Simulations of linear and branched actin networks with formin and arp 2/3 complex as mediators

Ondrej Maxian March 5, 2024

1 Introduction

A great mystery in cell biology is how cells mediate the assembly of different cytoskeletal structures at different times in the cell cycle. One important example of this is branched vs. linear actin filament networks. When cells migrate, they form a branched network at their leading edge, which can be important for stable adhesion to a substrate. On the other hand, when cells divide, they form a contractile ring of linear actin filaments. The unidirectional (non-branched) nature of the ring is vital to transmit force in the correct direction for the completion of cytokinesis [1].

While it is clear that different network architectures are required for different cellular functions, the molecular mechanisms by which cells control the assembly of branched vs. linear actin structures are still incompletely understood. At the most basic level, linear actin filaments can assemble in two ways: via formin-mediated assembly, and via spontaneous nucleation [1, 25]. Seeing as spontaneous nucleation is energetically unfavorable, formin can significantly speed up filament nucleation [28] to the point where it dominates assembly dynamics. This is also the case in the presence of profilin [29], which can block the spontaneous nucleation of filaments, thereby causing filament assembly to be dependent on formin [1, 26].

Branched actin networks are dependent on a different assembly factor called the arp 2/3 complex. While this complex can aid in nucleation of linear filaments, assembly is most efficient when it binds to an existing "mother" filament. The fixed binding angle of arp 2/3 complex (typically around 70 degrees) causes the formation of tree-like branched actin networks [25, 1]. Because the arp 2/3 complex fixes the minus end of the "daughter" filament onto the mother, the branches can only grow from their plus ends, and thus branched structures do not exhibit treadmilling and are typically more stable than their linear counterparts [1, 22, 17]. Unlike formin networks, whose

growth is aided by the presence of profilin, arp 2/3-complex-mediated growth is *inhibited* by profilin, which blocks the association of an arp 2/3 activator (WASP VCA) with monomeric actin, a step which is necessary for branch formation [32].

While the key molecular players in branched vs. linear network formation are by now well understood, there is still some debate over exactly how the cell might program the actin to polymerize into one type of structure vs. the other. In the simplest case when there is exactly one structure formed, the kinetics become quite simple: the elongation rate is proportional to a constant times the amount of available (free) actin monomers [33], and the dissociation rate is constant. Solving the simple equation $k_{\rm on}C = k_{\rm off}$ gives an equilibrium concentration (about 0.1 μ M) at which the structure assumes a steady state size. But introducing multiple structures (either of the same kind or of different types), throws off this balance, as actin assembly kinetics alone do not provide a mechanism for controlling the size (length) or type (linear vs. branched) of multiple actin structures [21]. To control this, the conventional viewpoint is that external or internal signals [6] activate signaling cascades which eventually cause the cell to produce more or less of a particular type of network [18, 25, 33]. However, recent experimental evidence has shown that a limiting actin pool must also play a part in determining the types of structures being produced [5].

The general evidence for the limiting pool model is as follows: if cells only use signaling to preferentially organize actin into one particular type of structure, then changes in the actin concentration should affect the amount, but not the relative proportion, of linear vs. branched networks. However, experiments in fission yeast show that underexpression of actin favors formin-mediated linear structures, while overexpression favors arp 2/3-complex-mediated actin patches [5]. Another important set of observations comes from inhibition of arp 2/3 complex. As long as there is sufficient formin, inhibition of arp 2/3 (via drug CK-666) leads to new assembly of filamentous structures, as might be expected (see [36] for a similar result in plant cells). However, subsequent wash out of CK-666 gives disassembly of the additional filamentous structures and subsequent reassembly of dense actin patches [5]. These results suggest that the release of F-actin via disassembly is important for the assembly of new structures [33], and that the occupation of monomers in existing branched structures limits the amount of linear filaments that can be created. Such a scenario was shown to be important in cytokinesis in C. elegans, where depleting the arp 2/3 complex gives an excess of filamentous actin and consequent delays in cell division [7]. It was also shown that depleting arp 2/3 decreases actin turnover times (by increasing the amount of available actin) in the cortex of drosophila embryos [35], demonstrating the versatility of these results across organisms.

These recent in vivo experimental studies have spawned a number of in vitro and in silico (modeling) studies that attempt to understand how a limiting actin pool might set the balance between linear and branched filament networks. Thus far, studies concerned with size control have ignored the precise geometry of branched networks, and have instead looked at filamentous actin structures only, finding that the steady state length distribution of individual filaments (or structures) is not well defined [21, 12] because actin kinetics only set the total amount of actin incorporated. In fact, the limiting pool hypothesis predicts that structures with slightly higher growth rates will capture all of the monomers at equilibrium [21]. But this is not the behavior we observe in vivo, where the balance between branched and filamentous actin is maintained at a relative steady state in time.

