Introduction to Simulation - Lecture 16

Methods for Computing Periodic Steady-State - Part II

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Thanks to Deepak Ramaswamy, Michal Rewienski, and Karen Veroy

Outline

- Three Methods so far
 - Time integration until steady-state achieved
 - Finite difference methods
 - Shooting Methods
- Shooting Methods
 - State transition function
 - Sensitivity matrix
 - Matrix-Free Approach
- Spectral Methods
 - Galerkin and Collocation Methods

Periodic Steady-State Basics

$$\frac{dx(t)}{dt} = F\left(\underbrace{x(t)}_{state}\right) + \underbrace{u(t)}_{input}$$

Suppose the system has a periodic input



• Many Systems eventually respond periodically

$$x(t+T) = x(t)$$
 for $t >> 0$

Periodic Steady-State Basics

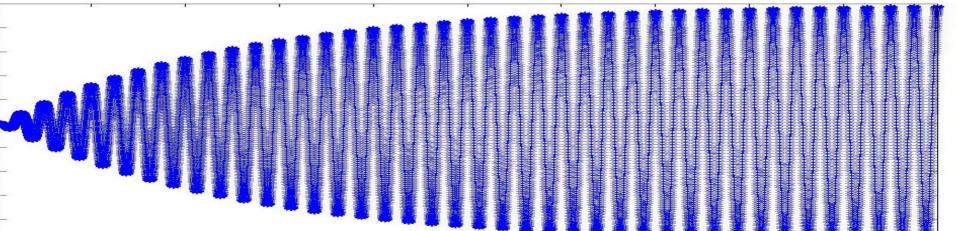
Computing Steady State

Time Integration Method

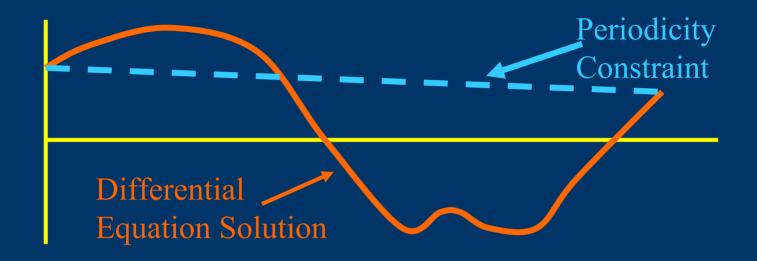
• Time-Integrate Until Steady-State Achieved

$$\frac{dx(t)}{dt} = F(x(t)) + u(t) \Rightarrow \hat{x}^{l} = \hat{x}^{l-1} + \Delta t \left(F(\hat{x}^{l}) + u(l\Delta t) \right)$$

Need many timepoints for lightly damped case!



Basic Formulation



N Differential Equations: $\frac{d}{dt}x_i(t) = F_i(x(t))$

N Periodicity Constraints: $x_i(T) = x_i(0)$

Finite Difference Methods

Nonlinear Problem

$$\frac{dx(t)}{dt} = F(x(t)) + \underbrace{u(t)}_{input} \quad t \in [0, T] \quad \underbrace{x(T) = x(t)}_{periodicity}$$

Discretize with Backward-Euler

$$H_{FD}\left(\begin{bmatrix} \hat{x}^1 \\ \hat{x}^2 \\ \vdots \\ \hat{x}^L \end{bmatrix}\right) = \begin{array}{c} \hat{x}^1 - \hat{x}^L - \Delta t \left(F\left(\hat{x}^1\right) + u\left(\Delta t\right)\right) \\ \hat{x}^2 - \hat{x}^1 - \Delta t \left(F\left(\hat{x}^2\right) + u\left(2\Delta t\right)\right) \\ \vdots \\ \hat{x}^L - \hat{x}^{L-1} - \Delta t \left(F\left(\hat{x}^L\right) + u\left(L\Delta t\right)\right) \end{array} = \mathbf{0}$$

Solve Using Newton's Method

Shooting Method

Basic Definitions

Start with
$$\frac{dx(t)}{dt} = F(x(t)) + u(t)$$

And assume x(t) is unique given x(0).

