

Soluciones

③ (24.29.)

①



La energía almacenada por un condensador es.

$$U = \frac{Q^2}{2C}$$

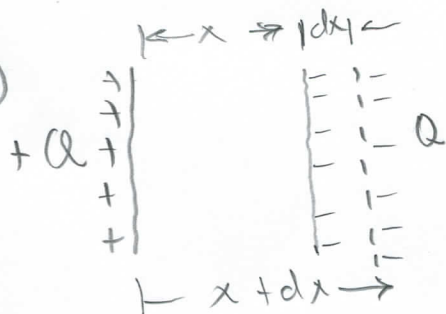
y la capacitancia para un condensador de placas paralelas de acuerdo a la geometría es. $C = \frac{\epsilon_0 A}{d} = \frac{\epsilon_0 A}{x}$.

luego de acuerdo con lo anterior la energía es.

$$U = \frac{Q^2 x}{2\epsilon_0 A}$$

cuyas variables son conocidos de acuerdo a los datos del problema.

②



luego. $U(x) = \frac{x Q^2}{2\epsilon_0 A} \Rightarrow \frac{dU(x)}{dx} = \frac{Q^2}{2\epsilon_0 A}$

$$\Rightarrow dU = \frac{Q^2}{2\epsilon_0 A} dx$$

③ $dU = dW = F dx \Rightarrow dU = \frac{Q^2}{2\epsilon_0 A} dx = F dx$

entonces $F = \frac{Q^2}{2\epsilon_0 A}$

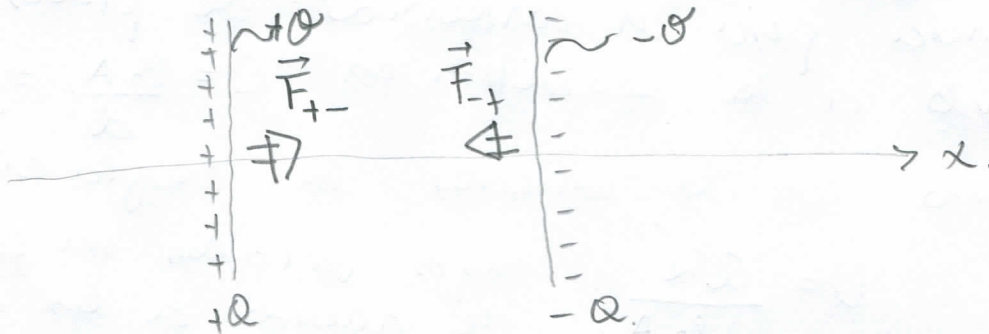
④ El campo debido a ambas placas para un condensador está dado por $E = \frac{\sigma}{\epsilon_0} = \frac{Q}{\epsilon_0 A}$. Para una sola placa la ley de Gauss establece que $E = \frac{\sigma}{2\epsilon_0} = \frac{Q}{2\epsilon_0 A}$ luego.

$F = \frac{Q^2}{2\epsilon_0 A}$ puesto que $F = QE$

↳ Fuerza que ejerce una placa sobre la otra.

⑥ De acuerdo con el problema ③ en la parte ①
 se tiene que la fuerza que ejerce una placa
 sobre la otra está dada por.

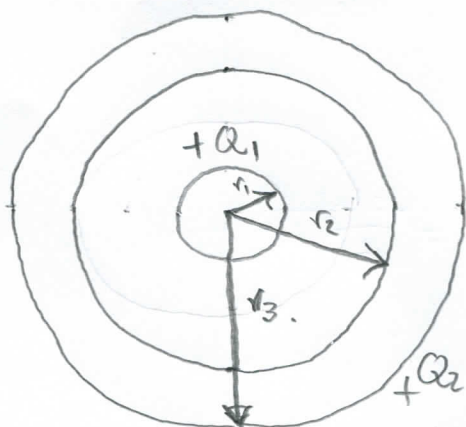
$$F = \frac{Q^2}{2\epsilon_0 A} = \frac{Q\sigma}{2\epsilon_0}$$



$$\vec{F}_{+-} = \frac{Q^2}{2\epsilon_0 A} \hat{i} = +\frac{Q\sigma}{2\epsilon_0} \hat{i}$$

$$\vec{F}_{-+} = -\frac{Q^2}{2\epsilon_0 A} \hat{i} = -\frac{Q\sigma}{2\epsilon_0} \hat{i}$$