There are a number of potential explanations for why the steady state theory does not match the experimental results. A first possibility is that each individual structure has an inherent mechanism of size control, in particular if there is a net negative feedback in the growth rate as a function of the structure size [3]. Such feedback could occur through fragmentation and annealing, where longer segments detach and attach to monomers, and fragmentation of longer filaments at steady states balances annealing of shorter ones [10]. Another mechanism for length-dependent attachment rate could occur in vivo from sensing of the cell boundary [19] or from the dynamics of linear filaments in bundles [27]. In fact, recent work [20] has suggested that size control can also emerge in a feedback-independent way if filaments form bundles (couples the dynamics of multiple filaments into one structure) and are removed randomly from the bundle with an independent probability (leads to a steady state because long filaments eventually get removed).

A second possibility for why the *in vivo* experiments do not match steady state predictions is more obvious: what if the cell simply is not operating at steady state? Indeed, it has been shown that intermediate filament lengths are quasi-stable on timescales of a few hours, and that the drift to steady state occurs on timescales as long as days [14, 34]. The intermediate-timescale regime was the focus of [2], where it was found that spontaneous nucleation of actin filaments gives rise to a quasi-stable, near-uniform length distribution, which can change to exponential in the presence of capping protein. This kind of length distribution does not match the expected steady state exponential one [9], because it is only stable on intermediate timescales. Nevertheless, these timescales, which are on the order of tens of minutes to a few hours, are the most relevant in *in vitro* experiments and in the cell itself, where cell-cycle cues typically take place every few hours to change the environment.

1.1 The focus of this study

It is therefore clear that further study is necessary to clear up how the balance of actin assembly changes on short to intermediate to long timescales. In addition, while previous size-control studies are pleasingly general [3, 21], the fact remains that branched actin networks require the presence of linear filaments to form, and so the specific modeling of branched and linear networks is necessary to truly understand how formin and arp 2/3 complex compete for actin monomers. Our goal in this study is to address these two points: by specifically modeling branched and linear networks, we can get an idea of how a limiting actin pool alone can give rise to different network architectures. And by simulating dynamically, we can understand how long the quasi-steady states with near uniform filament length distributions can persist in time [2].

Because some of the *in vitro* experiments on these systems pattern formin in one place and arp 2/3 complex in another, it will be important to look at the spatial dynamics of the filaments. Yet, because the nucleation reactions are reaction-limited, and because most reaction rates are reported in terms of bulk concentrations, it makes little sense to explicitly track actin monomers in space. Because of this, we settle on a hybrid approach where we treat actin nucleates as bulk quantities, which form fibers at random points in space. While we will not do so in this draft, our approach can be extended to incorporate diffusion, hydrodynamics, and explicit forces on the filaments which affect their dynamics.

The simulation approach we employ is an amalgamation of previous work [10, 22, 2]. As in [10, 2], we use a stochastic simulation algorithm to simulate a list of well-mixed reactions involving filaments and formins, but in this work we add the simulation of branched actin networks. Similar to [22, 17], we explicitly simulate branched actin networks, since their geometry (the number of branches and length of each branch) determines the rate of monomer binding and the possibility of arp 2/3 complex binding to form another branch. Unlike in [22], however, we do not explicitly track individual actin monomers in space, which allows us to reach longer time and lengthscales. Specifically, we do not introduce seed filaments to skip the lag phase of actin assembly [8], and are consequently able to explicitly simulate the actin nucleation reactions proposed in [28] to obtain a more accurate picture of how formin-based nucleation competes with spontaneous nucleation. Combining tools from these previous works allows us to study, for the first time to our knowledge, the spatial and temporal dynamics of how formins and arp 2/3 compete for actin monomers to form linear vs. branched networks.

Parameter	Description	Value	Units	
$k_0^{\mathrm{d}+}$	Dimer formation rate	3.5×10^{-6}	$\mu\mathrm{M}^{-1}\cdot\mathrm{s}^{-1}$	
$k_0^{ ext{d-}}$	Dimer dissociation rate	0.041	s^{-1}	
$k_0^{\mathrm{t+}}$	Trimer formation rate	13×10^{-5}	$\mid \mu \mathrm{M}^{-1} \cdot \mathrm{s}^{-1} \mid$	
$k_0^{ ext{t-}}$	Trimer dissociation rate	22	s^{-1}	
$k_0^{\rm b+}$	Barbed end addition rate	11.6	$\mu \mathrm{M}^{-1} \cdot \mathrm{s}^{-1}$	
$k_0^{ ext{b-}}$	Barbed end dissociation rate	1.4	s^{-1}	
$k_0^{\mathrm{p+}}$	Pointed end addition rate	1.3	$\mu \mathrm{M}^{-1} \cdot \mathrm{s}^{-1}$	
k_0^{p-}	Pointed end dissociation rate	0.8	s^{-1}	

Table 1: Parameter values for filament nucleation and growth [28].

