

# THE INCOMPLETENESS OF OBSERVATION

Why Quantum Mechanics and General Relativity  
Cannot Be Unified From Within

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## Abstract

The incompatibility between quantum mechanics and general relativity is widely viewed as a failing of modern physics. This paper proposes the opposite: the incompatibility is a structural consequence of embedded observation. Any observer embedded in the universe it measures must access reality through projections that discard degrees of freedom beyond spacetime’s causal boundaries.

Applying Wolpert’s (2008) physical limits of inference [4], we show that quantum and gravitational vacuum measurements are complementary projections of a shared hidden sector, reframing the  $10^{122}$  cosmological constant discrepancy as a measurement of  $\sim 10^{244}$  hidden-sector degrees of freedom. Tracing out this trans-horizon sector via the Nakajima-Zwanzig formalism yields dynamics that are provably non-Markovian, with multiple independent arguments—timescale collapse, conservation-law-enforced information backflow, sub-ohmic spectral density at strong coupling, and  $O(1)$  mode migration through the evolving horizon—arguing that the dynamics are CP-indivisible. By Barandes’ stochastic-quantum correspondence [14, 15], CP-indivisible dynamics recover the Schrödinger equation as the mandatory description of an embedded observer.

The framework yields falsifiable predictions: gravitational wave echoes ( $\sim 54$  ms for a  $30M_{\odot}$  remnant), a stochastic gravitational noise floor, and a slowly evolving effective cosmological constant in the Running Vacuum Model form  $\Lambda_{\text{eff}} = \Lambda_0 + 3\nu_{\text{OI}}(H^2 - H_0^2)/(8\pi G)$ , testable by DESI, Euclid, and the Vera Rubin Observatory.

# 1 INTRODUCTION

## 1.1 The Incompatibility Problem

Quantum mechanics and general relativity are extraordinarily successful yet structurally incompatible. Historically, physics has operated under the tacit assumption of a "God's-eye view"—that the universe can be fully described as if the observer were outside of it. However, physical observers and their measuring devices are strictly embedded subsystems within the universe they are measuring. This paper argues that the QM-GR incompatibility is the physical manifestation of the informational limits imposed on an embedded observer.

## 1.2 Embedded Inference and Complementary Projections

The sharpest consequence of embedded observation is that different physical subsystems, coupling to the same underlying degrees of freedom through different mechanisms, can produce irreconcilable measurements — not because one is wrong, but because each discards different information. We partition the total state space  $\Omega$  into degrees of freedom accessible to an observer (the visible sector) and degrees of freedom rendered inaccessible by spacetime's causal structure (the hidden sector) — primarily trans-horizon modes beyond the observable universe and the interiors of black holes.

Two physical couplings to the vacuum define two projection maps from  $\Omega$  to observables. Localized matter (e.g., an atom subject to the Lamb shift) measures the unsigned total vacuum activity through the squared field operator:

$$\pi_Q : \Omega \rightarrow \mathbb{R}, \quad \pi_Q(\omega) = \langle \hat{\phi}^2 \rangle_\omega \quad (1)$$

Spacetime geometry measures the net signed stress-energy density sourcing Einstein's equations:

$$\pi_G : \Omega \rightarrow \mathbb{R}, \quad \pi_G(\omega) = \langle T_{00} \rangle_\omega \quad (2)$$

Both maps are many-to-one: distinct full-universe configurations that differ only in the hidden sector yield identical outputs. But their kernels are non-identical —  $\pi_Q$  discards sign information while summing magnitudes, whereas  $\pi_G$  retains signs while summing algebraically. This is the origin of the discrepancy.

Wolpert's (2008) limits of inference [4] provide the formal no-go theorem confirming that this discrepancy is structural. Wolpert proved that any physical subsystem whose

state is a deterministic, many-to-one function of the total configuration constitutes an inference device subject to absolute limits: two such devices with non-identical surjective projections cannot fully reconstruct each other's conclusions from their own outputs alone (the Mutual Inference Impossibility). Both  $\pi_Q$  and  $\pi_G$  satisfy Wolpert's conditions — each is an embedded proper subsystem of  $\Omega$  with a surjective observation map — and their non-identical kernels trigger the impossibility result. No single embedded subsystem can simultaneously access both the unsigned total and the signed net of the vacuum. The  $10^{122}$  discrepancy between them is therefore a structural limit on embedded observation, not an error awaiting correction.

## 2 OBSERVATIONAL INCOMPLETENESS

### 2.1 Quantifying the Discrepancy

Having established that the  $10^{122}$  cosmological constant discrepancy [1, 2, 3] reflects a structural limit rather than a fine-tuning error (§ 1.2), we now extract its quantitative content. A thermodynamic analogy clarifies the mechanism. When measuring a glass of water, a thermometer records the total unsigned kinetic energy of all molecules (a variance-type measurement), producing a massive reading because nothing cancels out. Conversely, a suspended dust speck is only moved by the net canceled-out mechanical push of molecules striking it from all sides (a mean-type measurement), leaving a tiny statistical residual.

- **The Variance Target (Localized Matter):** Localized matter measuring the vacuum (e.g., via the Lamb shift) acts as the thermometer. Because the field operator is squared ( $\hat{\phi}^2$ ), its expectation value is strictly positive-definite for all modes. This measures the unsigned total activity:  $V \propto \sum_{i=1}^N \langle \hat{\phi}_i^2 \rangle = N \langle E \rangle$ .
- **The Mean Target (Spacetime):** The local gravitational field acts as the dust speck. It couples dynamically to the net signed sum of the stress-energy tensor. This is a mean-type projection:  $M = \langle T_{00} \rangle = \sum_{i=1}^N s_i |E_i|$ , where  $s_i \in \{+1, -1\}$  is the spin-statistics sign of the  $i$ -th mode.

Near a causal horizon, the Unruh temperature [18]  $T_U = \frac{\hbar a}{2\pi k_B c}$  diverges as proper acceleration  $a \rightarrow \infty$ . In this ultra-relativistic, infinite-temperature limit ( $T_U \gg m_i$ ), the active trans-horizon modes behave as a randomized, maximally mixed conformal fluid. Consequently, the spin-statistics sign  $s_i$  over the  $N$  modes functions as an independent, identically distributed random variable with a mean  $\langle s_i \rangle = 0$ .

By the Central Limit Theorem, the expected value of the gravitational sum is zero, but its root-mean-square fluctuation is strictly non-zero and scales with the square root of

the number of states:

$$M_{\text{rms}} = \sqrt{\sum_{i=1}^N (s_i |E_i|)^2} = \sqrt{N} \langle |E_i| \rangle \quad (3)$$

Therefore, the ratio of the variance projection to the mean-squared projection strictly scales as the square root of the hidden-sector dimensionality:

$$\frac{V}{M_{\text{rms}}} \approx \frac{N \langle |E_i| \rangle}{\sqrt{N} \langle |E_i| \rangle} = \sqrt{N} \quad (4)$$

**Observational Incompleteness Principle.** For any many-to-one partition of the universe into visible and hidden sectors, the variance-type projection  $V = \sum |E_i|$  and the mean-type projection  $M = \sum s_i |E_i|$  satisfy  $V/M_{\text{rms}} = \sqrt{N}$ , where  $N$  is the hidden-sector dimensionality. The discrepancy between them is a structural consequence of embedded observation, not a fine-tuning problem.

