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# Positive Geometry for Stringy Scalar Amplitudes

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Charles University

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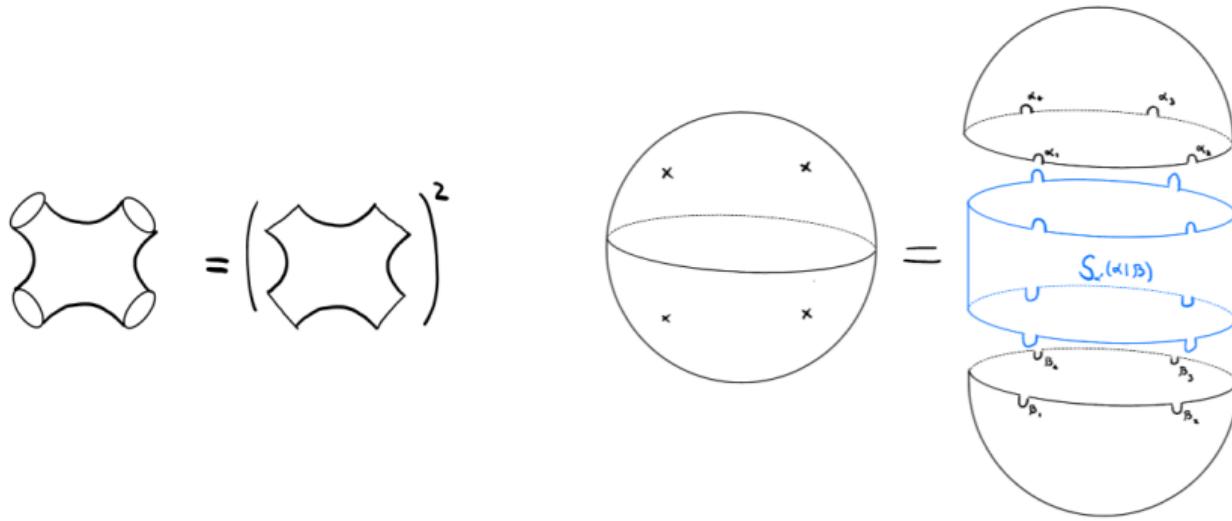
with C. Bartsch, K. Kampf, D. Podivín

February 4, 2026, University of Hertfordshire



# Introduction: KLT Double Copy

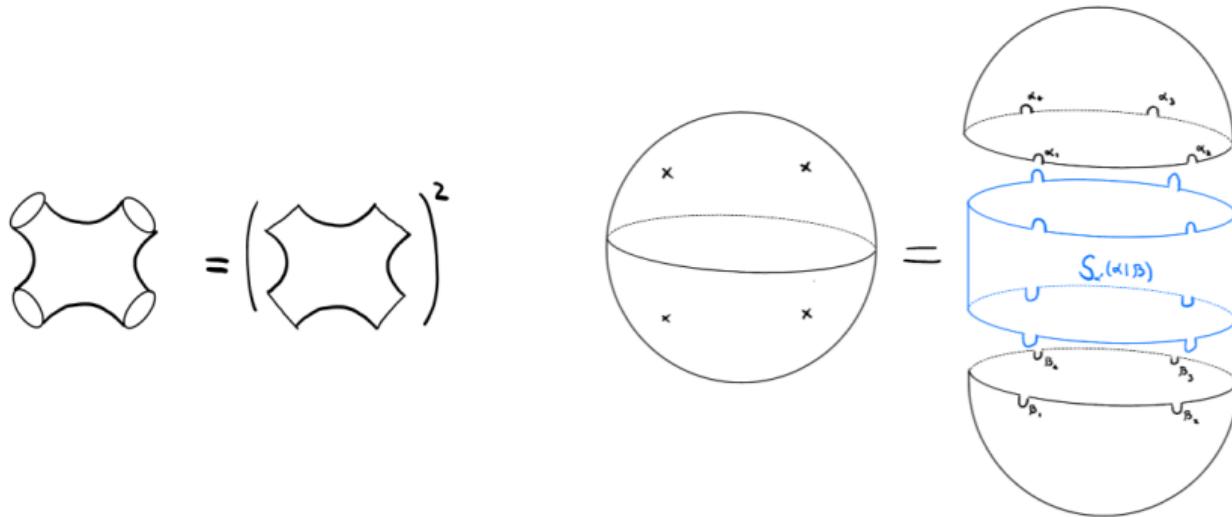
Kawai-Lewellen-Tay (1986): closed string = (open string)<sup>2</sup>



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# Introduction: Bi-adjoint $\phi^3$

Taking  $\alpha' \rightarrow 0$ : GR = (Yang-Mills)<sup>2</sup>

$$A^{\text{GR}} = \sum_{\alpha, \beta} A^{\text{YM}}(\alpha) S(\alpha|\beta) A^{\text{YM}}(\beta)$$

What is the kernel  $S(\alpha|\beta)$ ? Does it have a physical interpretation?

