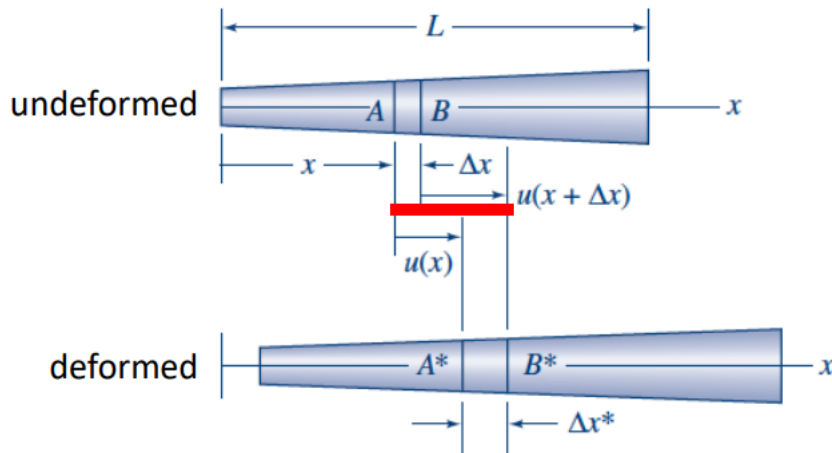


Lecture 21

Computer Project: Axially-deformable Bars

Kinematic assumptions: cross sections, which are plane and are perpendicular to the axis before deformation, remain plane and remain perpendicular to the axis after deformation. In addition, cross sections do not rotate about the axis.

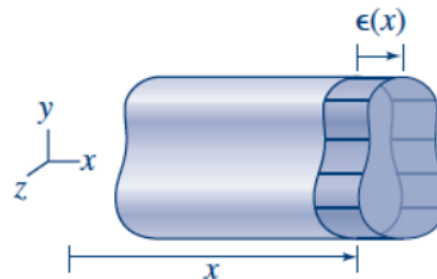


The strain-displacement relationship is ...

$$\epsilon_x(x) = \lim_{\Delta x \rightarrow 0} \frac{\Delta x^* - \Delta x}{\Delta x} = \frac{du(x)}{dx} = \epsilon(x)$$

$$\Delta x^* + u(x) = \Delta x + u(x + \Delta x)$$

$$\Delta x^* - \Delta x = u(x + \Delta x) - u(x)$$



HW8: Matlab Project

Consider an elastic bar with Young's modulus, $E = 10$, the cross section area $A = 1$, and the length of the bar $L = 1$. The bar has a built-in boundary condition at $x = 0$, i.e. $u(0) = 0$, and at $x = L$, the internal force $R(L) = 0$ as shown in Fig. 1.

The differential equation that governs the equilibrium of the bar has been derived as follows,

$$\frac{d}{dx} \left(EA \frac{du}{dx} \right) + b(x) = 0, \quad 0 < x < L,$$

where $u(x)$ is the displacement field.

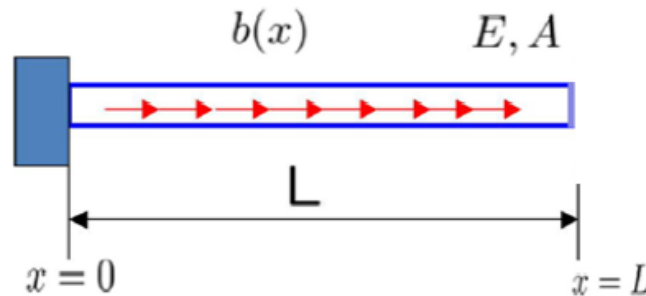
The bar is subjected a distributed load along its span, i.e.

$$b(x) = p \sin\left(\frac{2\pi x}{L}\right),$$

where $p = 1$ with a unit of force per unit length.

