delimited by a closed darker grey boundary, and we're indicating that the amount of energy within, the volume integral, is **8 J**.

As a visualization device, this representation can be useful. But let's straighten out some of its aspects:

- Recall that this is a snapshot at a given time instant. So there are 8 J of energy in the volume at that instant, but we don't know the situation earlier or later: there could be a different amount of energy, the region might be at a different position and have a different shape, or it might not even exist anymore.

- Recall that some scalar quantities, such as electric charge and in some situations matter (antimatter), can have negative amounts.
- We must not surmise that the amount of quantity is uniformly distributed within the volume. In fact there could be negative amounts of it in some subvolumes and positive in others.



A region with a negative amount of charge



Even if there is a zero amount of quantity in a volume, some subvolumes could have non-zero amounts: some positive and some negative, so that the total is zero.

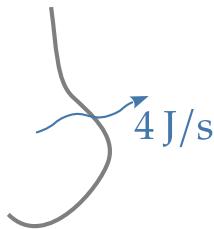
Fluxes

Suppose that at a given time instant we have this surface:



where we have again eliminated one spatial dimension for simplicity. We want to visualize and graphically represent the fact that there's a flux of **8 J/s** across this surface, from its left side to its right side.

One first visualization that comes to mind is this:



It conveys the correct information, but it's also misleading in several respects:

- This is a snapshot at a given time instant. An instant earlier or later, the surface might be at a different position and with a different shape. The **blue wavy arrow** might suggest that the surface is at rest, while something crosses it.

- A flux, properly speaking, says that after a very small time interval there's a lower amount of a quantity on one side of a surface, and a larger amount on the other. *But it doesn't say anything about a "left-to-right movement" of the quantity, or a right-to-left one.* In our example it could be that -4 J/s of energy "moved" from right to left; or even that $+7 \text{ J/s}$ moved from left to right, and -3 J/s from right to left, possibly across different parts of the surface. All these possibilities correspond to the same flux. Moreover, *the description of this situation depends on the coordinate system:* in one coordinate system the surface may be at rest, but in another it may be moving.

"Fechner [in 1845] supposed every current to consist in a streaming of electric charges, the vitreous charges travelling in one direction, and the resinous charges, equal to them in magnitude and number, travelling in the opposite direction with equal velocity." Whittaker 1951

Therefore we see that the flux of our example could be happening in several ways. Four equivalent ways are illustrated in fig. 4.1; combination of those are also possible.

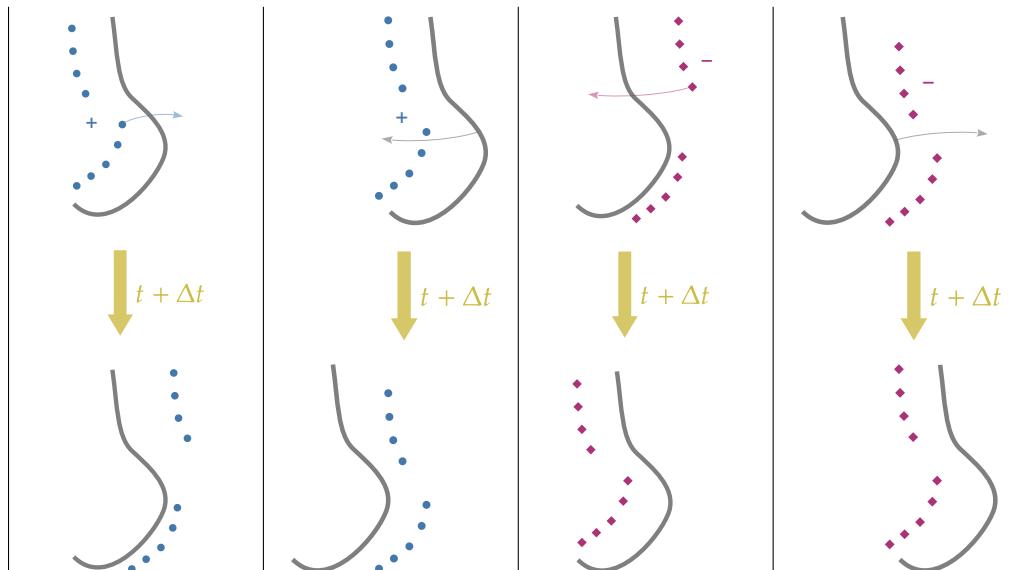
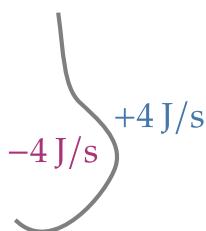


Figure 4.1 Four allusive depictions of the same flux

In view of these possibilities, a less misleading representation of the flux in our example could be as follows:



This picture says that there's an energy change of -4 J/s on the left side of

the surface, and of +4 J/s on the right. Note that **the amounts indicated on the two sides must have equal magnitude but opposite sign**, otherwise the picture doesn't make sense. The absence of arrows prevents unwarranted conclusions about "what's moving". If we know that the surface is moving in a particular way, then arrows could of course be added to indicate this movement.

Please feel free to use the graphical representation and mental visualization that you prefer. The ones above are just suggestions. The real goal of the discussion above was to make you aware of the subtleties regarding flux, so that you don't make wrong assumptions and jump to wrong conclusions when solving physics problems.

Exercise 4.1

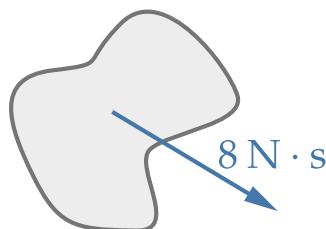
The volume integral of matter in a particular volume is equal to 36 mol. Can we conclude that the volume doesn't contain antimatter?

4.3 Intuition and visualization for vector quantities

Let's now consider the vector quantities momentum and angular momentum.

Volume integrals

A volume integral for, say, momentum could be represented as follows (we keep on simplifying our visualization to two dimensions):



Momentum is a vector quantity, so the total amount in the volume above is a vector. The picture shows the direction and orientation of this vector, and the magnitude of $8 \text{ N} \cdot \text{s}$ is explicitly reported.

This visual representation is useful, if we keep in mind remarks analogous to those made in the scalar case:



Remember that the *magnitude* of a vector is always positive, and that $\overleftarrow{\text{N}} = -1 \cdot \overrightarrow{\text{N}}$

- this is a time snapshot;
- the vector could have a different orientation or magnitude;
- different subvolumes could have amounts represented by different vectors; only the total vector is represented above.

This last remark will be especially important when we discuss some physical phenomena involving rotation. As an example, look at fig. 4.2: the volume integral for the whole is *zero*, but its left and right subregions have *non-zero and opposite* volume integrals.

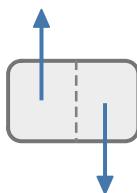


Figure 4.2 The volume integral for the whole region is zero, but the volume integrals for the left and right subregions are non-zero and opposite.

We add one more remark:

- the origin of the vector representing the volume integral is unimportant: it doesn't need to be placed at the centre of the volume, for instance. The vector refers to the volume as a whole, not to some specific point within.

Adding vectors in General Relativity

We are used to the idea of adding vectors placed at different points in space: we only have to first move each – keeping it parallel to itself – to a common point, and then add them all at that point with the usual rule.

This operation *cannot* be done in General Relativity: the notion of parallelism doesn't apply anymore in a simple way, owing to the curvature of spacetime. The addition would lead to different results depending on how we transported the vectors. So in General Relativity we can only sum vectors that are placed at the same spacetime point.

How can this operation be possible in Newtonian mechanics and in practical applications, then? After all, General Relativity surely applies here! The answer is that the discrepancies of vector transportation are small enough in the neighbourhood of the Earth, as the curvature of spacetime is very small here.

Fluxes

For the visualization and mental representation of the flux of a vector quantity we must pay attention to analogous warnings as for a scalar quantity. Consider again our example surface



and suppose we are told that there's a flux of momentum through it, from left to right, having the following direction and orientation:

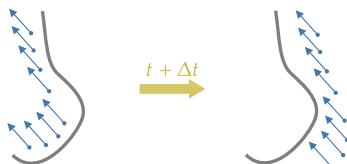


with magnitude **6 N**.