D.E. defines a State-Transition Function

$$\Phi(y,t_0,t_1) \equiv x(t_1)$$

where x(t) is the D.E. solution given $x(t_0) = y$

Shooting Method

Abstract Formulation

Solve
$$H(x(0)) = \Phi(x(0), 0, T) - x(0) = 0$$

$$x(T)$$

Use Newton's method

$$J_{H}(x) = \frac{\partial \Phi(x, 0, T)}{\partial x} - I$$

$$J_{H}\left(x^{k}\right)\left(x^{k+1}-x^{k}\right) = -H\left(x^{k}\right)$$

Shooting Method

Computing Newton

To Compute
$$\Phi(x(0), 0, T)$$
Integrate $\frac{dx(t)}{dt} = F(x(t)) + u(t)$ on [0,T]
What is $\frac{\partial \Phi(x, 0, T)}{\partial x}$? $x^{\varepsilon}(T)$
 $x(0) + \varepsilon$
 $x(0)$

Indicates the sensitivity of x(T) to changes in x(0)

Shooting Method

Sensitivity Matrix by Perturbation

$$\frac{\partial \Phi\left(x,0,T\right)}{\partial x} \approx$$

Shooting Method

Efficient Sensitivity Evaluation

Differentiate the first step of Backward-Euler

$$\frac{\partial}{\partial x(0)} \left(\hat{x}^{1} - x(0) - \Delta t \left(F(\hat{x}^{1}) + u(\Delta t) \right) = 0 \right)$$

$$\Rightarrow \frac{\partial \hat{x}^{1}}{\partial x(0)} - \frac{\partial x(0)}{\partial x(0)} - \Delta t \frac{\partial F(\hat{x}^{1})}{\partial x} \frac{\partial \hat{x}^{1}}{\partial x(0)} = 0$$

$$\Rightarrow \left(I - \Delta t \frac{\partial F(\hat{x}^{1})}{\partial x} \right) \frac{\partial \hat{x}^{1}}{\partial x(0)} = \frac{\partial x(0)}{\partial x(0)} I$$

Shooting Method

Efficient Sensitivity Matrix Cont

Applying the same trick on the 1-th step

$$\Rightarrow \left(I - \Delta t \frac{\partial F(\hat{x}^l)}{\partial x}\right) \frac{\partial \hat{x}^l}{\partial x(0)} = \frac{\partial \hat{x}^{l-1}}{\partial x(0)}$$

$$\frac{\partial \Phi(x,0,T)}{\partial x} \approx \prod_{l=1}^{L} \left(I - \Delta t \frac{\partial F(\hat{x}^{l})}{\partial x} \right)^{-1}$$

Shooting Method

Observations on Sensitivity Matrix

Newton at each timestep uses same matrices

$$\frac{\partial \Phi(x,0,T)}{\partial x} \approx \prod_{l=1}^{L} \left[I - \Delta t \frac{\partial F(\hat{x}^l)}{\partial x} \right]^{-1}$$

Timestep Newton Jacobian

Formula simplifies in the linear case

$$\frac{\partial \Phi(x,0,T)}{\partial x} \approx \left(I - \Delta tA\right)^{-L}$$

Matrix-Free Approach

Basic Setup

Start with
$$\frac{dx(t)}{dt} = F(x(t)) + u(t)$$

$$H(x(0)) = \Phi(x(0), 0, T) - x(0) = 0$$

Use Newton's method

$$J_{H}(x) = \frac{\partial \Phi(x, 0, T)}{\partial x} - I$$

$$J_{H}(x^{k})(x^{k+1} - x^{k}) = -H(x^{k})$$

Matrix-Free Approach

Matrix-Vector Product

Solve Newton equation with Krylov-subspace method

$$\underbrace{\left(\frac{\partial \Phi\left(x^{k},0,T\right)}{\partial x}-I\right)}_{A}\underbrace{\left(x^{k+1}-x^{k}\right)}_{X}=\underbrace{x^{k}-\Phi\left(x^{k},0,T\right)}_{b}$$

Matrix-Vector Product Computation

$$\left(\frac{\partial \Phi\left(x^{k},0,T\right)}{\partial x} - I\right) p^{j} \approx \frac{\Phi\left(x^{k} + \varepsilon p^{j},0,T\right) - \Phi\left(x^{k},0,T\right)}{\varepsilon} - p^{j}$$

