

The Tetrad Formulation of General Relativity

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Outline

- 1 Introduction
- 2 Vielbein Formalism
- 3 Tetrads in Special Relativity
- 4 Spin Connection
- 5 Differential Forms Approach
- 6 Tetrads in Curved Spacetime
- 7 Example Applications
- 8 Einstein-Cartan Theory

Motivation for Noncoordinate Bases

- Traditional approach uses coordinate bases $\hat{e}_{(\mu)} = \partial_\mu$
- Noncoordinate bases $\hat{e}_{(a)}$ are not derived from any coordinate system
- Advantages:
 - Reveals connection to gauge theories
 - Simplifies calculations in certain metrics
 - Enables description of spinor fields

Basis Vectors

At each point, introduce basis vectors $\hat{e}_{(a)}$ (Latin indices) satisfying:

$$g(\hat{e}_{(a)}, \hat{e}_{(b)}) = \eta_{ab}$$

where η_{ab} is the canonical form of the metric.

Terminology

- Tetrad (4D), vielbein ("many legs"), dreibein (3D), zweibein (2D)
- Relation to coordinate basis:

$$\hat{e}_{(\mu)} = e_\mu^a \hat{e}_{(a)}$$

$$\hat{e}_{(a)} = e_a^\mu \hat{e}_{(\mu)}$$

Metric Relations

$$g_{\mu\nu} e_a^\mu e_b^\nu = \eta_{ab}$$

$$g_{\mu\nu} = e_\mu^a e_\nu^b \eta_{ab}$$

Basis One-Forms

Orthonormal one-forms $\hat{\theta}^{(a)}$ compatible with basis vectors:

$$\hat{\theta}^{(a)}(\hat{e}_{(b)}) = \delta_b^a$$

Relations to coordinate basis:

$$\hat{\theta}^{(\mu)} = e_a^\mu \hat{\theta}^{(a)}$$

$$\hat{\theta}^{(a)} = e_\mu^a \hat{\theta}^{(\mu)}$$

Component Transformations

Vector Components

For vector $V = V^\mu \hat{e}_{(\mu)} = V^a \hat{e}_{(a)}$:

$$V^a = e_\mu^a V^\mu$$

$$V^\mu = e_a^\mu V^a$$

Tensor Components

Mixed components transform as:

$$V_b^a = e_\mu^a e_b^\nu V_\nu^\mu$$

Index Interpretation

- Greek indices: "curved" spacetime indices
- Latin indices: "flat" tangent space indices

Basic Concept

- Tetrad: Set of four orthonormal basis vectors $\{e_{(a)}^\mu\}$ at each spacetime point
- In Minkowski space: $\eta_{ab} = e_{(a)}^\mu e_{(b)}^\nu \eta_{\mu\nu}$
- Relates local frame to global coordinates

Inertial Observer Tetrad

For standard Minkowski coordinates (t, x, y, z) :

$$e_{(0)} = (1, 0, 0, 0), \quad e_{(1)} = (0, 1, 0, 0), \quad e_{(2)} = (0, 0, 1, 0), \quad e_{(3)} = (0, 0, 0, 1)$$

Satisfies $e_{(a)} \cdot e_{(b)} = \eta_{ab}$

Tetrad for Accelerated Observer

Rindler Coordinates

For observer with constant acceleration a in x -direction:

$$t = \frac{1}{a} e^{a\xi} \sinh(a\eta)$$
$$x = \frac{1}{a} e^{a\xi} \cosh(a\eta)$$

Rindler Tetrad

$$\mathbf{e}_{(0)} = \frac{1}{a\sqrt{x^2 - t^2}} (x\partial_t + t\partial_x)$$
$$\mathbf{e}_{(1)} = \frac{1}{a\sqrt{x^2 - t^2}} (t\partial_t + x\partial_x)$$
$$\mathbf{e}_{(2)} = \partial_y, \quad \mathbf{e}_{(3)} = \partial_z$$

Satisfies $\mathbf{e}_{(a)} \cdot \mathbf{e}_{(b)} = \eta_{ab}$ along hyperbolic trajectory

