



# An Experimental Study on the Aerodynamic and Aeroacoustic Performances of a Bio-Inspired UAV Propeller

Zhe Ning<sup>1</sup>, Hui Hu<sup>2</sup>(✉)

*Department of Aerospace Engineering, Iowa State University, Ames, Iowa, 50011*

Under rapid development, the small UAV is facing two essential problems: aerodynamics efficiency and noise emission. In this work, the aerodynamic and aeroacoustics characteristics of a novel bio-inspired UAV propeller with the unique planform shape inspired by the cicada wing and maple seed was experimentally investigated by comparing with conventional baseline propeller at hover flight condition. The design thrust (3N) and solidity (0.12) are kept same for both propellers. The aerodynamics forces and sound measurements revealed that, compared with baseline propeller, the bio-inspired propeller could generate equal thrust and emit lower noise under constant power input. In addition, the flow field measurements obtained from Partial Image Velocimetry (PIV) illustrated that the bio-inspired propeller produced a smaller wake region and demonstrated a fast decay rate of tip vortex strength compare to baseline propeller.

## Nomenclature

$A$	=	Swiping area of the propeller
$BPL$	=	Blade passing frequency
$C_0$	=	Sound speed
$c$	=	Chord length of the propeller
$c_{tip}$	=	Chord length at the tip of the propeller
$c_r$	=	Chord length at corresponding radius of the propeller
$D$	=	Diameter of the propeller
$Re_c$	=	Reynold number at corresponding chord length
$T$	=	Thrust
$p$	=	Sound pressure
$p_{ref}$	=	Sound pressure reference
$R$	=	Radius of the propeller
$r$	=	Non-dimensional radial distance
$SPL$	=	Sound pressure level
$U$	=	Relative velocity at certain chord
$U_T$	=	Rotational velocity at certain chord
$U_P$	=	Induced velocity at certain chord
$V_{tip}$	=	tip speed
$x$	=	Axial coordinate
$y$	=	Vertical coordinate
$z$	=	Transverse coordinate
$\Omega$	=	Rotational speed of the propeller
$\alpha$	=	Angle of attack
$\theta$	=	Twist angle
$\Phi$	=	Inflow angle
$\rho$	=	Density of air

<sup>1</sup> Graduate Student, Department of Aerospace Engineering.

<sup>2</sup> Martin C. Jischke Professor of Aerospace Engineering, AIAA Associate Fellow, Email: [huhui@iastate.edu](mailto:huhui@iastate.edu)

## I. Introduction

In the past twenty years, the Unmanned Aerial Vehicle (UAV) has developed rapidly due to the reduction of the sizes and costs of small electronic devices, such as processors, sensors, and batteries<sup>1</sup>. Specifically, the rotary-wing system has become a popular configuration due to its hovering ability and vertical take-off and landing (VTOL) motion<sup>2</sup>. Fitted with these attributes, the rotary-wing UAV has been put to use in a variety of applications (e.g., video taking & mapping, building & construction, delivery, rescue operation, and personal entertainment).

During the last two decades, the development of the rotary-wing UAV was focused on control theory and functional development. In a study performed by Achtek<sup>3</sup>, the flying motion of a quadcopter was controlled by visual feedback and sensor measurements to accomplish autonomous flight. Mellinger<sup>4</sup> studied real-time trajectories and controllers to drive a quadcopter to execute aggressive maneuvers. Lindsey<sup>5</sup> used teams of quadcopters to construct buildings, and Mirjan<sup>6</sup> used flying robots to build a bridge. Although the control theory and application has been well developed, less emphasis has been placed on the aerodynamic performance. Currently, the operational time of rotary-wing small UAV is lower than one hour, which is too short for many applications, such as delivery, inspection, and recording video. Therefore, the aerodynamic efficiency, especially the propeller efficiency, is an essential consideration for the design progress providing the known constraints of the battery charge and motor power. The numerical work done by Bristeau<sup>7</sup> illustrated that the flexibility of the propellers has significant effect on the quadcopter dynamic performance, under the assumption of constant local angle of attack (AOA) and induced velocity. Since the rotary-wing UAV always consists of multiple rotors, with small separation distances between rotors, numerical analyses have proved that the strong flow interaction will result in up to 5% thrust decrease<sup>8</sup>. During flight, the rotor will be affected by tip vortices generated from the rotor itself, as well as neighboring rotors<sup>9</sup>. For current small UAV propellers, a larger gap between rotors must be maintained in order to make up the thrust deficit due to the flow interaction. In other words, the size of the rotary-wing UAV will increase. A significant result is that the bending strength at the joint between arm and fuselage will increase. Furthermore, the weight of the UAV will slightly increase, which may balance out or reduce the total efficiency. Besides the aerodynamic constraints, the noise generated from the operating rotary-wing UAV is another essential problem, which will have significant impact on human health, as well as wild animals. In a study executed by Dittmer<sup>10</sup>, the noise generated from a UAV increases the heart rates of bears, which indicates raised levels of stress. According to Leslie<sup>11</sup> and Sinibaldi<sup>12</sup>, noise reduction of UAVs could broaden their mission ranges. Therefore, the noise generated from rotating propeller need to be considered during the design process.

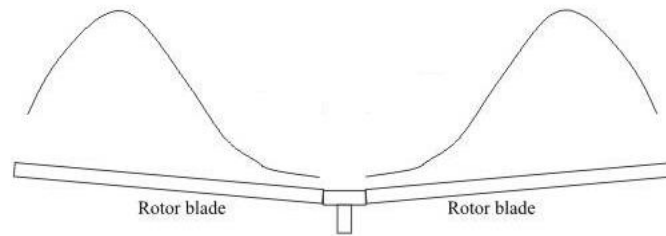
Understanding that the propeller is the main component of lift generation and aerodynamic noise, the focus of this study is to design a UAV propeller with the same power input that can achieve the same or improved aerodynamic performance, while reducing noise levels. In determining a potential solution, nature always provides inspiration to researchers and engineers. As shown in Figure 1, cicada wings and maple seeds have their planform different from traditional propellers. Instead of having a tapered planform shape, they have large chord length in the mid-span region. Most of the thrust for a rotary wing is known to be generated between the 50% to 90% radius, as shown in Figure 2. The planform of insect wings and maple seeds seems to have a better match with the lift distribution, where the chord length is large at high lift region. Beside the geometry, another important factor is the Reynolds number. At Reynolds numbers ranging from 10,000 to 100,000, the lift/drag ratio dramatically increases as the Reynolds number increases. As shown in figure 3, the lift/drag ratio at a Reynolds number of 100,000 is almost twice that of when the Reynolds number is 50,000 for the same E63 airfoil (a typical low Reynolds number airfoil)<sup>13</sup>. The increasing chord length will increase the chord Reynolds number. As a result, the propeller will operate in the regime with a better lift/drag ratio if this configuration of planform is used into design.

With this in mind, we conducted an explorative study to evaluate the aerodynamic and aeroacoustic performances of a novel bio-inspired UAV propeller with the unique planform shape inspired by the cicada wings and maple seeds. The bio-inspired UAV propeller was designed to have the same planform area and the same airfoil cross sectional shape, thus yielding the same weight as a conventional tapered UAV propeller. It was expected that the bio-inspired UAV propeller will have comparable or even better aerodynamic and aeroacoustic performances in comparison to the conventional tapered UAV propellers. In the present study, a conventional tapered UAV propeller was used as the baseline for the comparative study, and the aerodynamic forces generated by the bio-inspired UAV propeller under hovering conditions will be measured quantitatively by using a high-sensitive force/moment transducer (JR3 load cell). The aeroacoustic performance of the bio-inspired propeller was characterized in an anechoic chamber. Furthermore, a high-resolution digital Particle Image Velocimetry (PIV) system was also used to achieve detailed flow field measurements to quantify the evolution of the unsteady vortex structures in the wake of the propellers. The

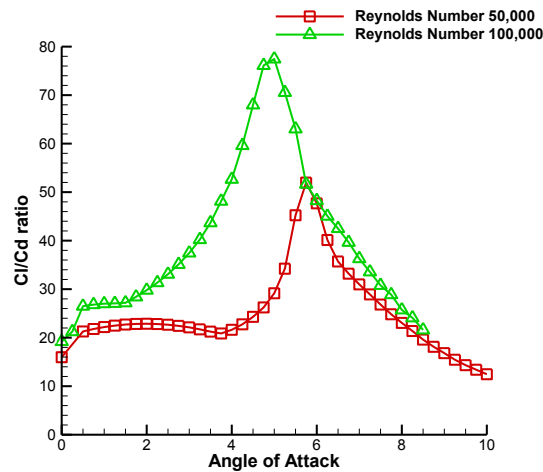
aerodynamic force and aeroacoustic measurement results were correlated to the detailed flow field measurements in order to elucidate the underlying physics.