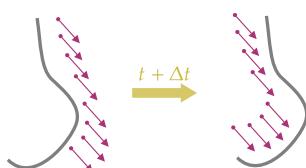
! Vector flux and crossing direction are two separate things

The vector that represents the flux, like the one above, and the surface-crossing direction, say left-to-right, are completely separate and independent things. In particular, they can have completely different directions.

We could visualize the flux above as follows:



with the familiar warning that the surface could be moving. But the following visualization also corresponds to exactly the same flux:



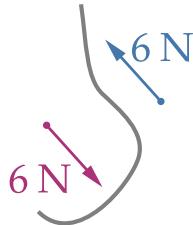
What happens in either illustration is that there's a change equal to a vectorial amount



of magnitude **6 N** on the *right* of the surface, and a

change equal to the *opposite* vectorial amount of the same magnitude on the *left* of the surface.

A less misleading representation of the flux in our example could therefore be like this:

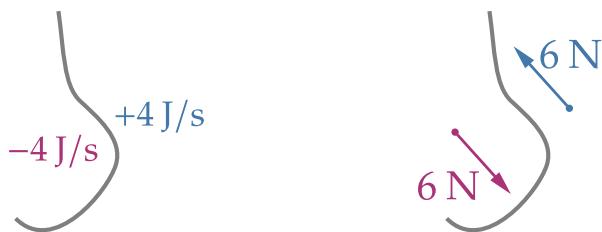


This picture says that there's a momentum change of a particular direction, orientation, magnitude on the left side of the surface, and an opposite momentum change on the right side of the surface. Note that **the vectors indicated on the two sides must be opposite and have equal magnitude**, otherwise the picture doesn't make sense.

An animated representation or mental visualization can also be quite illuminating. You can see one [at this link¹](#).

4.4 The symmetry of fluxes

Our discussion about fluxes of scalar and vector quantities, and their graphical representation, for instance



shows an important mathematical and physical symmetry of the flux of a quantity.

The *sign* of the flux is determined by the direction in which we imagine to cross the surface. In the first picture above, if we are imagining to cross the surface from left to right, then the energy flux is $+4 \text{ J/s}$; and if we are imagining to cross the surface from right to left, then the energy flux is -4 J/s . Analogously for the vector flux in the second picture above, where the minus sign corresponds to flipping the vector.

The crossing direction is arbitrary, completely left to us (just like the surface itself is arbitrary). It is therefore always important, when we report

a flux through a surface, to state which crossing direction we have agreed upon.

We can state this symmetry as follows:

Every flux is the same as the opposite flux in the opposite direction

A flux in a particular surface-crossing direction is equivalent to a flux of opposite sign in the opposite crossing direction.

This statement may sound familiar. Indeed it includes Newton's famous third law, the "principle of action and reaction", as a special case. Newton's third law is the expression of the symmetry of vector fluxes, but for the specific case of the flux of momentum, that is, force. Now we see that this property is more general: it applies not only to force, but also to the flux of angular momentum (torque), and to the flux of scalar quantities as well.

"LEX III. Actioni contrariam semper & æqualem esse reactionem: sive corporum duorum actiones in se mutuo semper esse æquales & in partes contrariae dirigi."

"LAW III. To every action there is always opposed an equal reaction: or, the mutual actions of two bodies upon each other are always equal, and directed to contrary parts."

Newton 1726

4.5 Force is momentum flux

We already mentioned in §§ 3.5 and 4.4 that *flux of momentum is what we call 'force'*. Owing to the importance of the notion of *force* in the many branches of physics which rely on Newtonian mechanics, it is appropriate to discuss this connection in more depth. This connection, as well as the connection to Newton's laws, will become even clearer when we discuss the balance of momentum in §X.

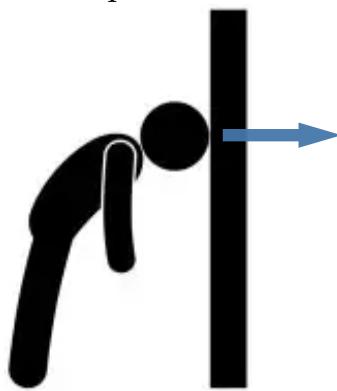
The notion of force is very intuitive. We associate it to the sensations that we feel in our skin, flesh, and even bones when, for instance, we push with our hand against a wall or against the other hand, or push on the ground with our feet to jump, or other similar actions. This force is typically represented by a vector, having the direction and orientation of the "push", and magnitude expressing its intensity. Such a force vector is exactly the vector expressing the flux of momentum; the two are the same.

A force can therefore also be visualized as a flow of momentum. This mental representation can be illuminating in some physical problems.

As a concrete example, imagine a person pushing against a wall. In terms of force, we say that

the person is exerting a force on the wall,
and the force vector has a person→wall orientation

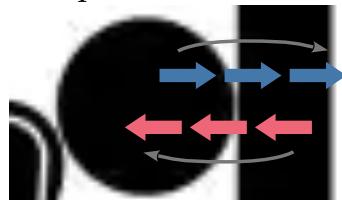
which can be depicted for example like this:



In terms of momentum flux, we imagine a surface separating the person and the wall, and we say that through this surface

there is a momentum flux from the person to the wall,
and the flux vector has a person→wall orientation

which, zooming in, can be depicted like this:

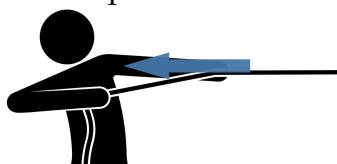


That is, an amount of momentum with [person→wall orientation](#) is flowing from the person to the wall; and, by symmetry of the flux, this is also equivalent to an amount of momentum with [opposite \(wall→person\) orientation](#) flowing from the wall to the person.

Let's take another concrete example, in order to distinguish the question of what exerts the force with the orientation of the force vector itself. Imagine a person pulling a rope fastened somewhere. In terms of force, we say that

the person is exerting a force on the rope,
and the force vector has a rope→person orientation

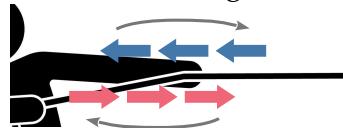
which can be depicted for example like this:



In terms of momentum flux we say that

there is a momentum flux from the person to the rope,
and the flux vector has a **rope→person orientation**

which can be depicted like this, zooming in:



That is, an amount of momentum with **rope→person orientation** is flowing from the person to the rope; and, by symmetry of the flux, this is also equivalent to an amount of momentum with **opposite (person→rope) orientation** flowing from the rope to the person.

A final example illustrates a situation in between the previous two. Consider the foot of a running person, as it pushes on the ground. In terms of momentum flux, there is a flux from the foot to the ground, and the flux vector is parallel to the ground, **oriented towards the back of the foot**. By flux symmetry, this is equivalent to an amount of momentum with **opposite orientation** (towards the running direction) flowing from the ground to the foot:



From the examples above, we see that thinking of force as momentum flux automatically leads to Newton's third law: if one side is gaining/losing momentum with a given orientation, by symmetry the other side is gaining/losing momentum with the opposite orientation. So if one side is experiencing a force with a given orientation, the other side must be experiencing a force with the opposite orientation.

Exercise 4.2

Using your intuition, try to guess the various momentum flows between this person and the walls:



(Buster Keaton in '*The Electric House*'²)

In the rest of these notes we shall use the terms ‘momentum flux’ and ‘force’ interchangeably, as synonymous words for the same thing. You can choose the mental representation that you prefer for this; but I recommend that you keep in mind the flux visualization. If a physics problem involving forces seems tricky, try to visualize them as momentum flow: suddenly the problem might become clearer.

4.6 Pressure, tension, shear force

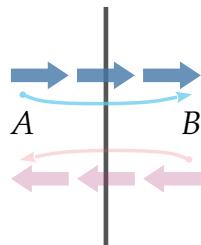
The examples of the previous section demonstrated a variety of possible orientations of the momentum-flux vector with respect to the surface through which it occurs. Obviously all orientations are possible. Special names are given, however, to three specific kinds of momentum flux: *pressure*, *tension*, and *shear force*.

Consider a surface and calls its sides *A* and *B*. Now consider a flux of momentum as seen from *A* to *B*.

Pressure

If the momentum-flux from *A* to *B* is a vector also oriented from *A* to *B*, then we call the momentum flux a **pressure**, or a *compressive momentum*

flux.



