

# The Tetrad Formulation of General Relativity

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- Traditional approach uses coordinate bases  $\hat{e}_{(\mu)} = \partial_\mu$
- Noncoordinate bases  $\hat{e}_{(a)}$  are not derived from any coordinate system
- Advantages:
  - Reveals connection to gauge theories
  - Simplifies calculations in certain metrics
  - Enables description of spinor fields

## Basis Vectors

At each point, introduce basis vectors  $\hat{e}_{(a)}$  (Latin indices) satisfying:

$$g(\hat{e}_{(a)}, \hat{e}_{(b)}) = \eta_{ab}$$

where  $\eta_{ab}$  is the canonical form of the metric.

## Terminology

- Tetrad (4D), vielbein ("many legs"), dreibein (3D), zweibein (2D)
- Relation to coordinate basis:

$$\hat{e}_{(\mu)} = e_{\mu}^a \hat{e}_{(a)}$$

$$\hat{e}_{(a)} = e_a^{\mu} \hat{e}_{(\mu)}$$

## Metric Relations

$$g_{\mu\nu} e_a^\mu e_b^\nu = \eta_{ab}$$

$$g_{\mu\nu} = e_\mu^a e_\nu^b \eta_{ab}$$

## Basis One-Forms

Orthonormal one-forms  $\hat{\theta}^{(a)}$  compatible with basis vectors:

$$\hat{\theta}^{(a)}(\hat{e}_{(b)}) = \delta_b^a$$

Relations to coordinate basis:

$$\hat{\theta}^{(\mu)} = e_a^\mu \hat{\theta}^{(a)}$$

$$\hat{\theta}^{(a)} = e_\mu^a \hat{\theta}^{(\mu)}$$

## Vector Components

For vector  $V = V^\mu \hat{e}_{(\mu)} = V^a \hat{e}_{(a)}$ :

$$V^a = e_\mu^a V^\mu$$

$$V^\mu = e_a^\mu V^a$$

## Tensor Components

Mixed components transform as:

$$V_b^a = e_\mu^a e_b^\nu V_\nu^\mu$$

## Index Interpretation

- Greek indices: "curved" spacetime indices
- Latin indices: "flat" tangent space indices

## Basic Concept

- Tetrad: Set of four orthonormal basis vectors  $\{e_{(a)}^\mu\}$  at each spacetime point
- In Minkowski space:  $\eta_{ab} = e_{(a)}^\mu e_{(b)}^\nu \eta_{\mu\nu}$
- Relates local frame to global coordinates

## Inertial Observer Tetrad

For standard Minkowski coordinates  $(t, x, y, z)$ :

$$e_{(0)} = (1, 0, 0, 0), \quad e_{(1)} = (0, 1, 0, 0), \quad e_{(2)} = (0, 0, 1, 0), \quad e_{(3)} = (0, 0, 0, 1)$$

Satisfies  $e_{(a)} \cdot e_{(b)} = \eta_{ab}$

# Tetrad for Accelerated Observer

## Rindler Coordinates

For observer with constant acceleration  $a$  in  $x$ -direction:

$$t = \frac{1}{a} e^{a\xi} \sinh(a\eta)$$
$$x = \frac{1}{a} e^{a\xi} \cosh(a\eta)$$

## Rindler Tetrad

$$e_{(0)} = \frac{1}{a\sqrt{x^2 - t^2}} (x\partial_t + t\partial_x)$$
$$e_{(1)} = \frac{1}{a\sqrt{x^2 - t^2}} (t\partial_t + x\partial_x)$$
$$e_{(2)} = \partial_y, \quad e_{(3)} = \partial_z$$

Satisfies  $e_{(a)} \cdot e_{(b)} = \eta_{ab}$  along hyperbolic trajectory



# Spin Connection Definition

## Covariant Derivative

For tensor with Latin indices:

$$\nabla_\mu X_b^a = \partial_\mu X_b^a + \omega_{\mu c}^a X_b^c - \omega_{\mu b}^c X_c^a$$