**Figure 1. Wing planform of cicada and maple seed**



**Figure 2. Aerodynamic force distribution of a rotor.**



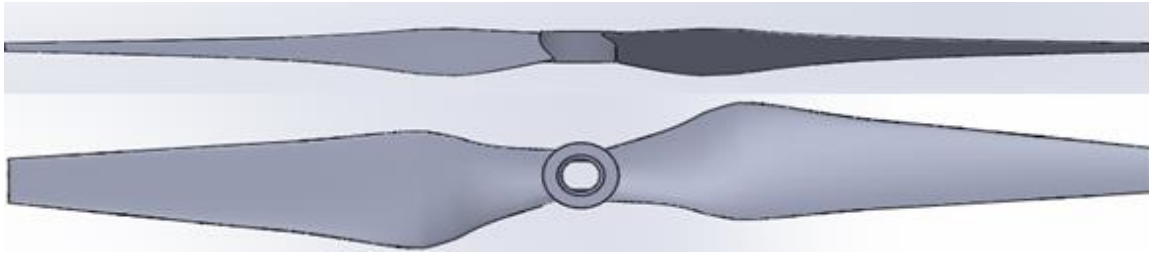
**Figure 3. lift/drag ratio of a low-Reynolds E63 airfoil vs AOA at Re=50,000 & Re=100,000**

## II. Propeller Design and Experimental setup

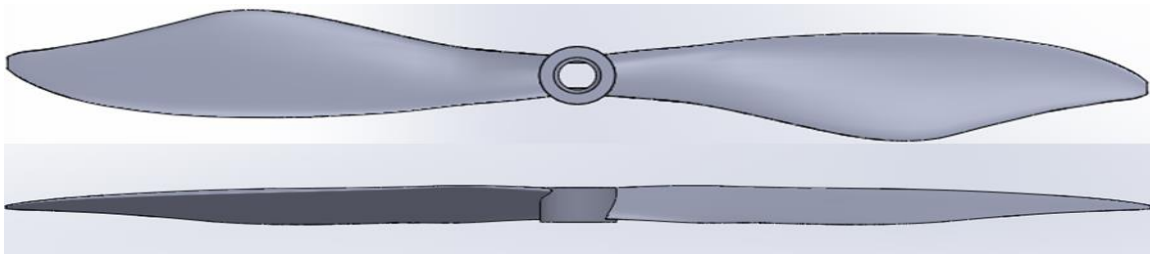
### 2.1 Propeller design

As shown in figure 4, in order to accomplish a designed thrust of 3N, an E63 airfoil shape was used to design the baseline propeller. The thickness of the airfoil was doubled based on the camber line due to a strength concern. The propeller was 240mm in diameter and 11mm at the chord tip. The chord lengths along the propeller radius from tip to 30% radius were calculated by the equation  $C_r = \frac{C_{tip}}{r}$ , where  $C_r$  is the chord length at corresponding radius location, and  $r$  represents a non-dimensional radial distance, which is 0 in the rotating center and 1 at tip. The blade twisted 11.6 degrees at the tip to 26.3 degrees at the 30% radius. The solidity was 0.12 which is widely used for conventional propeller. As shown in figure 5, the airfoil shape, designed thrust and rotational speed, and solidity used to design the bio-inspired propeller were kept the same as the baseline propeller design. This is with the notable exception of the design of the propeller planform, which was derived from a real maple seed configuration with a reduced aspect ratio.

Both propellers were made of a hard plastic material (i.e., VeroWhitePlus, RGD835, manufactured by Stratasys, Inc.) by using a rapid prototyping machine (i.e., Connex 3D printer).

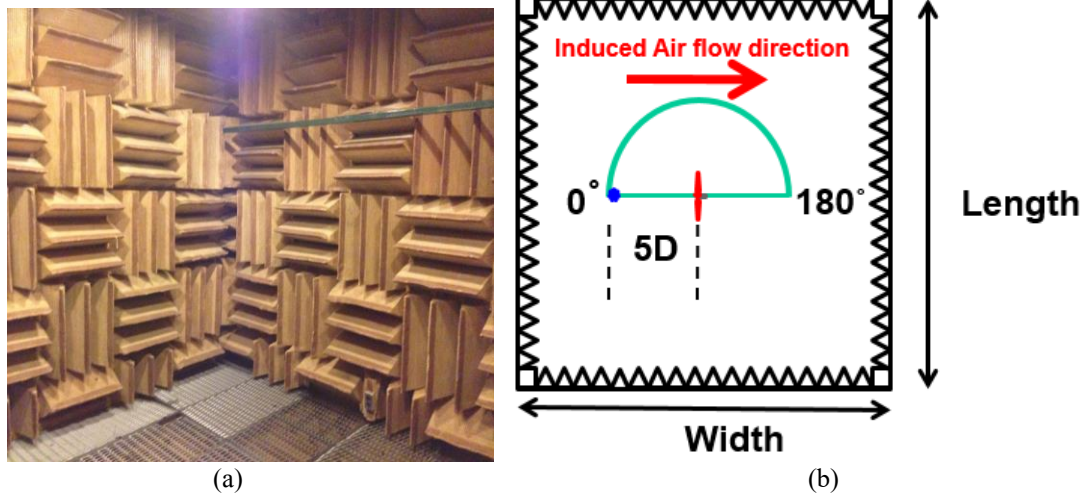


**Figure 4. Baseline propeller**



**Figure 5. Bio-inspired propeller**

## 2.2 Experimental setup



**Figure 6. Anechoic chamber and sound measurement sketch**

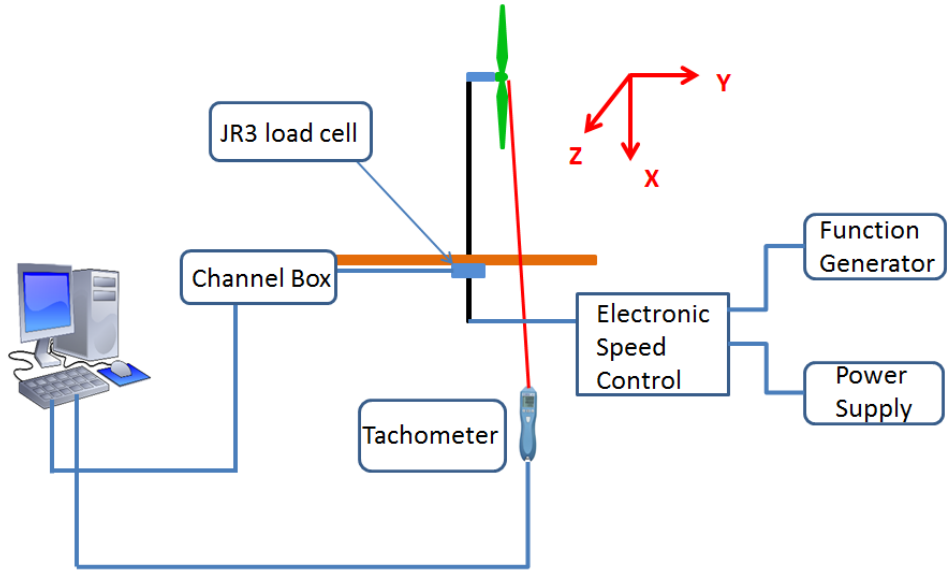
The noise at the hover flight condition was measured in an anechoic chamber (as shown in Figure 6 (a)) located in the Department of Aerospace Engineering at Iowa State University. The chamber has dimensions of 12\*12\*9 feet with a 100 Hz cut-off frequency. Figure 6 (b) shows the sketch of the sound measurement. The red arrow represents the induced airflow direction, and the blue points represent the test locations. The spectrum measurement was conducted at 0 degrees and 5D away from the propellers. In the near field, closer than 5D, the turbulence level is high, which would affect the sound measurement. 5D is the standard distance that can be considered a far field distance.