1.1.1 Program of study

While our main focus in this study is on branched networks, it is necessary to first establish a foundation with linear filaments, which we do in Section 2. We incorporate the most complete picture of spontaneous nucleation developed in [28] into a standard algorithm of filament assembly. We first simulate the dynamics of spontaneous filament assembly, then incorporate formins to speed the nucleation phase. Adding profilin slows the nucleation phase, but speeds up the growth phase; we consequently find an optimum profilin concentration at which the speed of polymerization is optimal.

After establishing foundations with linear filaments, we move on in Section 3 to our main focus in this study: branched networks. We find that a high arp 2/3 concentration can siphon actin monomers into branches, dropping the density of linear filaments in the process. In the equilibrium phase (after all the actin is polymerized), the branch density typically continues to increase as new branches form; this increase is blocked with profilin. Adding formin, we find that we cannot impact the density of branches unless we include profilin, because profilin-actin preferentially associates to formin-bound filaments.

2 Stochastic simulations of filamentous actin

The process of actin polymerization can be separated into a nucleation step and elongation step. In the nucleation step, two actin monomers come together to form a dimer, which then must be joined by another monomer to form a trimer. A trimer has a barbed and pointed end, and monomers can associate at those ends with the same rate as longer filaments. The reactions in the system are therefore: dimer formation and breakup, trimer formation and breakup, and polymerization and depolymerization from the barbed/pointed end of each filament. Table 1 lists the parameters, which come from the recent study [28].

As already mentioned in the introduction, we do not explicitly track "filaments" (really nucleates) of length 3 or smaller, instead only keeping a count of how many are in the system. Once a tetramer forms, we insert it randomly into the spatial domain and start to track its location explicitly. This allows us to retain some spatial dependence (which will be necessary later when we try to mimic *in vitro* experiments) while not wasting computational time on steps we know are well mixed. Thus, in our system, the key unknowns are the number of monomers, dimers, and trimers, and the length (and coordinates in space) of each filament.

We also incorporate actin binding proteins into our simulations. We allow for two types of proteins: one which binds to actin monomers in solution, and then is kicked off when the actin incorporates into filaments (e.g., profilin, thymosin- β 4, TTH-1), and another which binds to the barbed ends of existing filaments (e.g., formin, capping protein). These barbed-binding proteins can also nucleate new filaments. With this information, simulations become defined by a set of binding rates k_{ij} , which is the rate at which monomers bound to protein i bind to fibers whose barbed ends are bound by protein j, and another set of unbinding rates k_j , which is the rate at which fibers with barbed end bound by protein j depolymerize. Section A.1 gives more details on the implementation of this stochastic reaction scheme.

2.1 Simulations of actin polymerization without spatial dynamics

We now perform stochastic simulations of actin polymerization using the algorithm in Section A.1. In these simulations, we do not simulate the diffusion of the fibers, as there is no reason to do so if we are not interested in spatial information. We consider systems with various concentrations of monomers (2,5, and 10 μ M) in a box of size 125 μ m³ (5 μ m on each side).

Figure 1 shows the results of our simulations. The dynamics can be divided into a "lag phase" for nucleation (which is longer for smaller concentrations), followed by a "growth phase" (in which free actin is incorporated into the existing fibers), terminated by an "equilibrium phase." At the end of the lag phase, the number of fibers per volume reaches a steady state, so that in the growth phase the decrease in free actin is driven exclusively by incorporation into existing fibers. After

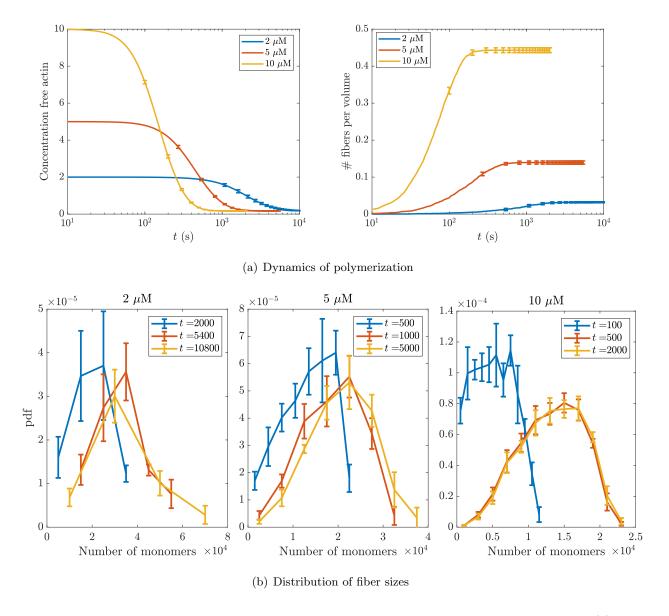


Figure 1: Stochastic simulations of polymerization dynamics at three different concentrations. (a) The dynamics of polymerization, with the concentration of free monomers shown at left and the number of fibers per unit volume $(1 \mu m^3)$ shown at right. The number of fibers comes to a steady state when about half the actin is depleted. (b) The distribution of fiber sizes (in terms of the number of monomers) at three different time points for each concentration.

the growth phase is complete, the concentration of free actin remains constant at 0.17 μ M, and the fibers must exchange monomers with other fibers to grow in length [14].