The internal force is defined as $R(x) = \sigma A = EA\epsilon$, i.e.

$$R(x) = EA \frac{du}{dx}$$

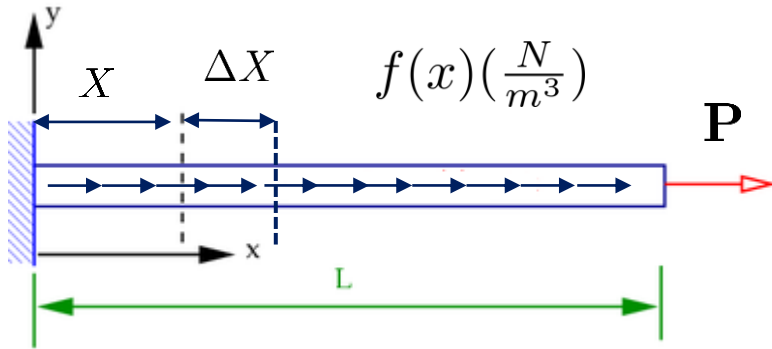


$$u(0) = 0, \quad R(L) = 0$$

Exact solution

$$R(x) = \frac{pL}{2\pi} \cos\left(\frac{2\pi}{L}x\right) - \frac{pL}{2\pi}; \quad u(x) = \frac{pL^2}{4\pi^2 EA} \sin\frac{2\pi x}{L} - \frac{pL}{2\pi EA}x$$

HW8: Matlab Project



B.C. : $u(0) = 0$, $R(L) = P$

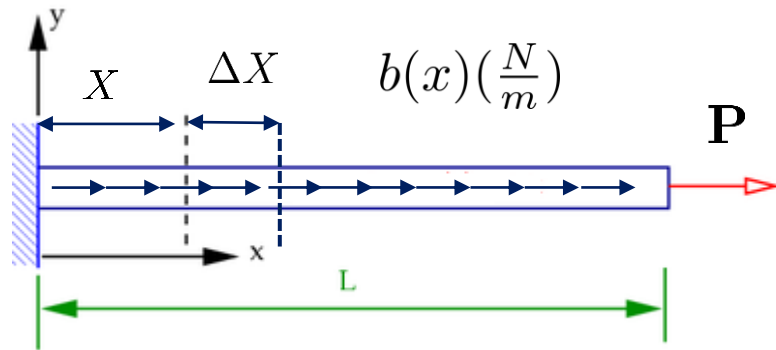
$$(1) \quad \frac{d\sigma}{dx} + f(x) = 0; \quad (2) \quad \frac{d}{dx} E \left(\frac{du}{dx} \right) + f(x) = 0 .$$

$$\text{Let } R(x) = EA \frac{du}{dx} \rightarrow \frac{d}{dx} R(x) + b(x) = 0 .$$

$$\frac{d}{dx} \left(AE \frac{du}{dx} \right) + b(x) = 0 , \text{ where } b(x) = Af(x) \sim (F/L)$$

Today's Lecture Attendance Password is: Matlab Project

HW8



$$\text{B.C. : } u(0) = 0, \quad R(L) = P$$

$$\sum F_x = 0 \rightarrow -R(x) + R(x + \Delta x) + b(x)\Delta x = 0$$

$$\frac{R(x + \Delta x) - R(x)}{\Delta x} + b(x) = 0$$

Let $R(x) = EA \frac{du}{dx} \rightarrow \frac{d}{dx} R(x) + b(x) = 0$.

$$\frac{d}{dx} \left(AE \frac{du}{dx} \right) + b(x) = 0; \quad \forall 0 < x < L, \quad \text{where } b(x) \sim (F/L)$$

How to solve this equation by using MatLab

In the lecture these four relations have been combined to obtain a single equation representing equilibrium in terms of the displacement,

$$\frac{d}{dx} \left[EA \frac{du}{dx} \right] + b = 0. \quad (5)$$

In order to utilize the solver in MATLAB, one must convert the governing equations into first-order form,

$$\boxed{\frac{dy}{dx} = f(y, x)}, \quad (6)$$

where y is a vector of unknown variables, and f is a vector of known functions depending on y and the position x .

