Experiments

Robin

February 16, 2021

Motivation: We wish to study the transport of M interacting particles within a square lattice of size N.

$$H(m) = -\sum_{(i,j)} J_{i,j} \ m_i \ m_j \qquad i, j \in \{1, ..N\} \qquad m_i \in \{0,1\} \qquad \sum_i m_i = M$$

From now on, let
$$J_{i,j} = \begin{cases} J & j \in \text{nbd}(i) \\ 0 & j \notin \text{nbd}(i) \end{cases}$$
, $\beta = \frac{J}{T}$ and $|\text{nbd}(i)| = q \ \forall i$.

Mean field theory: Suppose $m_i = \langle m_i \rangle + \delta m_i = m + \delta m_i$, where $\delta m_i = m_i - m$ is a small fluctuation, then we can simplify the Hamiltonian considerably.

$$H(m) = -J \sum_{(i,j)} m_i m_j$$

$$= -J \sum_{(i,j)} (m + \delta m_i)(m + \delta m_j)$$

$$\approx -J \sum_{(i,j)} (m^2 + m \ \delta m_i + m \ \delta m_j)$$

$$= -\frac{J}{2} \sum_{j \in \text{nbd}(i)} \sum_{i} (m^2 + 2m \ \delta m_i)$$

$$= -\frac{Jq}{2} \sum_{i} (m^2 + 2m \ (m_i - m))$$

$$= -\frac{Jq}{2} \left(2m \sum_{i} m_i - \sum_{i} m^2 \right)$$

$$= -\frac{Jq}{2} \left(2Mm - Nm^2 \right)$$

There is no interaction, which means that the only possible phase is a disordered phase. Similarly, adding a term $\propto \sum_i m_i$ to the Hamiltonian only changes it by a constant. Therefore, the Hamiltonian is invariant under application of a

constant external field.

For simulation: Choose an edge (i,j) randomly. Choose whether to exchange m_i and m_j using the Metropolis algorithm. The energy per temperature of an edge of the lattice is given by $-\frac{E_{(i,j)}}{T}=\beta(m_i\ m_j+\sum_{i'\neq j}m_i\ m_{i'}+\sum_{j'\neq i}m_j\ m_{j'}=m_i\sum_{i'}m_{i'}+m_j\sum_{j'}m_{j'}-m_i\ m_j).$

Let's compute the change in energy when m_i and m_j are exchanged:

$$-\frac{\Delta E_{(i,j)}}{T} = \beta((m_i + \Delta m_i) \left(\sum_{i'} m_{i'} + \Delta m_j\right) + (m_j + \Delta m_j) \left(\sum_{j'} m_{j'} + \Delta m_i\right)$$

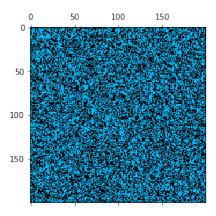
$$- (m_i + \Delta m_i) (m_j + \Delta m_j) - m_i \sum_{i'} m_{i'} - m_j \sum_{j'} m_{j'} + m_i m_j)$$

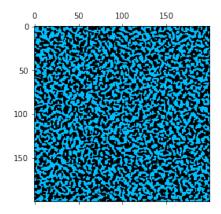
$$= \beta(\Delta m_i \sum_{i'} m_{i'} + \Delta m_j \sum_{j'} m_{j'} + \Delta m_i \Delta m_j)$$

$$= \beta \Delta m_i \left(\sum_{i'} m_{i'} - \sum_{j'} m_{j'} - \Delta m_i\right)$$

$$= \beta \Delta m_i \left(\sum_{i'} m_{i'} - m_j - \sum_{j'} m_{j'} + m_i\right)$$

$$= \beta(m_j - m_i) \left(\sum_{i' \neq j} m_{i'} - \sum_{j' \neq i} m_{j'}\right)$$





- (a) 200×200 lattice initialized uniformly at random with values $\in \{0, 1\}$.
- (b) 200×200 lattice after 15000000 iterations with $\beta = 100$.

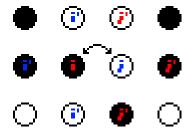


Figure 2: Sketch of the interaction between i, j and neighbouring sites. i interacts with the blue sites, and j with the red ones.

For exact solution: Work out the 1D problem.

$$H(m) = -J \sum_{i} m_i \ m_{i+1}$$
 if m is a cycle, $J \neq 0$ iff $j = i+1$

$$= -J \sum_{i} k_i$$
 where $k_i = 1$ if $m_i = m_{i+1} = 1$, $k_i = 0$ otherwise

This system has $\binom{N}{M}$ allowed configurations. To compute the partition function we must first find the degeneracy of all values of $K = \sum_i k_i$. This is equivalent to computing the number c(K, M, N) of induced subgraphs with M vertices and K edges in a cycle of length N.