**Clarification: independence of spin-statistics signs under trans-horizon entanglement.** The TFD entanglement structure invoked in § 3.2 entangles *amplitude* (occupation number) degrees of freedom between matched visible and trans-horizon mode pairs. The spin-statistics sign  $s_i$ , by contrast, is not a dynamical variable but a fixed topological property of each field species:  $s_i = +1$  for bosonic and  $s_i = -1$  for fermionic contributions at each energy scale. In the ultra-relativistic regime ( $T_U \gg m_i$ ), all Standard Model species contribute at comparable magnitudes, so the net sign at each energy bin is determined by species composition with no systematic bias. The CLT therefore applies to the *species-sign sector*, while entanglement resides in the *amplitude sector*; the two structures coexist without contradiction.

The iid assumption on the signs  $s_i$  can be relaxed. Even under residual correlations between nearby energy scales, the  $\sqrt{N}$  scaling is preserved provided the correlations decay sufficiently fast. Specifically, for a fast-scrambling horizon bath (§ 3.2), correlations decay exponentially on the scrambling timescale  $t_s \sim \ln S$ , satisfying the strong mixing ( $\alpha$ -mixing) condition:  $\alpha(k) \rightarrow 0$  as the separation  $k$  between energy bins grows. Bradley’s central limit theorem for  $\alpha$ -mixing sequences [40] guarantees that  $M/\sqrt{N} \rightarrow \mathcal{N}(0, \sigma^2)$  with  $\sigma^2 = \text{Var}(s_i) + 2 \sum_{k=1}^{\infty} \text{Cov}(s_0, s_k)$ , where the covariance sum converges due to exponential decay. The  $\sqrt{N}$  scaling is therefore robust to short-range correlations; only long-range power-law correlations with exponent  $\alpha < 1/2$  could alter the scaling to  $N^\alpha$ , which would change the predicted discrepancy ratio from  $10^{122}$  to  $10^{244\alpha}$ . Since the observed discrepancy provides a single data point, such a modification would constrain  $\alpha$

rather than falsify the framework—but would require a physical mechanism for long-range sign correlations across widely separated energy scales, which the scrambling dynamics disfavor.

## 2.2 Deriving the $10^{122}$ Discrepancy

The ratio of the two projections directly encodes the hidden sector's dimensionality. Setting this equal to the observed  $10^{122}$  discrepancy yields  $N \sim 10^{244}$  hidden degrees of freedom.

This specific number corroborates holographic principles [6]. The Bekenstein-Hawking entropy of the cosmological horizon [10] is  $S_{\text{dS}} \sim 10^{122}$ . Because every visible degree of freedom can correlate with every hidden degree of freedom, the total number of correlations is the product of their state spaces ( $S_{\text{visible}} \times S_{\text{hidden}}$ ). This saturated, all-to-all correlation network yields an effective trace-out dimensionality equal to the square of the boundary capacity:  $N = S_{\text{dS}}^2 \sim 10^{244}$ . The cosmological constant problem is therefore not an error; it is the most precise measurement of the dimensionality of the inaccessible universe.

# 3 THE EMERGENCE OF QUANTUM MECHANICS

## 3.1 Classical Axioms and the Trace-Out Operation

Having established the macroscopic scale of the hidden sector, we now ask how an embedded, localized particle behaves when subjected to this immense informational deficit. This framework derives quantum mechanics purely from three classical premises:

1. **Classical Continuous Dynamics:** The total universe evolves deterministically via the continuous Liouville equation:  $\frac{\partial \rho}{\partial t} = \{H, \rho\}$ .
2. **Classical General Relativity:** Einstein's field equations define the absolute information barriers of the hidden sector.
3. **Classical Probability Theory:** Observational predictions are classical expectation values.

Because the universe is fundamentally continuous and deterministic at the global level, an embedded observer's inability to track the  $10^{244}$  hidden states forces a massive mathematical data compression. The observer must "trace out" the hidden sector to predict the marginal dynamics of localized matter. Section 3.2 demonstrates that this trace-out is provably non-Markovian for IR horizon modes and gives strong physical arguments

that the dynamics are in fact CP-indivisible—the key condition required for quantum mechanics to emerge from Barandes’ stochastic-quantum correspondence.

This program shares a broad structural similarity with ’t Hooft’s deterministic quantum mechanics [46], which also derives quantum behavior from an underlying deterministic dynamics through information loss. The key difference is the mechanism: ’t Hooft posits information loss at the Planck scale through equivalence classes of cellular automaton states, whereas the present framework locates it at the cosmological horizon through the trace-out of causally inaccessible degrees of freedom. The two approaches make different empirical predictions—’t Hooft’s requires discreteness at the Planck scale, while the present framework predicts gravitational wave echoes and running vacuum evolution—and could in principle both be tested.

### 3.2 From Trace-Out to Quantum Dynamics

Tracing out the hidden sector via the Nakajima-Zwanzig formalism [12, 13, 20] yields a Generalized Langevin Equation with a non-local memory kernel  $\gamma(t - \tau)$ , where the  $10^{244}$  trans-horizon states act as a continuous background noise  $\xi(t)$ :

$$m\ddot{x}(t) = -\nabla V(x) - \int_0^t \gamma(t - \tau)\dot{x}(\tau)d\tau + \xi(t) \quad (5)$$

A standard objection is that tracing out a vast bath should immediately lead to classical decoherence—incoherent, memoryless noise that destroys quantum behavior. For a generic thermal bath, this is correct: the Markovian (memoryless) approximation holds whenever the bath correlation time is much shorter than the system’s dynamical timescale, and the resulting dynamics are CP-divisible—each time step is an independent, information-destroying channel. We now show that the cosmological horizon trace-out fails every prerequisite for this conclusion, through five independent and progressively stronger arguments.

**Argument 1: Timescale collapse.** The derivation of any Markovian master equation—Lindblad, Redfield, or Davies—requires a strict separation of three timescales [20]: the bath correlation time  $\tau_B$  must be much shorter than the system’s characteristic evolution timescale  $\tau_S$ , which must itself be much shorter than the relaxation time  $\tau_R$ :

$$\tau_B \ll \tau_S \ll \tau_R \quad (6)$$

For a localized qubit with transition frequency  $\omega$  coupled to a thermal bath, this hierarchy is readily satisfied:  $\tau_B \sim 1/T$ ,  $\tau_S \sim 1/\omega$ ,  $\tau_R \sim 1/(\lambda^2\omega)$ , and the ratios are controlled by

the small coupling  $\lambda$ . For the cosmological horizon, all three timescales collapse to the same order:

-  $\tau_B \sim 1/H$  — the bath correlation time for IR horizon modes, set by the Gibbons-Hawking temperature  $T_{\text{GH}} = H/2\pi$  -  $\tau_S \sim 1/H$  — the dynamical timescale of cosmological observables (the Hubble time) -  $\tau_R \sim 1/H$  — the de Sitter equilibration time

When  $\tau_B \approx \tau_S \approx \tau_R$ , the Markovian approximation fails by its own defining criterion. This is not a subtle argument about the structure of the bath state—it is a direct, quantitative failure of the most basic prerequisite. The dynamics of the visible sector at cosmological scales are provably non-Markovian for the IR modes that dominate the vacuum energy.

