

The Asymptotics of Localizing Entanglement

Leo L.-Y. Lee, Abigail Vaughan-Lee, Hanyang Sha, Akanksha Chablani, Christopher Vairogs, and Jacob Beckey

University of Illinois Urbana-Champaign

Introduction

Localization is a protocol that involves measuring a state and discarding a subset of subsystems.

Definition 1. The ***n*-tangle** is the entanglement measurement $\tau: (\mathbb{C}^2)^{\otimes n} \rightarrow [0, 1]$ given by $\tau_n(|\psi\rangle) := |\langle\psi|\tilde{\psi}\rangle|$ where $|\tilde{\psi}\rangle := \sigma_y^{\otimes n}|\psi^*\rangle$.

Definition 2. Suppose a von Neumann measurement Π_A on subsystem A of a state $|\Psi\rangle \in \mathcal{H}_A \otimes \mathcal{H}_B$ produces an ensemble $\{(p_i, |\psi\rangle_i)\}_i$ of states over subsystem B . The **multipartite entanglement assistance (MEA)** of $|\Psi\rangle$ with respect to τ is defined as the maximal average post-measurement entanglement:

$$L_{\text{global}}^\tau(|\Psi\rangle) := \max_{\Pi_A} \sum_i p_i \tau_{N_B}(|\psi_i\rangle).$$

In particular, if Π_A is *local* — i.e., the operators assume a tensor product form with respect to the subsystems in A — then such L_{global}^τ is alternatively called the **localizable multipartite entanglement (LME)** and denoted by L^τ .

Theorem 3. For a known $C > 0$, we have

$$\Pr_{|\Psi\rangle \sim \text{Haar}} \left(L_{\text{global}}^\tau(|\Psi\rangle) \leq 1 - \sqrt{2d_B/d_A} - \varepsilon \right) \leq 2 \exp(-Cd_A d_B \varepsilon^2),$$

which shows that when $d_A \gg d_B \gg 1$, the quantity L_{global}^τ is near maximal [1, Thm. 8].

We aim to obtain similar concentration results for L^τ and study its asymptotic behavior. Preliminary numerical estimation (see Fig. 1) suggests our hypothesis as follows.

Conjecture 4. For some function K of d_A and d_B , we have

$$\Pr_{|\Psi\rangle \sim \text{Haar}} (L^\tau(|\Psi\rangle) \geq K(d_A, d_B) + \varepsilon) \leq 2 \exp(-Cd_A d_B \varepsilon^2)$$

with $\lim_{d_B \rightarrow \infty} K(d_A, d_B) = 0$.

The above holds when $K = \mathbb{E}_{|\Psi\rangle \sim \text{Haar}} [L^\tau(|\Psi\rangle)]$ [1]. Yet, the explicit form of K is unknown.

Notation	Definition
\mathcal{H}_A	Hilbert space of subsystem A
N_A	Number of qubits in A
d_A	Dimension of \mathcal{H}_A . Equals 2^{N_A}
Haar	Haar measure
$\mathcal{C}(\mathcal{H}_{A_i})$	Collection of orthonormal bases in qubit A_i
$\mathcal{P}(\mathcal{H}_A)$	Collection of orthonormal bases that have a tensor product form with respect to the subsystems in A