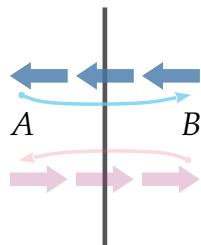
(click the picture for an animated version)³

By symmetry, there is an analogous relationship for the momentum flux as seen from *B* to *A*.

Pressure is the kind of momentum flux that we exert when we *push* on an object, and that air exerts on all objects it surrounds.

Tension

If the momentum-flux from *A* to *B* is a vector oriented in the opposite way, from *B* to *A*, then we call the momentum flux a **tension**, or a *tensile momentum flux*.



(click the picture for an animated version)⁴

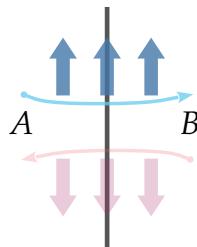
By symmetry, there is an analogous relationship for the momentum flux as seen from *B* to *A*.

Tension is the kind of momentum flux that we experience in our bones when we *pull* an object, and that occurs in any section of a stretched rubber band.

Shear force

If the momentum-flux from *A* to *B* is a vector oriented along the surface, then we call the momentum flux a **shear force**, or a *shearing momentum*

flux.



(click the picture for an animated version)⁵

By symmetry, there is an analogous relationship for the momentum flux as seen from **B** to **A**.

Shear force is the kind of momentum flux that we experience under our feet when we walk, and that occurs between a car's wheels and the ground.

In general, a momentum flux won't have any of the three special directions above, but rather a combination of them.

Exercise 4.3

Using your intuition, try to identify the various momentum fluxes that occur in the different parts of a tower crane. Which of the fluxes are (approximately) compressive, tensile, and shearing?



4.7 Total flows

Recall that a flux is defined as the amount of a quantity crossing a surface in a short time lapse Δt , divided by that time lapse. This is why the dimension of a flux is $\frac{\text{dimension of quantity}}{\text{time}}$. Denoting the flux by, say, J , this definition also means that the amount of quantity crossing the surface in Δt is equal to $J \Delta t$.

Take now a surface existing between two time instants t_1 and t_2 ; during this time lapse it could also be moving and changing shape. Choose a crossing direction through the surface. At each intermediate time instant t we can then measure the flux of a quantity crossing the surface in that direction, at that instant; denote it by $J(t)$.

The total amount of quantity that cross the surface between times t_1 and t_2 can be found by dividing the time interval into very short time

lapses of length Δt . For each short time lapse we know that the amount is $J(t) \Delta t$. The total is obtained by adding these small amounts. As Δt is considered shorter and shorter, this sum is by definition an integral:

Total flow of quantity through a surface between two times

The total amount of quantity crossing a surface (relative to a specified crossing direction) between times t_1 and t_2 is given by

$$\int_{t_1}^{t_2} J(t) dt . \quad (4.1)$$

The meaning of the integral above should be clear for any scalar quantity, for which the flux is also a scalar. In the case of a vector quantity, for instance momentum, the flux is also a vector. How is the integral of a vector calculated? The procedure is simple: a vector is represented by three components, and we calculate the integral for each component obtaining three results, which are the components of a new vector.

Take the case of momentum, whose flux (force) we denote $\mathbf{F} = (F_x, F_y, F_z)$. The integral of this flux is then

$$\int_{t_1}^{t_2} \mathbf{F}(t) dt := \left(\int_{t_1}^{t_2} F_x(t) dt , \int_{t_1}^{t_2} F_y(t) dt , \int_{t_1}^{t_2} F_z(t) dt \right) . \quad (4.2)$$

There can be zero total flow with non-zero flux

The result of the integral above could be zero – which means that no *net* amount of quantity crossed the surface between t_1 and t_2 . Yet the flux $J(t)$ can be non-zero, even at all times. Of course it needs to be positive at some times, and negative at others.

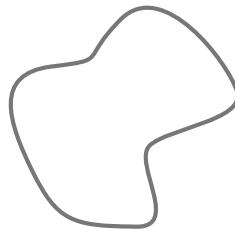
As a simple example, consider a room's door. During one minute, three people enter through the door; during the next minute, three people (not necessarily the same!) exit. The total flow of people is zero, but the flux was non-zero during the first minute, and non-zero during the second minute.

 **Exercise 4.4**

- What is the physical dimension of the total flow of a quantity?
- Suppose that we calculate the integral above in the case of matter, finding a total of $\int_{t_1}^{t_2} J(t) dt = 7 \text{ mol}$. Now we change our mind and choose the opposite crossing direction for that surface. How does the result above change?

4.8 Closed surfaces, influxes, effluxes

We shall often consider **closed** surfaces, that is, surfaces that don't have a boundary or border, like the surface of a sphere or of a cube. A closed surface delimits a specific three-dimensional volume, and we can naturally speak of its **interior** and its **exterior**. An example (simplified by removing one dimension as usual) is the surface delimiting the volume of our previous examples:



We can give two crossing directions to a closed surface: out-in, from exterior to interior; or in-out, from interior to exterior. A flux through the surface is usually called **influx** if we are considering the out-in crossing direction, and **efflux** or *outflux* if we are considering the in-out crossing direction. Obviously

$$\text{influx} \equiv -\text{efflux}$$



It makes sense to speak of influx or efflux only for a *closed* surface.

The influx and efflux, unless specified otherwise, are fluxes through the *whole* surface. In the picture above, the influx of 2 mol/s can be happening everywhere across the surface, or just across some portions of it.

Exercise 4.5

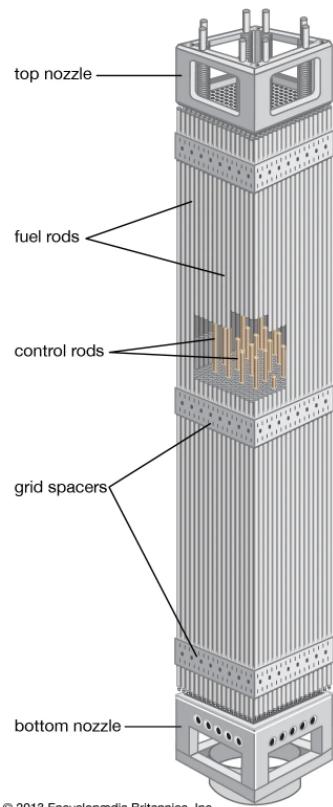
Take an imaginary cylindrical surface enclosing one [control rod⁶](#) in a [nuclear-fission reactor⁷](#). Let's say that in a reactor there are 20 such rods. Around 5×10^{19} neutrons are liberated in a second by the fission fuel, but 2/3 of these are absorbed by the control rods.

How much, on average, is the efflux of neutrons (matter) through the surface of one control rod?

Express the result first in neutrons/s, and then in mol/s, using the Avogadro constant

$$N_A = 6.022\,140\,76 \times 10^{23} \text{ particles/mol.}$$

Be careful about the signs!



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URLs for chapter 4

1. <https://pglpm.github.io/7wonders/media/vectorfluxanimation.webp>
2. <https://www.imdb.com/title/tt0013099/>
3. <https://pglpm.github.io/7wonders/media/pressureanimation.webp>
4. <https://pglpm.github.io/7wonders/media/tensionanimation.webp>
5. <https://pglpm.github.io/7wonders/media/shearforceanimation.webp>
6. https://energyeducation.ca/encyclopedia/Control_rod
7. <https://www.britannica.com/technology/nuclear-reactor>

Chapter 5

Physical laws

Every branch of physical science is based on two sets of fundamental equations. The first set is that of basic laws of physics, which are postulated to hold valid for all bodies under all conceivable circumstances [...]. The second set of fundamental equations are the constitutive equations: these are relationships which are not supposed to hold for all bodies, but only to describe the behavior of some restricted class of bodies, or possibly of a larger class of bodies for a more restricted class of phenomena.

G. Astarita 1990

5.1 Two kinds of laws

In ch. 3 we introduced the main quantities in terms of which we describe physical phenomena. How do we describe the patterns, regularities, and diversities that we observe in these phenomena? This is the role of physical laws. Different physics formalisms or languages express these laws in different ways, even if the actual physical phenomenon is exactly the same.