**Argument 2: Conservation-law-enforced correlations.** CP-divisible dynamics satisfies a monotonicity condition: the trace distance  $D(\rho_1(t), \rho_2(t))$  between any two system states can only decrease or remain constant over time. Physically, the system can only lose distinguishability to the environment, never recover it—information flows out but not back. However, exact conservation of stress-energy across the cosmological horizon forces specific correlations between the visible and hidden sectors to persist indefinitely. If two visible-sector configurations differ in total energy by  $\Delta E$ , the corresponding hidden-sector configurations must differ by  $-\Delta E$ . This correlation is enforced by the constraint equations of general relativity and cannot decay under any dynamics that respects diffeomorphism invariance. In a weakly coupled system, the conserved-quantity correlations are perturbatively small ( $\sim \lambda^2$ ) and compatible with approximately CP-divisible dynamics. But for the cosmological trace-out, the energy correlation between visible and hidden sectors is the dominant feature—it is what produces the  $10^{122}$  discrepancy. The persistent mutual information between sectors is not a small correction; it is the largest term in the problem. These conservation-law-enforced correlations periodically drive information backflow from hidden to visible sector as modes cross and recross the slowly evolving horizon, violating the monotonicity that CP-divisibility requires (Babu et al. [21]).

**Argument 3: No perturbative route to CP-divisibility.** Every existing proof that dynamics near a horizon can be rendered Markovian—the Davies-Lindblad weak-coupling limit, Merkli’s error bounds [41], the Kaplanek-Burgess open EFT resummation [37, 38]—explicitly requires that the system-bath coupling constant  $\lambda$  be small, with errors of order  $O(|\lambda|^{1/4})$  or  $O(g^2)$ . The gravitational coupling of the entire visible sector to the trans-horizon sector is not perturbatively small. It is what produces the  $10^{122}$  projection discrepancy of § 2.2. No perturbative Markovian result applies at this coupling strength.

This limitation is not merely practical but structural. The de Sitter Wightman function for a conformally coupled massless scalar,  $G^+(\tau) = H^2/(16\pi^2 \sinh^2(H\tau/2))$ , is strictly

positive and monotonically decreasing for all  $\tau > 0$ . At second order in the coupling  $g$ , the decoherence rates  $\gamma_k(t) = g^2 \int_0^t ds \operatorname{Re}[G^+(s)e^{-i\omega s}]$  are therefore positive for *all* system frequencies, including  $\omega \sim H$ —because a linear functional of a positive function is positive. Perturbation theory does not merely fail to prove CP-indivisibility; it actively yields CP-divisible dynamics regardless of the frequency regime. This confirms that CP-indivisibility, if it occurs, is a genuinely non-perturbative phenomenon that cannot be accessed by any finite-order expansion.

**Argument 4: Sub-ohmic spectral density at strong coupling.** The effective spectral density of the de Sitter bath, including thermal occupation at the Gibbons-Hawking temperature  $T_{\text{GH}} = H/2\pi$ , is  $J_{\text{eff}}(\omega) = J(\omega) \coth(\pi\omega/H)$ . For an ohmic bare density  $J(\omega) = \eta\omega$ , this gives  $J_{\text{eff}}(\omega) \rightarrow \eta H/\pi$  (constant) for  $\omega \ll H$  and  $J_{\text{eff}}(\omega) \rightarrow \eta\omega$  (linear) for  $\omega \gg H$ . The IR-relevant regime is therefore sub-ohmic with spectral exponent  $s = 0$ . This classification connects the cosmological trace-out to a large body of established results in the spin-boson literature. Non-perturbative calculations—tensor network simulations [42] and numerical renormalization group methods [43]—have demonstrated that sub-ohmic baths ( $s < 1$ ) at strong coupling produce CP-indivisible dynamics with decoherence rates that oscillate and go negative, exhibiting information backflow absent in the perturbative regime. The de Sitter bath falls squarely in this parameter class:  $s = 0$  at the relevant IR frequencies, with non-perturbative coupling. While these numerical results do not constitute a proof for the de Sitter case specifically, they establish that CP-indivisibility is the generic behavior in the relevant parameter regime.

**Argument 5: Information backflow from the evolving partition.** The preceding arguments analyze the trace-out dynamics for a fixed system-environment partition. But the cosmological horizon evolves in time, and modes continuously migrate between the visible and hidden sectors: the event horizon contracts in comoving coordinates (removing visible modes) while the particle horizon grows (adding new ones). This migration provides a structural mechanism for information backflow independent of the spectral density.

Consider a mode  $A_1$  with frequency  $\omega \sim H$  that is in the visible sector at time  $t_1$ . At time  $t_2 > t_1$ , this mode has crossed the event horizon and joined the hidden sector, carrying its vacuum correlations with the remaining visible modes  $\{A_2, A_3, \dots\}$ . Through the scrambling dynamics of the hidden sector, information about  $A_1$ 's state is distributed across all hidden-sector modes, including those causally connected to the visible sector. The stress-energy conservation constraint (Argument 2) forces this redistributed information to influence the visible sector's subsequent evolution. Information that was accessible at  $t_1$ , became inaccessible at  $t_2$ , and then influenced the visible sector again through the conservation-law channel. This is the definition of information backflow.



The rate of mode migration is set by  $\dot{H}/H^2 = -(1+q) \approx -0.45$  in the current epoch, where  $q \approx -0.55$  is the deceleration parameter. Roughly 45% of horizon-scale modes migrate per Hubble time—this is an  $O(1)$  effect, not a perturbative correction. The resulting information backflow drives the trace distance between visible-sector states to increase over intervals where partition changes dominate, violating the monotonicity required by CP-divisibility.

A formal subtlety must be acknowledged: the standard definition of CP-indivisibility assumes a fixed system-environment partition. Extending it to time-dependent partitions—which is the physically appropriate framework for cosmological dynamics—has not been formalized in the open quantum systems literature. This extension is itself a well-posed mathematical problem and a natural target for future work.

**Additional structural constraints.** Three further properties of the trans-horizon bath reinforce and sharpen the above arguments:

- *Maximal entanglement.* The vacuum state across a causal horizon takes the Thermofield Double (TFD) form. As established by Buscemi [23], initially entangled system-environment states violate the product-state assumption required for complete positivity. For almost any entangled initial state, the set of joint unitaries generating CP-divisible reduced dynamics is measure-zero in the space of all unitaries.
- *Spectral rigidity from fast scrambling.* Causal horizons scramble information exponentially fast (Sekino and Susskind [19]). The resulting spectral rigidity—manifesting as a non-decaying ramp in the Spectral Form Factor (Cotler et al. [22])—prevents bath correlation functions from decaying to zero. The non-decaying component is exponentially small in the entropy ( $\sim e^{-S}$ ), but persists to the Heisenberg time  $t_H \sim e^S$ , ensuring that the memory kernel never fully decays.
- *Born approximation failure.* Macroscopic conservation laws physically prevent the environment from resetting to equilibrium after each interaction, forcing persistent system-environment correlations ( $\chi(t) \neq 0$ ) that invalidate the Born approximation independently of coupling strength (Babu et al. [21]).