Krylov method search direction

Matrix-Free Approach

Convergence for GCR

Example

$$\frac{dx}{dt} - Ax = 0 \quad eig(A) \text{ real and negative}$$

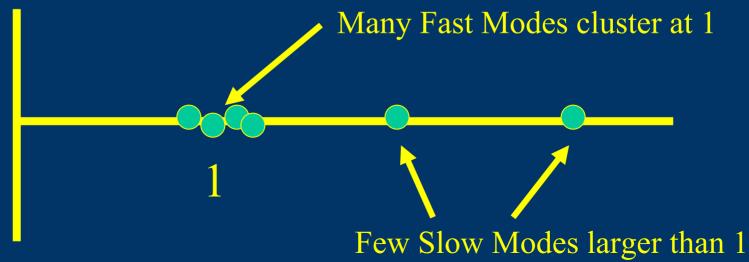
Shooting-Newton Jacobian

$$\frac{\partial \Phi(x,0,T)}{\partial x} - I = e^{AT} - I$$

Matrix-Free Approach

Convergence for GCR-evals

$$e^{AT} - I = S \begin{bmatrix} e^{\lambda_1 T} - 1 & & \\ & \ddots & \\ & e^{\lambda_N T} - 1 \end{bmatrix} S^{-1}$$



Fourier Representation

Truncation Approximation

• Periodic function \rightarrow fourier series

$$x(t) = \sum_{l=-\infty}^{\infty} X_l e^{-i2\pi l \frac{t}{T}}$$

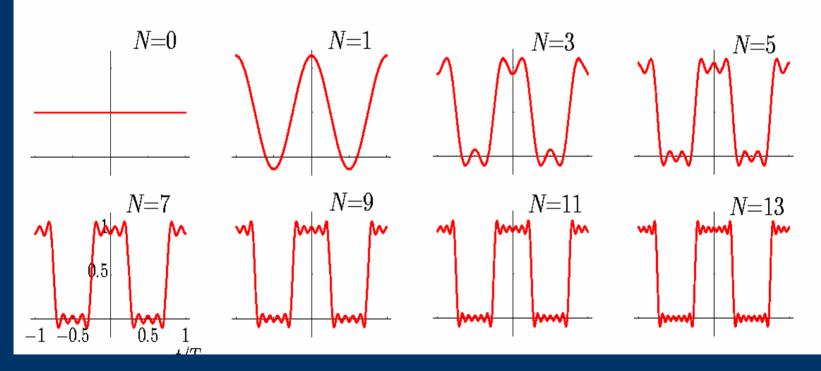
Approximate a function with truncated series

$$x(t) \approx \sum_{l=-L}^{L} X_l e^{-i2\pi l \frac{t}{T}}$$

Fourier Representation

Square Wave Example

$$x_N(t) = \frac{1}{2} + \sum_{n=1}^{N} \left(\frac{\sin(n\pi/2)}{n\pi/2} \right) \cos(2\pi nt/T).$$



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Fourier Representation

Annoyance for Real Functions

• Real x > Fourier Coeffs complex conjugate

$$X_{-l} = X_l^*$$

Can rewrite series with fewer unknowns

$$x(t) = \sum_{l=1}^{\infty} X_{l} e^{-i2\pi l \frac{t}{T}} + X_{l}^{*} e^{+i2\pi l \frac{t}{T}} + X_{0}$$
Real

Fourier Representation

Orthogonality

Terms in Fourier Series are orthogonal

$$\int_{0}^{T} e^{-i2\pi l \frac{t}{T}} e^{-i2\pi m \frac{t}{T}} dt = 0 \quad l \neq m$$

Simple formula for computing coefficients

$$\int_{0}^{T} e^{-i2\pi m \frac{t}{T}} x(t) dt = \int_{0}^{T} e^{-i2\pi m \frac{t}{T}} \sum_{l=-\infty}^{\infty} X_{l} e^{-i2\pi l \frac{t}{T}} dt = TX_{m}$$

Fourier Representation

Advantages

- For smooth functions (infinitely cont. diff)
 - Fourier Coefficients decay exponentially fast

$$\lim_{m\to\infty} \frac{1}{T} \int_{0}^{T} e^{-i2\pi m \frac{t}{T}} x(t) dt = \lim_{m\to\infty} X_m = O(c^m)$$

Automatically satisfies periodicity

$$x(t+T) = \sum_{l=-L}^{L} X_{l} e^{-i2\pi l \left(\frac{t+T}{T}\right)} = \sum_{l=-L}^{L} X_{l} e^{-i2\pi l \frac{t}{T}} = x(t)$$