By defining,

$$R = EA \frac{du}{dx}.$$

$$\frac{dR}{dx} + b = 0,$$

$$\frac{d}{dx} \begin{bmatrix} u \\ R \end{bmatrix} = \begin{bmatrix} R \\ \frac{R}{E(x)A(x)} \\ -b(x) \end{bmatrix}$$

$$\boxed{\begin{aligned} y &:= \begin{bmatrix} y_1 \\ y_2 \end{bmatrix} = \begin{bmatrix} u \\ R \end{bmatrix}, \\ f(y, x) &:= \begin{bmatrix} f_1(y, x) \\ f_2(y, x) \end{bmatrix} = \begin{bmatrix} \frac{y_2}{E(x)A(x)} \\ -b(x) \end{bmatrix}, \end{aligned}} \quad (7)$$

one obtains the desired first-order form,

$$\frac{d}{dx} \begin{bmatrix} y_1 \\ y_2 \end{bmatrix} = \begin{bmatrix} \frac{y_2}{E(x)A(x)} \\ -b(x) \end{bmatrix}.$$

2.2 Boundary condition

In order to apply boundary conditions in the solver in MATLAB, one must define a function which returns a residual of how much the boundary conditions are not satisfied; a residual of zero implies that the boundary conditions are satisfied exactly. The function has the form,

$$\boxed{g(y(a), y(b))}$$

where g is vector of functions depending on the value of y evaluated at the boundary points $x = a$ and $x = b$ (Here we assume the problem is defined on the interval (a, b)).

For example:

To clarify the form of the function, consider the boundary condition,

$$\begin{aligned} u(0) &= u_0, \\ R(L) &= R_L, \end{aligned}$$

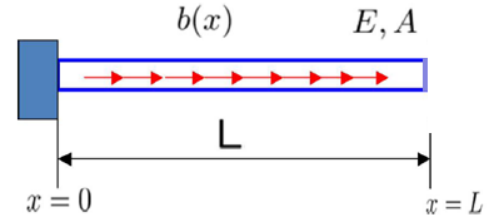
where the displacement is known as u_0 at the end point $x = 0$, and the force is known as R_L at the end point $x = L$. The vector defining the residual of how much the boundary condition is not satisfied is,

$$\begin{bmatrix} u(0) - u_0 \\ R(L) - R_L \end{bmatrix} \stackrel{=0}{\leftarrow} R(L) = EA \frac{du}{dx} \Big|_{x=L}$$

Using the correspondence between u, R and y defined in (7), one defines g as,

Res

$$\boxed{g(y(0), y(L)) := \begin{bmatrix} g_1(y(0), y(L)) \\ g_2(y(0), y(L)) \end{bmatrix} = \begin{bmatrix} y_1(0) - u_0 \\ y_2(L) - R_L \end{bmatrix}. \quad (8)}$$



Use Matlab to find the displacement field and internal force/stress field, and compare them with the statically indeterminate system that has the same dimensions and the same material properties, but with different boundary conditions: $u(0) = u(L) = 0$. Hint: Go to class Bcourses website and go to the lecture folder, and then download a Matlab-P1 folder that contains the file: bar1d.m. You start your solution there. For all details, please refer to Lecture20F.pdf slide.

bvp4c

Solve boundary value problem – fourth-order method

co

Syntax

```
sol = bvp4c(odefun,bcfun,solinit)
sol = bvp4c(odefun,bcfun,solinit,options)
```

Description

`sol = bvp4c(odefun,bcfun,solinit)` integrates a system of differential equations of the form $y' = f(x,y)$ specified by `odefun`, subject to the boundary conditions described by `bcfun` and the initial solution guess `solinit`. Use the `bvpinit` function to create the initial guess `solinit`, which also defines the points at which the boundary conditions in `bcfun` are enforced.