**Counterexamples and their scope.** An important body of work by Kaplanek and Burgess [37, 38] demonstrates that an accelerated Unruh-DeWitt detector can evolve via a Markovian Lindblad equation at late times near a Rindler or Schwarzschild horizon. This is fully consistent with the above arguments: their results apply to a localized qubit whose transition frequency  $\omega$  satisfies  $\omega \gg H$ , precisely the regime where the timescale separation of Argument 1 holds ( $\tau_S \sim 1/\omega \ll \tau_B \sim 1/H$ ), where the coupling is perturbatively weak (Argument 3), and where the spectral density is ohmic rather than sub-ohmic (Argument 4). Kaplanek and Burgess themselves identified that as  $\omega \rightarrow H$  (or equivalently  $m/H \rightarrow 0$ ), the Markovian validity window closes due to "critical slowing down" [39]—the thermalization timescale diverges and the Markovian resummation

breaks down. Their counterexamples therefore confirm rather than contradict the present analysis: Markovian dynamics emerge when the timescale hierarchy is satisfied and the spectral density is effectively ohmic, and fail when either condition is violated.

We state the central dynamical claim:

***The Trace-Out Conjecture.*** *Tracing out the trans-horizon sector in the cosmological setting forces the marginal dynamics of the visible sector into a CP-indivisible stochastic process. Any faithful mathematical representation of these dynamics is empirically equivalent to quantum mechanics.*

The status of this claim is stronger than a bare conjecture but short of a rigorous proof. What is *established* is: (a) the dynamics are provably non-Markovian for IR horizon modes, by the timescale collapse argument; (b) conservation laws force persistent system-environment correlations that are non-perturbatively large; (c) perturbation theory is structurally incapable of detecting CP-indivisibility in this setting, since the positive-definite de Sitter noise kernel guarantees positive perturbative decoherence rates at all orders in  $g^2$ ; (d) the effective spectral density is sub-ohmic ( $s = 0$ ) at IR frequencies, placing the system in a parameter class where non-perturbative calculations generically yield CP-indivisible dynamics; (e) the evolving cosmological partition drives  $O(1)$  information backflow per Hubble time through mode migration and conservation-law-enforced correlations; and (f) every additional structural property of the bath (TFD entanglement, spectral rigidity, Born approximation failure) independently makes CP-divisibility more fragile. What remains *open* is a non-perturbative numerical calculation—feasible with tensor network or functional renormalization group methods—of the decoherence rates  $\gamma_k(t)$  for the specific de Sitter spectral density, which would either confirm or refute CP-indivisibility directly. We note that the framework’s falsifiable predictions (§ 4.7, § 4.8) do not depend on the specific mechanism producing indivisibility, only on the fact of it; if the dynamics are CP-indivisible for any reason, Barandes’ correspondence [14, 15] guarantees the framework follows.

By Barandes’ stochastic-quantum correspondence [14, 15], any indivisible stochastic process on configuration space is mathematically equivalent to a unitarily evolving quantum system. The Schrödinger equation is therefore the unique description an embedded observer must use to predict the marginal behavior of a system coupled to this bath. (A technical caveat: Barandes’ proof is formulated for finite-dimensional configuration spaces. The cosmological trace-out acts on an infinite-dimensional visible sector. The extension is expected to follow from standard limiting arguments, but has not been carried out explicitly; we flag this as a mathematical gap distinct from the physical question

of CP-indivisibility.) The indivisible stochastic process recovers the quantum potential  $Q = -\frac{\hbar^2}{2m} \frac{\nabla^2 \sqrt{\rho}}{\sqrt{\rho}}$  as the effective force on configuration-space probability flow, with diffusion coefficient  $D = \hbar/2m$ . The value of Planck’s constant is consistent with the ratio of the holographic entropy to the hidden-sector dimensionality,  $\hbar \approx S_{\text{universe}}/N \approx 10^{-122}$  in dimensionless units—a non-trivial self-consistency check linking the macroscopic causal structure to the microscopic noise floor.

Le et al. [17] proved that CP-divisible dynamics satisfies temporal Tsirelson’s bound ( $B \leq 2\sqrt{2}$ ). This provides a sharp falsification criterion: if the trace-out of this constrained hidden sector violates temporal Tsirelson’s bound, it is provably CP-indivisible and inherently quantum. Conversely, if the cosmological dynamics can be shown to *satisfy* the bound, the conjecture would be significantly weakened.

## 4 IMPLICATIONS PREDICTIONS

### 4.1 The Double-Slit Experiment as Stochastic Guidance

If the dynamics are CP-indivisible as argued in §3.2, the double-slit experiment is explained through continuous, single-universe fluid dynamics. The framework is ontologically consistent with the particle remaining localized and traversing a single slit, though this is an interpretive commitment inherited from the stochastic picture rather than a distinguishing prediction—since the framework is empirically equivalent to standard quantum mechanics by construction (via Barandes’ correspondence), it cannot predict which-path information that standard QM forbids.

What the framework does explain is the mechanism behind interference. Opening or closing a second slit alters the boundary conditions of the local visible-sector field configuration. Because the quantum potential  $Q = -\frac{\hbar^2}{2m} \frac{\nabla^2 \sqrt{\rho}}{\sqrt{\rho}}$  is determined by the equilibrium distribution of the indivisible stochastic process, it inherits sensitivity to these boundary conditions through the Barandes correspondence—not through a direct causal link between laboratory geometry and the cosmological horizon. The trans-horizon sector provides the underlying noise that makes the dynamics indivisible; the specific interference pattern is then a consequence of the local potential landscape acting on that noise, just as in the Bohmian picture.

This picture shares structural features with de Broglie–Bohm pilot-wave theory, which also posits a localized particle guided by a quantum potential. The key difference is ontological: in Bohmian mechanics the guiding wave is a fundamental, ontic entity, whereas here the quantum potential is an emergent consequence of tracing out the trans-horizon bath. The framework therefore does not inherit Bohmian mechanics’ requirement for a

preferred foliation of spacetime, since the stochastic dynamics are defined on the configuration space of the visible sector rather than on a global pilot wave.

## 4.2 Bell’s Theorem and Dynamical Indivisibility

A standard objection to any stochastic underpinning of quantum mechanics is Bell’s Theorem [11], which proves that no theory satisfying certain assumptions can reproduce quantum correlations. The precise assumption is *factorizability*: conditioning on the complete hidden variable  $\lambda$  in the common past of two measurements makes their outcomes statistically independent,

$$P(a, b|x, y, \lambda) = P(a|x, \lambda) \cdot P(b|y, \lambda) \quad (7)$$

There are three logically distinct ways to violate this condition: (i) *nonlocality*—a direct causal influence between the distant measurements, (ii) *superdeterminism*—the measurement settings  $x, y$  are themselves correlated with  $\lambda$ , removing the experimenters’ freedom of choice, and (iii) *dynamical indivisibility*—the stochastic evolution connecting preparation to outcomes is irreducibly global and cannot be decomposed into independent local substeps.

This framework takes the third route. As shown in § 3.2, the trans-horizon bath forces non-Markovian, and on strong physical grounds CP-indivisible, dynamics. In Barandes’ formalization [14, 15, 16], the transition maps of an indivisible stochastic process from preparation time  $t_0$  to measurement time  $t_2$  cannot be factored through any intermediate time  $t_1$ : the global map is not the composition of two local maps. Consequently, the joint probability  $P(a, b|x, y, \lambda)$  inherits this irreducible global structure and does not factorize—even though each measurement is performed locally and the experimenters retain full freedom in choosing their settings. Measurement independence is preserved; what fails is the assumption that the dynamical law connecting the shared past to the two distant outcomes can be decomposed into independent, memoryless local processes. No faster-than-light signaling is required. This resolution of Bell’s theorem depends on the dynamics being CP-indivisible; the timescale collapse and conservation-law arguments of § 3.2 provide strong grounds for this, but a fully non-perturbative proof remains open.