(9)



$$r_1 = 10 \text{ cm}$$

$$r_2 = 30 \text{ cm}$$

$$r_3 = 40 \text{ cm}$$

$$Q_1 = 1 \times 10^{-5} \text{ C}$$

$$Q_2 = 8 \times 10^{-5} \text{ C}$$

Dentro de la bola metálica de radio r_1 el campo es nulo, luego $E(r < r_1) = 0 = -\frac{dV}{dr} \Rightarrow V = c + c!$

En el exterior de la bola metálica se tiene que

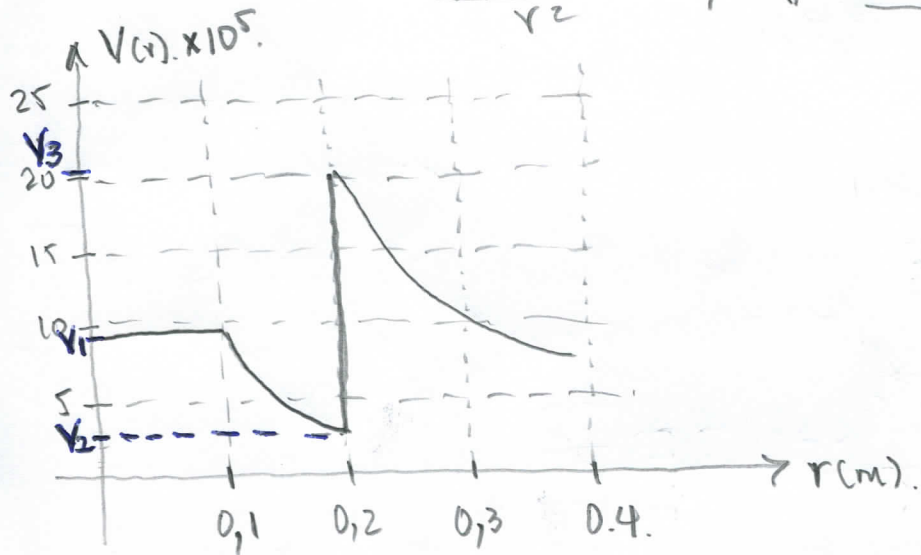
$$E(r_1 < r < r_2) = \frac{kQ_1}{r^2} \Rightarrow V = \frac{kQ_1}{r}$$

Entre el interior y el exterior del capacitor se tiene

$$E(r_2 < r < r_3) = 0 = -\frac{dV}{dr} \Rightarrow V = c + c!$$

En el exterior de la la coraza se tiene que

$$E(r > r_3) = \frac{k(Q_1 + Q_2)}{r^2} \Rightarrow V = \frac{k(Q_1 + Q_2)}{r}, \text{ luego.}$$

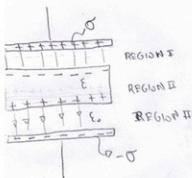


$$V_1 = \frac{kQ_1}{r_1} = 9 \times 10^5 \text{ V}$$

$$V_2 = \frac{k(Q_1 + Q_2)}{r_3} = 20,3 \times 10^5 \text{ V}$$

$$V_3 = \frac{kQ_1}{r_2} = 3 \times 10^5 \text{ V}$$

5



The electric field in the region I and II is

$$E = 2 \left(\frac{\sigma}{2\epsilon_0} \right) = \frac{\sigma}{\epsilon_0}$$

Now the field inside the metal slab is 0 therefore

$$\vec{E}(z) = \begin{cases} \sigma/\epsilon_0 & \text{if } 0 < z < \frac{1}{2}(d-a) \text{ or } \frac{(d+a)}{2} < z < d \\ 0 & \text{otherwise} \end{cases}$$