`sol = bvp4c(odefun,bcfun,solinit,options)` also uses the integration settings defined by `options`, which is an argument created using the `bvpset` function. For example, use the `AbsTol` and `RelTol` options to specify absolute and relative error tolerances, or the `FJacobian` option to provide the analytical partial derivatives of `odefun`.

$$\frac{dy}{dx} = f(y, x)$$

Res
y

```

1 %%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%
2 % -- Matlab code (BVP4C) to solve Tension-Compression Bar(1D)
3 %%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%%
4 function bar1d
5 % -- Define geometry
6 L = 1; % -- Length of bar
7
8 % -- Set solver parameters
9 nvar = 2; % -- Number of variables
10 np = 10; % -- Initial Number of points on [0,L]
11 xp = linspace(0,L,np); % -- Initial Points at which to satisfy ODE
12
13 % -- Set initial solution for the solver
14 solinit = bvpinit(xp,zeros(1,nvar));
15
16 % -- Set options on the tolerance for the accuracy of the solution
17 % Default: RelTol(1e-3), AbsTol(1e-6)
18 %
19 %  $y'(x) = f(x,y(x)) + \text{res}(x)$ 
20 %
21 % norm( res(i)/f(i) ) <= RelTol and
22 % norm( res(i) ) <= AbsTol
23 %
24 options = bvpset('RelTol',1e-3,'AbsTol',1e-6);
25
26 % -- Invoke solver bvp4c is a built-in Matlab function
27 sol = bvp4c(@bar1d_ode,@bar1d_bc,solinit,options);
28
29 % -- Plot solution
30 bar1d_plot(L,sol);
31
32 end

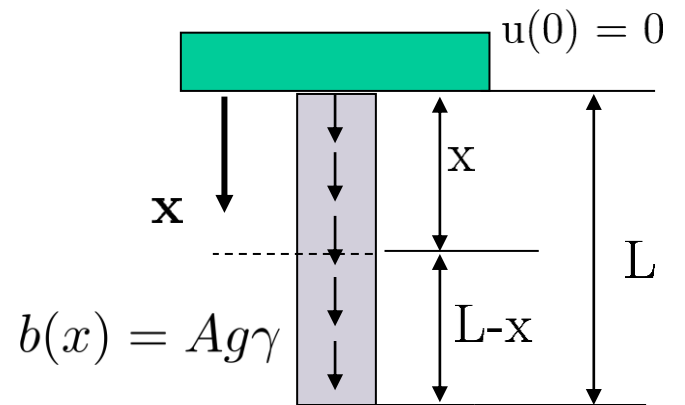
```

y


```

48 function [fxy] = barld_ode(x,y)
49
50 % -- Define material property and geometry
51 E = 1; % Young's Modulus
52 A = 1; % Cross-sectional area
53 L = 1; % Length of bar
54
55 g=9.8;
56 gamma=1;
57
58 % -- Define distributed load
59 b = gamma*g;
60
61 % -- Define function
62 fxy = [y(2)/(E*A);
63        -b];
64 end

```



$$\frac{d}{dx} \begin{bmatrix} y_1 \\ y_2 \end{bmatrix} = \begin{bmatrix} \frac{y_2}{E(x)A(x)} \\ -b(x) \end{bmatrix}.$$

```

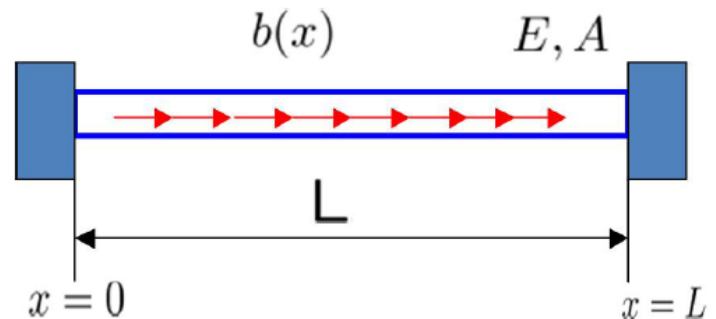
*****
function [fxy] = barld_ode(x,y)

% -- Define material property and geometry
E = 1; % Young's Modulus
A = 1; % Cross-sectional area
L = 1; % Length of bar

% -- Define distributed load
b = sin(pi*x/(2*L));

% -- Define function
fxy = [y(2)/(E*A);
       -b];
end