## 4.3 Wigner’s Friend and Relational Epistemology

The Wigner’s Friend paradox exposes the conceptual friction between objective physical events and subjective quantum states. By defining the wave function entirely as an algorithmic data compression tool, this framework natively resolves the paradox. Inside the isolated laboratory, the Friend performs a measurement, and a definite, objective physical event occurs—driven by the classical divergence of stochastic trajectories in the

local background noise. Stationed outside, Wigner lacks access to that local microscopic information. Because he cannot track the exact stochastic trajectory, Wigner is mathematically forced to apply the trace-out algorithm and assign a superposition to the lab. This superposition is strictly a measure of Wigner’s own epistemic deficit, not the physical state of the Friend.

#### 4.4 The Everettian Illusion and the Measure Problem

The Many-Worlds Interpretation (MWI) posits that the universal wave function is the fundamental, ontic physical reality, resolving quantum superpositions by continuously splitting the universe into deterministic branches. If the present framework is correct, MWI commits a structural category error: it mistakes a mandatory epistemic data compression algorithm for the physical universe itself.

Formulating a perfectly deterministic universal wave function implicitly requires an observer entirely outside the system. By applying Wolpert’s limits, the universal wave function is an informational impossibility for any embedded physical observer. In this framework, the total universe evolves deterministically as a single, continuous reality governed by the classical Liouville equation. What MWI interprets as newly created, parallel universes actually corresponds to the unmeasured, fluctuating configurations of the hidden sector within our single universe. The realization of an outcome is not the universe splitting, but simply the classical divergence of stochastic trajectories driven by the indivisible macroscopic background noise.

Furthermore, MWI struggles to organically derive the Born rule ( $p = |\psi|^2$ ) in a universe where all outcomes deterministically occur. In the embedded framework, this problem dissolves: because there is only a single continuous universe, probabilities are standard, classical stochastic probabilities arising directly from epistemic ignorance. Given the CP-indivisibility argued in § 3.2, the Born rule is natively recovered as the equilibrium distribution of the indivisible stochastic process generated by the trans-horizon bath.

A natural question is the relationship between this framework and decoherence theory, quantum Darwinism (Zurek [44]), and the decoherent histories approach (Griffiths, Gell-Mann and Hartle [45]). These programs explain how *classicality* emerges from quantum mechanics through environmental tracing-out—the inverse direction from the present work, which derives *quantum mechanics* from classical dynamics through the same operation. The two are complementary, not contradictory: in the present framework, the cosmological trace-out at the horizon scale generates quantum mechanics for the entire visible sector, while the standard decoherence program describes how *further* tracing-out of local environmental degrees of freedom within the visible sector recovers classical behavior for macroscopic subsystems. The hierarchy is: classical total dynamics  $\rightarrow$  quantum

visible-sector dynamics (via cosmological trace-out)  $\rightarrow$  classical macroscopic dynamics (via local decoherence).

## 4.5 Black Hole Singularities and the Information Paradox

The interior of a black hole is the paradigmatic example of the hidden sector: a region rendered strictly inaccessible by a causal horizon. In classical general relativity, extrapolating the field equations through this horizon predicts a singularity—a point of infinite curvature where the theory breaks down. The embedded observation framework reinterprets this breakdown. A singularity is what results from extending a visible-sector description past its domain of validity, into a region that belongs entirely to the hidden sector. It is a mathematical artifact of the God’s-eye-view assumption, not a physical event. Just as the  $10^{122}$  cosmological constant discrepancy is not a fine-tuning error but a measurement of hidden degrees of freedom (§ 2.2), the singularity is not a prediction of nature but a signal that the observer’s projection map  $\pi_G$  has been applied outside its domain.

This reframing is consistent with a near-universal consensus across quantum gravity programs that singularities must be resolved in any complete theory. Loop quantum gravity replaces the singularity with a quantum bounce at Planck density [24]. String theory regularizes it through non-perturbative corrections [25]. Recent work in higher-derivative gravity shows that singularity resolution can emerge from pure gravity alone, without additional matter fields [26]. The embedded observation framework provides a unifying explanation for *why* every approach resolves the singularity: all are, in different formalisms, recognizing that the classical description fails at the boundary of the hidden sector.

The black hole information paradox dissolves under the same logic. Hawking’s original argument [27] assumes that information falling past the horizon is destroyed at the singularity, leading to non-unitary evolution. But if the singularity is not physical and the interior is simply the hidden sector, then information is not destroyed—it is *inaccessible*. The exterior observer must trace out the interior degrees of freedom, producing exactly the thermal Hawking spectrum as their marginal prediction. The apparent information loss is an instance of observational incompleteness, not a violation of unitarity. The Page curve—the expected rise and fall of radiation entropy during evaporation—has recently been shown to emerge generically in open quantum systems at weak coupling [28], supporting the interpretation that it reflects the observer’s restricted access rather than a uniquely gravitational phenomenon.

## 4.6 String Theory, AdS/CFT, and the Limits of External Unification

This framework offers a structural explanation for the pattern of successes and difficulties in string theory. The AdS/CFT correspondence (Maldacena [7]) achieves an exact, mathematically rigorous unification of quantum mechanics and gravity by placing the observer on the *asymptotic boundary* of a hypothetical Anti-de Sitter space, looking inward. Because this observer is positioned on the outside looking in, they inherently possess a "God's-eye view" and are entirely immune to Wolpert's absolute limits for embedded observers. The unification is exact precisely because it is constructed from an external vantage point.

However, our actual universe is expanding (de Sitter space) and has no outer boundary. We are permanently embedded inside it, rendering the AdS/CFT exact unification structurally inapplicable to us. The embedded observation framework predicts that this is not a technical obstacle awaiting a clever de Sitter extension, but a fundamental limit: any exact unification requires access to the complete state space  $\Omega$ , which Wolpert's theorem prohibits for embedded observers.

This perspective also addresses string theory's landscape problem. String theory admits an enormous number of self-consistent vacuum solutions — estimated at  $10^{500}$  or more — with no known mechanism for selecting our universe from among them. The embedded observation framework suggests that this is not a failure of the theory but a consequence of attempting to construct a unique, complete description from within. An embedded observer, restricted to projections  $\pi_Q$  and  $\pi_G$ , cannot identify the correct vacuum from within — the resulting degeneracy is the vacuum-selection analogue of the  $10^{122}$  discrepancy.

## 4.7 Falsifiable Predictions

Because the quantum-gravity discrepancy is mapped to structural boundary limits, the global cosmological incompleteness theorem must scale down to local event horizons.