Similar to [2], we do not concern ourselves with the long-time "steady state" distribution of actin filament lengths, which can take days to reach [14, 10, 21], and instead focus on the length distributions on an intermediate timescale of a 30 minutes to 3 hours, which corresponds to a typical timescale for *in vitro* experiments. Our results for the filament length distributions, which are shown in Fig. 1(b), match those of [2] for spontaneous nucleation of filaments: intermediate timescales show a quasi-stable length distribution which is skewed to the left and quite broad (c.f. [2, Fig. 2c]), and higher concentrations have lower mean filament lengths (since more filaments are nucleated faster, the filaments that are nucleated first do not suck up as many monomers prior to nucleation of other filaments).

2.2 Incorporating formins

We now incorporate formin into the filament dynamics. We assume that formin can do three things,

- 1. Formin can nucleate filaments by binding two actin monomers in the reaction $F + A + A \rightarrow FA_2$, which is irreversible and has rate $k_{\text{for}}^{(n)}$ [23, 37]. Following previous studies [23, 37], we do not allow formin nucleates to break apart.
- 2. Fibers with formin at the barbed end add monomers with a modified rate $\alpha_{\rm for}k^{\rm b+}$ at the barbed end (the pointed end rate and unbinding rates are unchanged). In the absence of profilin, $\alpha_{\rm for}=0.5$.
- 3. Formin can bind and unbind from existing filament barbed ends, with rates k_{for}^+ and k_{for}^- respectively.

We choose to model the mammalian formins mDia1 and mDia2, which are able to nucleate filaments from free actin monomers. For these formins, association of free monomers to barbed ends with bound formin takes place at roughly 50% the rate it does without formin [15]. The key rates are reported in Table 2, and Appendix A.1.2 reports how we translate these rates into inputs to the simulation.

2.2.1 Results with formin

We now consider how adding formin affects the dynamics of actin polymerization at a fixed actin concentration of 5 μ M. Figure 2 shows the concentration of free monomers, fraction of free formins,

Parameter	Description	Value	Units	Ref
$k_{ ext{for}}^{(n)}$	Filament nucleation by formin	2×10^{-4}	$\mu\mathrm{M}^{-2}\cdot\mathrm{s}^{-1}$	[23, Tab. S1]
k_{for}^+	Formin association to barbed ends	29.1	$\mu M^{-1} \cdot s^{-1}$	[30]
k_{for}^-	Formin dissociation from barbed ends	5×10^{-4}	s ⁻¹	[15]
$\alpha_{ m for}$	Growth for formin barbed ends (no profilin)	0.5		[15]

Table 2: Parameter values for formin. We choose an average of values reported in literature for the mammalian formin mDia1 and mDia2 (with the exception of nucleation which comes from budding yeast data Bni1p.)

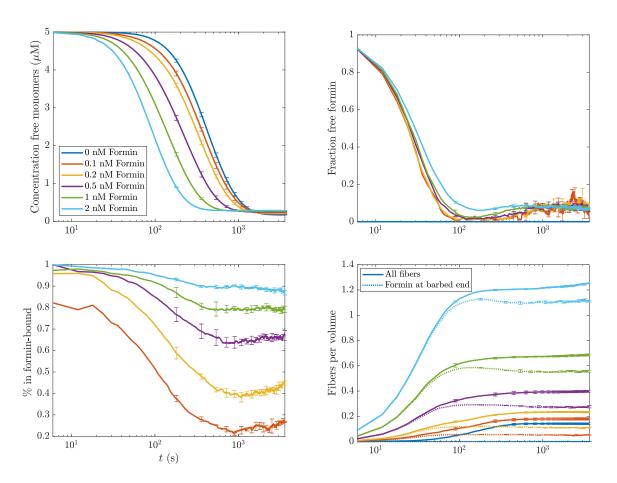


Figure 2: Dynamics of actin polymerization (at 5 μ M) when formin nucleates actin filaments, and barbed end growth with formin attached occurs at 50% of the free-ended rate. Top left: the concentration of free actin monomers over time. The colors correspond to different formin concentrations. Top right: the fraction of free formins over time. Bottom left: the percentage of fibers that have a formin-bound barbed end. Bottom right: the number of fibers per unit volume. The solid lines give the total density of filaments per volume, while the dotted lines give the density of formin-bound filaments.

and fiber density (with and without bound formins) from simulations lasting 3600 s, in which we include formin at concentrations varying from 0 to 2 nM. In all cases, we see how formin aids the process of nucleation; at time shorter than 60 seconds, almost all of the filaments created are formin-bound, regardless of the formin concentration. As time goes on, the percentage of formin-bound filaments decreases in proportion to the amount of formin in the system, as spontaneous nucleation gives filaments without bound ends. There is then a long-time steady state where formin assumes an equilibrium with about 7% unbound. The long time steady state demonstrates that the number of fibers is large enough in all cases for formin to assume an equilibrium state (there is never an excess of formins, because at higher formin concentrations more filaments are nucleated).