```



```

function [res] = barld_bc(ya,yb)

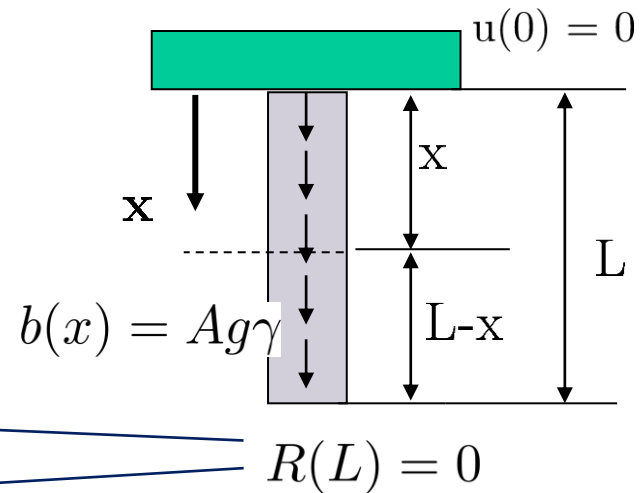
% -- Boundary Conditions (BC)
%   u: displacement
%   f: force

ua = 0;      % -- Fixed      at x=a
fb = 0;      % -- Zero force at x=b

res= [ya(1)-ua;
      yb(2)-fb];

end

```



```

function [res] = barld_bc(ya,yb)

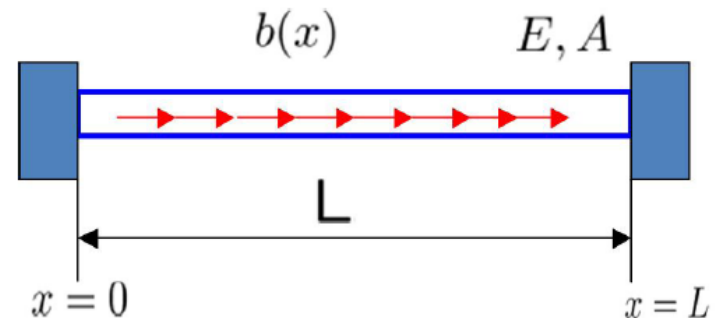
% -- Boundary Conditions (BC)
%   u: displacement
%   f: force

ua = 0;      % -- Fixed      at x=a
ub = 0;      % -- Fixed      at x=b
%fb = 0;      % -- Zero force at x=b

res= [ya(1)-ua;
      yb(1)-ub];

end

```



Statically indeterminate problem



HW8.pdf

9:58am 9:58am

86 KB



Matlab-P1.zip

9:53am 9:53am

75 KB

C30-2023 > Matlab-P1



Name



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Type

Size



bar1d

10/8/2023 9:50 AM

MATLAB Code

4 KB



bar1d_M

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MATLAB Live Script

42 KB



Fig1

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PNG File

34 KB

Matlab Live Editor

Live Editor - C:\Users\shaof\Desktop\C30-2023\Matlab-P1\bar1d.mlx *

LIVE EDITOR INSERT VIEW

Code Control Task Section Break Text Table of Contents Code Example Image Hyperlink Equation

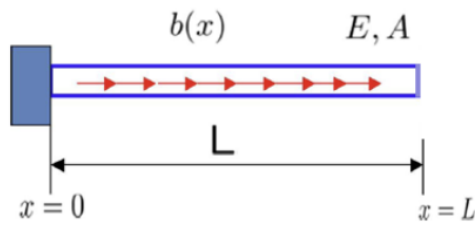
CODE SECTION TEXT IMAGE LINK EQUATION

bar1d.mlx * ✕

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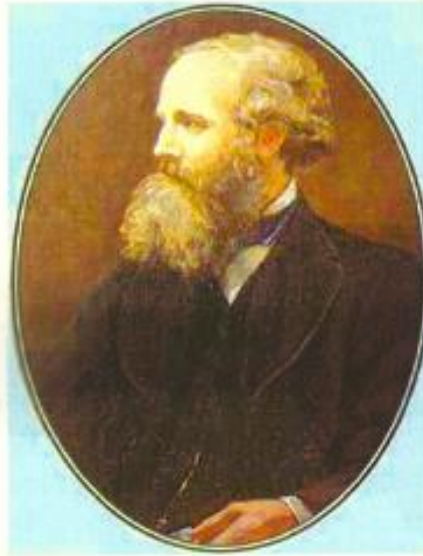
HW8 Assignment

Consider an elastic bar with Young's modulus, $E=10$, the cross section area $A=1$, and the length of the bar $L=1$. The bar has a built-in boundary condition at $x=0$, i.e. $u(0)=0$, and at $x=L$, the internal force $R(L)=0$ as shown in Fig. 1.