- **Gravitational Wave Echoes:** The classical event horizon is replaced by an informational boundary located a microscopic distance  $\epsilon \approx l_p$  outside the Schwarzschild radius  $r_h$ . Calculating the tortoise coordinate integral predicts a precise time delay for gravitational wave echoes:

$$\Delta t_{\text{echo}} \approx \frac{r_h}{c} \ln \left( \frac{r_h}{\epsilon} \right) \quad (8)$$

For a stellar-mass black hole remnant of  $M = 30M_\odot$ , the expected delay is approximately 54 ms. The identification  $\epsilon \approx l_p$  is physically motivated by the Planck-scale breakdown of

semiclassical geometry but is not derived from the framework’s internal logic; a different value of  $\epsilon$  would shift the predicted echo time logarithmically. To date, LIGO–Virgo–KAGRA searches for post-merger echoes have yielded null results [8], with the original claims of Abedi et al. remaining contested. A continued null result at the predicted timescale would constrain the location of the informational boundary but would not falsify the broader framework, since  $\epsilon$  is a free parameter; what *would* falsify it is the complete absence of any echo structure at all scales. Crucially, a *population* of echo detections across mergers with different remnant masses would overconstrain  $\epsilon$ . The predicted timescale  $\Delta t_{\text{echo}} \propto (GM/c^3) \ln(GM/c^2\epsilon)$  depends on mass logarithmically but on  $\epsilon$  through the same logarithm; echo detections at consistent timescales across, say, 10, 30, and 60  $M_\odot$  remnants would either confirm or rule out a Planck-scale  $\epsilon$  with no remaining freedom. A single detection has a free parameter; three or more do not.

- **Stochastic Gravitational Noise Floor:** Hidden-sector fluctuations must source a continuous stochastic background with an inverse-frequency-squared spectrum in the MHz–GHz band. Next-generation high-frequency gravitational wave detectors [9] could provide an independent test of this prediction.

## 4.8 Cosmological Evolution of Observational Incompleteness

The hidden-sector dimensionality  $N$  derived in § 2.2 is not a fixed constant. It is set by the Bekenstein-Hawking entropy of the cosmological horizon,  $S_{\text{ds}} \sim A/4l_p^2$ , which depends on the horizon area and therefore evolves with the expansion history of the universe. In an accelerating universe, this evolution has a specific and potentially observable consequence.

The cosmological event horizon—the boundary beyond which events occurring now will never be observable—has a finite comoving radius that is *shrinking* as the expansion accelerates. Each object pushed beyond this horizon permanently adds degrees of freedom to the hidden sector. In the far future, as the universe asymptotes to pure de Sitter space, essentially everything outside the gravitationally bound local group will have crossed beyond the horizon.

Because  $N = S_{\text{ds}}^2$  and the horizon entropy evolves with the Hubble parameter as  $S_{\text{ds}} \sim 1/H^2$ , the hidden-sector dimensionality inherits a specific time dependence. The ratio of the two projections (§ 2.1) becomes epoch-dependent:

$$\frac{V}{M_{\text{rms}}} = \sqrt{N(t)} = S_{\text{ds}}(t) \sim \frac{1}{H(t)^2} \quad (9)$$

An important subtlety must be addressed before extracting the observational prediction. A naive reading of this scaling might suggest  $\Lambda_{\text{eff}}(t) \propto H(t)^2$  as a standalone dark energy



model. However, as Hsu [30] demonstrated, a pure  $\rho_\Lambda \propto H^2$  scaling in a spatially flat FRW universe yields an effective equation of state  $w = 0$ : the dark energy mimics pressureless dust and cannot drive cosmic acceleration. This is because if  $\rho_\Lambda = \alpha H^2$ , the Friedmann equation  $H^2 \propto \rho_m + \rho_\Lambda$  forces  $\rho_\Lambda/\rho_{\text{total}}$  to remain constant throughout cosmic history, precluding a deceleration-to-acceleration transition.

The resolution lies in distinguishing the *present-day value* of the gravitational projection from its *evolution*. The observed cosmological constant  $\Lambda_0$  is the current epoch's gravitational residual  $M_{\text{rms}}(t_0)$ , set by today's hidden-sector dimensionality  $N(t_0)$ . This value is an empirical fact anchored to the present horizon area. What evolves is the *deviation* from this value as the horizon changes. Expanding the projection ratio around the present epoch:

$$\Lambda_{\text{eff}}(t) = \Lambda_0 + \left. \frac{d\Lambda_{\text{eff}}}{dH^2} \right|_{H_0} (H^2 - H_0^2) + \mathcal{O}((H^2 - H_0^2)^2) \quad (10)$$

The derivative is calculable from the  $\sqrt{N}$  scaling. Since  $M_{\text{rms}} \propto 1/\sqrt{N} \propto H^2$  and  $\Lambda_0 \gg \delta\Lambda$ , the leading correction takes the form:

$$\Lambda_{\text{eff}}(H) = \Lambda_0 + \frac{3\nu_{\text{OI}}}{8\pi G} (H^2 - H_0^2) \quad (11)$$

where  $\nu_{\text{OI}}$  is a dimensionless coefficient encoding the rate at which the hidden-sector dimensionality responds to changes in the Hubble parameter. This is precisely the form of the Running Vacuum Model (RVM) [31, 32], arrived at here from information-theoretic rather than renormalization-group arguments. The constant offset  $\Lambda_0$  is not imposed by hand but emerges naturally as the present-epoch gravitational residual; the  $H^2$ -dependent correction captures the slow migration of degrees of freedom across the cosmological horizon.

**Relationship to the standard RVM and coefficient comparison.** In the standard RVM, the coefficient  $\nu$  arises from QFT renormalization group running of the vacuum energy and is predicted to scale as  $\nu_{\text{RVM}} \sim \frac{1}{12\pi} \sum_i M_i^2/M_{\text{Pl}}^2$  [31], giving  $\nu_{\text{RVM}} \sim 10^{-3}$  for GUT-scale particles. In the present framework,  $\nu_{\text{OI}}$  encodes the logarithmic derivative of the hidden-sector dimensionality with respect to the Hubble parameter. A direct estimate from  $N = S_{\text{dS}}^2 \sim H^{-4}$  gives  $\nu_{\text{OI}} \sim \Lambda_0/(M_{\text{Pl}}^2 H_0^2) \sim \Omega_\Lambda \sim 0.7$ , which is far too large — it would predict order-unity deviations from  $\Lambda\text{CDM}$ , which are excluded. This discrepancy indicates that the linear expansion above is oversimplified: the full relationship between  $\Lambda_{\text{eff}}$  and  $N$  must include a suppression mechanism.

There are two natural candidates for this suppression within the framework, and we can

estimate their effects to bound the problem from both sides.

*Upper bound (naive estimate).* The naive  $\nu_{\text{OI}} \sim \Omega_\Lambda \sim 0.7$  derived above treats every change in  $H$  as producing a proportional change in  $\Lambda_{\text{eff}}$ . This is excluded by observation (DESI + Planck constrain  $\nu < 10^{-2}$ ), establishing that the full mapping from  $N$  to  $\Lambda_{\text{eff}}$  must include suppression.

*Lower bound (thin-shell estimate).* The total hidden-sector mode count  $N$  receives contributions across a vast frequency range, from  $\omega \sim H$  (horizon scale) up to  $\omega \sim M_{\text{Pl}}$  (Planck scale). As  $H$  changes by  $\delta H$  over a Hubble time, only modes in a thin shell near the horizon frequency  $\omega \sim H$  actively migrate between visible and hidden sectors; the deep UV modes at  $\omega \gg H$  are insensitive to the slowly evolving horizon. The fraction of active modes is of order  $\delta N/N \sim H/M_{\text{Pl}}$ , since the mode density scales as  $\omega^3$  and the shell width is  $\sim H$ , while the total extends to  $M_{\text{Pl}}$ . This gives  $\nu_{\text{OI}} \sim (H_0/M_{\text{Pl}})^2 \sim 10^{-122}$ , which is far too small — it would render the evolution undetectable by any foreseeable experiment.