## Tetrad Postulate

$$\nabla_\mu e_\nu^a = \partial_\mu e_\nu^a - \Gamma_{\mu\nu}^\lambda e_\lambda^a + \omega_{\mu b}^a e_\nu^b = 0$$

## Relation to Christoffel Symbols

$$\Gamma_{\mu\lambda}^\nu = e_a^\nu \partial_\mu e_\lambda^a + e_a^\nu e_\lambda^b \omega_{\mu b}^a$$

$$\omega_{\mu b}^a = e_\nu^a e_b^\lambda \Gamma_{\mu\lambda}^\nu - e_b^\lambda \partial_\mu e_\lambda^a$$

## Under Local Lorentz Transformations

$$\omega_{\mu}^a{}_b \rightarrow \Lambda_{a'}^a \Lambda_b^{b'} \omega_{\mu}^{a'}{}_{b'} - \Lambda_c^{b'} \partial_{\mu} \Lambda_c^{a'}$$

where  $\Lambda_{a'}$  satisfies:

$$\Lambda_{a'}^a \Lambda_b^b \eta_{ab} = \eta_{a'b'}$$

## Under General Coordinate Transformations

Transforms as a one-form in the Greek index:

$$\omega_{\mu}^a{}_b \rightarrow \frac{\partial x^{\nu}}{\partial x^{\mu'}} \omega_{\nu}^a{}_b$$

# Cartan Structure Equations

## Basis One-Forms and Connection

Define:

$$e^a = e^a_\mu dx^\mu$$

$$\omega^a_b = \omega^a_{\mu b} dx^\mu$$

## Torsion and Curvature

$$T^a = de^a + \omega^a_b \wedge e^b$$

$$R^a_b = d\omega^a_b + \omega^a_c \wedge \omega^c_b$$

## Bianchi Identities

$$dT^a + \omega^a_b \wedge T^b = R^a_b \wedge e^b$$

$$dR^a_b + \omega^a_c \wedge R^c_b - R^a_c \wedge \omega^c_b = 0$$

# Cartan Structure Equations

## Geometric Interpretation

- Torsion measures failure of infinitesimal parallelograms to close
- First structure equation defines torsion as "twisting" of frame fields
- Second structure equation relates curvature to connection rotation

## Bianchi Identities from Cartan Formalism

$$DT^a = R^a_b \wedge e^b \quad (\text{First identity})$$

$$DR^a_b = 0 \quad (\text{Second identity})$$

where  $D$  is exterior covariant derivative:

$$DT^a = dT^a + \omega^a_b \wedge T^b$$

## General Definition

At each point  $p$  in spacetime:

$$g_{\mu\nu} e_{(a)}^\mu e_{(b)}^\nu = \eta_{ab}$$

Inverse relation:

$$g_{\mu\nu} = \eta_{ab} e_\mu^{(a)} e_\nu^{(b)}$$

## Physical Interpretation

- $e_{(a)}^\mu$  gives local inertial frame at each point
- $e_\mu^{(a)}$  converts between coordinate and local frames
- Spin connection  $\omega_\mu^{ab}$  describes how frames rotate during parallel transport

# Example: FLRW Metric in Vielbein Form

## Metric and Vielbein Choice

$$ds^2 = -dt^2 + a^2(t)\delta_{ij}dx^i dx^j$$

Choose orthonormal basis:

$$e^0 = dt, \quad e^i = a(t)dx^i$$

## Solving for Spin Connection

From torsion-free condition  $\omega^a_b \wedge e^b = -de^a$ :

$$de^0 = 0$$

$$de^i = \dot{a}dt \wedge dx^i$$

With antisymmetry:  $\omega^0_0 = 0, \omega^0_j = \omega^j_0, \omega^i_j = -\omega^j_i$

## Solution

$$\omega^0_j = \dot{a} dx^j, \quad \omega^i_0 = \dot{a} dx^i, \quad \omega^i_j = 0$$