Table 1: Notations used in this poster

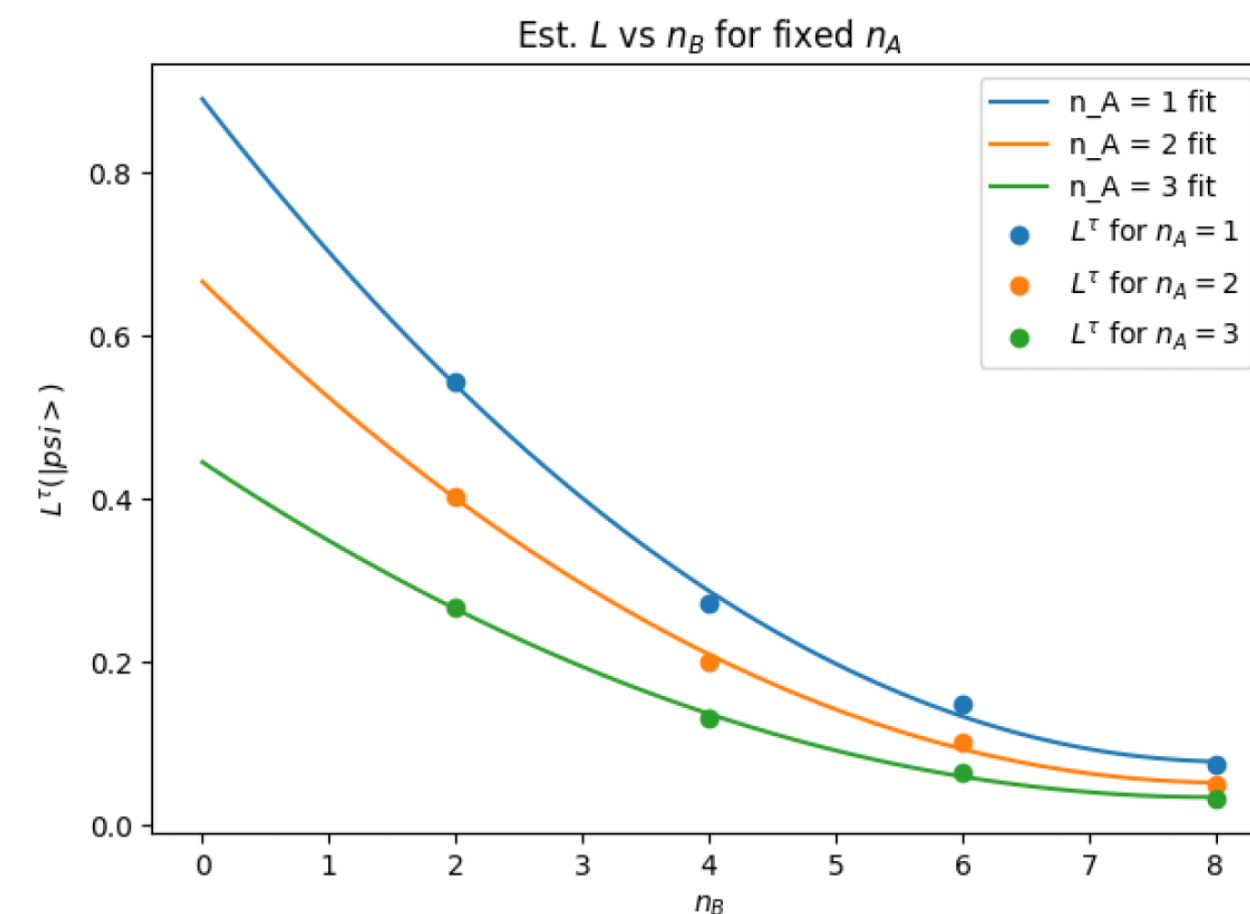


Figure 1: Exponential fitting of $\mathbb{E}_{|\Psi\rangle \sim \text{Haar}} [L^\tau(|\Psi\rangle)]$ for $N_A = 1, 2, 3$ as N_B increases

Methods

Given a basis $\beta = \{|\varphi_i\rangle\}_i \in \mathcal{P}(\mathcal{H}_A)$, the set $\{|\varphi_i\rangle\langle\varphi_i| \otimes I_B\}_i$ containing operators associated with a Π_A , and a state $|\Psi\rangle \in \mathcal{H}_A \otimes \mathcal{H}_B$, we denote by $\bar{\tau}_\beta(|\Psi\rangle)$ the average N_B -tangle. We also denote $\beta_{\text{max}} := \arg\max_\beta \bar{\tau}_\beta(|\Psi\rangle)$. Note that by Def. 2, $\bar{\tau}_{\beta_{\text{max}}}(|\Psi\rangle) = L^\tau(|\Psi\rangle)$

Recall that our goal is to bound $\Pr(L^\tau(|\Psi\rangle) \geq K(d_A, d_B) + \varepsilon)$. Lemma 22 of [1] suggests that

$$K(d_A, d_B) = \sqrt{\frac{2}{d_B + 1}}$$

is a decent choice as it is an upper bound of $\mathbb{E}[\bar{\tau}_\beta(|\Psi\rangle)]$. We can then prove the following using Levi's lemma [2, Lem. 53].

Lemma 5. Given a fixed basis $\gamma \in \mathcal{P}(\mathcal{H}_A)$, for all $\varepsilon, \delta > 0$ such that $\varepsilon - \sqrt{2}d_A d_B \delta > 0$, we have

$$\Pr_{|\Psi\rangle \sim \text{Haar}} \left(L^\tau(|\Psi\rangle) \geq \sqrt{2/(d_B + 1)} + \varepsilon \text{ and } \|\beta_{\text{max}} - \gamma\|_B \leq \delta \right) \leq 2 \exp(-2d_A d_B (\varepsilon - \sqrt{2}d_A d_B \delta)^2 / 9\pi^3 (2 + 4\sqrt{2})^2).$$

Inspired by [3, Lem. II.4], we define Defs. 6 and 7.