Answer:

$$S(\alpha|\beta)^{-1} = m_n(\alpha|\beta)$$

Scattering amplitudes in *bi-adjoint  $\phi^3$*

Diagonal elements:  $m_n \equiv m_n(\alpha|\alpha)$  are  $\text{Tr}(\phi^3)$  amplitudes

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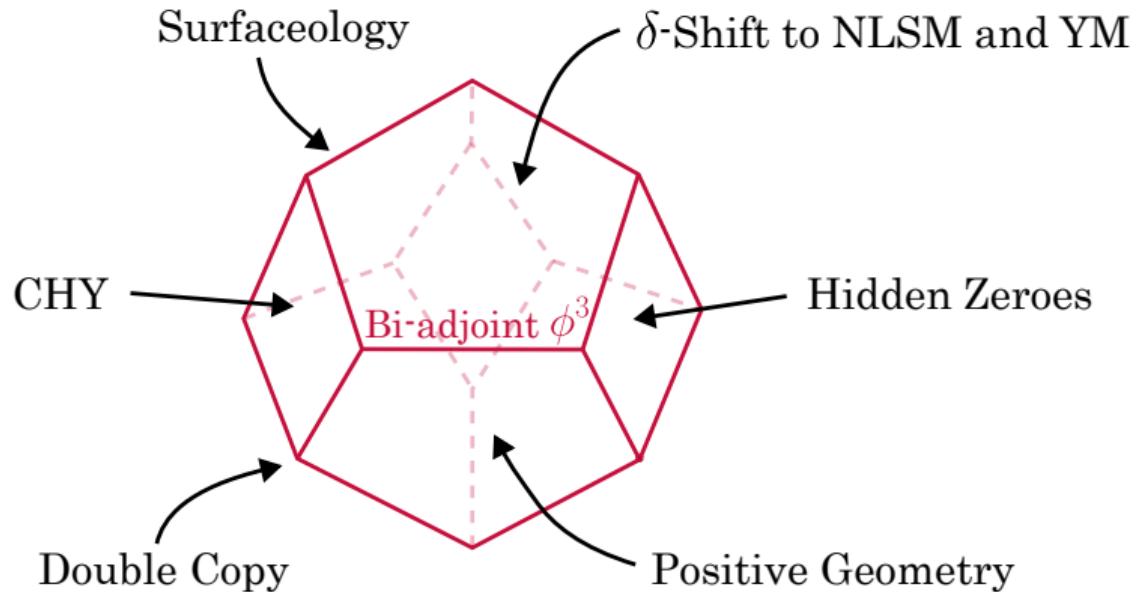
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# Introduction: Bi-adjoint $\phi^3$

Bi-adjoint  $\phi^3$ : A theory with many facets



# Introduction: Inverse KLT Kernel

What about the inverse *string theory* KLT kernel?

$$S_{\alpha'}(\alpha|\beta)^{-1} = m_n^{\alpha'}(\alpha|\beta)$$

A lot less is known, studied by Mizera (2017)

Some ‘stringy’ version of bi-adjoint  $\phi^3$

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# Outline

- Review:  $\text{Tr}(\phi^3)$  Amplitudes
- Review: The ABHY Associahedron
- Geometry of Stringy  $\text{Tr}(\phi^3)$
- Pions and Mixed Amplitudes
- Off-Diagonal Amplitudes

# $\text{Tr}(\phi^3)$ Amplitudes

Cubic scalar theory, only planar Feynman diagrams

Propagators in tree-level  $\text{Tr}(\phi^3) \longleftrightarrow$  Planar Mandelstam variables