In the formalism we're using we distinguish between two kinds of physical laws:

- Balance laws
- Constitutive relations

Balance laws, or simply balances, have this name because they express a sort of trade-off or “budget” about physical quantities. **Conservation laws** are a special and important kind of balance laws.

Constitutive equations

We shall see that there's a set of balances which apply to *all* physical phenomena, without exceptions; they represent the basic universal patterns that we observe everywhere. Constitutive relations instead only apply to particular phenomena and only in particular theories. They express the diversity that we observe around us.

In this chapter we

5.2 Conservation laws

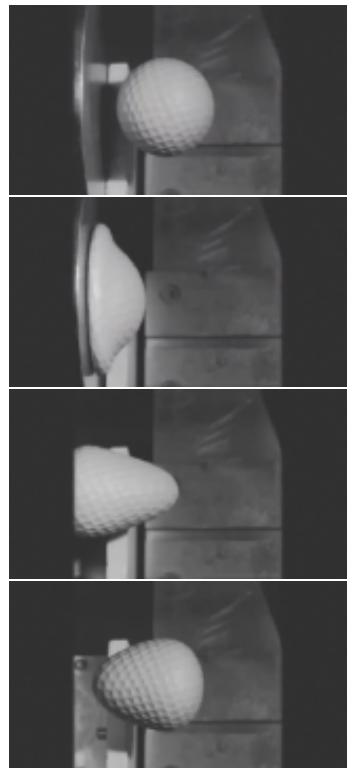
Conservation laws tightly connect the *volume integrals* and *fluxes* discussed in ch. 4. They are intuitively and visually easy to grasp. The archetype is as follows:

Consider a *closed* (§ 4.8) control surface at a (coordinate) time t_1 , delimited by a closed control surface. A ball-shaped region, delimited by a spherical surface, is an example. (Actually the control volume doesn't need to be connected: it could consist of several separate volumes, each delimited by a distinct closed surface.)

Starting from that particular time, imagine the closed control surface moving and deforming, changing shape and size but remaining closed – no holes or cuts – through time, until a later time t_2 . At this time we shall have a final closed control surface that can be different from the initial one, and at a different position. The temporal sequence of closed control surface delimits a sequence of control volumes; see the sequence on the right as an example. We focus in particular on the initial and the final control volumes. Note that, if we wanted, we could start with a spherical surface delimiting a region the size of a tennis ball, and end with a cubical surface delimiting a region the size of the Sun.

Now consider any one of the seven primitive quantities of § 3.1, except magnetic flux. Remember that we can measure the amount of this quantity in any control volume, and its flux across any control surface. For concreteness let's take *matter*, for which we denote volume integrals by N and flux by J . We make the following measurements:

- $N(t_1)$: total amount of matter contained within the control volume at time t_1 .
- $J(t)$: *influx* (flux flowing inwards, § 4.8) of matter (and antimatter!) through the closed surface at time t , for all times t between t_1 and t_2 .
- $N(t_2)$: total amount of matter contained within the control volume at time t_2 .



Temporal sequence of closed control surfaces, represented by the golf ball ([click here for an animated version¹](#))

Each of the three amounts above is, in this example, a scalar that can be positive, negative, or zero. For instance we could have:

$$\begin{aligned} t_1 &= 00:00:00Z & t_2 &= 00:00:18Z \\ N(t_1) &= 3.50 \text{ mol} & J(t) &= -0.06 t \text{ mol/s}^2 & N(t_2) &= -6.22 \text{ mol} \end{aligned}$$

note that $J(t)$ is given as a function of the time t .

With this setup we can introduce two kinds of balance laws.

Conservation law

A quantity with volume integral $N(t)$ and flux $J(t)$ is said to be **conserved**, or to satisfy a **conservation law**, if the following equality holds always:

$$N(t_2) - N(t_1) - \int_{t_1}^{t_2} J(t) dt = 0 \quad (5.1)$$

The meaning of the equation above is intuitive, especially if we rewrite it like this:

$$N(t_2) = N(t_1) + \int_{t_1}^{t_2} J(t) dt$$

which says that the amount of quantity in the final control volume is equal to the amount in the initial control volume, plus the total amount (which could be negative) that flowed in between those times. Said otherwise, any amount of quantity that appears in (or disappears from) the volume, must come in (or go out) through the surface; it can't appear out of nowhere.

A conservation law allows us to make several kinds of predictions or deductions. For example:

- If we know the amount of quantity within the control surface at t_1 and the amount that flowed in through the surface between t_1 and t_2 , then we can predict the amount within the control surface at t_2 : this is given by the last equation above.
- If we know the amount of quantity within the control surface at t_1 and the one at t_2 , then we can deduce the total amount that flowed in through the surface between the times t_1 and t_2 :

$$\int_{t_1}^{t_2} J(t) dt = N(t_2) - N(t_1).$$

- If we know the amount of quantity within the surface at t_2 and the amount that flowed in between t_1 and t_2 , we can deduce the amount within the surface at t_1 :

$$N(t_1) = N(t_2) - \int_{t_1}^{t_2} J(t) dt .$$

These kinds of predictions and deductions are very powerful, because we are fully free to decide the shapes and motion of the control surfaces and volumes, and even the times t_1 and t_2 .

We unconsciously use several conservation laws all the time in our everyday life. If a bike tyre is suddenly deflated, we conclude that that air must have got out of it through its surface, which must therefore have a hole or a defective valve; we don't conclude that air "just disappeared".



Deflated tyre: did the air within just vanish out of existence?

Differential form of a conservation law

The conservation law

$$N(t_2) - N(t_1) - \int_{t_1}^{t_2} J(t) dt = 0$$

is said to be written in *integral form*, because it contains the integral of the flux, necessary to calculate the total amount of quantity that flowed through the control surface.

Change the name of time t_1 simply to t , and take t_2 to come after a *very short* lapse of time Δt after t :

$$t_1 \text{ becomes } t , \quad t_2 \text{ becomes } t + \Delta t$$

Then the conservation law above can be approximated as

$$N(t + \Delta t) - N(t) - J(t) \Delta t \approx 0$$

If we divide the whole expression by Δt and consider smaller and smaller Δt , the fraction $\frac{N(t+\Delta t)-N(t)}{\Delta t}$ becomes, by definition, a derivative, and we obtain a new form for the conservation law:



Some mathematical conditions must be satisfied to make these steps; or otherwise the derivative below must be carefully defined in a generalized sense.

Conservation law in differential form

A quantity is said to satisfy a **conservation law** if the following equality holds always:

$$\frac{dN(t)}{dt} - J(t) = 0 \quad (5.2)$$

This form says that the rate of change of the volume integral of a quantity must be equal to the *influx* of that quantity.

Both forms, integral and differential, of a conservation law are useful, in different situations. The integral form is obviously useful if we know the total flow during a time lapse but we don't know the flux at each instant. The differential form is often useful for simulating physical processes and make predictions.

5.3 Balance laws

Our intuitive understanding of a conserved quantity is that it cannot be "created" or "destroyed". This leads to the definition of a more general law, where this condition is *not* satisfied.

The setup is just like before: we consider an initial and final control volumes at t_1 , t_2 , and a sequence of closed control surfaces in between. This time let's take *energy-mass* as a concrete example, with initial volume integral $E(t_1)$, final volume integral $E(t_2)$, and flux $\Phi(t)$.

Balance law

A quantity with volume integral $E(t)$ and flux $\Phi(t)$ is said to be only **balanced**, or to satisfy a **balance law** or simply **balance**, if the following equality holds always:

$$E(t_2) - E(t_1) - \int_{t_1}^{t_2} \Phi(t) dt = \int_{t_1}^{t_2} R(t) dt \quad (5.3)$$

where $R(t)$, called the **source**, is the amount of quantity created within the closed control surface at time t .

If the quantity is a scalar, then $R(t)$ can also be negative, in which case we say that some of the quantity has been destroyed.

A balance law is in some respects more trivial than a conservation law. If we measure that $E(t_1) - E(t_2) + \int_{t_1}^{t_2} \Phi(t) dt$ is not zero (so there's no conservation law), we can always say that some amount of quantity must have been created or destroyed between t_1 and t_2 . So a quantity can always be said to satisfy a balance. But a balance is not trivial if we have some specific physical law that tells us in advance how the amount created or destroyed at each instant, $R(t)$, can be calculated.