Computing Coefficients

Residual

Plug representation into differential equation

$$R(\vec{X},t) = \frac{d}{dt} \left(\sum_{l=-L}^{L} X_l e^{-i2\pi l \frac{t}{T}} \right) - F\left(\sum_{l=-L}^{L} X_l e^{-i2\pi l \frac{t}{T}} \right) - u(t)$$
Residual

Simplify by differentiating representation

$$R(\vec{X},t) = \sum_{l=-L}^{L} \frac{-i2\pi l}{T} X_l e^{-i2\pi l \frac{t}{T}} - F\left(\sum_{l=-L}^{L} X_l e^{-i2\pi l \frac{t}{T}}\right) - u(t)$$
Residual

Computing Coefficients

Collocation and Galerkin

Collocation – Residual = 0 at test points

$$\underbrace{R(\vec{X}, t_l)}_{\text{Residual}} = 0 \quad l = \{1, ..., 2L + 1\}$$

Galerkin – Residual orthog to Fourier Terms

$$\int_{0}^{T} e^{-i2\pi m \frac{t}{T}} \underbrace{R(\vec{X}, t)}_{\text{Residual}} dt = 0 \quad m \in \{-L, ..., 0, ...L\}$$

Computing Coefficients

Galerkin Equation

Galerkin – Residual orthog to Fourier Terms

$$-\left(\int_{0}^{T} e^{-i2\pi m\frac{t}{T}} \left(\sum_{l=-L}^{L} \frac{-i2\pi l}{T} X_{l} e^{-i2\pi l\frac{t}{T}} - F\left(\sum_{l=-L}^{L} X_{l} e^{-i2\pi l\frac{t}{T}}\right) - u(t)\right) dt\right) = 0$$

$$i2\pi mX_{l} + \int_{0}^{T} e^{-i2\pi m\frac{t}{T}} F\left(\sum_{l=-L}^{L} X_{l} e^{-i2\pi l\frac{t}{T}}\right) dt + \int_{0}^{T} e^{-i2\pi m\frac{t}{T}} u(t) dt = 0$$

$$m \in \{-L, ..., 0, ...L\}$$

Computing Coefficients

Linear Galerkin F(x)=Ax

$$i2\pi mX_{l} + \int_{0}^{T} e^{-i2\pi m\frac{t}{T}} A\left(\sum_{l=-L}^{L} X_{l} e^{-i2\pi l\frac{t}{T}}\right) dt + \int_{0}^{T} e^{-i2\pi m\frac{t}{T}} u(t) dt = 0$$

$$U_{m}$$

$$\begin{bmatrix} \frac{i2\pi L}{T} + A & 0 & 0 & 0 \\ 0 & \frac{i2\pi(L-1)}{T} + A & 0 & 0 \\ 0 & 0 & \ddots & 0 \\ 0 & 0 & 0 & -\frac{i2\pi L}{T} + A \end{bmatrix}$$

$$\begin{bmatrix} X_{-L} \\ X_{-(L-1)} \\ \vdots \\ \vdots \\ X_L \end{bmatrix} = - \begin{bmatrix} U_{-L} \\ U_{-(L-1)} \\ \vdots \\ \vdots \\ U_L \end{bmatrix}$$

Diagonal

Computing Coefficients

Collocation Equations

Collocation – Residual zero at test times

$$R(\vec{X}, t_l) = 0 = \sum_{l=-L}^{L} \frac{-i2\pi l}{T} X_l e^{-i2\pi l \frac{t_l}{T}} - F\left(\sum_{l=-L}^{L} X_l e^{-i2\pi l \frac{t_l}{T}}\right) - u(t_l)$$
Residual

$$l = \{1, ..., 2L + 1\}$$

Computing Coefficients

Discrete Fourier Transform

$$\begin{bmatrix} e^{i2\pi \frac{L}{T}t_{1}} & \cdots & e^{-i2\pi \frac{L}{T}t_{1}} \\ \vdots & \ddots & \ddots & \vdots \\ \vdots & \ddots & \ddots & \vdots \\ \vdots & \ddots & \ddots & \vdots \\ e^{i2\pi \frac{L}{T}t_{(2L+1)}} & \cdots & e^{-i2\pi \frac{L}{T}t_{(2L+1)}} \end{bmatrix} \begin{bmatrix} X_{-L} \\ X_{-(L-1)} \\ \vdots \\ X_{L} \end{bmatrix} = \begin{bmatrix} x(t_{1}) \\ x(t_{2}) \\ \vdots \\ x(t_{(2L+1)}) \end{bmatrix}$$