Figure 7 shows the schematic of the experimental setup to be used in the present study to measure the dynamic force of the both propellers. A 0.2 diameter clearance gap was maintained between propeller and support tower. The power used to drive the motor was tabulated by a direct current power supply. The UAV propeller was connected to a function generator via an electronic speed control. The rotational speed of the propeller was controlled by adjusting

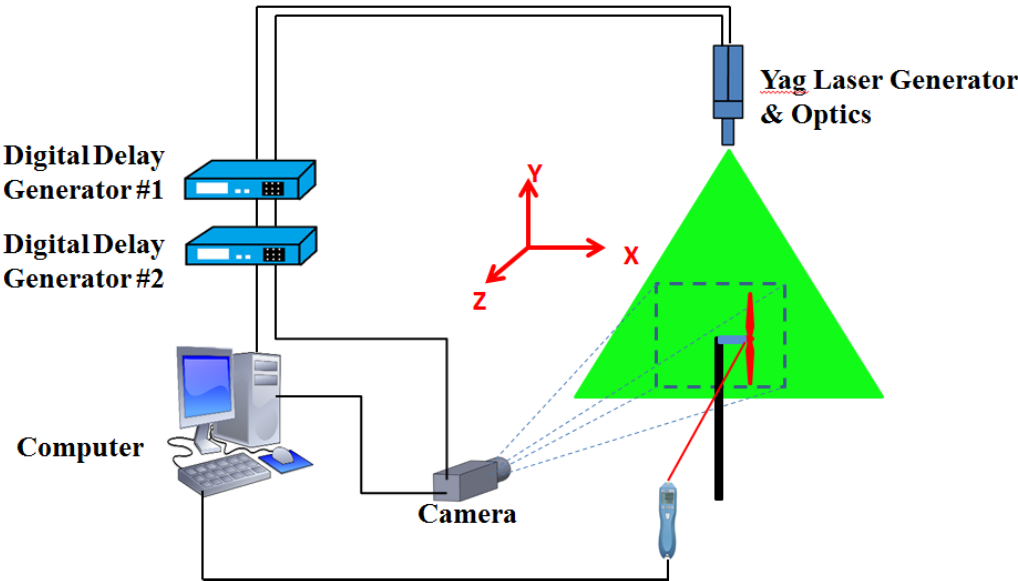
the duty cycle of the output signal from the function generator. During the experiment test, the rotational speed of the propeller was measured by tachometer. A high-sensitivity force-moment sensor (JR3 load cell) was used to measure the force and moment, and the precision of this sensor for force measurements was  $\pm 0.1\text{N}$  ( $\pm 0.25\%$  of the full range). Different types of propellers (baseline & Bio-inspired) were tested under the conditions shown in Table 1.

**Table1. Test conditions**

Blade type	Rotational direction	Voltage Condition (V)	Rotational speed (RPM)
Baseline propeller	CCW (Counter Clock-wise)	11.1	0-6000
Bio-inspired propeller	CCW (Counter Clock-wise)	11.1	0-6000



**Figure 7. Experimental setup for the dynamic measurements of propeller**



**Figure 8. Experimental setup for the flow field measurement**

In addition to the aerodynamic force measurements, a high-resolution digital PIV system was used to measure the flow field in order to quantify the detailed flow structures around the propeller. An interrogation window with  $32 \times 32$  pixels and a 50% effective overlap were used to calculate the instantaneous velocity vectors. In order to determine the ensemble-average wake flow statistics and detailed flow structure at certain phase angles, both “free-run” and “phase-locked” PIV measurements were conducted in this study. The ensemble-average results were obtained from 1000 frames of instantaneous PIV measurements for free-run test. For the “phase-locked” PIV measurements, the phase-averaged flow velocity at certain phase angle range from 0 to 330 degrees subdivided into 30-degree increments were calculated from 255 frames of the instantaneous PIV measurements. Figure 8 shows the schematic of the PIV system used in this study. The air was seeded with  $\sim 1 \mu\text{m}$  water based droplets by using fog generation. The field illumination was provided by a Nd:YAG laser generator with power 200 mJ per pulse at the wavelength of 532 nm. The laser sheet thickness was set to around 1 mm at test section. A high-resolution ( $2048 \times 2048$  pixels) charge-coupled device (CCD) camera with axis perpendicular to the laser sheet was used to capture the raw image. The camera and laser generator were connected to the computer via a digital delay generator, which controlled the timing of the image capture and laser illumination. A second digital delay generator was used to modify the signal, in order to achieve the “phase-locked” PIV measurement.

### III. Results and Discussion

#### 3.1 Aerodynamics performance results

The aerodynamic force measurements were conducted at different rotational speeds ranging from 0 to 6000 RPM for both propellers. Figure 9 presents the comparative results, where the design thrust is denoted by dashed line, and the results of baseline and bio-inspired are shown in red and black lines, respectively. This color line notation is also used in later comparisons. The symbols indicate the experimental data, and the lines are the best curve fits. Each data point is the average value of 30 second measurement with 1000Hz sampling frequency. In figure 9 (a), it should be noted that the thrust grows with the input power increasing. Compared with the baseline propeller, the bio-inspired propeller required same power to generate the design thrust. However, the rotational speed of the bio-inspired propeller was 3.5Hz lower than the baseline propeller at hover motion (thrust = 3N). This indicates that the thrust coefficient of bio-inspired propeller is higher than the baseline propeller. The thrust coefficient is defined by

$$C_T = \frac{T}{\rho A V_{tip}^2} \quad (1)$$

It can be found in figure 9 (c), the thrust coefficient of bio-inspired propeller is about 10% higher than the baseline propeller.

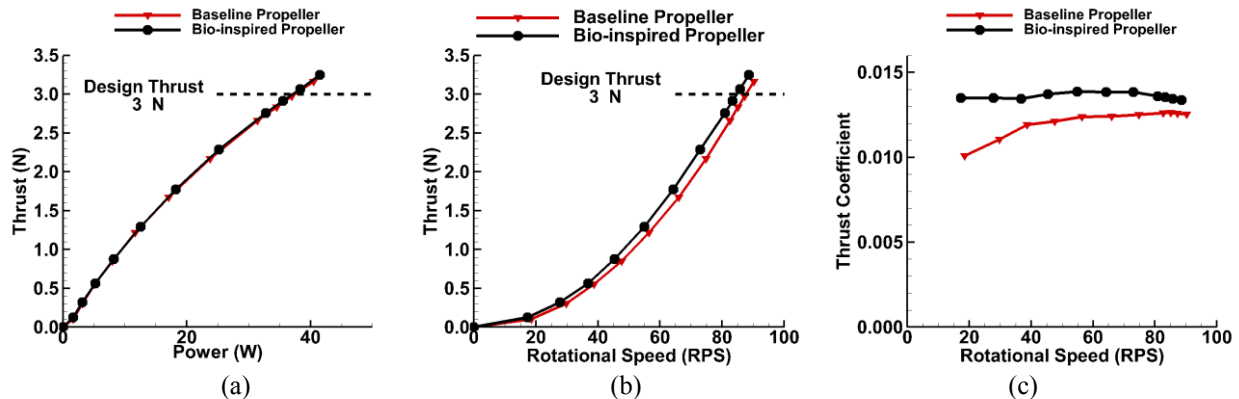


Figure 9. Force measurement results of baseline and bio-inspired propellers

#### 3.2 Sound measurement results

As mentioned before, the sound measurements were conducted at hover motion, thus the thrust was kept constant (3N) for both propellers. Hover is unique flight condition for vertical take-off and landing aircraft, where the thrust generated from propeller equal to the aircraft weight. At this flight condition, the forward and vertical flight speed is



zero. The experimental conditions at hover are listed in table 2. The aeroacoustic performance was characterized by sound pressure level (SPL) at different frequency, which calculated by

$$SPL (dB) = 20 \log_{10} \left( \frac{p}{p_{ref}} \right) (dB) \quad (2)$$

The reference pressure is  $2 \times 10^{-5}$  Pa. Compared with baseline propeller, the bio-inspired propeller could reduce noise up to 4 dB in the human hearing domain (20Hz~20,000Hz). As shown in figure 9 (a), the bio-inspired propeller generates lower SPL than the baseline propeller, which means it has the potential to reduce tonal noise as well as broadband noise. Figure 9 (b) reveals the sound spectra comparison results at the low frequency domain. The frequency axial is non-dimensionalized by blade passing frequency (BPF), which is defined by