One could object: "why don't you just add the two integrals in the equation above, writing

$$E(t_2) - E(t_1) - \int_{t_1}^{t_2} [\Phi(t) - R(t)] dt = 0$$

and get a conservation law?"

The point is this: the flux Φ involves only what's happening *on the control surface*; the source R , on the other hand, involves what's happening *in the control volumes, at each time*.

So, whereas a conservation law allows us to predict the amount of quantity in the final volume just by making measurements on the surface during the time lapse, a general balance law requires us to make measurements at every point within the volume during the time lapse.

Thus a conservation law can be seen as a special and powerful case of a balance law, one for which $R(t) = 0$ *always and everywhere*.

Differential form of a balance law

By considering a very short time lapse Δt , with reasoning analogous to the case of a conservation law, we can write a balance law in a different form:

Balance law in differential form

A quantity is said to satisfy a **balance law** if the following equality holds always:

$$\frac{dN(t)}{dt} - J(t) = R(t) \quad (5.4)$$

Balance laws in General Relativity

Conservation and balance laws appear even simpler from the point of view of Relativity Theory. From a four-dimensional spacetime perspective, a 3D volume is a region, called *hypersurface*, having one less dimension than spacetime. But a sequence of 2D surfaces through time is also just a region having one less dimension than spacetime – two spatial dimensions and one temporal one. Thus the distinction between the 3D region at t_1 , the sequence of 2D surfaces between t_1 and t_2 , and the 3D region at t_2 disappear: they are seen to be just different parts of the same three-dimensional hypersurface. We perceive some parts of this hypersurface as belonging “to the same time”, showing their three dimensions all at once; and other parts as extending through time, showing only two dimensions at any time. In fact, different observers make this division in different ways.

And from a spacetime perspective, the amount of a quantity $E(t_1)$ or $E(t_2)$ within a 3D region is seen as a flux through time; so its apparent difference from the flux Φ also disappear. Their only difference is that one (C) points exclusively in the time direction, while the other (Φ) also points partially in a spatial direction.

What about the magnetic flux?

In the case of magnetic flux, the idea of a conservation law is analogous, but is formulated with one less spatial dimension: we consider a closed 1D curve at t_1 , one at t_2 , and a sequence in between these times. The magnetic flux turns out to be a quantity for which it's possible to ask how much is “linked” to a closed curved, and how much is crossing a closed curve. One way to understand this is to imagine magnetic flux as a bundle of tubes or lines that are either closed or extend to infinity. It is a very fascinating quantity, and one may wonder if other quantities exist which satisfy similar balances even with one less dimension. But we shall not pursue magnetic flux or similar topics in these notes.

5.4 Constitutive relations

In the previous section we saw that a balance law can be used to predict how the density of a quantity changes with time. But we also saw that in order to make such prediction we need to already know what the *future* fluxes of the quantity will be, throughout the region of interest. This

doesn't sound like a great achievement: to predict the future, we need the future. The situation doesn't improve if we have two or more quantities that satisfy balance laws: we shall need to know in advance the fluxes of each.

It turns out, however, that there are physical laws which *relate the densities of some quantities to the fluxes of others*. Such relations can potentially allow us to calculate the time evolution without knowing the fluxes in advance. Roughly, the basic idea is as follows: From the density of one quantity, we calculate the fluxes of the second quantity, and use them to predict its density at a later time. Once we know this density, we use it to calculate the new fluxes of the first quantity, and therefore its density at a later time; and so on. This cross-prediction does not fully eliminate the necessity of specifying some future quantities in advance; but the remaining ones that we need to specify are often in our control. Indeed this is how we can predict how a physical system will respond to influences controlled by us.

Physical laws of this kind turn out not to be universal: they depend on, or are '[constitutive](#)'² of, the particular physical phenomenon and the physical theory being used. For this reason they are called **constitutive relations** or **constitutive equations**. In some fields they are called **closure equations**, because they allow us to "close" the system of balance equations in such a way that it can be used for future predictions without knowing in advance the future value of some quantities.

Constitutive relations express the diversity that we observe around us, for example the different behaviours of a drop of water, which obeys some constitutive relations, as compared with a block of wood, which obeys others. They also mark the difference between specialized or approximate physical theories, for example between Newtonian mechanics, which is based on particular approximate constitutive relations, and general relativity, which is based on different and more exact constitutive relations. Depending on the specific scientific field you'll work in, you'll learn some constitutive relations in more detail than others.

Constitutive relations come in a great variety of mathematical forms. Some of them are simple algebraic relations between the density of one quantity and the flux of another. Others involve spatial or time derivatives. Other still involve integrals in space or in time.

❖ Add examples of laws that reduce to the balances: Bernoulli, Poynting?, mechanical-energy-work theorem

URLs for chapter 5

1. <https://pglpm.github.io/7wonders/media/surfacemotion.mp4>
2. <https://www.merriam-webster.com/dictionary/constitutive>

Chapter 6

The Seven Wonders of the World

We already remarked that the seven primitive quantities introduced in §3.1 are common to all our main physical theories. Even more remarkable is that each of them obeys a balance law like (5.3); three actually obey a strict conservation law like (5.4).

These seven balance laws are known, so far, to be satisfied by *all* phenomena and in all our main physical theories, approximate or not: from subatomic scales to cosmological scales, from Newtonian mechanics to general relativity and quantum theory. No exceptions are known.

The seven fundamental balances

- Conservation of matter
- Conservation of electric charge
- Conservation of magnetic flux
- Balance of energy
- Balance of momentum
- Balance of angular momentum
- Balance of entropy

These balances are truly the “Seven Wonders of the World”¹.

Once you learn these balances, you can apply them to any kind of physical phenomenon or problem you work with: construction of bridges, control of chemical reactions, operation of GPS navigation and satellites,

¹ Compare with the traditional seven wonders: <https://www.britannica.com/topic/Seven-Wonders-of-the-World>, <https://education.nationalgeographic.org/resource/seven-wonders-ancient-world/>.

monitoring of nuclear power plants, sending robots to Mars, or collisions of subatomic particles.

Let us say something more about each.

6.1 Conservation of matter

More precisely: matter, electric charge, magnetic flux always satisfy strict conservation laws. For matter we actually have several conservation laws, one for each kind of matter. Energy, momentum, angular momentum, entropy satisfy balance laws, which become conservation laws only in special circumstances, but not always.

Chapter 7

Densities and flux-densities

How do we mathematically represent the seven main quantities? We said that, for each, we can ask *how much in this region?* and *how much through this surface during this time?* We must therefore use mathematical objects that allow us to answer these two questions. These objects are the *density* and the *flux vector*. The idea behind them is quite simple.

Density

Take a point in spacetime with coordinates (t_0, x_0, y_0, z_0) . At the instant t_0 , imagine a very small cuboid centred at (x_0, y_0, z_0) , with sides aligned with the coordinate axes x, y, z and of lengths $\Delta x, \Delta y, \Delta z$. Then we have

$$\text{amount of quantity in cuboid} = n(t_0, x_0, y_0, z_0) \Delta x \Delta y \Delta z . \quad (7.1)$$

where $n(t_0, x_0, y_0, z_0)$ is the **density** at the spacetime point (t_0, x_0, y_0, z_0) .

The density tells us how much quantity there's in a unit of volume. As the notation suggests, it is a function of the coordinates, that is, of time and spatial position. This functional dependence reflects the fact that we can have a larger amount of a quantity concentrated in some regions at some times, than in other regions or at other times.

In order to calculate the total amount of the quantity in an arbitrary 3D region, we simply divide it into very small cuboids similar to the one above. For each cuboid we calculate the respective amount; this amount will generally be different from cuboid to cuboid, because the value of the

density $n(t, x, y, z)$ will be different at each cuboid's centre. Then we sum up all these amounts.

Isn't this an integral?

You probably recognize this procedure as the description of integration. Indeed we can write:

$$\text{total amount of quantity in 3D region} = \iiint_{\text{3D region}} n(t, x, y, z) dx dy dz . \quad (7.2)$$

Clearly we must know the value of the density n at all points within the region, in order to calculate the total amount. Note that if the quantity is absent in some subregion, then $n(\dots) = 0$ there.

