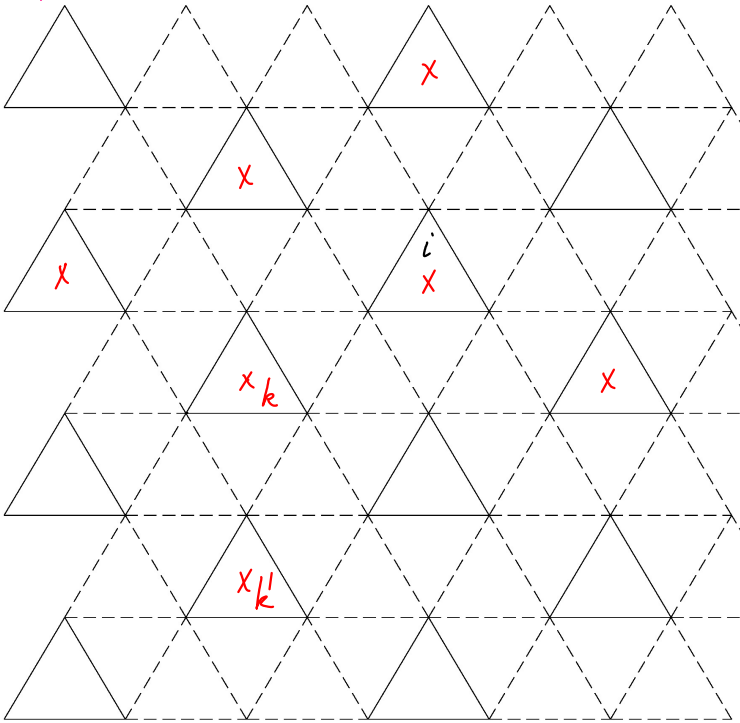


Real Space Renormalisation for the Ising Model on a Triangular Lattice



$$s_i = \pm 1.$$

$$-\beta H(\{s_i\}) = J \sum_{\langle ij \rangle} s_i s_j + B \sum_i s_i$$

$$t_k = \pm 1$$

$$\sum_i s_i^{(k)} > 0 \rightarrow t_k = 1$$

$$< 0 \quad t_k = -1$$

We carry out the standard renormalisation procedure

$$A e^{J' t_k t_{k'}} = \sum_{\{s_i^{(k)}, s_j^{(k')}\}} e^{-\beta \mathcal{H}(s_i^{(k)}, s_j^{(k')})} W(t_k, s_i^{(k)}) W(t_{k'}, s_j^{(k')}).$$

k, k' nearest neighbours; $J > 0$ Ferromagnetic

Take $t_k = t_{k'} = +1$: 4 possibilities for each triangle.

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+	+	+	+	+	+
+ $\times k$ +	+ +	+ +	- +	- +	- +
+	+	-	+	+	-
+ $\times k'$ +	+ \rightarrow -	+ +	+ +	+ \rightarrow -	+ +
e^{8J+6B}	$2 \times$ $2 e^{4J}$	1	$2 \times$ $2 e^{2J}$	$4 \times$ $4 e^{-2J}$	$2 \times$ $2 e^{-2J}$
-	-	-			
+ +	+ +	+ +			
+	+	-			
+ +	- \rightarrow +	+ +			
e^{4J}	$2 \times$ 2	e^{-4J}			