At low (\leq 0.2 nM) formin concentrations, the dynamics of polymerization closely resemble the spontaneous nucleation case; we see that it takes about 1500 s to reach the steady state concentration of free monomers. As we increase the formin concentration, the nucleation phase is faster (bottom right plot), going from about 300 s for no formin to 100 s with formin. The faster nucleation phase occurs because more filaments are nucleated. The nucleation of additional filaments leads to a faster growth phase, since there are more barbed ends available to bind actin monomers. Thus, despite the fact that formin-bound barbed ends grow more slowly than free barbed ends, the additional barbed ends still lead to a faster growth phase. If we define the growth phase as the time it takes the free monomer concentration to go from 90% to 10%, we obtain a time of $252 - 30 \approx 200$ s with 2 nM formin, compared to $1080 - 144 \approx 1000$ s without formin. With 2 nM formin, there are ten times as many filaments, but each filament grows at half the speed, which gives a factor of five difference in the dynamics.

How do these dynamics affect the length distribution of fibers? As shown in Fig. 3, and as previously reported in [2], an increase in the formin concentration leads to the nucleation of more filaments at early times, which translates to a smaller mean and standard deviation in the number of monomers per filament. The slower speed of formin-filament growth leads to a bimodality in the distribution of filament lengths; this is especially clear in the case of 0.1 nM formin, where there is a peak at 2×10^4 which matches the non-formin-bound fibers, along with a second peak around 1×10^4 monomers for formin-bound fibers. When the formin concentration increases, we typically observe a large peak associated with formin-bound fibers. This peak tends to lose mass over time, with growth occurring for larger numbers of monomers. This reflects the theoretical expectation that the slower growing fibers (the formin-bound ones) will give up their actin over long times to the faster-growing (non formin-bound) ones [21].

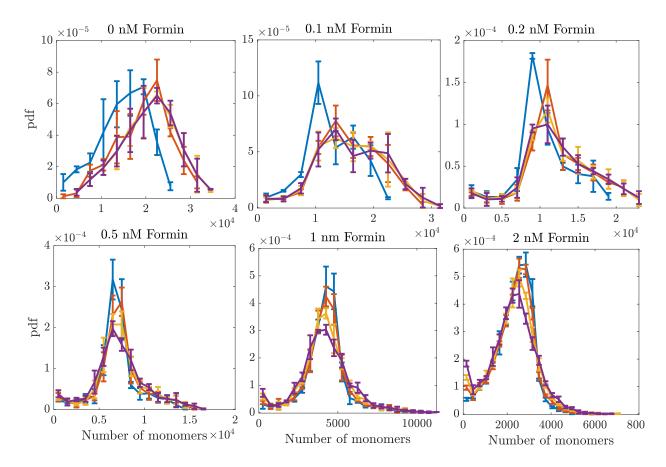


Figure 3: Length distribution for actin filaments in simulations with formin of varying concentrations. We show the distributions of fiber sizes (in terms of the number of monomers) at four time points (t = 600 s in blue, t = 1200 s in red, t = 2400 s in yellow, and t = 3600 s in purple). The distributions become more uniform (smaller range) as the formin concentration increases.

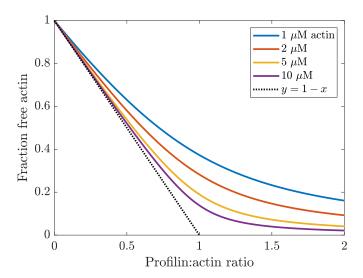


Figure 4: How the profilin: actin equilibrium affects the fraction of free actin. Here the x axis shows the ratio of profilin to actin, while the y axis shows the fraction of free monomers that result. Because of nonlinearities in the equilibrium, the curves are different for different actin concentrations. The dotted black curve shows the result if all the profilin binds to actin.

2.3 Explicitly incorporating profilin

Profilin is a protein which binds actin monomers in solution. The binding reaction has forward rate $45 \,\mu\text{M}^{-1}\cdot\text{s}^{-1}$ and reverse rate $10 \,\text{s}^{-1}$. At typical profilin concentrations of order $1 \,\mu\text{M}$, this reaction is much faster than the other reactions in the system, which means we can approximate it as in instantaneous equilibrium. Appendix A.1 discusses how we do this; in Fig. 4 we show the resulting fraction of free actin as a function of the profilin actin ratio. At low profilin concentrations, all of the profilin binds to actin, while at higher concentrations the equilibrium is more complex.