Do not confuse flux with movement

Flux

For the flux of the quantity through a surface we consider three different cases:

x -Flux

Take again a point in spacetime with coordinates (t_0, x_0, y_0, z_0) . Keeping x_0 fixed, imagine a very small rectangular surface centred at (x_0, y_0, z_0) , with sides aligned with the coordinate axes y, z and of lengths $\Delta y, \Delta z$. Imagine that this surface exists for a lapse of time Δt around the time t_0 . Then we have

$$\left. \begin{array}{l} \text{amount of quantity} \\ \text{through rectangle } \Delta y \Delta z \\ \text{towards positive } x \\ \text{during time lapse} \end{array} \right\} = j_x(t_0, x_0, y_0, z_0) \Delta t \Delta y \Delta z . \quad (7.3)$$

where $j_x(t_0, x_0, y_0, z_0)$ is the **x -flux** at the spacetime point (t_0, x_0, y_0, z_0) .

The x -flux tells us how much quantity is flowing in a unit of time through a unit of surface parallel to y, z . Also the x -flux is a function of the coordinates, because the flux could be larger through some surfaces at some times, than through other surfaces at other times. If the quantity is

flowing in the negative x direction, then j_x will be negative; and clearly if no quantity is flowing through the small rectangle, then $j_x = 0$.

In an analogous way we define a y -flux and a z -flux:

y -Flux

$$\left. \begin{array}{l} \text{amount of quantity} \\ \text{through rectangle } \Delta z \Delta y \\ \text{towards positive } y \\ \text{during time lapse} \end{array} \right\} = j_y(t_0, x_0, y_0, z_0) \Delta t \Delta z \Delta y . \quad (7.4)$$

z -Flux

$$\left. \begin{array}{l} \text{amount of quantity} \\ \text{through rectangle } \Delta x \Delta y \\ \text{towards positive } z \\ \text{during time lapse} \end{array} \right\} = j_z(t_0, x_0, y_0, z_0) \Delta t \Delta x \Delta y . \quad (7.5)$$

The three fluxes can be grouped into the **flux vector**

$$\mathbf{j}(t, x, y, z) := (j_x(t, x, y, z), j_y(t, x, y, z), j_z(t, x, y, z)) \quad (7.6)$$

In order to calculate the total amount of the quantity through an arbitrary sequence of 2D surfaces during a time lapse, we simply divide the total time into very brief time lapses, and for each of these we approximate the surface with very small rectangles aligned along the three axes. For each such lapse and rectangle we calculate the respective flux amount, using j_x , j_y , or j_z depending on the orientation of the rectangle. Then we sum up all these amounts.

T

This is also the description of an integration, which can be written

total amount of quantity through sequence of surfaces =

$$\begin{aligned} & \iiint j_x(t, x, y, z) dt dy dz + \\ & \iiint j_y(t, x, y, z) dt dz dy + \\ & \iiint j_z(t, x, y, z) dt dx dy \quad (7.7) \end{aligned}$$

We shall not have to calculate density integrals (7.2) and flux integrals (7.7) in these notes. In concrete physics problems these integrals are

often difficult to calculate, and require advanced mathematical techniques and specialized software. But if you'll ever end up working in fields such as computational fluid dynamics, atmospheric or ocean modelling, or numerical relativity, then you shall probably encounter them again.

The definitions above of density and flux are appropriate if the quantity in question is a scalar. For vector quantities such as momentum and angular momentum, the density is a vector, and each of the three fluxes is a vector as well.

7.1 Balance laws expressed with derivatives

In ch. 7 we saw that each of the main seven quantities is represented by a density and by a flux vector, which allow us to calculate the amount of quantity within a 3D region at a particular time, and flowing through a 2D surface during an time interval. These amounts are exactly what $N(t_1)$, $N(t_2)$, J express. Therefore, a conservation or balance law must also give some relationship between density and flux vector.

The way to find such relationship is intuitive and mathematically simple, if only a bit lengthy. The idea is to formulate the balance law for a very simple case, the same simple set-up that we used in ch. 7 to introduce density and flux.

Take a point in spacetime with coordinates (t, x, y, z) . We apply the conservation law (5.4) to the following situation:

$N(t_1)$: Choose an initial time $t_1 = t - \frac{\Delta t}{2}$, slightly before t , and a cuboid region centred at (x, y, z) with sides $\Delta x, \Delta y, \Delta z$. The total amount of the quantity within this region is given, adapting formula (7.1), by

$$N(t_1) = n\left(t - \frac{\Delta t}{2}, x, y, z\right) \Delta x \Delta y \Delta z . \quad (7.8)$$

$N(t_2)$: Choose a final time $t_2 = t + \frac{\Delta t}{2}$, slightly after t , and a cuboid region centred at (x, y, z) with sides $\Delta x, \Delta y, \Delta z$. The total amount of the quantity within this region is given, adapting formula (7.1), by

$$N(t_2) = n\left(t + \frac{\Delta t}{2}, x, y, z\right) \Delta x \Delta y \Delta z . \quad (7.9)$$

J : We calculate the flux separately through the six rectangular surfaces bounding the small cuboid region:

J_x : First choose a rectangular surface centred slightly to one side of the point (x, y, z) : at $(x - \frac{\Delta x}{2}, y, z)$, with sides parallel to y, z and lengths $\Delta y, \Delta z$. We keep this surface constant for the small duration Δt . The x -flux through this rectangle, towards positive x , during Δt , is according to formula (7.3)

$$J_x^{\text{in}} = j_x(t, x - \frac{\Delta x}{2}, y, z) \Delta t \Delta y \Delta z .$$

Now choose the rectangular surface parallel to the previous one, but on the other side of the point (x, y, z) , centred at $(x + \frac{\Delta x}{2}, y, z)$. The x -flux through this rectangle, towards positive x , during Δt , is

$$J_x^{\text{out}} = j_x(t, x + \frac{\Delta x}{2}, y, z) \Delta t \Delta y \Delta z .$$

Note that the flux J_x^{in} points to the interior of the 3D region, towards its centre (x, y, z) ; whereas J_x^{out} points to the exterior, away from the centre. The total flux into the region is therefore given by their difference:

$$\begin{aligned} J_x &= J_x^{\text{in}} - J_x^{\text{out}} \\ &= [j_x(t, x - \frac{\Delta x}{2}, y, z) - j_x(t, x + \frac{\Delta x}{2}, y, z)] \Delta t \Delta y \Delta z \end{aligned} \quad (7.10)$$

J_y, J_z : An analogous reasoning can be made to find the total flux through the other four surfaces, two parallel to z, y and two to x, y . We find

$$\begin{aligned} J_y &= J_y^{\text{in}} - J_y^{\text{out}} \\ &= [j_y(t, x, y - \frac{\Delta y}{2}, z) - j_y(t, x, y + \frac{\Delta y}{2}, z)] \Delta t \Delta z \Delta x \end{aligned} \quad (7.11)$$

$$\begin{aligned} J_z &= J_z^{\text{in}} - J_z^{\text{out}} \\ &= [j_z(t, x, y, z - \frac{\Delta z}{2}) - j_z(t, x, y, z + \frac{\Delta z}{2})] \Delta t \Delta x \Delta y \end{aligned} \quad (7.12)$$

The total flux into the region during the time interval Δt is finally the sum of the three above:

$$J = J_x + J_y + J_z \quad (7.13)$$

Up to now we have only chosen arbitrary surfaces, regions, and a time interval. Now let's assume that the corresponding amounts and fluxes

satisfy a balance law like (5.4). Putting together the puzzle pieces we find

$$N(t_2) - N(t_1) - J = 0 \quad \implies$$

$$\left. \begin{aligned} & n\left(t + \frac{\Delta t}{2}, x, y, z\right) \Delta x \Delta y \Delta z - n\left(t - \frac{\Delta t}{2}, x, y, z\right) \Delta x \Delta y \Delta z \\ & - [j_x(t, x - \frac{\Delta x}{2}, y, z) - j_x(t, x + \frac{\Delta x}{2}, y, z)] \Delta t \Delta y \Delta z \\ & - [j_y(t, x, y - \frac{\Delta y}{2}, z) - j_y(t, x, y + \frac{\Delta y}{2}, z)] \Delta t \Delta z \Delta x \\ & - [j_z(t, x, y, z - \frac{\Delta z}{2}) - j_z(t, x, y, z + \frac{\Delta z}{2})] \Delta t \Delta x \Delta y \end{aligned} \right\} = 0$$

This expression is quite long, but it should be intuitively understandable if you try to identify the individual summed terms.