Let's conceptualize a cycle of length N-1 as an open string of length N with periodic boundary conditions. We'll start with an empty string and calculate the number of ways to put M particles and K interacting pairs in it. Now let's take a vertex at one end of the string. We can either fill in with a particle or leave it empty. This observation gives us a simple way to divide the length N problem into two smaller problems. One of them is the calculation of the number of ways to put M-1 particles and K interacting pairs in a length N string with both ends filled. The other one is the calculation of the number of ways to put M particles and K interacting pairs in a length N string with one end filled and the one end empty. Repeating the process, we find that $c(K, M, N) = f(K, M-1, N-1) + \sum_{O=1}^{N-2} g(K, M-1, O)$, where f(K, M, N) and g(K, M, N) are the calculations on smaller strings we mentioned. Similarly to how we decomposed c(K, M, N) in a combination of f(K, M, N) and g(K, M, N), f(K, M, N) and g(K, M, N) can be shown to obey the following recurrences:

$$\begin{split} f(K,M,N) &= f(K-1,M-1,N-1) + g(K,M,N-1) \\ g(K,M,N) &= f(K,M-1,N-1) + g(K,M,N-1) \\ \Rightarrow f(M,N,K) - g(M,N,K) &= f(M-1,N-1,K-1) - f(M-1,N-1,K) \\ \Rightarrow g(M,N-1,K) &= f(M,N-1,K) - f(M-1,N-2,K-1) \\ &+ f(M-1,N-2,K) \\ \Rightarrow f(M,N,K) &= f(M-1,N-1,K-1) + f(M,N-1,K) \\ &+ f(M-1,N-2,K) - f(M-1,N-2,K-1) \end{split}$$

We'll use the method of characteristic functions to find f and g. Let $F(x,y,z) = \sum_{K,M,N} f(K,M,N) x^K y^M z^N$, $G(x,y,z) = \sum_{K,M,N} g(K,M,N) x^K y^M z^N$ and f(1,0,0) = g(0,0,0) = 1, f(K,M,0) = g(K,M,0) = 0 for all other K,M. We need to calculate F(x,y,z) on subsets with one or two arguments equal to 0 before we're able to fully calculate F(x,y,z) and f(K,M,N). Directly from initial conditions, we get g(0,0,N) = f(0,0,N) = 1 and f(K,0,N) = 1

$$g(K,0,N) = \begin{cases} 1 & K=0 \\ 0 & K \neq 0 \end{cases}$$
 when $N \neq 0$. This result allows us to calculate

F(x, 0, z) and G(x, 0, z).

$$F(x, 0, z) = \sum_{K,N} f(K, 0, N) x^{K} z^{N}$$

$$= x + \sum_{N \neq 0} z^{N}$$

$$= x + \frac{z}{1 - z}$$

$$G(x, 0, z) = \sum_{K,N} g(K, 0, N) x^{K} z^{N}$$

$$= 1 + \frac{z}{1 - z}$$

Calculating F(0, y, z) and G(0, y, z) is a bit more complicated.

$$f(0, M, N) = g(0, M, N - 1)$$

$$g(0, M, N) = f(0, M - 1, N - 1) + g(0, M, N - 1)$$

$$\begin{split} \sum_{M,N} f(0,M,N+1)y^M z^N &= \sum_{M,N} g(0,M,N)y^M z^N \\ \frac{1}{z} \Big(\sum_{M,N} f(0,M,N)y^M z^N - \sum_{M} f(0,M,0)y^M \Big) &= \sum_{M,N} g(0,M,N)y^M z^N \\ \frac{1}{z} \sum_{M,N} f(0,M,N)y^M z^N &= \sum_{M,N} g(0,M,N)y^M z^N \\ F(0,y,z) &= zG(0,y,z) \end{split}$$

$$\sum_{M,N} g(0,M+1,N+1)y^M z^N &= \sum_{M,N} f(0,M,N)y^M z^N + \sum_{M,N} g(0,M+1,N)y^M z^N \\ \frac{1}{yz} \Big(\sum_{M,N} g(0,M,N)y^M z^N - \frac{1}{1-z} \Big) &= \sum_{M,N} f(0,M,N)y^M z^N + \frac{1}{y} \Big(\sum_{M,N} g(0,M,N)y^M z^N - \frac{1}{1-z} \Big) \\ G(0,y,z) - \frac{1}{1-z} &= yzF(0,y,z) + zG(0,y,z) - \frac{z}{1-z} \end{split}$$