The physically relevant coefficient therefore lies between these bounds:  $10^{-122} \ll \nu_{\text{OI}} \ll 1$ . The standard RVM derives  $\nu_{\text{RVM}} \sim M_i^2/M_{\text{Pl}}^2 \sim 10^{-3}$  through adiabatic regularization of the vacuum energy, in which the dangerous quartic divergences ( $\sim m^4$ ) cancel and only the mild quadratic ( $\sim m^2 H^2$ ) terms survive [33]. An analogous mechanism may operate here: the  $\sqrt{N}$  scaling of § 2.1 relates the *ratio* of the two projections to  $N$ , not the absolute value of  $\Lambda_{\text{eff}}$ , and the mapping from projection ratio to effective cosmological constant involves the spectral density of the trans-horizon bath, which may introduce a suppression analogous to the  $M_i^2/M_{\text{Pl}}^2$  factor. If the heaviest species contributing to the active migration shell has mass  $M_*$ , one expects  $\nu_{\text{OI}} \sim M_*^2/M_{\text{Pl}}^2$  on dimensional grounds, recovering the RVM scaling from a different derivation path. A first-principles calculation from the trans-horizon spectral density is required to determine whether this heuristic is correct; the prediction at this stage is the *functional form*  $\Lambda_{\text{eff}} = \Lambda_0 + 3\nu_{\text{OI}}(H^2 - H_0^2)/(8\pi G)$  with  $\nu_{\text{OI}} > 0$  and bounded as  $10^{-122} < \nu_{\text{OI}} < 1$ , but the precise numerical value remains an open problem that constitutes the most important calculational target for future work within this framework.

**Observational status.** Dark energy surveys — DESI, Euclid, and the Vera Rubin Observatory — are measuring the equation of state parameter  $w$  with increasing precision. DESI’s 2024 Data Release 1 [29] and 2025 Data Release 2 report evidence for evolving dark energy at  $2.8\sigma$ – $4.2\sigma$  depending on dataset combination, with the signal strengthening as data has doubled. De Cruz Pérez, Gómez-Valent, and Solà Peracaula [34] tested RVM variants against DESI DR2 + Planck PR4 + supernovae and found preference for dynamical vacuum ( $\nu > 0$ ) at  $2.7\sigma$ – $3.1\sigma$  over  $\Lambda\text{CDM}$ , with the best-fit sign consistently indicating energy flow from matter into the vacuum. If confirmed by independent

probes, this would be consistent with the framework’s prediction that the hidden-sector dimensionality is slowly growing as the cosmological horizon evolves.

**Differentiation from competing models.** The functional form  $\Lambda_{\text{eff}} = \Lambda_0 + \beta H^2$  is shared by several independent theoretical programs: the QFT-derived RVM [31], holographic dark energy with Granda-Oliveros cutoff [35], and the stringy RVM from Chern-Simons gravitational anomalies [36]. These models make identical background-cosmology predictions and cannot be distinguished by expansion-history measurements alone. However, the present framework differs from all of them in a structurally important way: it predicts gravitational wave echoes (§ 4.7) and a stochastic gravitational noise floor from the same hidden-sector physics that sources the cosmological evolution. No purely cosmological RVM or holographic model makes such predictions. The conjunction of confirmed dark energy evolution *and* detected GW echoes would uniquely favor an information-theoretic origin over renormalization-group or holographic derivations.

In the deep future, as  $H$  asymptotes to a constant (set by the dark energy density),  $N$  saturates at its maximum value and  $\Lambda_{\text{eff}}$  stabilizes. An observer in that epoch would inhabit a nearly empty de Sitter void — possessing an extraordinarily precise quantum mechanics (because the trace-out is nearly total and the compression nearly perfect) but describing almost nothing, since essentially the entire universe would have migrated into the hidden sector.

## 5 CONCLUSION

The incompatibility between quantum mechanics and general relativity is not a bug to be fixed. It is the physical analogue of Gödel incompleteness [5]—the universe demonstrating, through the  $10^{122}$  cosmological discrepancy, that observers are inside the system they are trying to describe.

This paper has developed that intuition into a concrete framework. Wolpert’s limits on embedded inference, formalized as complementary projections  $\pi_Q$  and  $\pi_G$  (§ 1.2), reinterpret the cosmological constant discrepancy as a direct measurement of  $\sim 10^{244}$  hidden degrees of freedom (§ 2.2). Tracing out this trans-horizon sector via the Nakajima-Zwanzig formalism yields dynamics that are provably non-Markovian for IR horizon modes—the timescale separation required for memoryless evolution collapses when bath, system, and relaxation timescales all equal the Hubble time (§ 3.2). Conservation-law-enforced correlations, a sub-ohmic effective spectral density at strong coupling, and  $O(1)$  mode migration through the evolving horizon provide strong physical arguments that the dynamics are further CP-indivisible; through Barandes’ stochastic-quantum correspondence, this recovers the Schrödinger equation as the mandatory description of an embedded observer’s

marginal predictions. The framework addresses Bell’s theorem through dynamical indivisibility rather than nonlocality or superdeterminism (§ 4.2), natively resolves the Wigner’s Friend paradox (§ 4.3), dissolves the Everettian measure problem (§ 4.4), reinterprets black hole singularities and the information paradox as consequences of observational incompleteness (§ 4.5), and explains why string theory’s AdS/CFT correspondence succeeds mathematically while remaining structurally inapplicable to embedded de Sitter observers (§ 4.6). Because the hidden-sector dimensionality is set by the evolving horizon area, the framework further predicts a slowly time-dependent effective cosmological constant in the Running Vacuum Model form,  $\Lambda_{\text{eff}} = \Lambda_0 + 3\nu_{\text{OI}}(H^2 - H_0^2)/(8\pi G)$ , derived from information-theoretic rather than renormalization-group arguments and testable by current dark energy surveys (§ 4.8).

Several open problems remain. The dynamics of the cosmological trace-out are provably non-Markovian, and the physical arguments for full CP-indivisibility—timescale collapse, sub-ohmic spectral density at strong coupling, conservation-law-enforced information backflow, and  $O(1)$  mode migration through the evolving horizon—are individually strong and mutually reinforcing (§ 3.2). What remains open is a non-perturbative numerical calculation of the decoherence rates  $\gamma_k(t)$  for the de Sitter spectral density, feasible with existing tensor network methods, which would confirm or refute CP-indivisibility directly. Additionally, the value of Planck’s constant is recovered as a dimensional consistency condition, but a first-principles derivation from the trans-horizon bath’s spectral density has not yet been achieved. The coefficient  $\nu_{\text{OI}}$  in the cosmological evolution prediction requires a first-principles calculation from the trans-horizon spectral density to determine whether the framework’s predicted value matches or differs from the standard RVM coefficient of  $\nu \sim 10^{-3}$  (§ 4.8). On the experimental side, the gravitational wave echo prediction (§ 4.7) is testable with current and near-future detectors; continued null results would progressively constrain  $\epsilon$ , while a population of detections across different remnant masses would overconstrain the free parameter entirely.

Recognizing the Schrödinger equation as the macroscopic shadow of  $10^{244}$  missing causal variables reframes the unification problem: the goal is not to force quantum mechanics and general relativity into a single formalism, but to understand why an embedded observer necessarily sees them as separate.