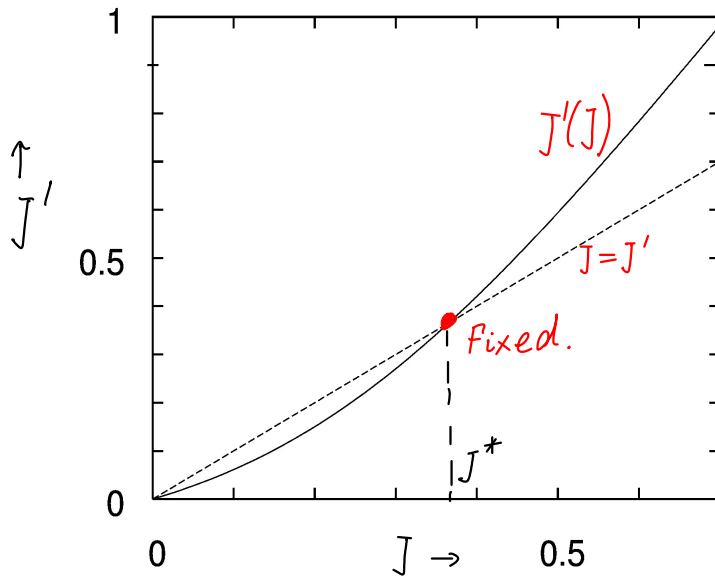
$$A e^{J'} = e^{8J} + 3 e^{4J} + 2 e^{2J} + 3 + 6 e^{-2J} + e^{-4J}.$$

We can do the same for opposite coarse spins, and

find: $A e^{-J'} = e^{4J} + 2 e^{2J} + 4 + 6 e^{-2J} + 2 e^{-4J}$

$$\frac{A e^{J'}}{A e^{-J'}} = e^{2J'} = \frac{e^{8J} + 3 e^{4J} + 2 e^{2J} + 3 + 6 e^{-2J} + e^{-4J}}{e^{4J} + 2 e^{2J} + 4 + 6 e^{-2J} + 2 e^{-4J}}.$$

$$J' = \frac{1}{2} \log \frac{A e^{J'}}{A e^{-J'}} = \frac{1}{2} \log \left(\frac{e^{8J} + 3 e^{4J} + 2 e^{2J} + 3 + 6 e^{-2J} + e^{-4J}}{e^{4J} + 2 e^{2J} + 4 + 6 e^{-2J} + 2 e^{-4J}} \right)$$



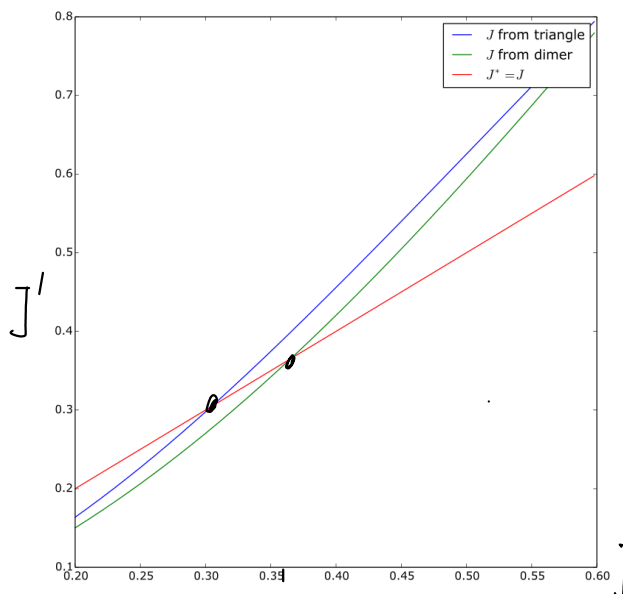
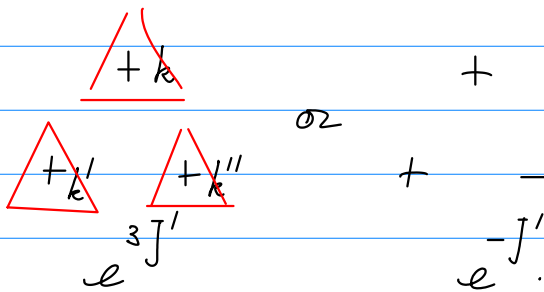
We find $J^* = 0.365$
and $\frac{dJ^*(J)}{dJ} = 1.544 = l^2$

$$l = \sqrt{3} \rightarrow z \approx 0.79$$

$$v = \frac{1}{z} \approx 1.26$$

$$v_{\text{exact}} = 1 \quad (\text{Houtappel 1950})$$

Add a magnetic field: $v = 2.02$ $v_{\text{exact}} = 1.875$.