In our simulations, monomers bound to profilin cannot nucleate new filaments (dimers), regardless of whether or not they are assisted by formin [24]. Without formin, profilin-bound monomers bind barbed ends at a rate 80% of the free monomer value, while with formin the binding rate is 300% of the free monomer value (these represent average rates for mDia1 and mDia2 from [15]). Binding at the pointed end with profilin-bound monomers is small (10% of the free monomer value). See Appendix A.1.2 for how these rates are incorporated into simulations.

We begin our study by looking at how profilin affects polymerization without formin. Figure 5 compares the dynamics of polymerization of 5 μ M actin with varying profilin concentrations in the range 0 to 5 μ M. Low profilin concentrations do not show too much of a change from the spontaneous

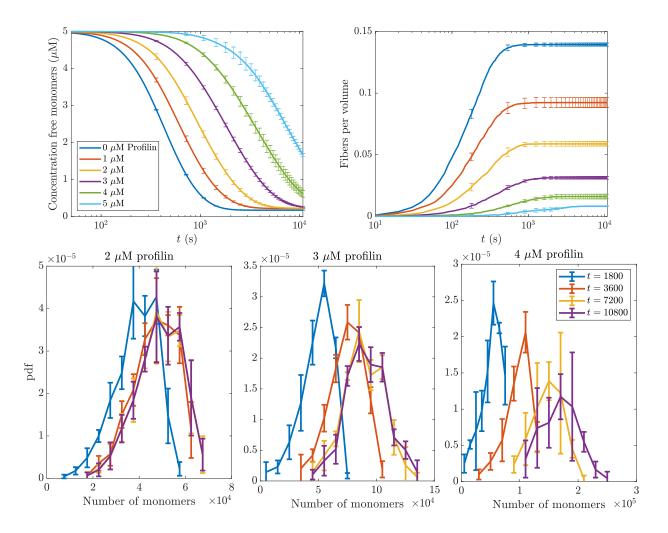


Figure 5: Dynamics of actin polymerization (at 5 μ M) with profilin. Top left: the concentration of free actin monomers over time. The colors correspond to different profilin concentrations. Top right: the number of fibers per volume over time. Bottom: fiber length distributions over time for 2 μ M profilin (left), 3 μ M profilin (middle), and 4 μ M profilin (right).

nucleation curves, since at 1 μ M profilin concentration we only expect about 20% of the monomers to be bound to profilin. Increasing the profilin concentration has drastic effects; once we reach 5 μ M profilin, the growth phase is not complete by t=3 hours. The slowdown in polymerization is driven by two factors: first, profilin has a reduced rate of binding for both pointed and barbed ends, which slows down the rate of polymerization. Second and more importantly, increases in profilin lead to less filaments being nucleated, which means there are less barbed ends available to bind monomers. The length distributions for each profilin concentration show another effect of having less filaments; each filament tends to be longer, and the distributions are more spread out. The reason for the latter is that there is no longer a big nucleation event that spawns many filaments. There are only a few filaments in the system, and the length of each filament is directly correlated with the time since it was nucleated.

2.3.1 Formin (high concentration) with profilin

We now consider the effect of adding formin at concentration 1 nM. At this formin concentration, we previously saw (Fig. 2) that about 80% of the fibers are bound to formin in the absence of profilin. Figure 6 shows the results of polymerization with formin in the presence of profilin. There is an interesting trade-off in the polymerization dynamics: on the one hand, profilin inhibits nucleation by formin, so the nucleation phase is slower with profilin. On the other hand, once filaments are nucleated, profilin-actin binds at a faster rate to the barbed end than actin without profilin. This process is auto-catalytic: as more profilin-actin binds to fibers, profilin is released, and there is more of it in solution, which means more of the actin monomers are profilin bound. As a result, low profilin concentrations of 1–3 μ M have dynamics which are faster than actin without profilin.

Unlike for the case without formin, with formin the behavior with increasing profilin concentration is non-monotonic: at higher profilin concentrations (4–5 μ M), the increased growth rate at barbed ends cannot compensate for the smaller number of fibers, and the dynamics are slower. The smaller number of fibers occurs because formin has to nucleate filaments from *free* actin monomers, which are less abundant as the profilin concentration increases. Indeed, when profilin concentrations are high, there are substantially less filaments, all of which are formin-bound (and therefore elongate rapidly). Unlike for the case without profilin, with profilin formin is in excess, and the number of free formins increases as the profilin concentration increases. Decreasing the formin concentration to 0.1 nM gives the same trends as for 1 nM formin concentration (results not shown).