Discrete Fourier Transform(DFT)

If
$$t_l = \frac{l}{2L+1}T$$
 then DFT Matrix has orthog columns

Spectral Differentiation

Computing Coefficients

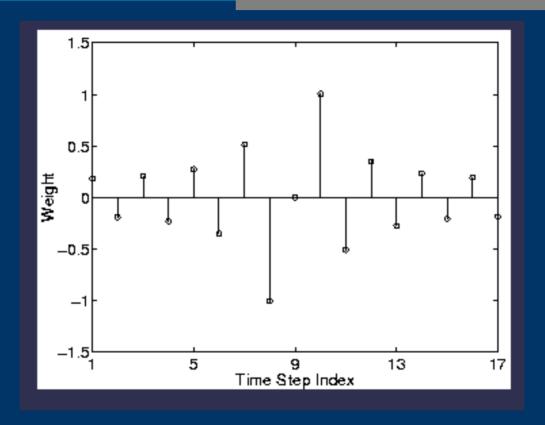
Collocation using timepoints

$$DFT \begin{bmatrix} \frac{i2\pi L}{T} & 0 & 0 & 0 \\ 0 & \frac{i2\pi(L-1)}{T} & 0 & 0 \\ 0 & 0 & \ddots & 0 \\ 0 & 0 & 0 & -\frac{i2\pi L}{T} \end{bmatrix} (DFT)^{-1} \begin{bmatrix} x(t_1) \\ x(t_2) \\ \vdots \\ x(t_{(2L+1)}) \end{bmatrix} - \begin{bmatrix} F(x(t_1)) \\ F(x(t_2)) \\ \vdots \\ F(x(t_{(2L+1)})) \end{bmatrix} = \begin{bmatrix} u(t_1) \\ u(t_2) \\ \vdots \\ u(t_{(2L+1)}) \end{bmatrix}$$

Converting timepoint into Fourier Coeffs, Differentiating, and then returning to time

Computing Coefficients

Spectral Differentiation Example



Middle row, T = 17 and 2L+1 = 17

Computing Coefficients

Spectral Colloc vs. F-D

$$\begin{bmatrix} \frac{1}{\Delta t} & 0 & 0 & -\frac{1}{\Delta t} \\ -\frac{1}{\Delta t} & \frac{1}{\Delta t} & 0 & 0 \\ 0 & \ddots & \ddots & 0 \\ 0 & 0 & -\frac{1}{\Delta t} & \frac{1}{\Delta t} \end{bmatrix} \begin{bmatrix} \hat{x}^1 \\ \hat{x}^2 \\ \vdots \\ \hat{x}^{2L+1} \end{bmatrix} + \begin{bmatrix} F(x(t_1)) \\ F(x(t_2)) \\ \vdots \\ F(x(t_{(2L+1)})) \end{bmatrix} = \begin{bmatrix} u(t_1) \\ u(t_2) \\ \vdots \\ \vdots \\ u(t_{(2L+1)}) \end{bmatrix}$$

$$DFT \begin{bmatrix} \frac{i2\pi L}{T} & 0 & 0 & 0 \\ 0 & \frac{i2\pi(L-1)}{T} & 0 & 0 \\ 0 & 0 & \ddots & 0 \\ 0 & 0 & 0 & -\frac{i2\pi L}{T} \end{bmatrix} (DFT)^{-1} \begin{bmatrix} x(t_1) \\ x(t_2) \\ \vdots \\ x(t_{(2L+1)}) \end{bmatrix} - \begin{bmatrix} F(x(t_1)) \\ F(x(t_2)) \\ \vdots \\ F(x(t_{(2L+1)})) \end{bmatrix} = \begin{bmatrix} u(t_1) \\ u(t_2) \\ \vdots \\ u(t_{(2L+1)}) \end{bmatrix}$$

Summary

Four Methods

- Time integration until steady-state achieved
- Finite difference methods
- Shooting Methods
- Spectral Methods

Shooting Methods

- State transition function
- Sensitivity matrix
- Matrix-Free Approach

Spectral Methods

Galerkin and Collocation Methods