Covariant Derivative

For tensor with Latin indices:

$$\nabla_{\mu} X_b^a = \partial_{\mu} X_b^a + \omega_{\mu c}^a X_b^c - \omega_{\mu b}^c X_c^a$$

Tetrad Postulate

$$\nabla_{\mu} e_{\nu}^a = \partial_{\mu} e_{\nu}^a - \Gamma_{\mu\nu}^{\lambda} e_{\lambda}^a + \omega_{\mu b}^a e_{\nu}^b = 0$$

Relation to Christoffel Symbols

$$\Gamma_{\mu\lambda}^{\nu} = e_{a}^{\nu} \partial_{\mu} e_{\lambda}^a + e_{a}^{\nu} e_{\lambda}^b \omega_{\mu b}^a$$

$$\omega_{\mu b}^a = e_{\nu}^a e_b^{\lambda} \Gamma_{\mu\lambda}^{\nu} - e_b^{\lambda} \partial_{\mu} e_{\lambda}^a$$

Transformation Properties

Under Local Lorentz Transformations

$$\omega_{\mu}^a b \rightarrow \Lambda_{a'}^a \Lambda_b^{b'} \omega_{\mu}^{a'} b' - \Lambda_c^{b'} \partial_{\mu} \Lambda_c^{a'}$$

where $\Lambda_{a'}^a$ satisfies:

$$\Lambda_{a'}^a \Lambda_{b'}^b \eta_{ab} = \eta_{a'b'}$$

Under General Coordinate Transformations

Transforms as a one-form in the Greek index:

$$\omega_{\mu}^a b \rightarrow \frac{\partial x^{\nu}}{\partial x^{\mu'}} \omega_{\nu}^a b$$

Cartan Structure Equations

Basis One-Forms and Connection

Define:

$$e^a = e_\mu^a dx^\mu$$

$$\omega^a{}_b = \omega_\mu{}^a{}_b dx^\mu$$

Torsion and Curvature

$$T^a = de^a + \omega^a{}_b \wedge e^b$$

$$R^a{}_b = d\omega^a{}_b + \omega^a{}_c \wedge \omega^c{}_b$$

Bianchi Identities

$$dT^a + \omega^a{}_b \wedge T^b = R^a{}_b \wedge e^b$$

$$dR^a{}_b + \omega^a{}_c \wedge R^c{}_b - R^a{}_c \wedge \omega^c{}_b = 0$$

Geometric Interpretation

- Torsion measures failure of infinitesimal parallelograms to close
- First structure equation defines torsion as "twisting" of frame fields
- Second structure equation relates curvature to connection rotation

Bianchi Identities from Cartan Formalism

$$DT^a = R^a{}_b \wedge e^b \quad (\text{First identity})$$

$$DR^a{}_b = 0 \quad (\text{Second identity})$$

where D is exterior covariant derivative:

$$DT^a = dT^a + \omega^a{}_b \wedge T^b$$

General Definition

At each point p in spacetime:

$$g_{\mu\nu} e_{(a)}^{\mu} e_{(b)}^{\nu} = \eta_{ab}$$

Inverse relation:

$$g_{\mu\nu} = \eta_{ab} e_{\mu}^{(a)} e_{\nu}^{(b)}$$

Physical Interpretation

- $e_{(a)}^{\mu}$ gives local inertial frame at each point
- $e_{\mu}^{(a)}$ converts between coordinate and local frames
- Spin connection ω_{μ}^{ab} describes how frames rotate during parallel transport

Example: FLRW Metric in Vielbein Form

Metric and Vielbein Choice

$$ds^2 = -dt^2 + a^2(t)\delta_{ij}dx^i dx^j$$

Choose orthonormal basis:

$$e^0 = dt, \quad e^i = a(t)dx^i$$

Solving for Spin Connection

From torsion-free condition $\omega^a{}_b \wedge e^b = -de^a$:

$$de^0 = 0$$

$$de^i = \dot{a}dt \wedge dx^i$$

With antisymmetry: $\omega^0{}_0 = 0$, $\omega^0{}_j = \omega^j{}_0$, $\omega^i{}_j = -\omega^j{}_i$