$$X_{ij} := (p_i + p_{i+1} + \dots + p_{j-1})^2$$

Masslessness:  $X_{ii+1} = X_{1n} = 0$

Momentum conservation:  $X_{ij} = X_{ji}$

At a pole  $X_{ij} = 0$  the amplitude factorizes: **Locality & Unitarity**



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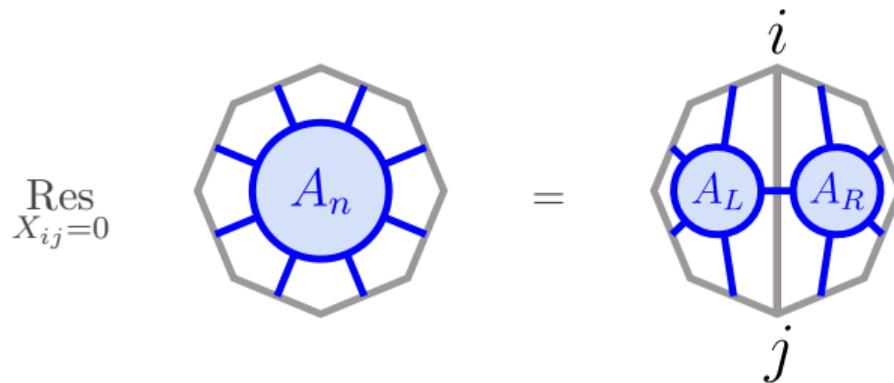
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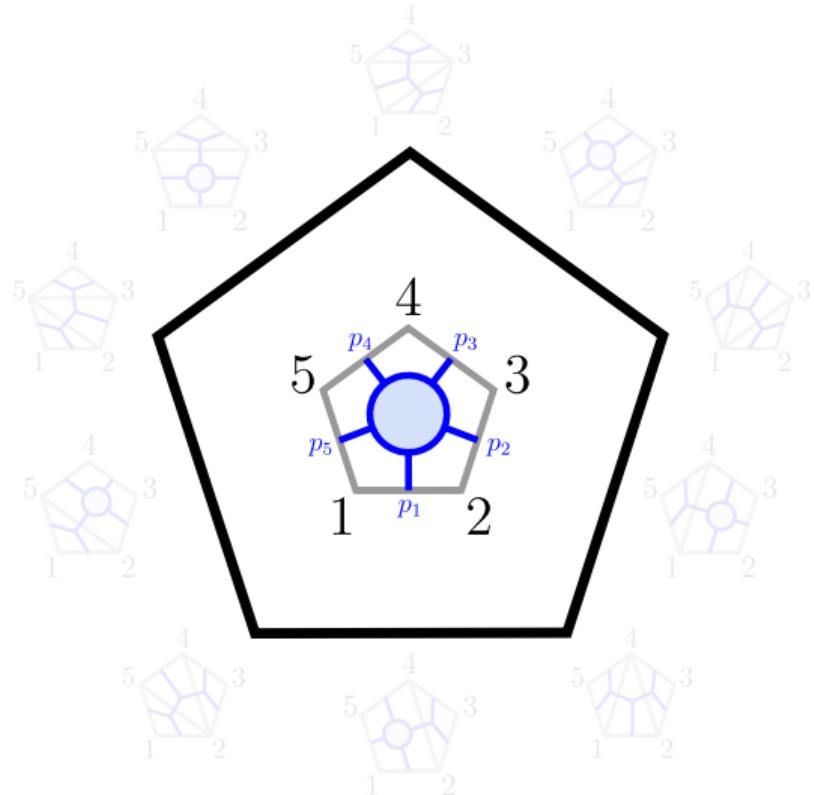
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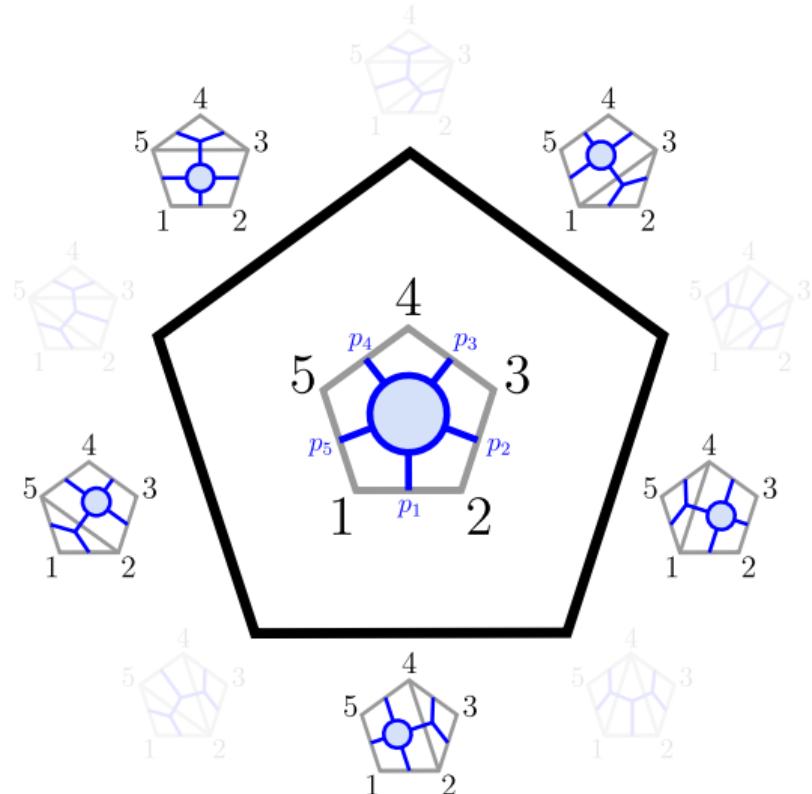
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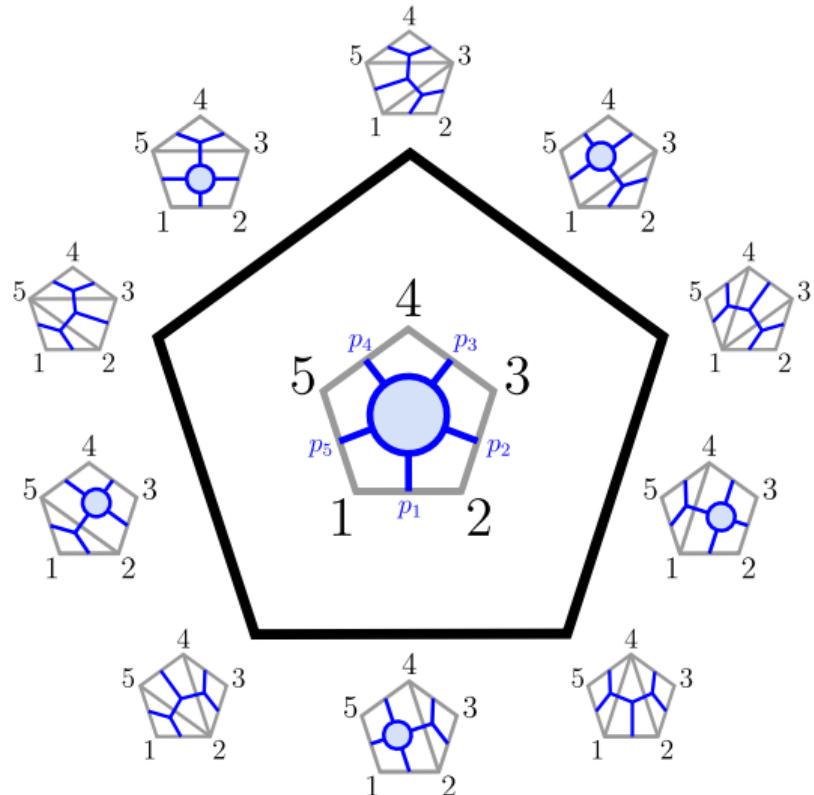
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# $\text{Tr}(\phi^3)$ Amplitudes: Examples