The diagram shows a horizontal elastic bar of length L . The left end is fixed at $x=0$, and the right end is at $x=L$. A distributed load $b(x)$ is applied along the length of the bar, represented by red arrows pointing to the right. The bar has Young's modulus E and cross-sectional area A .

The person who invented force method



(13 June 1831 – 5 November 1879)

James Clerk Maxwell

Scottish mathematician and physicist, is known for his contribution in electromagnetics theory, (the celebrated Maxwell equations), and his contribution in statistical physics (Maxwellian distribution).



412

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J. Clerk Maxwell F.R.S.

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[294]

L. On the Calculation of the Equilibrium and Stiffness of Frames.
By J. CLERK MAXWELL, F.R.S., Professor of Natural Philo-
sophy in King's College, London.*

THE theory of the equilibrium and deflections of frameworks
subjected to the action of forces is sometimes considered

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L. on the calculation of the equilibrium and stiffness of frames

JC Maxwell - The London, Edinburgh, and Dublin Philosophical ..., 1864 - Taylor & Francis

THE theory of the equilibrium and deflections of frameworks subjected to the action of forces is sometimes considered as more complicated than it really is, especially in cases in which the framework is not simply stiff, but is strengthened (or weakened as it may be) by additional connecting pieces. I have therefore stated a general method of solving all such questions in the least complicated manner. The method is derived from the principle of ...

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L. *On the Calculation of the Equilibrium and Stiffness of Frames.*
By J. CLERK MAXWELL, F.R.S., Professor of Natural Philo-
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THE theory of the equilibrium and deflections of frameworks subjected to the action of forces is sometimes considered as more complicated than it really is, especially in cases in which the framework is not simply stiff, but is strengthened (or weakened as it may be) by additional connecting pieces.

I have therefore stated a general method of solving all such questions in the least complicated manner. The method is derived from the principle of Conservation of Energy, and is referred to in Lamé's *Leçons sur l'Elasticité*, Leçon 7^{me}, as Clapeyron's Theorem; but I have not yet seen any detailed application of it.

If such questions were attempted, especially in cases of three dimensions, by the regular method of equations of forces, every point would have three equations to determine its equilibrium, so as to give $3s$ equations between e unknown quantities, if s be the number of points and e the number of connexions. There are, however, six equations of equilibrium of the system which must be fulfilled necessarily by the forces, on account of the equality of action and reaction in each piece. Hence if

$$m \sim e, \quad s \sim n$$

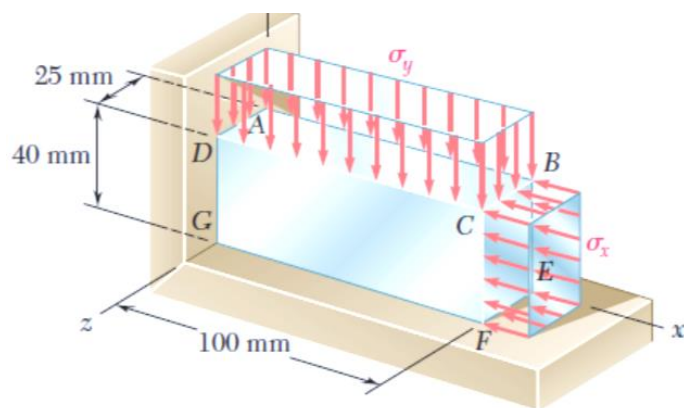
$$m + r = 3n \rightarrow$$

$$e = 3s - 6,$$

the effect of any external force will be definite in producing tensions or pressures in the different pieces; but if $e > 3s - 6$, these forces will be indeterminate. This indeterminateness is got rid