$$G(0, y, z) = yz^{2}G(0, y, z) + zG(0, y, z) + 1$$

$$G(0, y, z) = \frac{1}{1 - z - yz^{2}}$$

$$F(0, y, z) = \frac{z}{1 - z - yz^{2}}$$

$$\begin{split} G(0,y,z) &= \sum_{n=0}^{\infty} (z+yz^2)^n \\ &= \sum_{n=0}^{\infty} \sum_{k=0}^{n} \binom{n}{k} z^{n-k} y^k z^{2k} \\ &= \sum_{n=0}^{\infty} \sum_{k=0}^{n} \binom{n}{k} y^k z^{n+k} \\ &= \sum_{M,N} g(0,M,N) y^M z^N \text{ where } g(0,M,N) = \begin{cases} \binom{N-M}{M} & M \in \{0,N/2\} \\ 0 & M \notin \{0,N/2\} \end{cases} \end{split}$$

$$F(0,y,z) = \sum_{M,N} g(0,M,N) y^M z^N \text{ where } g(0,M,N) = \begin{cases} \binom{N-M-1}{M} & M \in \{0,N/2-1\} \\ 0 & M \notin \{0,N/2-1\} \end{cases}$$

It takes a bit more work to calculate G(x, y, z).

$$\begin{split} \sum_{K,M,N} f(K+1,M+1,N+1) x^K y^M z^N &= \sum_{K,M,N} f(K,M,N) x^K y^M z^N + \sum_{K,M,N} g(K+1,M+1,N) x^K y^M z^N \\ &\frac{1}{xyz} \bigg(F(x,y,z) - \frac{z}{1-z-yz^2} - x \bigg) = F(x,y,z) + \frac{1}{xy} \bigg(G(x,y,z) - \frac{1}{1-z-yz^2} \bigg) \\ &F(x,y,z) - x = xyz F(x,y,z) + z G(x,y,z) \end{split}$$

$$\begin{split} \sum_{K,M,N} g(K,M+1,N+1) x^K y^M z^N &= \sum_{K,M,N} f(K,M,N) x^K y^M z^N + \sum_{K,M,N} g(K,M+1,N) x^K y^M z^N \\ &\frac{1}{yz} \bigg(G(x,y,z) - 1 - \frac{z}{1-z} \bigg) = F(x,y,z) + \frac{1}{y} \bigg(G(x,y,z) - 1 - \frac{z}{1-z} \bigg) \\ &G(x,y,z) - 1 - \frac{z}{1-z} = yz F(x,y,z) + z G(x,y,z) - z - \frac{z^2}{1-z} \\ &G(x,y,z) = yz F(x,y,z) + z G(x,y,z) + 1 \end{split}$$

$$F(x, y, z) = \frac{z}{1 - xyz}G(x, y, z) + \frac{x}{1 - xyz}$$

$$G(x, y, z) = \frac{yz^2}{1 - xyz}G(x, y, z) + \frac{xyz}{1 - xyz} + zG(x, y, z) + 1$$

$$\frac{1 - xyz - yz^2 - z + xyz^2}{1 - xyz}G(x, y, z) = \frac{1}{1 - xyz}$$

$$G(x, y, z) = \frac{1}{1 - xyz - yz^2 - z + xyz^2}$$

$$= \sum_{x=0}^{\infty} (z + yz^2 + xyz - xyz^2)^n$$

$$G(x,y,z) = \sum_{n=0}^{\infty} (z + yz^{2} + xyz - xyz^{2})^{n}$$

$$= \sum_{n=0}^{\infty} \sum_{k=0}^{n} \binom{n}{k} (z + yz^{2})^{k} (xyz - xyz^{2})^{n-k}$$

$$= \sum_{n=0}^{\infty} \sum_{k=0}^{n} \binom{n}{k} \sum_{l=0}^{k} \binom{k}{l} z^{k-l} y^{l} z^{2l}$$

$$= \sum_{m=0}^{n-k} \binom{n-k}{m} (-1)^{m} x^{n-k-m} y^{n-k-m} z^{n-k-m} x^{m} y^{m} z^{2m}$$

$$= \sum_{n=0}^{\infty} \sum_{k=0}^{n} \binom{n}{k} \sum_{l=0}^{k} \binom{k}{l} y^{l} z^{k+l} \sum_{m=0}^{n-k} \binom{n-k}{m} (-1)^{m} x^{n-k} y^{n-k} z^{n-k+m}$$