$$m_3 = 1$$

$$m_4 = \begin{array}{c} 4 \\ | \\ \square \text{ (blue lines)} \\ | \\ 1 \end{array} + \begin{array}{c} 4 \\ | \\ \square \text{ (blue lines)} \\ | \\ 1 \end{array}$$

$$= \frac{1}{X_{13}} + \frac{1}{X_{24}}$$

$$m_5 = \begin{array}{c} 5 \\ | \\ \square \text{ (blue lines)} \\ | \\ 1 \end{array} + \begin{array}{c} 5 \\ | \\ \square \text{ (blue lines)} \\ | \\ 1 \end{array} + \begin{array}{c} 5 \\ | \\ \square \text{ (blue lines)} \\ | \\ 1 \end{array} + \begin{array}{c} 5 \\ | \\ \square \text{ (blue lines)} \\ | \\ 1 \end{array} + \begin{array}{c} 5 \\ | \\ \square \text{ (blue lines)} \\ | \\ 1 \end{array}$$

$$= \frac{1}{X_{13}X_{14}} + \frac{1}{X_{14}X_{24}} + \frac{1}{X_{24}X_{25}} + \frac{1}{X_{25}X_{35}} + \frac{1}{X_{35}X_{13}}$$

# The ABHY Associahedron

Positive Geometry:

- Geometry  $\mathcal{A}$  with boundaries  $\longleftrightarrow$  poles of amplitude
- With a unique canonical form  $\Omega(\mathcal{A})$  (“=” scattering amplitude)
- $\Omega(\mathcal{A})$  has log singularities *at and only at* boundaries of  $\mathcal{A}$

For  $\text{Tr}(\phi^3)$ :

- Boundary structure: triangulations of an  $n$ -gon  $\implies$  Associahedron
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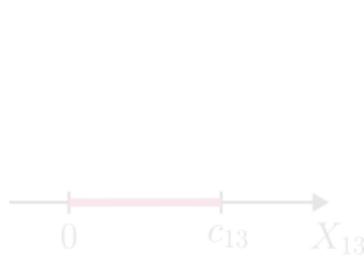
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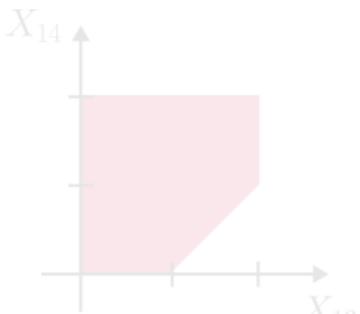
- All  $X_{ij} \geq 0 \implies$  correct facets
- Not full  $X_{ij}$  space: keep  $(n - 3)$   $X$ 's and  $(n - 2)(n - 3)/2$  constraints

$$c_{ij} = X_{ij} + X_{i+1j+1} - X_{ij+1} - X_{i+1j} = -(p_i + p_j)^2 = \text{positive const.}$$