$$BPF = \text{rotational frequency} \times \text{blade number} \quad (3)$$

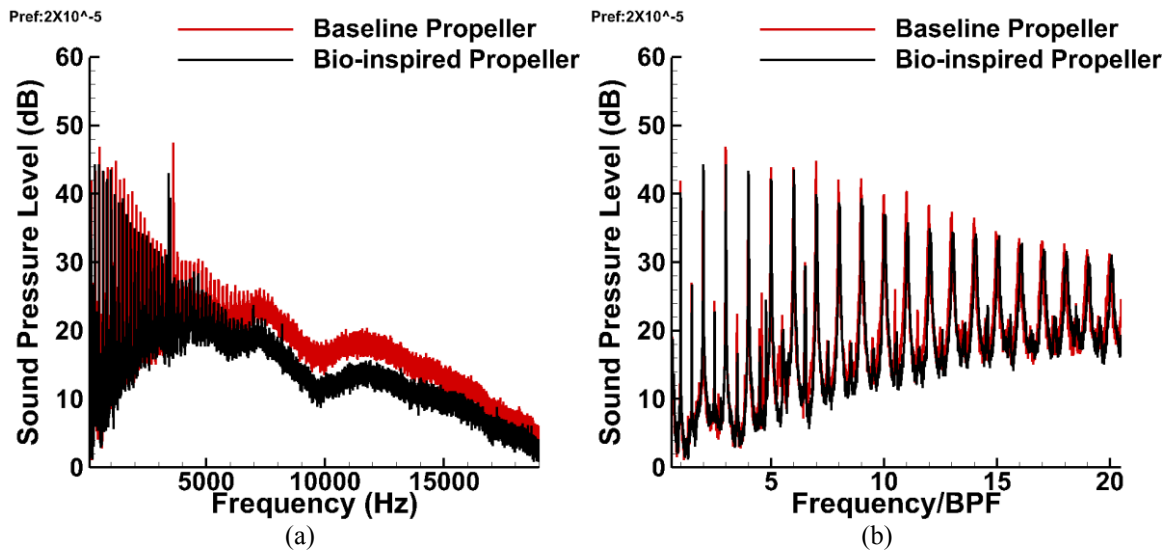
All of the peak values shown in spectra are the tonal noise components, which are associated with BPF. The tonal noise consists of loading noise and thickness noise. The loading noise is generated from the force variation and the thickness noise comes from mass variation, which has been illustrated in terms 1 and 2 on right hand side of Lighthill's acoustic analogy equation

$$\left( \frac{\partial^2}{\partial t^2} - C_0^2 \frac{\partial^2}{\partial x_i^2} \right) \rho = \frac{\partial m}{\partial t} - \frac{\partial f_i}{\partial x_i} + \frac{\partial^2 T_{ij}}{\partial x_i \partial x_j} \quad (4)$$

At Mach number lower than 1, the loading noise dominates the tonal noise. Compared with baseline propeller, the reduction of the tonal noise of the bio-inspired propeller could be easily observed. This noise attenuation could be attributed to the small force variation of the bio-inspired propeller. As shown in table 2, the thrust standard deviation of the bio-inspired propeller is 24% lower than the baseline propeller. Another possible reason is the lower rotational speed. By applying the free-space Green's function to equation 3, and do a scale analysis, it was found that the thickness noise is proportional to the  $U^2$ , the loading noise is proportional to the  $U^3$ , and the noise due to Lighthill's stress tensor is proportional to the  $U^4$ . Here, the  $U$  is the flow velocity. At hover motion, the rotational speed of bio-inspired propeller is low, which benefits the noise reduction.

**Table 2. Flight condition at hover motion**

	Rotational speed (RPS)	Mean Thrust (N)	Thrust standard deviation (N)
Baseline propeller	86.5	3	0.278
Bio-inspired propeller	83	3	0.211



**Figure 11. Sound spectra comparison between the baseline propeller and bio-inspired propeller**

### 3.3 “Free-run” PIV measurement results

In addition to the force and sound measurements, the flow field measurements downstream of the propellers were achieved through a high-resolution PIV system. Akin to the sound measurements, the flow structure measurements were also conducted at hover motion, where the flow field is azimuthally axisymmetric. The ensemble-averaged velocity distribution results of a typical “free-run” PIV measurement of both the baseline and bio-inspired propellers are shown in figure 12. The term “free-run” means the images taken were not harmonic with the propeller rotation. The X and Y axes were normalized based on the propeller radius (120mm). The induced airflow direction is denoted by the red arrow. The large blue region is the quiescent flow region, where the flow velocity is low. It indicates the rotating propeller did not affect this area significantly. The region above propeller is the inflow region. Due to the propeller rotation, the flow was sucked into the rotating plane. Another region is the induced flow region, also known as slipstream<sup>14</sup>. At there, the velocity is higher compared to other two regions. This is due to the propeller pushing the air. The pressure difference between inflow and induced flow region resulted in a thrust generated from propeller. The black lines shown in the plot are the streamlines. The densely streamlined area is the flow boundary between quiescent and induced flow region.

Compared with the baseline propeller, the flow speeds at quiescent and inflow region of bio-inspired propeller were almost the same. However, the velocity distribution of these two propellers were different at the induced flow region. The bio-inspired propeller had larger velocity gradient from flow boundary towards rotating axis compared with baseline propeller. Figure 13 shows the velocity profile at different downstream locations (0.166R, 0.5R, and 1R). The induced flow velocities downstream were almost same from tip to 0.75R for both propellers. Compared with baseline propeller, the bio-inspired propeller generated higher induced flow velocity in the region between 0.5R and 0.75R, where it has larger chord length. Meanwhile, the induced velocity of the bio-inspired propeller was lower than baseline propeller from center to 0.5R. Although the velocity distributions are different, the mass flow rate of both propellers were same, which are confirmed by the force measurement results.

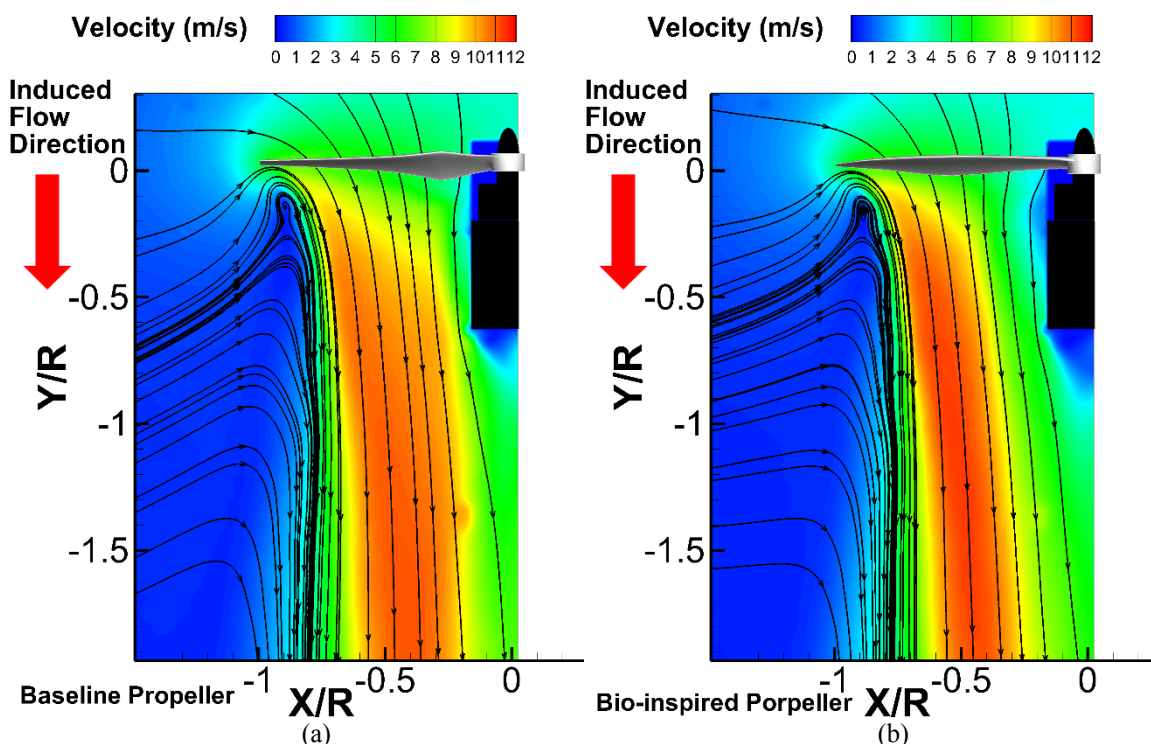


Figure 12. “Free-run” PIV measurement results\_ensemble-average velocity distribution

