

GR - Cosmology TDs

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Chapter 1

TD 1

1.1 Transverse Doppler effect and aberration.

1. ...
2. ...
3. ...

1.2 Relativistic invariants and conserved quantities.

1. (a) The four position is given by $\vec{x} = \begin{pmatrix} ct \\ \mathbf{x} \end{pmatrix}$ and we have that the proper time differential is given by:
 $dt = \gamma d\tau$. Hence the four velocity is given by:

$$\vec{u} = \frac{d\vec{x}}{dt} = \gamma \begin{pmatrix} c \\ \mathbf{u} \end{pmatrix}$$

- (b) We have the simple relation $p^\mu = mu^\mu$. Which gives:

$$\vec{p} = \begin{pmatrix} \gamma mc \\ \gamma m\mathbf{u} \end{pmatrix} = \begin{pmatrix} \frac{E}{c} \\ \gamma m\mathbf{u} \end{pmatrix}$$

- (c) The invariant built from the momentum is given by:

$$p_\mu p^\mu = \eta_{\mu,\nu} p^\mu p^\nu = -m^2 c^2$$

Which gives:

$$m^2 c^4 = E^2 - |\vec{p}|^2 c^2$$

2. (a) The conserved physical quantities are the energy of the whole system and the momentum of the center of mass. Hence in other words the four momentum is conserved.
- (b) The photon will lose energy and hence it will have a lower frequency $\nu' < \nu$. The four-momenta are given by:

$$p_{e0} = (m, 0, 0, 0), p_{p0} = (E_0, 0, 0, E_0)$$

And after the collision by:

$$p_{e1} = (E'', 0, p_y, p_z), p_{p1} = (E', 0, -E' \sin \theta, -E' \cos \theta)$$

- (c)

Chapter 2

Chapter 3

3.1 Covariant derivatives and Christoffel symbols

1. (a) We have that:

$$J_{\mu'}^{\mu} J_{\nu}^{\mu'} = \frac{\partial x^{\mu}}{\partial x^{\mu'}} \frac{\partial x^{\mu'}}{\partial x^{\nu}} = \frac{\partial x^{\mu}}{\partial x^{\nu}} = \delta_{\nu}^{\mu}$$

- (b) We have that:

$$\left(J_{\sigma'}^{\mu} J_{\mu}^{\mu'} \right)_{,\gamma} = (\delta_{\mu}^{\mu})_{,\gamma} = 0 \Rightarrow J_{\sigma',\gamma}^{\mu} J_{\mu}^{\mu'} = -J_{\sigma'}^{\mu} J_{\mu,\gamma}^{\mu'}$$

2. (a) We have that:

$$\partial_{\mu'} \phi'(x') = \frac{\partial}{\partial x^{\mu'}} \phi(x) = \frac{\partial x^{\mu}}{\partial x^{\mu'}} \frac{\partial}{\partial x^{\mu}} \phi(x) = J_{\mu'}^{\mu} \partial_{\mu} \phi(x)$$

Hence it is a tensor.

- (b) No.

- (c) We have:

$$\begin{aligned} \nabla_{\mu'} A^{\nu'} &= \partial_{\mu'} A^{\nu'} = J_{\mu'}^{\mu} J_{\nu'}^{\nu} (\partial_{\mu} A^{\nu} + \Gamma_{\mu\rho}^{\nu} A^{\rho}) \Leftrightarrow J_{\mu'}^{\mu} J_{\nu'}^{\nu} \partial_{\mu} A^{\nu} + J_{\mu'}^{\mu} J_{\nu',\mu'}^{\nu} A^{\nu} = J_{\mu'}^{\mu} J_{\nu'}^{\nu} \partial_{\mu} A^{\nu} + J_{\mu'}^{\mu} J_{\nu'}^{\nu} \Gamma_{\mu\rho}^{\nu} A^{\rho} \\ &\Leftrightarrow (J_{\mu'}^{\mu} J_{\nu',\mu'}^{\nu} - J_{\mu'}^{\mu} J_{\sigma}^{\nu'} \Gamma_{\mu\nu}^{\sigma}) A^{\nu} = 0 \end{aligned}$$

Which gives the desired result.

- (d) The exchange of the lower indices changes nothing.

- (e) We consider the invariant quantity $B^{\mu} A_{\mu}$ then:

$$\nabla_{\rho} (A_{\mu} B^{\mu}) = \partial_{\rho} (A_{\mu} B^{\mu}) = (\partial_{\rho} A_{\mu}) B^{\mu} + A_{\mu} \partial_{\rho} B^{\mu}$$

However we also have that:

$$\nabla_{\rho} (A_{\mu} B^{\mu}) = (\nabla_{\rho} A_{\mu}) B^{\mu} + A_{\mu} \nabla_{\rho} B^{\mu} = (\nabla_{\rho} A_{\mu}) B^{\mu} + A_{\mu} \partial_{\rho} B^{\mu} + A_{\mu} \Gamma_{\rho\nu}^{\mu} B^{\nu}$$

Now equating both sides we obtain:

$$(\partial_{\rho} A_{\mu}) B^{\mu} = \Gamma_{\rho M}^{\sigma} B^{\mu} A_{\sigma} + \nabla_{\rho} A_{\mu} B^{\mu}$$

Which gives:

$$\nabla_{\rho} A_{\mu} = \partial_{\rho} A_{\mu} - \Gamma_{\rho\mu}^{\sigma} A_{\sigma}$$

3. (a) We have that:

$$\Gamma_{\nu'\mu'}^{\rho'} = \frac{\partial x^{\rho'}}{\partial X^{\alpha}} \frac{\partial^2 X^{\alpha}}{\partial x^{\mu'} \partial x^{\nu'}} =$$

- (b)

3.2 Geodesics and the shortest path

1. (a) We have that:

$$S = -m \int |ds| = -m \int \sqrt{ds_\mu ds^\mu} = -m \int \sqrt{g_{\mu\nu} ds^\mu ds^\nu}$$

Now notice that:

$$ds = \frac{dx}{d\lambda} d\lambda$$

Hence re-writing on top we obtain:

$$S = -m \int \sqrt{-g_{\mu\nu} \dot{x}^\mu \dot{x}^\nu} d\lambda$$

- (b) The Euler-Lagrange equations are given by:

$$\frac{d}{d\lambda} \frac{\partial \mathcal{L}}{\partial \dot{x}} - \frac{\partial \mathcal{L}}{\partial x} = \frac{d}{d\lambda} \frac{\partial \mathcal{L}}{\partial \dot{x}} = \frac{d}{d\lambda} \frac{-g_{\mu\nu} \dot{x}^\mu}{2\mathcal{L}} =$$