For  $n = 4$ :  $X_{13}, X_{24} = c_{13} - X_{13}$



$\mathcal{A}_4$



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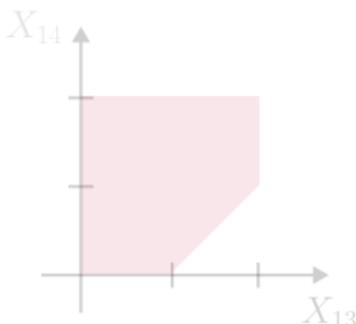
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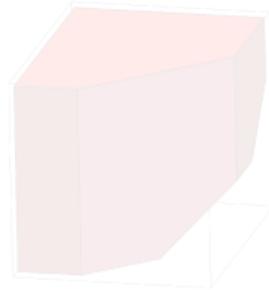
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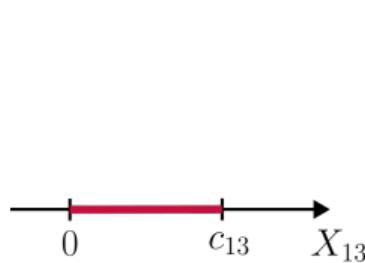
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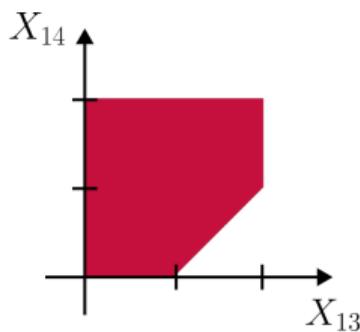
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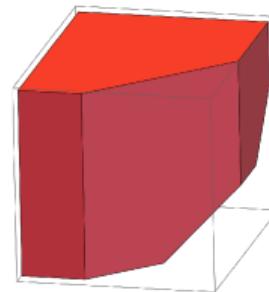
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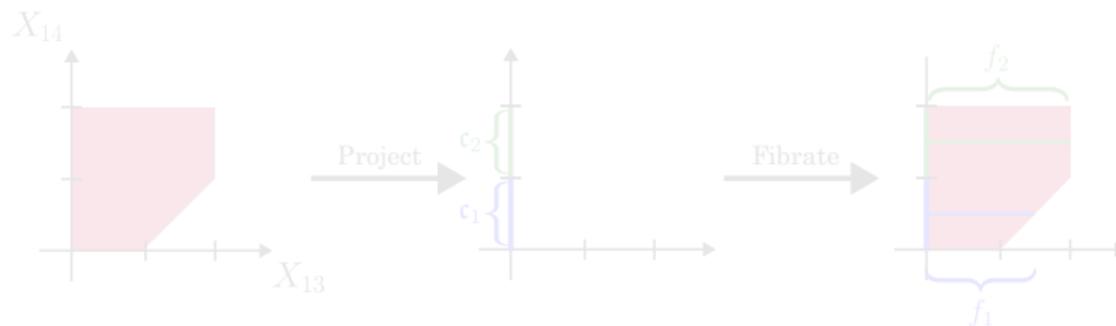
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# The ABHY Associahedron: Canonical Form

Canonical form:

$$\Omega \left( \begin{array}{c} \xrightarrow{\mathcal{A}_4} \\ 0 \quad c_{13} \quad X_{13} \end{array} \right) = \frac{dX_{13}}{X_{13}} + \frac{dX_{13}}{c_{13} - X_{13}} = \left( \frac{1}{X_{13}} + \frac{1}{X_{24}} \right) dX_{13}$$

For  $n \geq 5$ , we can use *chambers* and *fibers*



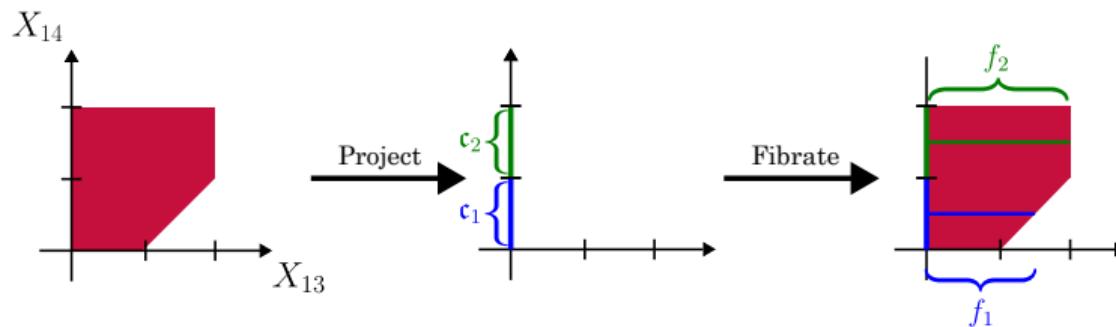
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In general:

- Chambers = product of lower-point associahedra
- Fibers = line-segment ( $= \mathcal{A}_4$ )

*Geometric Recursion*

$$m_n = \sum_{i=4}^n m(1, 2, 3, i)m(2, 3, \dots, i)m(i, i+1, \dots, n, 1, 2) \Big|_{X_{2j} \rightarrow X_{2j} - X_{2i}}$$

Example:

$$\begin{aligned} m_5 &= \left( \frac{1}{X_{13}} + \frac{1}{X_{24}} \right) \left( \frac{1}{X_{14}} + \frac{1}{X_{25} - X_{24}} \right) \\ &\quad + \left( \frac{1}{X_{13}} + \frac{1}{X_{25}} \right) \left( \frac{1}{X_{35}} + \frac{1}{X_{24} - X_{25}} \right) \end{aligned}$$

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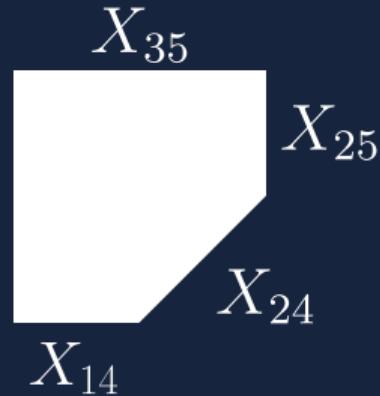
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# Summary

- $\text{Tr}(\phi^3)$  Amplitudes  $m_n$ 
  - Diagonal elements of  $S(\alpha|\beta)^{-1} = m(\alpha|\beta)$
  - Poles as  $X_{ij} = 0$
  - Factorizes on poles
- The ABHY Associahedron  $\mathcal{A}_n$ 
  - Polytope in kinematic space
  - Boundaries at  $X_{ij} = 0$
  - $\Omega(\mathcal{A}_n) = m_n$
  - Once we know  $\Omega(\mathcal{A}_4)$ , we can find  $m_n$  through geometric recursion