Now we take the last, long expression – which is an equality – and divide its left and right sides by $\Delta t \Delta x \Delta y \Delta z$. Note that many of the “ Δ ” factors simplify out. We arrive at this:

$$\left. \begin{aligned} & \frac{n\left(t + \frac{\Delta t}{2}, x, y, z\right) - n\left(t - \frac{\Delta t}{2}, x, y, z\right)}{\Delta t} \\ & + \frac{j_x\left(t, x + \frac{\Delta x}{2}, y, z\right) - j_x\left(t, x - \frac{\Delta x}{2}, y, z\right)}{\Delta x} \\ & + \frac{j_y\left(t, x, y + \frac{\Delta y}{2}, z\right) - j_y\left(t, x, y - \frac{\Delta y}{2}, z\right)}{\Delta y} \\ & + \frac{j_z\left(t, x, y, z + \frac{\Delta z}{2}\right) - j_z\left(t, x, y, z - \frac{\Delta z}{2}\right)}{\Delta z} \end{aligned} \right\} = 0 .$$

Still a long expression; but examine each fraction: we have a difference between a term calculated at some point, and one calculated at some point plus an increment “ Δ ”; this difference is then divided by that increment. This is the definition of *partial derivative*, that is, the derivative with respect to one variable, while keeping the other variables fixed. The short symbol for this is “ $\frac{\partial \dots}{\partial \dots}$ ”, so we can simply write

Conservation law in derivative form

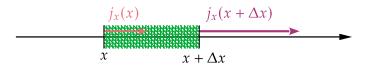
$$\frac{\partial n}{\partial t} + \frac{\partial j_x}{\partial x} + \frac{\partial j_y}{\partial y} + \frac{\partial j_z}{\partial z} = 0 \quad (7.14)$$

Let's make a simple check to see if this formula makes sense, in a simple

one-dimensional case where we disregard the y, z coordinates (let's say their fluxes are zero), so that the balance law becomes

$$\frac{\partial n}{\partial t} + \frac{\partial j_x}{\partial x} = 0.$$

Suppose that the derivative of the x -flux is positive: $\frac{\partial j_x}{\partial x} > 0$. This means that the x -flux increases as we move from x to $x + \Delta x$, as shown in the side figure. The x -flux at x is bringing some amount of quantity into the region of interest, the x -flux at $x + \Delta x$ is taking a larger amount out of that region. Therefore, during the short time Δt in which this flux takes place, the total amount of quantity within the region will decrease. But the total amount is given by $n(t)$, which is therefore decreasing with time. This means that its time derivative is negative: $\frac{\partial n}{\partial t} < 0$. This agrees with balance law above: the sum of the two derivatives must be zero, so if one is positive, the other must be negative.



Balance laws like (7.14) are the basis of many important computational and simulation methods in a huge variety of physical applications: simulation of the ocean around an offshore oil platform, of wind in a wind farm, of air around an aeroplane's wing, of earthquakes, of electromagnetic-wave propagation, of oscillations in a suspension bridge, of percolation of fluids through terrain, of chemical reactions, of energy transfer... The list could go on for pages!

The basic idea for using a balance law for a simulation is easy to understand. Take again the simpler one-dimensional formula without y, z coordinates, and rewrite it remembering the basic meaning of the derivative:

$$\frac{n(t + \Delta t, x) - n(t, x)}{\Delta t} + \frac{j_x(t, x + \frac{\Delta x}{2}) - j_x(t, x - \frac{\Delta x}{2})}{\Delta x} = 0$$

Suppose we know all quantities at time t ; we can then find the density n at time $t + \Delta t$. Leave it on the left side and bring all other terms on the right side, multiplying them by Δt :

$$n(t + \Delta t, x) = \Delta t \left[n(t, x) - \frac{j_x(t, x + \frac{\Delta x}{2}) - j_x(t, x - \frac{\Delta x}{2})}{\Delta x} \right]$$

which is essentially [Euler's method](#)¹.

This formula can then be used iteratively to find the density n at later times. As a simple simulation set-up, we consider a region of space of interest and divide it into cells of width Δx . We also treat time in steps of Δt . Then the iterative procedure goes as follows:

1. *Initial values:* Assign the value of n at each x -cell at the initial time.
2. *For-loop in x :* for each x -cell, calculate the value of n at $t + \Delta t$, using the formula above.
3. Increase time by Δt .
4. Go to 2.

For example, suppose we have two x -cells of width $\Delta x = 1$: one centred at x_1 and the other at $x_2 = x_1 + 1$. There are three boundaries: the leftmost at $x_1 - \frac{1}{2}$, one between the cells at $x_1 + \frac{1}{2} \equiv x_2 - \frac{1}{2}$, and the rightmost at $x_2 + \frac{1}{2}$. Then we have the following steps:

1. Assign the values $n(t_0, x_1), n(t_0, x_2)$ of n at x_1 and x_2 at time t_0 .
2. Assign the values

$$j_x(t_0, x_1 - \frac{1}{2}), j_x(t_0, x_1 + \frac{1}{2}), j_x(t_0, x_2 + \frac{1}{2})$$

of the fluxes at the three boundaries at time t_0 .

3. Calculate the values of n at x_1 and x_2 at time $t_1 = t_0 + \Delta t$, for instance

$$n(t_1, x_1) = \Delta t \left[n(t_0, x_1) - \frac{j_x(t_0, x_1 + \frac{1}{2}) - j_x(t_0, x_1 - \frac{1}{2})}{\Delta x} \right].$$

Note that we have all quantities required on the right side.

4. Assign the values

$$j_x(t_1, x_1 - \frac{1}{2}), j_x(t_1, x_1 + \frac{1}{2}), j_x(t_1, x_2 + \frac{1}{2})$$

of the fluxes at the three boundaries at time t_1 .

5. Calculate the values of n at x_1 and x_2 at time $t_2 = t_1 + \Delta t$, for instance

$$n(t_2, x_1) = \Delta t \left[n(t_1, x_1) - \frac{j_x(t_1, x_1 + \frac{1}{2}) - j_x(t_1, x_1 - \frac{1}{2})}{\Delta x} \right].$$

Note that we have again all quantities required on the right side.

6. And so on.

This routine requires us to know the flux at each x -cell and at each time. But we shall soon see that this requirement can be bypassed.

The rudimentary routine above is the basis from which more precise and complex simulation routines are developed.

Exercise 7.1

Write a script, in your preferred programming language, that implements the simulation routine above in an x -grid with four cells $x = 1, \dots, x = 4$. Take a grid size $\Delta x = 1$ and time step $\Delta t = 1$. We can imagine that the quantity in question is electric charge. Use the following initial values for $n(t, x)$ at $t = 0$:

$$n(t=0, x=1) = 7, \quad n(t=0, x=2) = 0, \quad n(t=0, x=3) = 0, \quad n(t=0, x=4) = 7$$

and assume the flux is, at each time t :

$$\begin{aligned} j_x(t, x=0.5) &= 0, \\ j_x(t, x=1.5) &= +2, \quad j_x(t, x=2.5) = 0, \quad j_x(t, x=3.5) = -2, \\ j_x(t, x=4.5) &= 0. \end{aligned}$$

Questions and tasks:

1. Which time evolution for n do you observe? Does it make sense, given the fluxes above?
2. Feel free to try out or generalize your script with:
 - different initial values
 - different fluxes
 - an explicit formula for the flux, for instance

$$j_x(t, x) = 8x^3 - 60x^2 + 118x - 45$$

- time-dependent fluxes
- more cells
- two or three dimensions, including y - or z -fluxes

URLs for chapter 7

1. <https://mathworld.wolfram.com/EulerForwardMethod.html>

Chapter 8

Matter

Matter is probably the easiest quantity to grasp intuitively; it's what we can really call "stuff". It is important to clearly distinguish matter from *mass*: mass can be considered a property of matter, but the two are different: for example, the mass of some amount of matter may change, even if the amount of matter stays the same.