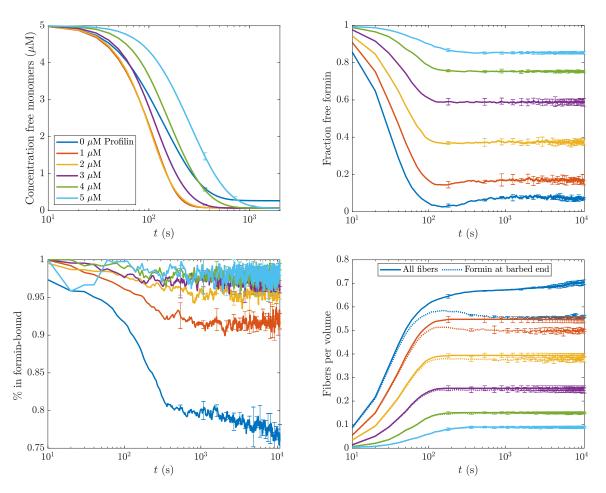


Figure 6: Dynamics of actin polymerization (at 5 μ M) with 1 nM formin and various profilin concentrations. Top left: the concentration of free actin monomers over time. The colors correspond to different formin concentrations. Note the non-monotonic dependence of overall monomer depletion on profilin concentration. Top right: the fraction of free formins over time. Bottom left: the percentage of fibers that have a formin-bound barbed end. Bottom right: the number of fibers per unit volume. In addition to the colors representing different concentrations, the solid lines give the total density of filaments per volume, while the dotted lines give the density of formin-bound filaments.

2.4 Summary of results

Summing up our results from this section on linear filaments, we found the following

- For spontaneous nucleation, the number of fibers increases (roughly quadratically) with the actin concentration. Because of the limited pool, higher concentrations have shorter fibers (since there are more fibers to suck up the monomers).
- Adding formin accelerates nucleation and polymerization by producing more filaments. Without profilin, the number of additional fibers is linearly proportional to the formin concentration. Consequently, the fraction of free formin is independent of the formin concentration.
- With profilin, we slow down nucleation and (with formin attached) accelerate the growth phase. This gives rise to non-monotonic dynamics: at low profilin concentrations, the nucleation phase is not substantially affected, but the growth phase is substantially faster, and polymerization is overall faster. At higher profilin concentrations, the accelerated growth phase cannot make up for a slow nucleation phase, and the dynamics are overall slower. At high profilin concentrations, formin is in excess, as it cannot nucleate filaments when profilin is bound to actin.

In general, we see that free monomers at concentration 5 μ M take about 10³ s to reach the equilibrium phase. Adding 4 μ M profilin, we do not reach the equilibrium phase until beyond 10⁴ s. When we add formin, the equilibrium phase is reached in about 500 s (rescuing the original timeline). These qualitative results agree with older bulk-actin polymerization curves in [16, Fig. 2] and [13, Fig. 6].

3 Stochastic simulations of branched filaments

We now discuss our model of filament branching. Similar to previous work [22], we consider the binding of arp 2/3 complex to a linear filament, which results in the formation of a branch whose tangent vector makes a seventy degree angle with the mother filament [25]. Previous work [4] has shown that the kinetics of actin monomers binding VCA-bound arp 2/3 complex are much faster than this complex (with or without a monomer attached) binding to an existing filament. Therefore, letting B be a VCA-bound arp 2/3 complex and N_j a filament of any size, the reaction

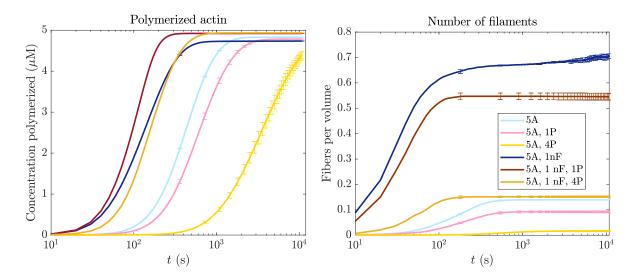


Figure 7: Summary of the role of formin and profilin for linear fiber growth. We show the base 5 μ M actin in light blue, with 1 μ M profilin in pink and 4 μ M profilin in yellow. The darker colors then show the same concentrations, but with 1 nM formin.

proceeds in two steps

$$B + N_1 \stackrel{K_{\text{fast}}}{\longleftrightarrow} BN_1 \qquad \text{(fast)}$$
 (1a)

$$BN_1 + N_j \underset{k_{\text{slow}}}{\overset{k_{\text{slow}}^+}{\rightleftharpoons}} BN_1 N_j \qquad \text{(slow)}$$

$$(1b)$$

To simplify this, we assume the fast reaction is in equilibrium and coarse grain this process into a single reaction involving three reactants: (VCA-bound) arp 2/3 complex B, an actin monomer N_1 , and an actin filament N_j ,

$$B + N_1 + N_j \underset{k^{\text{br-}}}{\overset{k^{\text{br+}}}{\rightleftharpoons}} BN_1 N_j. \tag{2}$$

