

Closing Lemmas 2.6–2.9 (Dec 11 version): suggested completion  
 December 12, 2025

**To:** Milan Zlatanović  
**Cc:** Prof. Elshad Allahyarov  
**Project:** NS Overleaf ([link](#))

Dear Milan,

Thank you for clarifying. You are right: my earlier response was tied to the older (Dec 8) “single Lemma 2.6” iteration. I have now reviewed your Dec 11 restructuring into Lemmas 2.6–2.9 (singular point; normalization; rescaled domain; ancient limit) in `new-version-12-11.tex`. The structure is good and much easier to check.

**What is still needed to fully close the chain:**

1. **Lemma 2.6 (existence of a CKN singular point):** the logic is correct (if no singular point exists at  $T^*$ ,  $\varepsilon$ -regularity gives local boundedness near  $t = T^*$ , hence continuation past  $T^*$ , contradiction). This is the right “anchor” for the blow-up.
2. **Lemma 2.7 (blow-up normalization):** as currently written it chooses  $x_k$  from the *global* vorticity supremum at times  $t_k \uparrow T^*$ . This leaves an implicit issue:  $x_k$  could in principle drift (the noncompactness worry you flagged in the Dec 8 comments).

The clean fix is to *anchor the blow-up near a fixed singular point  $x^*$  from Lemma 2.6*. Concretely, replace the choice of  $x_k$  by either:

- (*Local vorticity choice*) choose  $x_k \in B_1(x^*)$  with  $|\omega(x_k, t_k)| = \|\omega(\cdot, t_k)\|_{L^\infty(B_1(x^*))}$ , and note this local supremum must diverge as  $t_k \uparrow T^*$  if  $(x^*, T^*)$  is singular; or
- (*CKN-normalization choice, recommended*) choose scales  $r_k \downarrow 0$  such that the CKN functional at  $(x^*, T^*)$  satisfies

$$r_k^{-2} \iint_{Q_{r_k}(x^*, T^*)} (|u|^3 + |p|^{3/2}) dx dt \geq \varepsilon_{\text{CKN}},$$

and define the blow-up by  $\lambda_k := r_k$  and center  $x_k := x^*$  (so no drift can occur).

The second option makes the nontriviality in Lemma 2.9 essentially automatic by semicontinuity.

3. **Lemma 2.8 (domain exhaustion):** looks correct.
4. **Lemma 2.9 (ancient limit):** this is the main place where the proof needs to be written out. There are two distinct sub-steps:
  - (a) *Compactness / passage to the limit.* State explicitly the uniform estimates on each fixed cylinder  $Q_R = B_R \times (-R^2, 0)$  inherited from the local energy inequality under scaling, e.g.

$$\sup_{s \in (-R^2, 0)} \int_{B_R} |u^{(k)}(s)|^2 + \int_{Q_R} |\nabla u^{(k)}|^2 \leq C(R), \quad \|p^{(k)}\|_{L^{3/2}(Q_R)} \leq C(R),$$

plus a standard bound on  $\partial_s u^{(k)}$  in a negative Sobolev space (so Aubin–Lions applies). This yields a subsequence converging strongly in  $L_{\text{loc}}^p$  (for  $p < 3$ ) and is enough to pass the nonlinear term and the local energy inequality, giving a suitable weak limit  $(u^\infty, p^\infty)$  on  $\mathbb{R}^3 \times (-\infty, 0)$ .

- (b) *Nontriviality of the limit (your item (iii)).* This is easiest if the blow-up is anchored by a *scale-invariant lower bound* (CKN functional) rather than a pointwise vorticity normalization. With the CKN-normalization choice in Lemma 2.7, we get on  $Q_1$ :

$$\iint_{Q_1} (|u^{(k)}|^3 + |p^{(k)}|^{3/2}) dx dt \geq \varepsilon_{\text{CKN}}.$$

By strong/weak convergence and lower semicontinuity, the same lower bound holds for the limit, hence  $u^\infty \not\equiv 0$ , which implies your desired statement  $\int_{Q_r} |u^\infty|^3 \geq c > 0$  for some  $r, c$ .

**One stylistic point:** I recommend removing any claim that the ancient limit is locally  $L^\infty$  in space; from local energy bounds one naturally inherits  $L_t^\infty L_x^2 \cap L_t^2 H_x^1$  and the scale-invariant  $L^3/L^{3/2}$  controls, which are what the later “geometric depletion” steps use.

Thanks again for doing the hard work of restructuring this. If you agree, I suggest implementing the CKN-normalized blow-up (anchored at  $x^*$ ) as the main route, and keeping the vorticity normalization as an optional remark (it is conceptually nice, but technically less robust for proving nontriviality of the ancient limit).

Sincerely,  
Jonathan Washburn