potential is $V = - \int E dz$

$$\therefore = \begin{cases} - \int \sigma/\epsilon_0 dz & \text{if } z < \frac{1}{2}(d-a) \text{ or } \frac{(d+a)}{2} < z < d \\ \text{CONST} & \end{cases}$$

$$V\left(\frac{1}{2}(d-a)\right) = \frac{\sigma}{\epsilon_0} \left(\frac{d-a}{2} \right) + C_1 = \text{const} = \frac{\sigma}{\epsilon_0} \left(\frac{d+a}{2} \right) + C_2 = V\left(\frac{1}{2}(d+a)\right)$$

$$\text{Hence } C_1 - C_2 = \frac{\sigma}{\epsilon_0} \left(\frac{d+a}{2} \right) - \left(\frac{d-a}{2} \right) \frac{\sigma}{\epsilon_0} = \frac{\sigma}{\epsilon_0} a$$

∴ result

$$= V(0) - V(d) = (C_1) - \left(\frac{d\sigma}{\epsilon_0} + C_2 \right) = (C_1 - C_2) - \frac{d\sigma}{\epsilon_0} = \frac{\sigma(a-d)}{\epsilon_0} = V_{12}$$

capacitance is

$$= \frac{Q}{V_{12}} = \left| \frac{\sigma A}{\sigma/\epsilon_0 (a-d)} \right| \Rightarrow$$

$$C = \epsilon_0 \left(\frac{A}{d-a} \right)$$

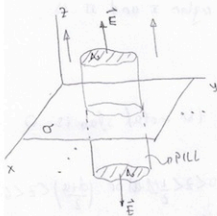
$$\text{if } a=0 \Rightarrow C=C_0 = \epsilon_0 \frac{A}{d}$$

$$\therefore C = \frac{\epsilon_0 A}{d} \left(\frac{d}{d-a} \right) \Rightarrow$$

$$C = C_0 \left(\frac{d}{d-a} \right)$$

$$+ \frac{1}{4\pi\epsilon_0} \frac{q^2}{r^2} \frac{1}{r^2} (2x-x) + \frac{1}{4\pi\epsilon_0} \frac{q^2}{r^2} B = \frac{1}{4\pi\epsilon_0} B$$

EM 8
finite sheet



Using the Gauss law

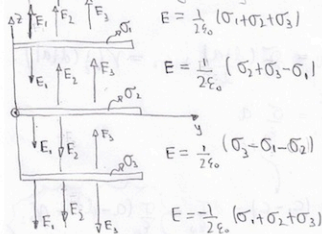
$$\oint_{\text{PILL}} \vec{E} \cdot d\vec{S} = \frac{Q_{\text{enc}}}{\epsilon_0}$$

$$2EA_0 = \frac{\sigma A_0}{\epsilon_0} \Rightarrow E = \frac{\sigma}{2\epsilon_0}$$

$$\vec{E}(z) = \text{sgn}(z) \frac{\sigma}{2\epsilon_0} \hat{z}$$

Electric field of a single sheet
placed at $z=0$

current problem is the following



$$E = \frac{1}{2\epsilon_0} (\sigma_1 + \sigma_2 + \sigma_3)$$

$$E = \frac{1}{2\epsilon_0} (\sigma_2 + \sigma_3 - \sigma_1)$$

$$E = \frac{1}{2\epsilon_0} (\sigma_3 - \sigma_1 - \sigma_2)$$

$$E = -\frac{1}{2\epsilon_0} (\sigma_1 + \sigma_2 + \sigma_3)$$

The field in the region $z \in [-h_2, h_1]$
is given by

$$\vec{E} = \frac{\hat{z}}{2\epsilon_0} \cdot \begin{cases} \sigma_2 + \sigma_3 - \sigma_1 & \text{if } z > 0 \\ \sigma_3 - \sigma_1 - \sigma_2 & \text{if } z < 0 \end{cases}$$