PROBLEM 9.81

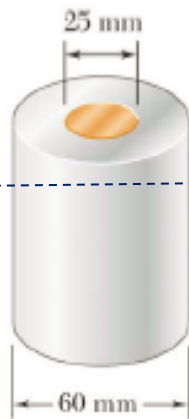
The block shown is made of a magnesium alloy, for which $E = 45 \text{ GPa}$ and $\nu = 0.35$. Knowing that $\sigma_x = -180 \text{ MPa}$, determine (a) the magnitude of σ_y for which the change in the height of the block will be zero, (b) the corresponding change in the area of the face $ABCD$, (c) the corresponding change in the volume of the block.



$$\begin{aligned} (a) \quad & \delta_y = 0 \quad \varepsilon_y = 0 \quad \sigma_z = 0 \\ & \varepsilon_y = \frac{1}{E}(\sigma_x - \nu\sigma_y - \nu\sigma_z) \\ & \sigma_y = \nu\sigma_x = (0.35)(-180 \times 10^6) \\ & = -63 \times 10^6 \text{ Pa} \end{aligned}$$

$$\sigma_y = -63.0 \text{ MPa} \quad \blacktriangleleft$$

$$\begin{aligned} \varepsilon_z &= \frac{1}{E}(\sigma_z - \nu\sigma_x - \nu\sigma_y) = -\frac{\nu}{E}(\sigma_x + \sigma_y) = -\frac{(0.35)(-243 \times 10^6)}{45 \times 10^9} = +1.890 \times 10^{-3} \\ \varepsilon_x &= \frac{1}{E}(\sigma_x - \nu\sigma_y - \nu\sigma_z) = \frac{\sigma_x - \nu\sigma_y}{E} = -\frac{157.95 \times 10^6}{45 \times 10^9} = -3.510 \times 10^{-3} \end{aligned}$$



Brass core
 $E = 105 \text{ GPa}$
 $\alpha = 20.9 \times 10^{-6}/^\circ\text{C}$

Aluminum shell
 $E = 70 \text{ GPa}$
 $\alpha = 23.6 \times 10^{-6}/^\circ\text{C}$

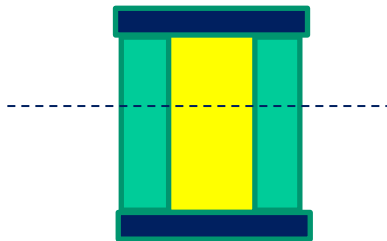
PROBLEM 9.38

$$\Delta T = 195 - 15 = 180^\circ\text{C}$$

The aluminum shell is fully bonded to the brass core and the assembly is unstressed at a temperature of 15°C . Considering only axial deformations, determine the stress in the aluminum when the temperature reaches 195°C .

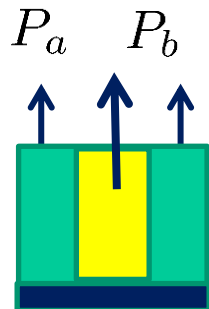
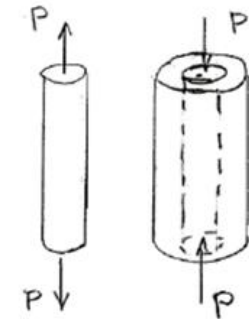
$$\Delta_a^T = \alpha_a \Delta T L > \Delta_b^T = \alpha_b \Delta T L$$

$$\sum F_y = 0 \rightarrow P_b + P_a = 0 \rightarrow P_b = -P_a =: P$$



$$f_a = \frac{L}{E_a A_a} \quad f_b = \frac{L}{E_b A_b}$$

Constrain condition



$$\Delta_a^T + \Delta_a^P = \Delta_b^T + \Delta_b^P$$

$$\rightarrow \alpha_a \Delta T L - f_a P = \alpha_b \Delta T L + f_b P$$

$$P = \frac{(\alpha_a - \alpha_b) \Delta T L}{(f_a + f_b)} \rightarrow \sigma_b = \frac{P}{A_b}$$