Based on the rotational speeds and chord lengths at different span-wise locations of both propellers, the chord Reynolds number distributions along the blade could be obtained. As shown in figure 15 (a), the bio-inspired propeller operated at higher AOA from 0.4R to 0.9 R compare to baseline propeller, which is mainly due to the large chord length at these locations. Because of the small chord length at the tip, the bio-inspired propeller operated at a lower AOA compare to baseline propeller. With the induced velocities provided by PIV measurements and the rotational



speeds at hover motion, the operating AOA at different span-wise locations could be calculated. It was discovered in figure 15 (b) that the bio-inspired propeller operated at almost constant AOA, which was around 6 degrees. The operating AOA of baseline model varied from 11 to 4 degrees from the root to 0.85R. For both propellers, there is a large jump from 0.85R to 0.9R, which is mainly due to the large gradient of the induced velocity. From 0.85R to 0.9R, the induced velocities dropped from around 8.4 m/s to about 3.7 m/s for both propellers, and the rotational speed increased from 54 m/s to 60 m/s. According to blade element diagram (figure 14), the inflow angle of attack ( $\Phi$ ) decreased from 9 degrees to 3.5 degrees. Since the twist angle ( $\Theta$ ) varied smoothly, the sudden change of the inflow AOA would lead to the operating AOA ( $\alpha$ ) jump. The  $U$ ,  $U_P$ , and  $U_T$  represent the relative velocity, induced velocity, and rotational velocity, respectively.

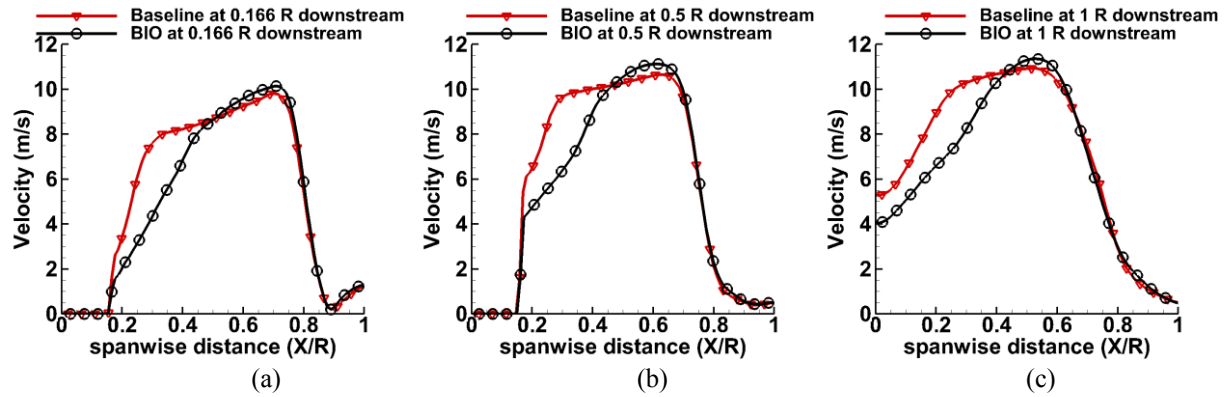


Figure 13. Velocity distribution comparison at different downstream location

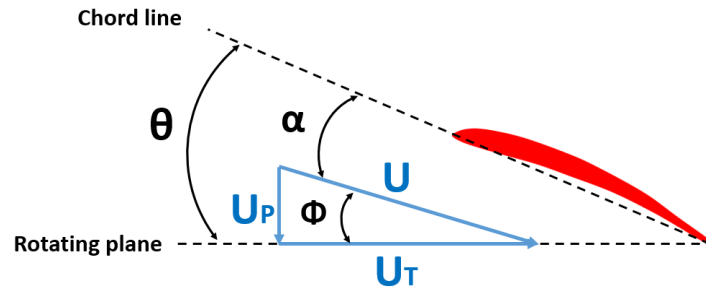


Figure 14. blade element

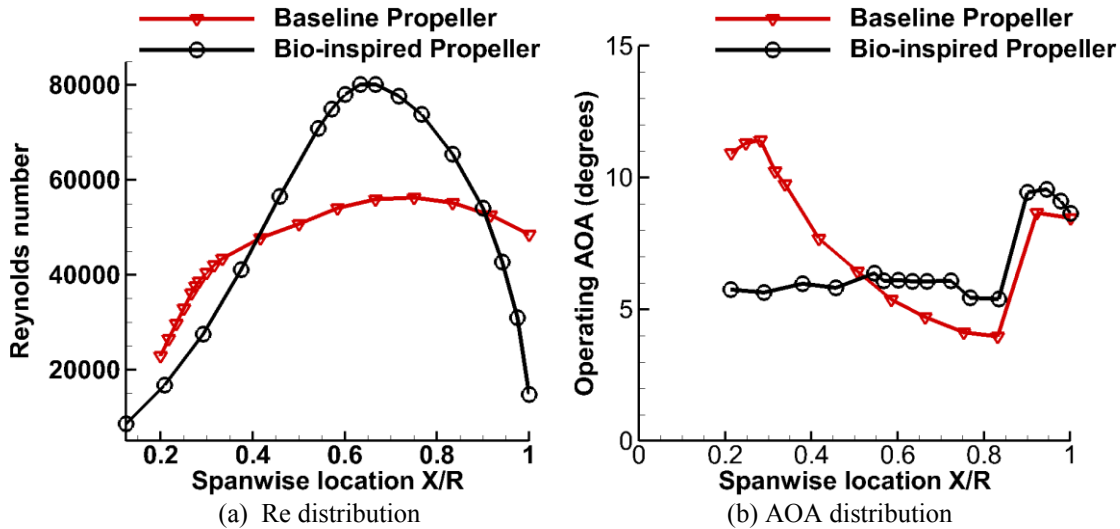


Figure 15. Reynolds number and operating AOA along the propellers

### 3.4 “Phase-locked” PIV measurement results

A “phase-locked” PIV measurement was also conducted at hover motion to illustrate the detailed flow characteristics of the downstream flow for the two different types of propellers. The term “phase-locked” means the images taken were harmonic with propeller rotation, so that the propeller appears in the same orientation every time an image is obtained. The phase angle was defined as the angle between position of a pre-marked blade and the measurement plane. In this study, the phase angles ranged from 0 to 330 degrees with incremental changes of 30 degrees, which could provide detailed information of the wake generation and translation within a full cycle of propeller rotation.

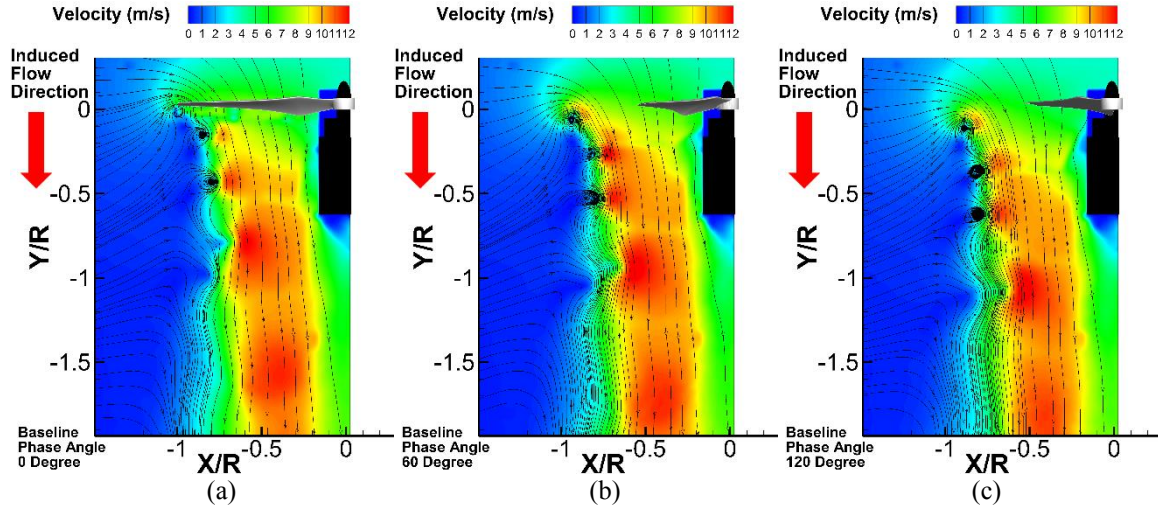


Figure 16. “Phase locked” PIV measurement results\_velocity distribution of baseline propeller