## Stringy $\text{Tr}(\phi^3)$

Recall:

$$S_{\alpha'}^{-1} = m_n^{\alpha'}(\alpha|\beta)$$

“Stringy”  $\text{Tr}(\phi^3)$ :  $m_n^{\alpha'} \equiv m_n^{\alpha'}(\alpha|\alpha)$

$$m_3^{\alpha'} = 1$$

$$m_4^{\alpha'} = \frac{1}{\tan(\pi\alpha'X_{13})} + \frac{1}{\tan(\pi\alpha'X_{24})}$$

$$\begin{aligned} m_5^{\alpha'} &= 1 + \frac{1}{\tan(\pi\alpha'X_{13})\tan(\pi\alpha'X_{14})} + \frac{1}{\tan(\pi\alpha'X_{14})\tan(\pi\alpha'X_{24})} \\ &+ \frac{1}{\tan(\pi\alpha'X_{24})\tan(\pi\alpha'X_{25})} + \frac{1}{\tan(\pi\alpha'X_{25})\tan(\pi\alpha'X_{35})} + \frac{1}{\tan(\pi\alpha'X_{35})\tan(\pi\alpha'X_{13})} \end{aligned}$$

# Geometry of Stringy $\text{Tr}(\phi^3)$

$m_n^{\alpha'}$  is ‘stringy’:

- Infinite resonance structure: Poles at  $X_{ij} = k/\alpha'$ ,  $k \in \mathbb{Z}$
- Factorizes on the poles
- Satisfies monodromy relations
- Reduces to  $m_n$  as  $\alpha' \rightarrow 0$

For stringy  $\text{Tr}(\phi^3)$  amplitudes the geometric recursion *still holds*:

$$m_n^{\alpha'} = \sum_{i=4}^n m^{\alpha'}(1, 2, 3, i) m^{\alpha'}(2, 3, \dots, i) m^{\alpha'}(i, i+1, \dots, n, 1, 2) \Big|_{X_{2j} \rightarrow X_{2j} - X_{2i}}$$

We can recurse down to products of 4-point  $m_4^{\alpha'}$

Thus: A geometry for  $m_4^{\alpha'}$   $\implies$  a geometry for  $m_n^{\alpha'}$ !

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Want: a one-dimensional geometry with canonical form

$$\omega_4^{\alpha'} = d \log \frac{\sin(\pi \alpha' X_{13})}{\sin(\pi \alpha'(c_{13} - X_{13}))} = \left( \frac{1}{\tan(\pi \alpha' X_{13})} + \frac{1}{\tan(\pi \alpha' X_{24})} \right) dX_{13}$$

Using Euler's infinite product formula for  $\sin(\pi x)$ :

$$\omega_4^{\alpha'} = \sum_{k \in \mathbb{Z}} d \log \frac{X_{13} + k/\alpha'}{X_{13} - c_{13} + k/\alpha'}$$

Which is the canonical form of an infinite sum of line segments!



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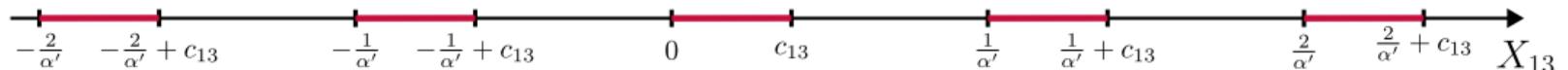
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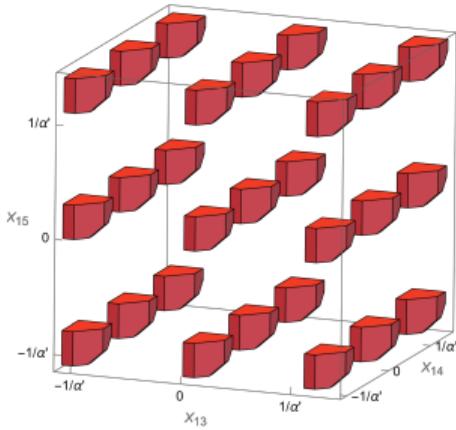
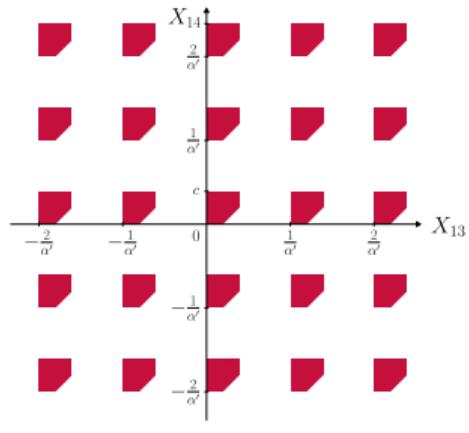
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# Geometry of Stringy $\text{Tr}(\phi^3)$

For higher  $n$ : Infinite grid of ABHY associahedra!