There must be a
3-point interaction.

Qu: How to find better values for κ and ν ?

Answer: cumulant expansion.

$$H = H_0 + \bar{V} \quad H_0: \text{'simple'}$$
$$\bar{V}: \text{small correction.}$$

$$\mathcal{Z} = \sum_r e^{H_0 + \bar{V}}$$

$$\langle e^{\bar{V}} \rangle_0 = \frac{1}{\mathcal{Z}_0} \sum_r e^{H_0} e^{\bar{V}} = \frac{\mathcal{Z}}{\mathcal{Z}_0} \rightarrow \mathcal{Z} = \mathcal{Z}_0 \langle e^{\bar{V}} \rangle_0$$

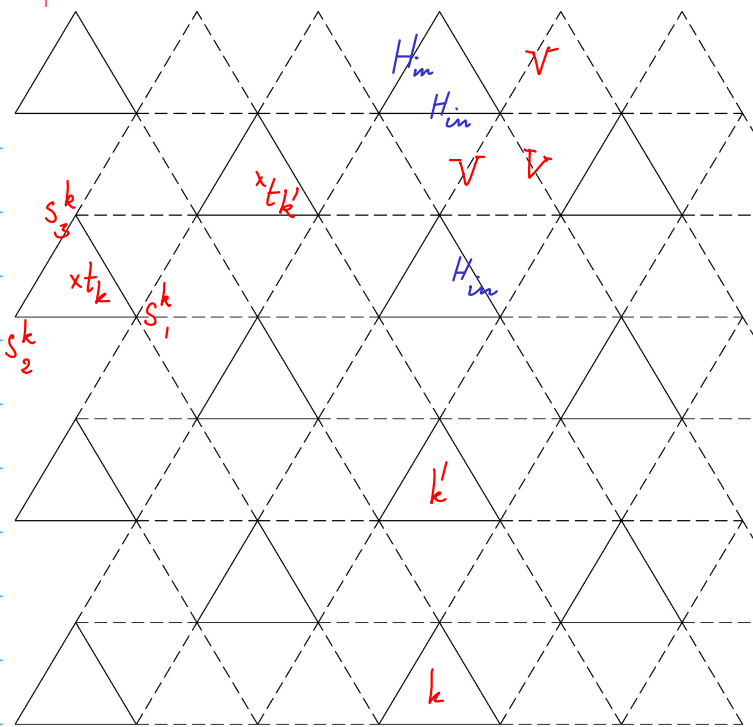
$$\langle e^{\bar{V}} \rangle_0 = 1 + \langle \bar{V} \rangle_0 + \frac{1}{2} \langle \bar{V}^2 \rangle_0 + \frac{1}{3!} \langle \bar{V}^3 \rangle_0 + \dots$$

\Rightarrow Perturbation expansion of \mathcal{Z} .

$$F = -k_B T \ln \mathcal{Z} = -k_B T \ln \mathcal{Z}_0 - k_B T \ln \left(1 + \underbrace{\langle \bar{V} \rangle_0 + \frac{1}{2} \langle \bar{V}^2 \rangle_0}_{\kappa} + \dots \right)$$

$$\ln(1+x) = x - \frac{x^2}{2} + \dots$$

$$F = F_0 - k_B T \left(\langle \bar{V} \rangle_0 + \frac{1}{2} \left(\langle \bar{V}^2 \rangle_0 - \langle \bar{V} \rangle_0^2 \right) + \frac{1}{3} \dots \right). \quad \text{Cumulant expansion.}$$



$$H_0 = H_{in}$$

H_{in} : couplings across solid lines

V : couplings across dashed lines.