There are different kinds of matter, each of which satisfies a balance relation. The distinction into different kinds depends on the physical theory. In most everyday situations, the distinction corresponds to different [chemical elements¹](#), and each satisfies its own balance. These balances are the basis of [stoichiometry²](#). If we observe phenomena such as nuclear fission or fusion, however, we notice that the balances of chemical elements are not really satisfied. With such phenomena we make a different distinction of types of matter, for instance *baryons* and *leptons*, and each satisfies again its own balance. It is unclear whether these balances might be broken in other physical phenomena at smaller scales. In these notes we shall usually consider chemical elements as the different kinds of matter, making some exceptions in discussion of nuclear phenomena.

8.1 Electric charge and magnetic flux**8.2 Magnetic flux****8.3 Energy****8.4 Momentum****8.5 Angular momentum****8.6 Entropy**

URLs for chapter 8

1. <https://doi.org/10.1351/goldbook.C01022>
2. <https://doi.org/10.1351/goldbook.S06026>

Chapter 9

Thermodynamics

9.1 Notes on “quasi-static” processes

In general, none of the statements "reversible \Rightarrow quasi-static" or "quasi-static \Rightarrow reversible" is true.

A counterexample to the second implication are systems with internal state variables, which cannot be made non-dissipative, no matter how slowed-down they are. See the discussion and mathematical analysis in Astarita § 2.5.

A counterexample to the first implication is a system of spins in a crystal lattice. It is possible to *reversibly* bring the system from an equilibrium state to another with opposite temperature by reversing the external magnetic field *as fast as possible* – and therefore *not* through a quasi-static process. In fact it is key here that the process be *not-quasi-static*, but as fast as possible, because a slow change of the external magnetic field would lead to an irreversible process with dissipation. For more details see the discussion in Buchdahl, Lecture 20.

The point is that for some systems a *fast* change can actually prevent the onset of dissipative phenomena, and so the process needs to be fast if we want it to be reversible. Adiabatic processes often also need to be fast (as a curious historical fact, Truesdell & Bharatha, Preface p. xii, remark that “In introducing what we today call an ‘adiabatic process’, Laplace called it ‘a sudden compression’, in which he was followed by Carnot”).

In fact, clearly non-quasi-static phenomena such as *explosions* can in some circumstances be described by *reversible* processes! This is possible if the explosion involves many shock waves, as explained by Oppenheim, chap. 1 p. 63:

If there is more than one shock, the losses in available energy are diminished, so that in the limit, with an infinite number of shocks, they become negligible, and the process acquires the character of a thermodynamically optimal, i.e., reversible, change of state. The study of explosion processes reveals that, indeed, they are associated not with one but with a multitude of shocks.

For explosions see also the mathematical analysis by Dunwoody: *Explosion and implosion in a mixture of chemically reacting ideal gases*, where again reversible-process equations are used.

A caveat about reversible and quasi-static associations is given by Erickson (§ 1.2):

Some tend to associate nearly reversible processes with those taking place very slowly – the "quasi-static" processes. This probably stems, at least in part, from experience with classical theories of heat conduction, viscosity, and so on. However, a ball made of silly putty behaves almost reversibly when bounced rapidly and various other high polymers have similar predilections. So, it seems prudent to be open-minded in considering what may be reversible processes for particular systems.

He later discusses (§ 3.1) the case of bars subjected to dead loads, for which we can have reversible processes under sudden jumps in elongation. He concludes (p. 46) that "the sudden jump provides an example of a process that is reversible but not reasonably considered to be quasi-static".

But there's an important question that underlies our discussion: what do we actually mean by "quasi-static"? We need to specify a time scale, otherwise the term is undefined. For example, a geological process (say, tectonic motion) can be considered as quasi-static – or even completely static – on time scales of minutes or days; but it is not quasi-static on time scales of millions of years.

Whether a process is reversible or not, within any tolerance needed, is an experimental question. We can measure any relevant quantities, say pressure p and exchanged heat q , under the process, and compare them with those, p^* and q^* , determined by the equations for a reversible process. We may find for example that at all times

$$\left| \frac{p - p^*}{p^*} \right| < 0.001 , \quad \left| \frac{q - q^*}{q^*} \right| < 0.001$$

and conclude that the process is reversible, if relative discrepancies of 0.1% or less are negligible in our concrete application.

But suppose that someone tells us "if you want the process to be reversible, you must make sure that it is quasi-static". Alright, but how much is "quasi-static"? Is it OK if the piston moves with a speed of 1 cm/s? or is that too much? How about 1 mm/s? – In fact we may find that for some kind of fluid 1 cm/s is absolutely acceptable for the process to be reversible, whereas for another kind of fluid that speed would lead (at the same temperature) to an irreversible process.

You see how this imprecise situation can lead to circular definitions: "if the process is irreversible, then it means it isn't quasi-static" – but then we are actually *defining* "quasi-static" in terms of "reversible"! Any statement of the kind "reversible \Rightarrow quasi-static" or "quasi-static \Rightarrow reversible" then becomes not a matter of experimental verification, but of pure *semantics*. At this point we can simply get rid of "quasi-static" terminology since it doesn't bring any new physics to the table. This circularity is admitted for example by Callen in discussing irreversible gas expansion (Problem 4.2-3 p. 99):

The fact that $dS > 0$ whereas $dQ = 0$ is inconsistent with the presumptive applicability of the relation $dQ = T dS$ to all quasi-static processes. We define (by somewhat circular logic!) the continuous free expansion process as being "essentially irreversible" and *non-quasi-static*.

A similar criticism can be read in Astarita, § 2.9, p. 62, where he also provides a mathematical quantification of quasi-static, similar to the one given above for reversibility:

Often this point is circumvented by bringing in another difficult concept, that of a quasi-static transformation, which proceeds "through a sequence of equilibrium states." Quasi-static is an impressive word, but the only meaning which can be attached to it is the less impressive word "slow" – and how can one speak of slowness without implying the concept of time? How slow is slow enough? If one chooses to develop a thermodynamic theory (rather than a thermostatic one), the answer is easy. For instance, in the case of a system where the state is V, T, \dot{V} [the latter is the rate of change of V], one needs to assume

that [the non-equilibrium pressure] $p(V, T, \dot{V})$ is a Taylor-series expandable at $\dot{V} = 0$ to obtain [that

$$p = p^* + \frac{\partial p}{\partial \dot{V}} \Big|_{\dot{V}=0} \dot{V} + O(\dot{V}^2),$$

where $p^* = p(V, T, 0)$ is the pressure at equilibrium]. One then reaches the conclusion that if the condition

$$\dot{V} \ll \frac{p^*}{\partial p / \partial \dot{V} \Big|_{\dot{V}=0}}$$

is satisfied, then indeed the difference between p and p^* is negligibly small as compared to p^* , and thus the process can be regarded as a quasi-static one.

Criticisms against the fuzzy notion of "quasi-static" have appeared in many other works. Truesdell & Bharatha (Preface p. xii), make the historical remark that "the 'quasi-static process' was barely mentioned for the first time in 1853 and was altogether foreign to the early work [in thermodynamics]". See also the mathematical analysis by Serrin: *On the elementary thermodynamics of quasi-static systems and other remarks*.

I also want to point out that "quasi-static" in some works has specific meanings somewhat unrelated to the discussion above. For example that the rate of increase of the total kinetic energy K of the system is negligible, so that the law of energy balance, which in its full generality is

$$\frac{d(U + K)}{dt} = Q + W$$

(that is, the rate of increase of internal energy U and kinetic energy is equal to the heat rate Q and work rate W provided to the system) can be approximated by

$$\frac{dU}{dt} = Q + W.$$

Or that similar inertial terms in the motion of the system are negligible. See for example the book by Day, chap. 2.

But note that such definitions of "quasi-static" have, again, *no* a-priori relation with reversibility.

Finally, the equation $dS = Q/T$ is only valid for a process that is:

- reversible (by definition),

- closed (no exchange of mass),
- with a homogeneous surface temperature,
- without bulk heating (such as instead happens in a microwave oven).

Under the last three conditions we have in general that $dS \geq Q/T$; when the equality sign is satisfied, then the process is *defined* as reversible. See Astarita, § 1.5, or Müller & Müller, for the different forms of the second law under different circumstances. This equation may be valid in quasi-static and non-quasi-static processes, as explained above.

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(“de X” is listed under D, “van X” under V, and so on, regardless of national conventions.)

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