Definition 6 (Basis-norm). For bases $\beta = \{|\varphi_i\rangle\}, \gamma = \{|\eta_i\rangle\}$ in \mathcal{H}_A ,

$$\|\beta - \gamma\|_B := \max_i |||\varphi_i\rangle\langle\varphi_i| - |\eta_i\rangle\langle\eta_i|||_1.$$

Definition 7 (Basis ε -net). For $\varepsilon > 0$, if a set $\mathcal{N} \subseteq \mathcal{P}(\mathcal{H}_A)$ satisfies that for all $\beta \in \mathcal{P}(\mathcal{H}_A)$ there exists $\gamma \in \mathcal{N}$ such that $\|\beta - \gamma\|_B \leq \varepsilon$, then we call \mathcal{N} a **basis ε -net** on $\mathcal{P}(\mathcal{H}_A)$.

Theorem 8. There exists a basis ε -net \mathcal{N} on $\mathcal{P}(\mathcal{H}_A)$ with

$$|\mathcal{N}| \leq [5(1 + 2\sqrt{2})^2 N_A^2 / \varepsilon^2]^{8N_A}.$$

The above theorem allows us to find a set of bases such that β_{max} is close to at least one of them. Combining the result with Lem. 5, we finally arrive at the Thm. 9.

Theorem 9. For $\varepsilon > 0$,

$$\Pr_{|\Psi\rangle \sim \text{Haar}} \left(L^\tau(|\Psi\rangle) \geq \sqrt{2/(d_B + 1)} + \varepsilon \right) \leq 2 \left(\frac{40(1 + 2\sqrt{2})^2 N_A^2 d_A^2 d_B^2}{\varepsilon^2} \right)^{8N_A} \exp\left(-\frac{d_A d_B \varepsilon^2}{18\pi^3 (2 + 4\sqrt{2})^2}\right).$$

Results

Observe the inequality displayed in Thm. 9: Suppose ε is fixed. We can notice that for a fixed d_A , the right-hand side can be written as

$$c_1 d_B^{16N_A} \exp(-c_2 d_B)$$

for some known positive reals c_1 and c_2 , which obviously tends to 0 as d_B approaches infinity. Likewise, we can discover that for a fixed d_B , the right-hand side vanishes as d_A approaches infinity. In conclusion, we may state the following.

Corollary 10. Let $\varepsilon, \delta > 0$ be arbitrary. Then for any $d_B \geq 2$, there exists an $N_0 \in \mathbb{N}$ such that for all $d_A \geq 2^{N_0}$, we have

$$\Pr_{|\Psi\rangle \sim \text{Haar}} \left(L^\tau(|\Psi\rangle) \geq \sqrt{2/(d_B + 1)} + \varepsilon \right) \leq \delta.$$

Likewise, for any $d_A \geq 2$, there exists an $N'_0 \in \mathbb{N}$ such that for all $d_B \geq 2^{N'_0}$, the above bound holds.

Discussion

Theorem 3 and Cor. 10 gives the following table.

	d_A fixed, $d_B \rightarrow \infty$	$d_A \rightarrow \infty$, d_B fixed
L^τ	$< \sqrt{2/(d_B + 1)}$	$< \sqrt{2/(d_B + 1)}$
L_{global}^τ	?	$> 1 - \sqrt{2d_B/d_A}$

Table 2: Bounds for typical values of L^τ and L_{global}^τ as d_A or d_B approaches infinity

Our next and perhaps last step is to find the value around which L_{global}^τ concentrates when d_A is fixed and d_B approaches infinity.

Acknowledgements

I would like to thank Ph.D. candidate Christopher Vairogs for his mentorship, and IQIST Postdoctoral Scholar Jacob Beckey for advising this project. Also, I would like to thank Illinois Mathematics Lab for supporting this project in the 2025 Spring semester.

References

- [1] C. Vairogs, S. Hermes, and F. Leditzky. *Localizing multipartite entanglement with local and global measurements*. 2024.
- [2] A. A. Mele. Introduction to Haar Measure Tools in Quantum Information: A Beginner's Tutorial. *Quantum*, **8**: (May 2024), p. 1340.
- [3] P. Hayden, D. Leung, P. W. Shor, and A. Winter. Randomizing Quantum States: Constructions and Applications. *Communications in Mathematical Physics*, **250**:2 (July 2004), pp. 371–391.