$$= \sum_{n=0}^{\infty} \sum_{k=0}^{n} \sum_{l=0}^{k} \sum_{m=0}^{n-k} \binom{n}{k} \binom{k}{l} \binom{n-k}{m} (-1)^{m} x^{n-k} y^{n-k+l} z^{n+m+l}$$

$$K = n - k$$

$$M = n - k + l$$

$$N = n + m + l$$

$$\begin{split} &= \sum_{K,M,N} \sum_{m=0}^{K} \binom{N-m-M+K}{N-m-M} \binom{N-m-M}{M-K} \binom{K}{m} (-1)^m x^K y^M z^N \\ &= \sum_{K,M,N} \sum_{m=0}^{K} \frac{(N-m-M+K)!}{(N-m-2M+K)!} \binom{K}{m} (-1)^m \frac{1}{K!} \frac{1}{(M-K)!} x^K y^M z^N \\ &= \sum_{K,M,N} \sum_{m=0}^{K} (-1)^m \binom{K}{m} \binom{N-m-M+K}{N-m-2M+K} \frac{M!}{K!(M-K)!} x^K y^M z^N \\ &= \sum_{K,M,N} \sum_{m=0}^{K} (-1)^m \binom{K}{m} \binom{N-m-M+K}{M} \binom{M}{K} x^K y^M z^N \\ &= \sum_{K,M,N} \binom{N-M}{M-K} \binom{M}{M} x^K y^M z^N \text{ by Wolfram} \end{split}$$

$$g(K,M,N) = \begin{cases} \binom{N-M}{M-K} \binom{M}{K} & M \in \{0,N\} \\ 0 & M \notin \{0,N\} \end{cases}$$

We can easily calculate f(K, M, N) with the help of the recurrence relations we wrote down at the beginning.

$$\begin{split} f(K,M,N) &= g(K,M+1,N+1) - g(K,M+1,N) \\ &= \binom{N-M}{M-K+1} \binom{M+1}{K} - \binom{N-M-1}{M-K+1} \binom{M+1}{K} \\ &= \binom{N-M-1}{M-K} \binom{M+1}{K} \text{by Pascal's identity} \end{split}$$

$$c(K, M, N) = f(K, M - 1, N - 1) + \sum_{O=1}^{N-2} g(K, M - 1, O)$$

$$= {N - M - 1 \choose M - K - 1} {M \choose K} + \sum_{O=M-1}^{N-2} {O - M + 1 \choose M - K - 1} {M - 1 \choose K}$$

$$\sum_{O=M-1}^{N-1} \binom{O-M+1}{M-K-1} = \sum_{O=M-1}^{N-2} \binom{O-M+2}{M-K} - \sum_{O=M-1}^{N-2} \binom{O-M+1}{M-K}$$

$$= \sum_{O=M}^{N-1} \binom{O-M+1}{M-K} - \sum_{O=M-1}^{N-2} \binom{O-M+1}{M-K}$$

$$= \binom{N-M}{M-K}$$

$$c(K, M, N) = \binom{N-M-1}{M-K-1} \binom{M}{K} + \binom{N-M}{M-K} \binom{M-1}{K}$$

We take the thermodynamics limit $M \to \infty$, $N \to \infty$, $\frac{M}{N} = \rho$. We'll assume the maximum of c(K, M, N) along K is attained when $K \propto M$ and find the constant of proportionality η . η can be thought of as the number of links per particle.

$$c(K,M,N) \approx \sqrt{\frac{N-M-1}{2\pi(M-K-1)(N-2M+K)}} \\ \frac{(N-M-1)^{N-M-1}}{(M-K-1)^{M-K-1}(N-2M+K)^{N-2M+K}} \\ \sqrt{\frac{M}{2\pi K(M-K)}} \frac{M^M}{K^K(M-K)^{M-K}} \\ + \sqrt{\frac{N-M}{2\pi(M-K)(N-2M+K)}} \\ \frac{(N-M)^{N-M}}{(M-K)^{M-K}(N-2M+K)^{N-2M+K}} \\ \sqrt{\frac{M-1}{2\pi K(M-K-1)}} \frac{(M-1)^{M-1}}{K^K(M-K-1)^{M-K-1}}$$