- Correct poles as  $X_{ij} = k/\alpha'$
- Factorises on poles
- $\Omega(\mathcal{A}_n^{\alpha'}) = m_n^{\alpha'}$

# Pions and Mixed Amplitudes

$m_n^{\alpha'}$  contains  $\text{Tr}(\phi^3)$  as  $\alpha' \rightarrow 0$

Surprisingly, it also contains pions in the NLSM! [Bartsch, Kampf, Novotný, Trnka]

- Rescale  $\alpha' \rightarrow \alpha'/2$
- Shift certain  $X_{ij} \rightarrow X_{ij} \pm 1/\alpha'$
- NLSM amplitudes as  $\alpha' \rightarrow 0$

$$\frac{1}{\tan(\pi\alpha'X_{ij})} \rightarrow \frac{1}{\tan(\pi\alpha'X_{ij}/2 + \pi/2)} = -\tan(\pi\alpha'X_{ij}/2)$$

Example:

$$m_4^{\alpha'} = \frac{1}{\tan(\pi\alpha'X_{13})} + \frac{1}{\tan(\pi\alpha'X_{24})} \rightarrow -\tan(\pi\alpha'X_{13}/2) - \tan(\pi\alpha'X_{24}/2)$$

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Surprisingly, it also contains pions in the NLSM! [Bartsch, Kampf, Novotný, Trnka]

- Rescale  $\alpha' \rightarrow \alpha'/2$
- Shift certain  $X_{ij} \rightarrow X_{ij} \pm 1/\alpha'$
- NLSM amplitudes as  $\alpha' \rightarrow 0$

$$\frac{1}{\tan(\pi\alpha'X_{ij})} \rightarrow \frac{1}{\tan(\pi\alpha'X_{ij}/2 + \pi/2)} = -\tan(\pi\alpha'X_{ij}/2)$$

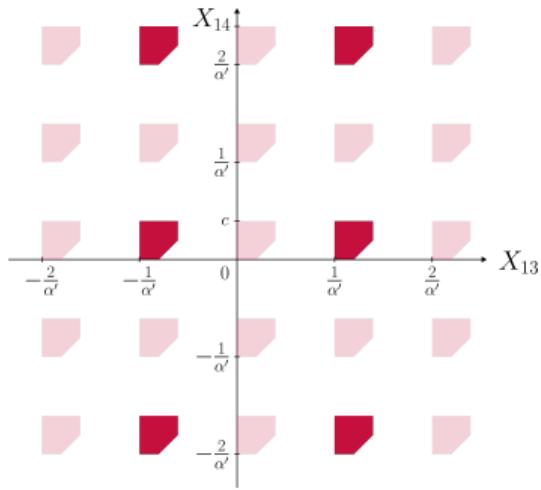
Example:

$$m_4^{\alpha'} = \frac{1}{\tan(\pi\alpha'X_{13})} + \frac{1}{\tan(\pi\alpha'X_{24})} \rightarrow -\tan(\pi\alpha'X_{13}/2) - \tan(\pi\alpha'X_{24}/2)$$

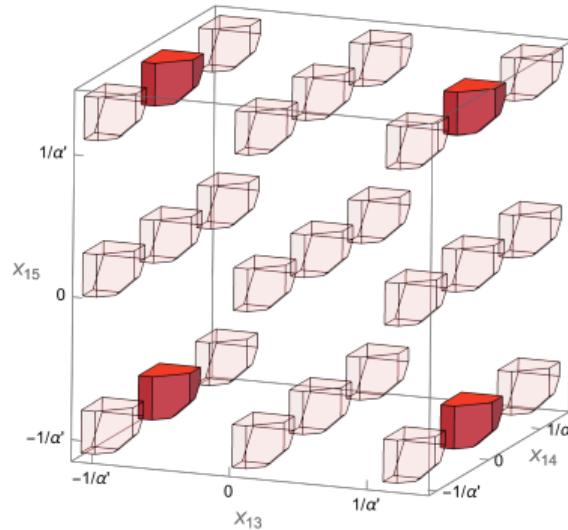
As  $\alpha' \rightarrow 0$  we recover  $A_4^{\text{NLSM}} = -X_{13} - X_{24}$

# Pions and Mixed Amplitudes

Geometrically: this isolates an **infinite subgrid** of associahedra



$$m_5^{\alpha'}(\pi\pi\phi\phi\phi)$$



$$m_6^{\alpha'}(\pi\pi\pi\pi\pi\pi)$$

# Pions and Mixed Amplitudes

Field theory amplitude in  $\alpha' \rightarrow 0$ :

$$\begin{aligned} \lim_{\alpha' \rightarrow 0} & \text{---} \left[ \begin{array}{ccccccccc} -\frac{2}{\alpha'} & -\frac{2}{\alpha'} + c_{13} & -\frac{1}{\alpha'} & -\frac{1}{\alpha'} + c_{13} & 0 & c_{13} & \frac{1}{\alpha'} & \frac{1}{\alpha'} + c_{13} & \frac{2}{\alpha'} & \frac{2}{\alpha'} + c_{13} \end{array} \right] \text{---} \\ &= \lim_{\alpha' \rightarrow 0} \text{---} \left[ \begin{array}{cc} \frac{1}{\alpha'} & \frac{1}{\alpha'} + c_{13} \end{array} \right] \text{---} \\ &= \lim_{\alpha' \rightarrow 0} \left( \frac{1}{X_{13} + 1/\alpha'} + \frac{1}{X_{24} - 1/\alpha'} \right) = -X_{13} - X_{24} \end{aligned}$$

This is the same as the  $\delta$ -shift if we equate  $\delta = 1/\alpha'$ !