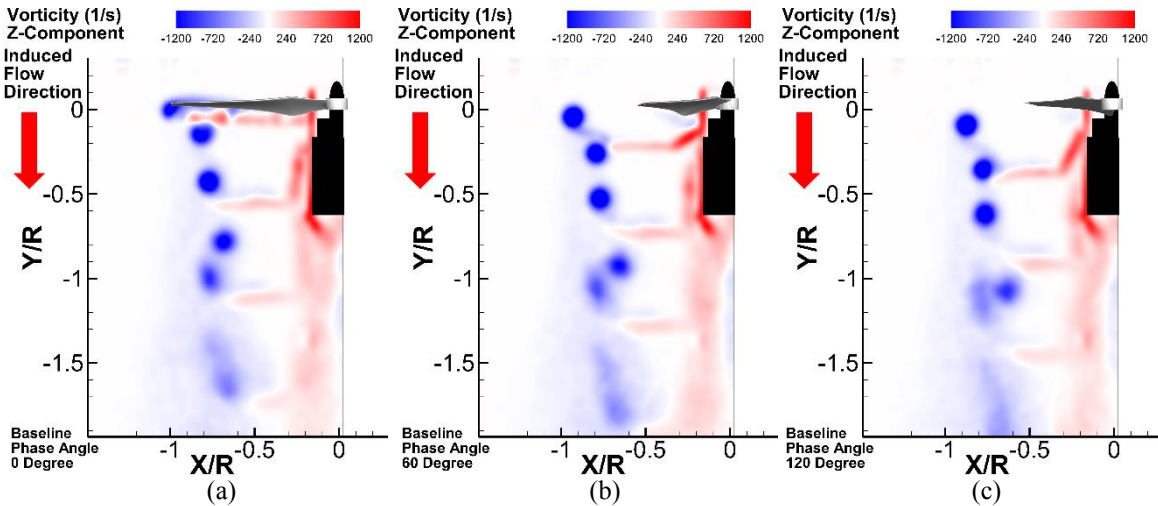


Figure 17. Phase locked PIV measurement results\_vorticity distribution of baseline propeller

Figure 16 shows the velocity distribution of the “phase-locked” PIV measurement results of the baseline propeller at 0, 60 and 120 degree phase angles. The 0 degree phase angle indicates that the propeller was just crossing the measurement plane. The three flow regions (inflow, induced flow and quiescent) can be visualized clearly in these plots, which are similar to the results achieved from the “free-run” PIV measurement. One phenomenon which was not captured by the “free-run” PIV measurements was the periodic velocity deficit close to the boundary between the induced flow and quiescent regions, which is due to the blade tip cutting flow and results in the formation of the tip vortices. Several strip shape, low-speed areas within the induced flow region also were not observed in the “free-run” PIV measurement. This phenomenon is due to the propeller cut-in. The circular motion of the streamline represents the tip vortex structure. The vorticity distribution is displayed in figure 17. The tip vortex structure is shown clearly in figures (a) through (f). With the phase angle increasing, the tip vortex structure shrinks inward and travels towards the downstream direction. The positive and negative vorticity region located in the induced flow area is due to the

flow passing the trailing edge. Due to the velocity difference, these vortices travel faster compared to tip vortices. As a result, the coming tip vortex starts to interact with the previous one at around the 70% radius downstream. This phenomenon is generally observed in rotor hover condition. Through flow visualization, Leishman<sup>14</sup> also observed this phenomenon in the wake of hovering helicopter blade. After the interaction, a pair of tip vortices will form one. The strength of it dropped dramatically and dissipated quickly.

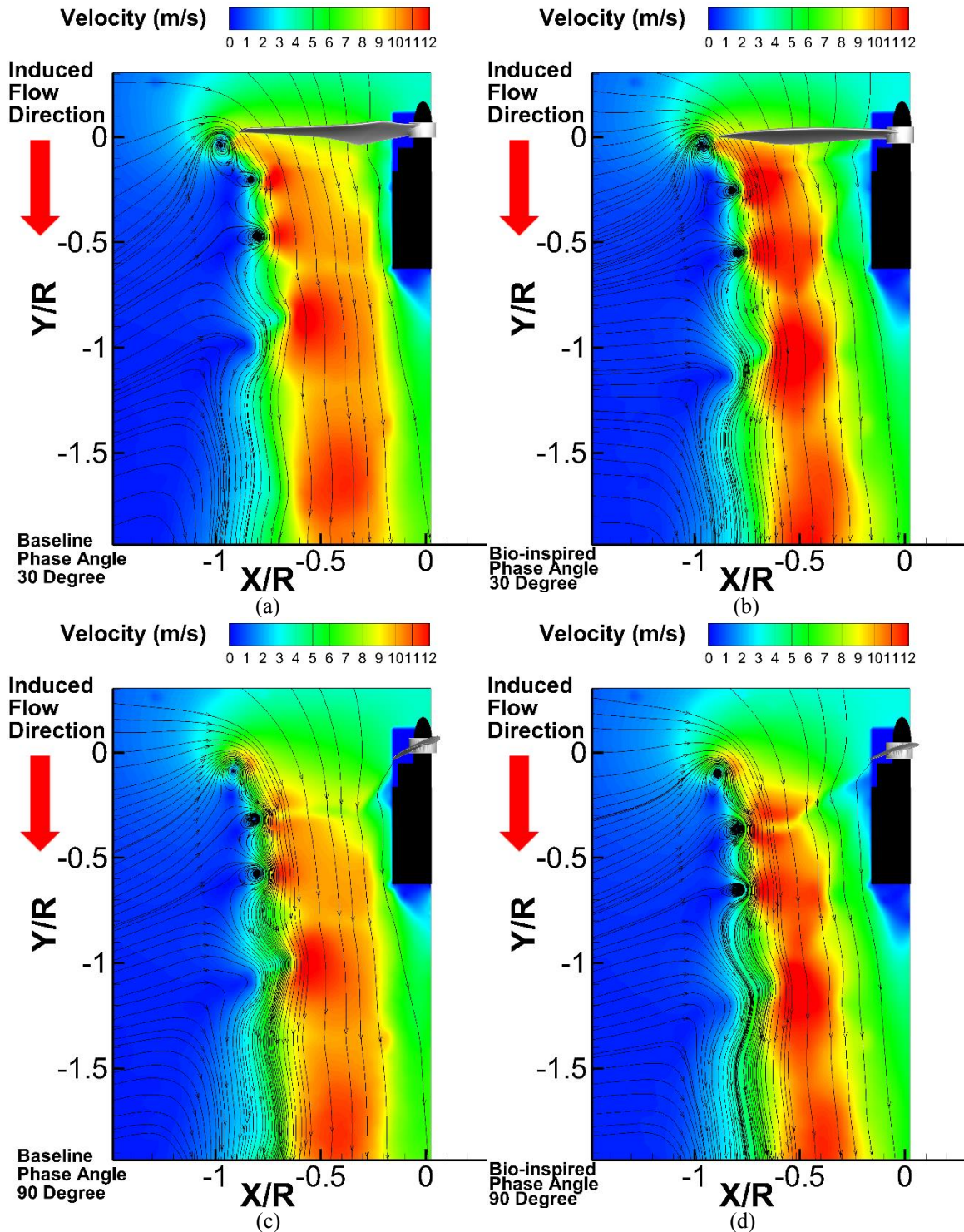


Figure 18. Velocity distribution comparison results at 30 and 90 degree phase angle