Fine grained sites: k_i .

$$H_{in} = \sum_{\langle k_i, k_i' \rangle} s_i^k s_{i'}^k$$

$$V = \sum_{\substack{\langle k_i, k_i' \rangle \\ k \neq k'}} s_i^k s_{i'}^{k'}$$

Recall:

$$e^{\mathcal{H}(\{t_k\})} = \sum_{\{s_i^k\}} e^{\mathcal{H}(\{s_i^k\})} \prod_k W(t_k, s_i^{(k)}) = \sum_{\{s_i^k\}} e^{\mathcal{H}_{in} + V} \prod_k W(t_k, s_i^{(k)})$$

We define $e^{\overline{\mathcal{H}}(\{t_k\}, \{s_i^{(k)}\})} = e^{\mathcal{H}_{in}(\{s_i^{(k)}\})} \prod_k W(t_k, s_i^{(k)})$,

$$\text{So } e^{\mathcal{H}(\{t_k\})} = \sum_{\{s_i^k\}} e^{\overline{\mathcal{H}}(\{t_k\}, \{s_i^{(k)}\})} e^V$$

$$\text{We have } e^{\mathcal{H}(\{t_k\})} = Z_0 \langle e^V \rangle_0 \Rightarrow$$

$$\mathcal{H}(\{t_k\}) = A + \langle V \rangle_0 + \frac{1}{2} (\langle V^2 \rangle_0 - \langle V \rangle_0^2) + \frac{1}{3} \dots$$

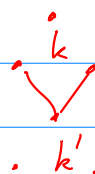
$$\langle V \rangle_0 = \frac{1}{\Omega} \sum_{\langle i, i' \rangle} \langle s_i^k s_{i'}^{k'} \rangle = \frac{1}{\Omega} \sum_{\langle i, i' \rangle} \langle s_i^k \rangle \langle s_{i'}^{k'} \rangle$$

Take $t_k = +1$.

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$$\langle s_i^k \rangle = \frac{e^{3J} + e^{-J} - e^{-J} + e^{-J}}{e^{3J} + 3e^{-J}} = \frac{e^{3J} + e^{-J}}{e^{3J} + 3e^{-J}}$$

$$t_k = -1 : \langle s_i^k \rangle = - \frac{e^{3J} + e^{-J}}{e^{3J} + 3e^{-J}}$$



$$\chi^p(t_k, t_{k'}) = J' t_k t_{k'} = 2J \left(\frac{e^{3J} + e^{-J}}{e^{3J} + 3e^{-J}} \right)^2 t_k t_{k'}$$

Fixed point: $J' = J \Rightarrow \frac{e^{3J^*} + e^{-J^*}}{e^{3J^*} + 3e^{-J^*}} = \frac{1}{\sqrt{2}} \rightarrow J^* = 0.3356$

$$\left. \frac{dJ'}{dJ} \right|_{J=J^*} = 2 \left(\frac{e^{4J^*} + 1}{e^{4J^*} + 3} \right)^2 + 4J^* \left(\frac{e^{4J^*} + 1}{e^{4J^*} + 3} \right) \left(\frac{4e^{4J^*}}{e^{4J^*} + 3} - \frac{e^{4J^*} + 1}{(e^{4J^*} + 3)^2} 4e^{4J^*} \right)$$

$$= 1.624 = 1^{\frac{2}{3}} = \sqrt{3}^{\frac{2}{3}} \rightarrow \frac{2}{3} = 0.667$$

Before $\frac{2}{3} = 0.79$

exact: $\frac{2}{3} = 1.0$

magnetic scaling dim.

$$\nu = 1.37 < \nu_{\text{exact}} = 1.875$$

$$\lambda_T = (\sqrt{3})^{\frac{2}{3}} \rightarrow \text{converges towards } 1.782 \text{ (5-cells)}$$

Cumulant expansion

$$\lambda_T = 1.784$$