Other subgrids give amplitudes inaccessible by  $\delta$ -shift

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# Off-Diagonal Amplitudes

So far: stringy  $\text{Tr}(\phi^3) \leftarrow$  Diagonal  $m_{\alpha'}(\mathbb{I}|\mathbb{I})$

How about stringy  $m_{\alpha'}(\alpha|\beta)$ ?

They can be written as some  $m_n^{\alpha'}$ 's times products of  $1/\sin(\pi\alpha' X_{ij})$ . Example:

$$m_{\alpha'}(\mathbb{I}|13245) = \frac{1}{\sin(\pi\alpha' X_{13})} \left( \frac{1}{\tan(\pi\alpha' X_{14})} + \frac{1}{\tan(\pi\alpha' X_{35})} \right)$$

Geometry for  $1/\sin(\pi\alpha' X_{ij}) \implies$  Geometry for all off-diagonal amplitudes

Canonical form:

$$d \log \tan(\pi\alpha' X_{ij}/2) = \frac{dX_{ij}}{\sin(\pi\alpha' X_{ij})}$$

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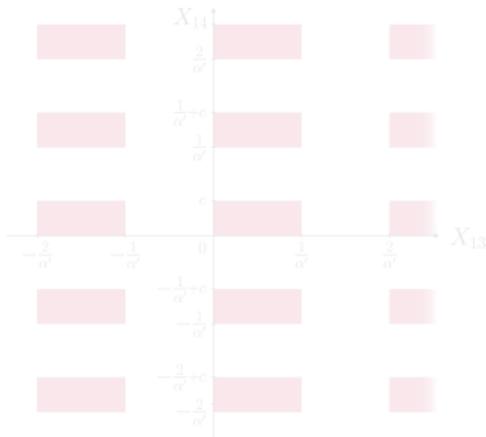
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Playing the same game as before: we again find infinitely many line segments!



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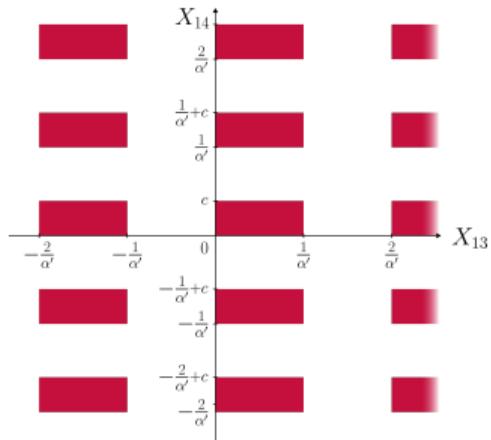
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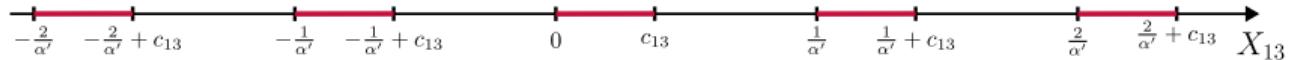
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# Summary

- “Stringy bi-adjoint  $\phi^3$  amplitudes”  $m_n^{\alpha'}(\alpha|\beta)$ 
  - Central in KLT double copy
  - Natural stringification of the well-studied  $\text{Tr}(\phi^3)$
- Geometric description in terms of the *Associahedral grid*
- Geometry contains all NLSM and mixed  $\pi/\phi$  amplitudes
- Positive geometry containing stringy features, beyond rational functions



# *Q&A*

*Thank you for listening!*

# Bonus Slides

$$m_5^{\alpha'} = \left( \frac{1}{\tan(\pi\alpha'X_{13})} + \frac{1}{\tan(\pi\alpha'X_{24})} \right) \left( \frac{1}{\tan(\pi\alpha'X_{14})} + \frac{1}{\tan(\pi\alpha'(X_{25} - X_{24}))} \right) \\ + \left( \frac{1}{\tan(\pi\alpha'X_{13})} + \frac{1}{\tan(\pi\alpha'X_{25})} \right) \left( \frac{1}{\tan(\pi\alpha'X_{35})} + \frac{1}{\tan(\pi\alpha'(X_{24} - X_{25}))} \right)$$

## Bonus Slides

The 'spurious boundaries' give all the correct contact terms:

$$\frac{1}{\tan(\pi\alpha'(X_a - X_b))} = \frac{1 + \tan(\pi\alpha'X_a)\tan(\pi\alpha'X_b)}{\tan(\pi\alpha'X_a) - \tan(\pi\alpha'X_b)}$$

Which leads to the correct

$$m_5^{\alpha'} = 1 + \frac{1}{\tan(\pi\alpha'X_{13})\tan(\pi\alpha'X_{14})} + \frac{1}{\tan(\pi\alpha'X_{14})\tan(\pi\alpha'X_{24})} \\ + \frac{1}{\tan(\pi\alpha'X_{24})\tan(\pi\alpha'X_{25})} + \frac{1}{\tan(\pi\alpha'X_{25})\tan(\pi\alpha'X_{35})} + \frac{1}{\tan(\pi\alpha'X_{13})\tan(\pi\alpha'X_{35})}$$