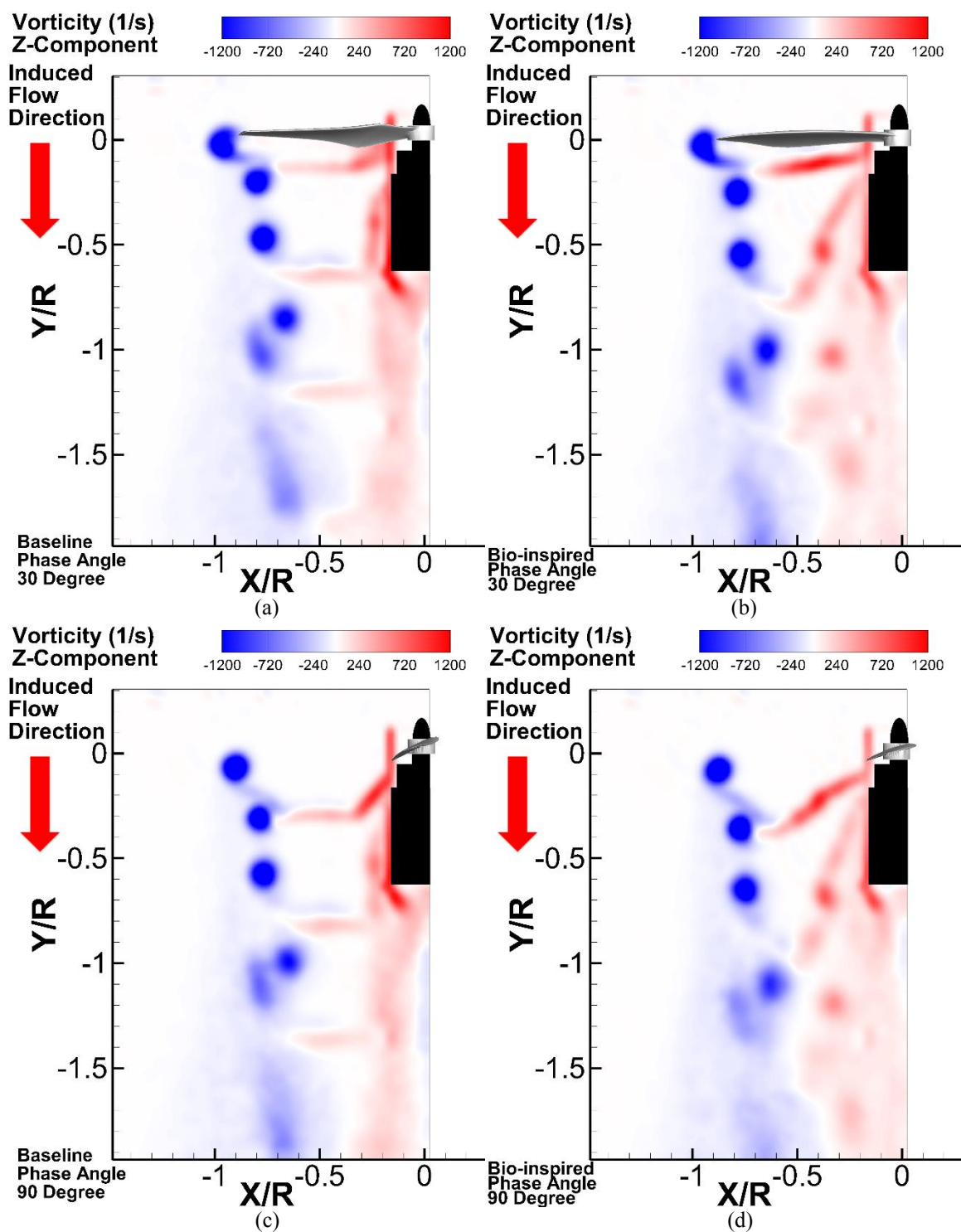


Figure 19. Vorticity distribution comparison results at 30 and 90 degree phase angle

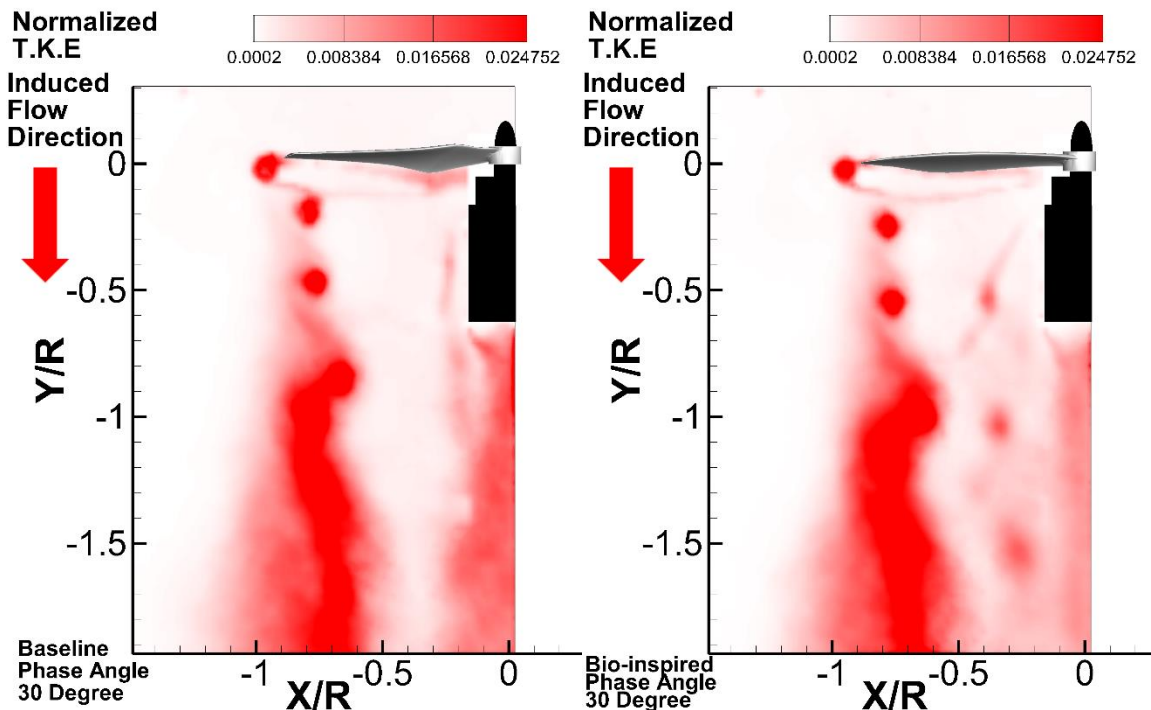


Figure 20. Normalized T.K.E distribution comparison results at 30 and 90 degree phase angle

The comparative results of the velocity distributions of baseline and bio-inspired propellers at 30 and 90 degrees phase angle are shown in figure 18. Compared with the baseline propeller, the induced velocity of bio-inspired propeller was higher at the outer part of the induced flow region, and lower at the inner part. The velocity gradient from the flow boundary towards rotating axis of the bio-inspired propeller was larger than that of the baseline propeller. Figure 19 shows the comparative results of the vorticity distribution of two propellers. One significant difference between these two propellers was the shape of the vortices due to the flow passing trailing edge. As mentioned before, the velocity gradient of the baseline propeller along radial direction was small. As a result, the shape of these vortices were almost parallel to the rotating disk. However, the shape of the vortices of bio-inspired propeller tilted downward in the wake while traveling downstream. Another difference was the vorticity distribution in the induced flow region after  $0.5R$  downstream. The vorticity distribution of bio-inspired propeller was less organized than that of the baseline propeller. This phenomenon can be observed in figure 20, which represents the normalized turbulent kinetic energy (TKE). Here, the TKE is normalized by the average induced velocity. The high TKE level appears at the tip vortex region, and tip vortex interaction region.

### 3.5 Comparison of the characteristics of the wake flow for both propellers

The wake structure traveling speed was acquired by tracking the location of it at different wake ages. In the wake of the propellers, the translation speed of tip vortex was almost half of the tip of the shear layer. This is clearly shown in Figure 21 for both propellers. As mentioned before, this is the reason for the tip vortex interaction in the downstream. By tracking the center of the tip vortex at different wake age, the flow boundary could be acquired. It can be seen in figure 22 (a), the wake size of the bio-inspired propeller is smaller compared with the baseline model's. The tip vortex interaction happened at an earlier stage compared bio-inspired propeller to the baseline propeller. The tip vortex interaction of the bio-inspired propeller appeared at around a 700-degree wake age, and this phenomenon appeared at about an 800-degree wake age for baseline propeller. The figure 22 (b) reveals this comparison of the tip vortex strengths. At early wake age, two propeller had similar strength. After 500 degrees, the vorticity value of the tip vortices was lower for the bio-inspired propeller, as compared with the baseline propeller. This indicates the tip vortices decay rate of the bio-inspired propeller was faster than baseline.