Solution

$$\omega^0{}_j = \dot{a}dx^j, \quad \omega^i{}_0 = \dot{a}dx^i, \quad \omega^i{}_j = 0$$

Verification:

$$\omega^0{}_j \wedge e^j = \dot{a}dx^j \wedge adx^j = 0$$

$$\omega^i{}_0 \wedge e^0 + \omega^i{}_j \wedge e^j = \dot{a}dx^i \wedge dt + 0 = -\dot{a}dt \wedge dx^i$$

Curvature Two-Form (Key Components)

$$R^0{}_j = \ddot{a}dt \wedge dx^j$$

$$R^i{}_0 = \ddot{a}dt \wedge dx^i$$

$$R^i{}_j = \dot{a}^2 dx^i \wedge dx^j$$

Vielbein Components

$$e_\mu^a = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & a & 0 & 0 \\ 0 & 0 & a & 0 \\ 0 & 0 & 0 & a \end{pmatrix}, \quad e_b^\nu = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & \frac{1}{a} & 0 & 0 \\ 0 & 0 & \frac{1}{a} & 0 \\ 0 & 0 & 0 & \frac{1}{a} \end{pmatrix}$$

Riemann Tensor Components

Using $R^\rho_{\sigma\mu\nu} = e_a^\rho e_\sigma^b R^a_{b\mu\nu}$:

$$R^0_{j0i} = a\ddot{a}\delta_{ji}$$

$$R^i_{0k0} = -\frac{\ddot{a}}{a}\delta_k^i$$

$$R^i_{jkl} = \dot{a}^2(\delta_k^i\delta_{jl} - \delta_l^i\delta_{jk})$$

Canonical Tetrad

For Schwarzschild metric:

$$ds^2 = -(1 - 2M/r)dt^2 + \frac{dr^2}{1 - 2M/r} + r^2d\Omega^2$$

Choose:

$$e_{(0)} = \sqrt{1 - 2M/r} \partial_t$$

$$e_{(1)} = \frac{1}{\sqrt{1 - 2M/r}} \partial_r$$

$$e_{(2)} = \frac{1}{r} \partial_\theta$$

$$e_{(3)} = \frac{1}{r \sin \theta} \partial_\phi$$

Non-zero Spin Connection Components

$$\omega^0{}_1 = \frac{M}{r^2} dt$$

$$\omega^1{}_2 = -\sqrt{1 - 2M/r} d\theta$$

$$\omega^1{}_3 = -\sqrt{1 - 2M/r} \sin \theta d\phi$$

$$\omega^2{}_3 = -\cos \theta d\phi$$

Advantages of Tetrad Approach

- Separates gravitational and coordinate effects
- Reveals physical structure of spacetime
- Simplifies calculation of curvature invariants

- Extension of GR allowing for torsion $T^a \neq 0$
- Spin connection ω_μ^{ab} becomes independent dynamical field
- Field equations:

$$G^{ab} = 8\pi T^{ab}$$

$$T^a{}_{bc} = 8\pi \tau^a{}_{bc}$$

where $\tau^a{}_{bc}$ is spin density tensor

Tetrad Action

Action in terms of tetrad and spin connection:

$$S = \frac{1}{16\pi} \int e^a \wedge e^b \wedge R^{cd} \epsilon_{abcd} + S_{\text{matter}}$$

where $e = \det(e_\mu^a)$ and ϵ_{abcd} is Levi-Civita tensor

Field Equations Derivation

Varying with respect to tetrad:

$$\epsilon_{abcd} e^b \wedge R^{cd} = 16\pi \tau^{(a)}$$

where $\tau^{(a)}$ is energy-momentum 3-form

Varying with respect to spin connection:

$$\epsilon_{abcd} T^c \wedge e^d = 16\pi \sigma_{ab}$$

where σ_{ab} is spin current 3-form

- Torsion propagates spin density through spacetime
- Avoids singularities in some cases
- Reduces to GR in vacuum (no spin sources)

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