$$\begin{split} Z &= \sum_{K} c(K, M, N) e^{\beta K} \\ &\approx \sum_{K} \sqrt{\frac{N - M - 1}{2\pi (M - K - 1)(N - 2M + K)}} \\ &\frac{(N - M - 1)^{N - M - 1}}{(M - K - 1)^{M - K - 1}(N - 2M + K)^{N - 2M + K}} \\ &\sqrt{\frac{M}{2\pi K(M - K)}} \frac{M^{M}}{K^{K}(M - K)^{M - K}} e^{\beta K} \\ &+ \sum_{K} \sqrt{\frac{N - M}{2\pi (M - K)(N - 2M + K)}} \\ &\frac{(N - M)^{N - M}}{(M - K)^{M - K}(N - 2M + K)^{N - 2M + K}} \\ &\sqrt{\frac{M - 1}{2\pi K(M - K - 1)}} \frac{(M - 1)^{M - 1}}{K^{K}(M - K - 1)^{M - K - 1}} e^{\beta K} \end{split}$$

$$\log(c(K, M, N)e^{\beta K}) \sim (N - M)\log(N - M) - (M - K)\log(M - K)$$

$$- (N - 2M + K)\log(N - 2M + K) + M\log(M)$$

$$- K\log(K) - (M - K)\log(M - K) + \beta K$$

$$\frac{\partial}{\partial K}\log(c(K, M, N)e^{\beta K}) = 2\log(M - K) + 2 - \log(N - 2M + K)$$

$$- 1 - \log(K) - 1 + \beta$$

$$= 2\log(M - K) - \log(N - 2M + K) - \log(K) + \beta$$

$$0 = 2\log(M - K) - \log(N - 2M + K) - \log(K) + \beta$$

$$e^{-\beta} = \frac{(M - K)^2}{(N - 2M + K)K}$$

$$e^{-\beta} = \frac{(1 - \eta)^2}{(1/\rho - 2 + \eta)\eta}$$

$$K = \frac{\sqrt{(e^{-\beta} - 1)M^2 + (e^{-\beta}M - \frac{1}{2}e^{-\beta}N - M)^2} + (e^{-\beta}M - \frac{1}{2}e^{-\beta}N - M)}{e^{-\beta} - 1}$$

$$\eta = \frac{\sqrt{(e^{-\beta} - 1) + (e^{-\beta} - \frac{1}{2\rho}e^{-\beta} - 1)^2} + (e^{-\beta} - \frac{1}{2\rho}e^{-\beta} - 1)}{e^{-\beta} - 1}$$

$$\eta = \sqrt{\frac{1}{|e^{-\beta} - 1|} + \left(1 - \frac{1}{2\rho|e^{\beta} - 1|}\right)^2} + 1 - \frac{1}{2\rho|e^{\beta} - 1|}$$

$$C = -\frac{\partial \eta}{\partial \beta} = \frac{\left(2\left(\rho - 1\right)\rho\left(e^{\beta} - 1\right) + \sqrt{1 - 4\left(\rho - 1\right)\rho\left(e^{\beta} - 1\right)} - 1\right)e^{\beta}}{2\sqrt{1 - 4\left(\rho - 1\right)\rho\left(e^{\beta} - 1\right)}\left(e^{\beta} - 1\right)^{2}}$$

Some interesting limits include $\beta \to \infty$, $\beta \to 0$ and $\beta \to -\infty$.

$$\beta \to \infty \Rightarrow \eta \to \sqrt{-1+1}+1$$

= 1 particles are clustering

$$\begin{split} \beta \to 0 \Rightarrow \eta &= \frac{|e^{-\beta} - \frac{1}{2\rho}e^{-\beta} - 1|\sqrt{1 + \frac{e^{-\beta} - 1}{(e^{-\beta} - \frac{1}{2\rho}e^{-\beta} - 1)^2}} + (e^{-\beta} - \frac{1}{2\rho}e^{-\beta} - 1)}{e^{-\beta} - 1} \\ &\to \frac{\frac{1}{2\rho}\sqrt{1 + 4(e^{-\beta} - 1)\rho^2} - \frac{1}{2\rho}}{e^{-\beta} - 1} \\ &\to \frac{\frac{1}{2\rho} + (e^{-\beta} - 1)\rho - \frac{1}{2\rho}}{e^{-\beta} - 1} \end{split}$$

 $\rightarrow \rho$ particles are distributed randomly

$$\begin{split} \beta \Rightarrow -\infty \Rightarrow \eta \to \left| 1 - \frac{1}{2\rho} \right| + 1 - \frac{1}{2\rho} \\ &= \begin{cases} 2 - \frac{1}{\rho} & \rho \ge \frac{1}{2} \\ 0 & \rho < \frac{1}{2} \end{cases} \text{ particles avoid contact} \end{split}$$

The discontinuity indicates a phase transition when $\beta \to -\infty$ and $\rho \to \frac{1}{2}$.