$$\text{Now } \vec{E} = -\frac{\partial V}{\partial z} \hat{z} \text{ hence}$$

$$V(z) = -\int \vec{E} dz$$

so

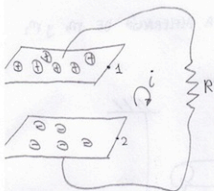
$$V(z) = -\frac{z}{2\epsilon_0} \cdot \begin{cases} \sigma_2 + \sigma_3 - \sigma_1 & \text{if } 0 < z < h_1 \\ \sigma_3 - \sigma_1 - \sigma_2 & \text{if } h_2 < z < 0 \end{cases} + \text{constant}$$

so $V(h_1) = -\frac{h_1}{2\epsilon_0} (\sigma_2 + \sigma_3 - \sigma_1)$ and $V(h_2) = +\frac{h_2}{2\epsilon_0} (\sigma_3 - \sigma_1 - \sigma_2)$. The potential difference between the upper and lower sheet is

$$V(h_1) - V(h_2) = V_{12} = -\frac{h_1}{2\epsilon_0} (\sigma_2 + \sigma_3 - \sigma_1) - \frac{h_2}{2\epsilon_0} (\sigma_3 - \sigma_1 - \sigma_2) = \frac{(\sigma_1 - \sigma_3)(h_1 + h_2)}{2\epsilon_0} + \frac{\sigma_2(h_2 - h_1)}{2\epsilon_0}$$

FM 24.27 WELL, the solution is too simple $W = \frac{1}{2} C V_0^2$ with $C = 450 \mu\text{F}$ and $V_0 = 245 \text{ V}$. However, there is more physics in this problem. Let us see

The conservation of energy implies



$$V_{12} = \frac{q}{C} + Ri = V_0$$

The current is $i = \frac{dq}{dt}$ then

$$\frac{dq}{dt} + \frac{1}{RC} q = \frac{V_0}{R}$$

$q(t) = Q_A e^{\alpha t} + Q_0$, then $\frac{dq}{dt} = \alpha Q_A e^{\alpha t}$ Hence (1) takes the form
with $Q_A, \alpha \in \mathbb{R}$

$$Q_A e^{\alpha t} + \frac{1}{RC} (Q_A e^{\alpha t} + Q_0) = \frac{V_0}{R}$$

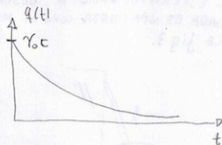
$$\left(1 + \frac{1}{RC}\right) Q_A e^{\alpha t} + \frac{Q_0}{RC} = \frac{V_0}{R} \quad \text{if we choose } \frac{Q_0}{RC} = \frac{V_0}{R} \Rightarrow \alpha + \frac{1}{RC} = 0$$

$$\text{Hence } \boxed{Q_0 = V_0 C} \quad \text{AND} \quad \boxed{\alpha = -1/RC}$$

a result, the charge in one of the plates of the condenser is

$$\boxed{q = V_0 C e^{-t/RC}}$$

this charge is decreasing exponentially with time



potential is dissipated by the resistance (R) of wire. This is

$$s = R \frac{di}{dt} = R \left(-\frac{V_0 C}{RC} e^{-t/RC} \right) \Rightarrow \boxed{V_{dis} = -V_0 e^{-t/RC}}$$

$$= dq V_{dis} = \underbrace{\left[V_0 C e^{-t/RC} \left(-\frac{dt}{RC} \right) \right]}_{dq} \underbrace{\left(-V_0 e^{-t/RC} \right)}_{V_{dis}} = \frac{V_0^2}{R} e^{-2t/RC} dt$$

$$J = \frac{V_0^2}{R} \left(-\frac{RC}{2} \right) e^{-t/RC} \Big|_0^t \Rightarrow \boxed{W(t) = \frac{1}{2} C V_0^2 \left[1 - e^{-t/RC} \right]}$$

ENERGY DISSIPATED
IN THE WIRE