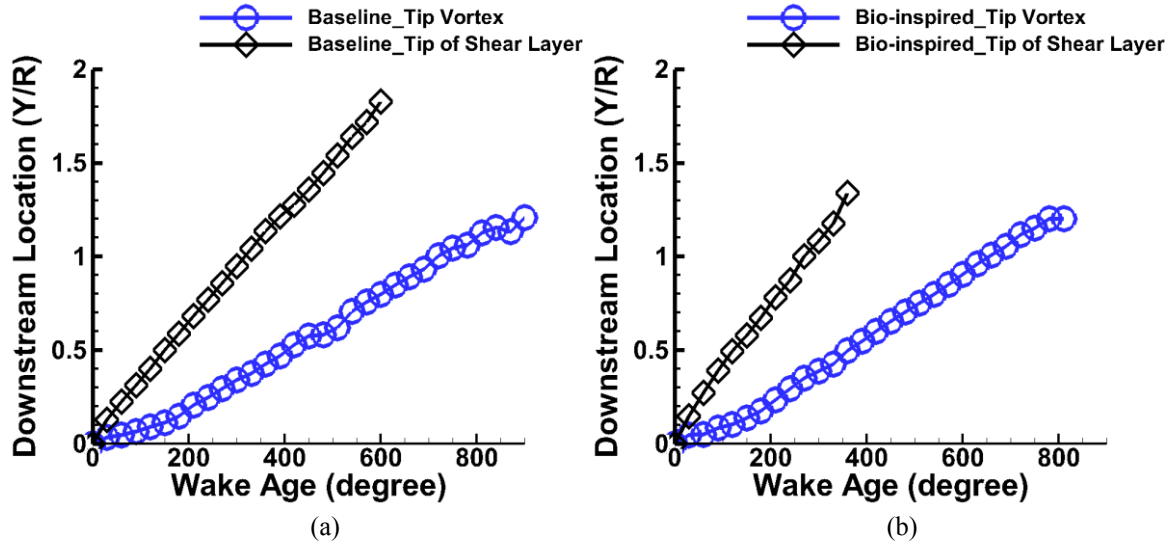


Figure 21. Vortex location at different wake age

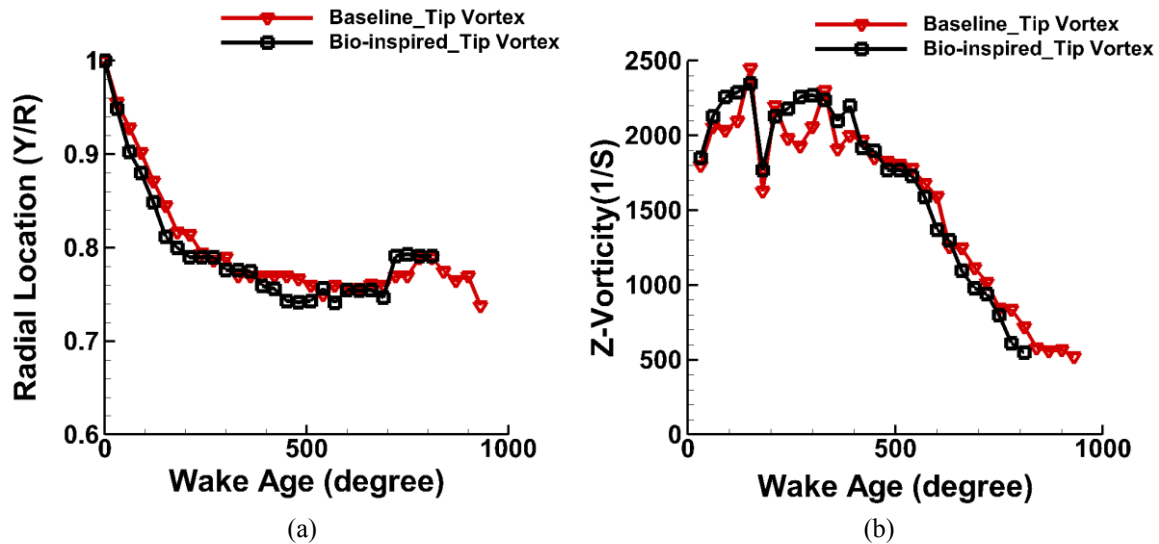


Figure 22. Tip vortex location and strength comparison

#### IV. Conclusion

In conclusion, a comparative experimental study on the aerodynamics and aeroacoustics characteristic between a baseline and a bio-inspired propeller was conducted through force measurements, as well as sound and flow field measurements. In this study, the design thrust and aerodynamic area were kept the same for both propellers. The force measurement results revealed that the bio-inspired propeller could generate same thrust as the baseline propeller under same power input. However, the rotational speed of bio-inspired propeller was slower than baseline at hover flight condition. The sound measurement results indicated that the bio-inspired propeller could reduce noise up to 4dB compared to baseline propeller at hover motion. Through PIV measurements, the velocity and vorticity distribution in the downstream are obtained. Compared with baseline, the bio-inspired propeller generated a smaller wake region and demonstrated a faster decay rate of the tip vortex strength.

#### Acknowledgments

The research work was supported by National Science Foundation (NSF) under award number of IRES-1064235. The authors want to thank Dr. Kai Zhang, Dr. Wenwu Zhou of Iowa State University for their help in conducting the experiments.



## References

- <sup>1</sup> Floreano, D., and Wood, R., "Science, technology and the future of small autonomous drones," *Nature*, vol. 521(7553), 2015, pp. 460–466.
- <sup>2</sup> Hoffmann, G., Huang, H., Waslander, S., and Tomlin, C., "Quadrotor Helicopter Flight Dynamics and Control: Theory and Experiment," *AIAA Guidance, Navigation and Control Conference and Exhibit*, 2007.
- <sup>3</sup> Achtehlik, M., Tianguang Zhang, Kuhnlenz, K., and Buss, M., "Visual tracking and control of a quadcopter using a stereo camera system and inertial sensors," *2009 International Conference on Mechatronics and Automation*, IEEE, 2009, pp. 2863–2869.
- <sup>4</sup> Mellinger, D., Michael, N., and Kumar, V., "Trajectory generation and control for precise aggressive maneuvers with quadrotors," *The International Journal of Robotics Research*, vol. 31, Apr. 2012, pp. 664–674.
- <sup>5</sup> Lindsey, Q., Mellinger, D., and Kumar, V., "Construction with quadrotor teams," *Autonomous Robots*, vol. 33, Oct. 2012, pp. 323–336.
- <sup>6</sup> Mirjan, A., Augugliaro, F., D'Andrea, R., Gramazio, F., and Kohler, M., "Building a Bridge with Flying Robots," *Robotic Fabrication in Architecture, Art and Design 2016*, Cham: Springer International Publishing, 2016, pp. 34–47.
- <sup>7</sup> Bristeau, pierre-jean, Martin, P., Salaun, E., and Petit, N., "the role of propeller aerodynamics in the model of a quadcopter UAV," *European control conference*, vol. Vol.2009, 2009.
- <sup>8</sup> Yoon, S., Lee, H. C. H., and Pulliam, T. H., "Computational Analysis of Multi-Rotor Flows," *54th AIAA Aerospace Sciences Meeting*, 2016.
- <sup>9</sup> Hwang, J. Y., Jung, M. K., and Kwon, O. J., "Numerical Study of Aerodynamic Performance of a Multicopter Unmanned-Aerial-Vehicle Configuration," *Journal of Aircraft*, vol. 52, May 2015, pp. 839–846.
- <sup>10</sup> Ditmer, M., Vincent, J., Werden, L., Tanner, J., Laske, T., Iaizzo, P., Garshelis, D., and Fieberg, J., "Bears Show a Physiological but Limited Behavioral Response to Unmanned Aerial Vehicles," *Current Biology*, vol. 25, Aug. 2015, pp. 2278–2283.
- <sup>11</sup> Leslie, A., Wong, K. C., and Auld, D., "Broadband noise reduction on a mini-UAV Propeller," *29th AIAA Aeroacoustics Conference*, 2008, pp. 5–7.
- <sup>12</sup> Sinibaldi, G., and Marino, L., "Experimental analysis on the noise of propellers for small UAV," *Applied Acoustics*, vol. 74, Jan. 2013, pp. 79–88.
- <sup>13</sup> Eppler, R., *Airfoil design and data*, Springer Berlin Heidelberg, 1990.
- <sup>14</sup> Leishman, J. G., *Principles of helicopter aerodynamics*, Cambridge University Press, 2006.