Verification:

$$\begin{aligned}\omega^0_j \wedge e^j &= \dot{a} dx^j \wedge a dx^j = 0 \\ \omega^i_0 \wedge e^0 + \omega^i_j \wedge e^j &= \dot{a} dx^i \wedge dt + 0 = -\dot{a} dt \wedge dx^i\end{aligned}$$

## Curvature Two-Form (Key Components)

$$R^0_j = \ddot{a} dt \wedge dx^j$$

$$R^i_0 = \ddot{a} dt \wedge dx^i$$

$$R^i_j = \dot{a}^2 dx^i \wedge dx^j$$

# Converting FLRW Curvature to Coordinate Basis

## Vielbein Components

$$e_{\mu}^a = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & a & 0 & 0 \\ 0 & 0 & a & 0 \\ 0 & 0 & 0 & a \end{pmatrix}, \quad e_b^{\nu} = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & \frac{1}{a} & 0 & 0 \\ 0 & 0 & \frac{1}{a} & 0 \\ 0 & 0 & 0 & \frac{1}{a} \end{pmatrix}$$

## Riemann Tensor Components

Using  $R^{\rho}_{\sigma\mu\nu} = e_a^{\rho} e_{\sigma}^b R^a_{b\mu\nu}$ :

$$R^0_{j0i} = a\ddot{a}\delta_{ji}$$

$$R^i_{0k0} = -\frac{\ddot{a}}{a}\delta_k^i$$

$$R^i_{jkl} = \dot{a}^2(\delta_k^i\delta_{jl} - \delta_l^i\delta_{jk})$$



## Canonical Tetrad

For Schwarzschild metric:

$$ds^2 = -(1 - 2M/r)dt^2 + \frac{dr^2}{1 - 2M/r} + r^2 d\Omega^2$$

Choose:

$$e_{(0)} = \sqrt{1 - 2M/r} \partial_t$$

$$e_{(1)} = \frac{1}{\sqrt{1 - 2M/r}} \partial_r$$

$$e_{(2)} = \frac{1}{r} \partial_\theta$$

$$e_{(3)} = \frac{1}{r \sin \theta} \partial_\phi$$

## Non-zero Spin Connection Components

$$\omega^0_1 = \frac{M}{r^2} dt$$

$$\omega^1_2 = -\sqrt{1 - 2M/r} d\theta$$

$$\omega^1_3 = -\sqrt{1 - 2M/r} \sin \theta d\phi$$

$$\omega^2_3 = -\cos \theta d\phi$$

## Advantages of Tetrad Approach

- Separates gravitational and coordinate effects
- Reveals physical structure of spacetime
- Simplifies calculation of curvature invariants

- Extension of GR allowing for torsion  $T^a \neq 0$
- Spin connection  $\omega_\mu^{ab}$  becomes independent dynamical field
- Field equations:

$$G^{ab} = 8\pi T^{ab}$$
$$T^a{}_{bc} = 8\pi \tau^a{}_{bc}$$

where  $\tau^a{}_{bc}$  is spin density tensor

## Tetrad Action

Action in terms of tetrad and spin connection:

$$S = \frac{1}{16\pi} \int e^a \wedge e^b \wedge R^{cd} \epsilon_{abcd} + S_{\text{matter}}$$

where  $e = \det(e_\mu^a)$  and  $\epsilon_{abcd}$  is Levi-Civita tensor

## Field Equations Derivation

Varying with respect to tetrad:

$$\epsilon_{abcd} e^b \wedge R^{cd} = 16\pi \tau^{(a)}$$

where  $\tau^{(a)}$  is energy-momentum 3-form

Varying with respect to spin connection:

$$\epsilon_{abcd} T^c \wedge e^d = 16\pi \sigma_{ab}$$

where  $\sigma_{ab}$  is spin current 3-form

- Torsion propagates spin density through spacetime
- Avoids singularities in some cases
- Reduces to GR in vacuum (no spin sources)

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