## Declaration of AI-Assisted Technologies

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## References

- [1] S. Weinberg, "The cosmological constant problem," *Rev. Mod. Phys.* **61**, 1 (1989).
- [2] J. Martin, "Everything you always wanted to know about the cosmological constant problem (but were afraid to ask)," *C. R. Phys.* **13**, 566–665 (2012).
- [3] S. M. Carroll, "The Cosmological Constant," *Living Rev. Relativ.* **4**, 1 (2001).
- [4] D. H. Wolpert, "Physical limits of inference," *Physica D* **237**, 1257–1281 (2008).
- [5] K. Gödel, "Über formal unentscheidbare Sätze der Principia Mathematica und verwandter Systeme I," *Monatsh. Math. Phys.* **38**, 173–198 (1931).
- [6] G. 't Hooft, "Dimensional Reduction in Quantum Gravity," arXiv:gr-qc/9310026 (1993).
- [7] J. Maldacena, "The Large-N Limit of Superconformal Field Theories and Supergravity," *Int. J. Theor. Phys.* **38**, 1113–1133 (1999).
- [8] J. Abedi, H. Dykaar, and N. Afshordi, "Echoes from the Abyss," *Phys. Rev. D* **96**, 082004 (2017).
- [9] A. Arvanitaki and A. A. Geraci, "Detecting High-Frequency Gravitational Waves with Optically Levitated Sensors," *Phys. Rev. Lett.* **110**, 071105 (2013).
- [10] G. W. Gibbons and S. W. Hawking, "Cosmological event horizons, thermodynamics, and particle creation," *Phys. Rev. D* **15**, 2738 (1977).
- [11] J. S. Bell, "On the Einstein Podolsky Rosen paradox," *Physics Physique Fizika* **1**, 195–200 (1964).
- [12] S. Nakajima, "On Quantum Theory of Transport Phenomena," *Prog. Theor. Phys.* **20**, 948–959 (1958).
- [13] R. Zwanzig, "Ensemble Method in the Theory of Irreversibility," *J. Chem. Phys.* **33**, 1338–1341 (1960).
- [14] J. A. Barandes, "The Stochastic-Quantum Theorem," arXiv:2309.03085 (2023).
- [15] J. A. Barandes, "The Stochastic-Quantum Correspondence," *Philosophy of Physics* **3**(1):8 (2025).
- [16] J. A. Barandes, S. Hasan, and J. Kagan, "The CHSH Game, Tsirelson's Bound, and Causal Locality," arXiv:2512.18105 (2025).

- [17] T. Le, F. A. Pollock, T. Paterek, M. Paternostro, and K. Modi, "Divisible quantum dynamics satisfies temporal Tsirelson's bound," *J. Phys. A* **50**, 055302 (2017).
- [18] W. G. Unruh, "Notes on black-hole evaporation," *Phys. Rev. D* **14**, 870 (1976).
- [19] Y. Sekino and L. Susskind, "Fast Scramblers," *JHEP* **2008**, 065 (2008).
- [20] H.-P. Breuer and F. Petruccione, *The Theory of Open Quantum Systems* (Oxford University Press, 2002).
- [21] K. Babu et al., "Unfolding system-environment correlation in open quantum systems: Revisiting master equations and the Born approximation," *Phys. Rev. Research* **6**, 013243 (2024).
- [22] J. Cotler et al., "Black Holes and Random Matrices," *JHEP* **2017**, 118 (2017).
- [23] F. Buscemi, "Complete positivity, Markovianity, and the quantum data-processing inequality, in the presence of initial system-environment correlations," *Phys. Rev. Lett.* **113**, 140502 (2014).
- [24] C. Rovelli and F. Vidotto, "Planck Stars," *Class. Quantum Grav.* **31**, 045003 (2014).
- [25] A. Wu, Y. Yan, and L. Ying, "Revisiting Schwarzschild black hole singularity through string theory," *Eur. Phys. J. C* **85**, 168 (2025).
- [26] P. Bueno, P. A. Cano, and R. A. Hennigar, "Regular Black Holes from Pure Gravity," *Phys. Lett. B* **861**, 139260 (2025).
- [27] S. W. Hawking, "Breakdown of predictability in gravitational collapse," *Phys. Rev. D* **14**, 2460 (1976).
- [28] J. Glatthard, "Page curves and typical entanglement in linear optics," *Phys. Rev. D* **109**, L081901 (2024).
- [29] DESI Collaboration, "DESI 2024 VI: Cosmological Constraints from the Measurements of Baryon Acoustic Oscillations," arXiv:2404.03002 (2024).
- [30] S. D. H. Hsu, "Entropy bounds and dark energy," *Phys. Lett. B* **594**, 13–16 (2004).
- [31] J. Solà Peracaula, "The cosmological constant problem and running vacuum in the expanding universe," *Phil. Trans. R. Soc. A* **380**, 20210182 (2022).
- [32] J. Solà Peracaula, A. Gómez-Valent, and J. de Cruz Pérez, "Running vacuum in the Universe: phenomenological status in light of the latest observations," *Universe* **9**, 262 (2023).

- [33] C. Moreno-Pulido and J. Solà Peracaula, "Renormalizing the vacuum energy in cosmological spacetime: implications for the cosmological constant problem," *Eur. Phys. J. C* **82**, 551 (2022).
- [34] J. de Cruz Pérez, A. Gómez-Valent, and J. Solà Peracaula, "Dynamical Dark Energy models in light of the latest observations," arXiv:2512.20616 (2025).
- [35] L. N. Granda and A. Oliveros, "Infrared cut-off proposal for the holographic density," *Phys. Lett. B* **669**, 275–277 (2008).
- [36] N. E. Mavromatos, S. Basilakos, and J. Solà Peracaula, "Stringy running vacuum model and current tensions in cosmology," *Class. Quantum Grav.* **41**, 015026 (2024).
- [37] G. Kaplanek and C. P. Burgess, "Hot accelerated qubits: decoherence, thermalization, secular growth and reliable late-time predictions," *JHEP* **2020**, 008 (2020).
- [38] G. Kaplanek and C. P. Burgess, "Qubits on the horizon: decoherence and thermalization near black holes," *JHEP* **2021**, 098 (2021).
- [39] G. Kaplanek and C. P. Burgess, "Hot cosmic qubits: late-time de Sitter evolution and critical slowing down," *JHEP* **2020**, 053 (2020).
- [40] R. C. Bradley, "Basic properties of strong mixing conditions. A survey and some open questions," *Prob. Surveys* **2**, 107–144 (2005).
- [41] M. Merkli, "Dynamics of Open Quantum Systems I, Oscillation and Decay," *Quantum* **6**, 615 (2022).
- [42] A. W. Chin, J. Prior, S. F. Huelga, and M. B. Plenio, "Generalized polaron ansatz for the ground state of the sub-ohmic spin-boson model: An analytic theory of the localization transition," *Phys. Rev. Lett.* **107**, 160601 (2011).
- [43] H.-P. Breuer, E.-M. Laine, and J. Piilo, "Measure for the degree of non-Markovian behavior of quantum processes in open systems," *Phys. Rev. Lett.* **103**, 210401 (2009).
- [44] W. H. Zurek, "Quantum Darwinism," *Nature Physics* **5**, 181–188 (2009).
- [45] R. B. Griffiths, "Consistent histories and the interpretation of quantum mechanics," *J. Stat. Phys.* **36**, 219–272 (1984).
- [46] G. 't Hooft, *The Cellular Automaton Interpretation of Quantum Mechanics* (Springer, 2016).