The negative rate for this reaction is obtained from [4],

$$k^{\text{br-}} = k_{\text{slow}}^- = 3.4 \times 10^{-3} \,\text{s}^{-1},$$

while the positive rate can be obtained directly from the *in vitro* TIRF data in [31]. In that study, which uses budding yeast (for which branching is possible without bound VCA), the rate of branching is measured in units of $\mu M^{-1} \cdot s^{-1}$ per μm of mother filament. Because the branching occurs in a background of 1 μM monomeric actin, we can extrapolate the rate to the correct units by dividing by the background concentration to obtain

$$k^{\rm br+} \approx 5 \times 10^{-3} \,\mu \mathrm{M}^{-2} \cdot \mathrm{s}^{-1} / (\mu \mathrm{m} \text{ of mother filament}).$$

This rate is an approximation, and represents a midpoint between the rate in the presence of VCA (9.7×10^{-3}) and the rate in its absence (2.5×10^{-3}) . Because each monomer occupies an average of $2 \times 10^{-3} \mu m$, the rate we use is $10^{-5} \mu M^{-2} \cdot s^{-1}/(monomer of mother filament)$.

To simplify the process of polymerization and depolymerization, we implement the following rules for branching:

- 1. A filament is eligible to form a branch when it has $N_{\rm mb} = 4$ monomers or more. Once a branch forms, it forms on the fourth monomer or higher (the first three monomers cannot form branches).
- 2. The pointed end of the mother cannot unbind if a branch is sitting on the fourth monomer.

 This always leaves two monomers between the pointed end and the first branch.
- 3. The barbed end (of a mother or a branch) cannot unbind if another branch is attached to it.
- 4. Branches can only unbind if they have length 1, in which case arp 2/3 also detaches from the filament (reverse of reaction (2)). "Barbed ends" alone cannot unbind if branches have length 1 (the only unbinding is through the reverse of reaction (2)).

This set of rules ensures simplicity in the construction and deconstruction of branched structures (in essence, the structure has to depolymerize in the order it was constructed, and cannot break into multiple branched structures). While these rules are certainly a simplification, our goal here is to study the number of branches per unit length, so whether one structure with 40 branches should in fact be two structures with 20 branches is of little consequence to us.

3.1 Branching with spontaneous nucleation

We begin by looking at the branching process with spontaneous nucleation of actin filaments only (no formins). We fix the actin concentration at 5 μ M and vary the concentration of arp 2/3 across two orders of magnitude (2 nM, 20 nM, and 200 nM). The dynamics of polymerization and assembly are shown in Fig. 8, and the filament length distributions are shown in Fig. 9.

To understand these data, it is important to remember that actin filament polymerization with spontaneous nucleation proceeds in two phases: nucleation of filaments (slow), and then elongation (fast). Figure 8 shows that the speed of branching relative to these two timescales determines the behavior. In the case of 2 nM arp 2/3, significant branching does not appear until $t \approx 400$ s, at which time the nucleation phase is already complete. Thus in this case most of the actin

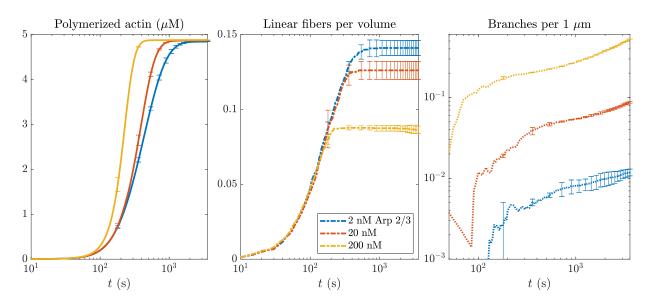


Figure 8: Course of polymerization and assembly with 5 μ M actin concentration and various concentrations of arp 2/3 (2 nM in blue, 20 nM in red, and 200 nM in yellow). From left to right, these plots show the total amount of polymerized actin, the number of linear fibers (including mother filaments) per volume, and the number of branches per length.

sits in linear filaments, and there is a small number of branches (0.01) per μ m. Increasing to 20 nM arp 2/3 and then 200 nM arp 2/3, the timescale of branching gets shorter, until it starts to compete with nucleation. At the highest arp 2/3 concentration, there are substantially fewer linear filaments, because branching on the filaments that do form takes up the available actin.

Further understanding of the dynamics can be gained by looking at the length distributions in Fig. 9, where we have made separate distributions for linear fibers (including mothers of branches) in the top set of plots, and branches in the bottom set. For 2 nM Arp 2/3 concentration, the length distributions look similar to the case without arp 2/3, where a quasi-steady state is achieved at roughly 1000 s that has a mean of 2×10^4 monomers. The branches sit in a rough steady state here, reflecting an equilibrium where monomers typically exchange between existing filaments. For 20 nM and 200 nM Arp 2/3 concentration, the quasi-steady state in the linear fibers occurs at an early time, and they do not elongate past 2×10^4 monomers. It is instead branches which form, in particular new branches, as the distribution of branch lengths becomes more and more concentrated on smaller branches over time. This is in line with our previous observation that branching is faster than elongation for these concentrations. The formation of new branches at later times demonstrates that there is an excess of arp 2/3 which looks for available monomers to form new branches.

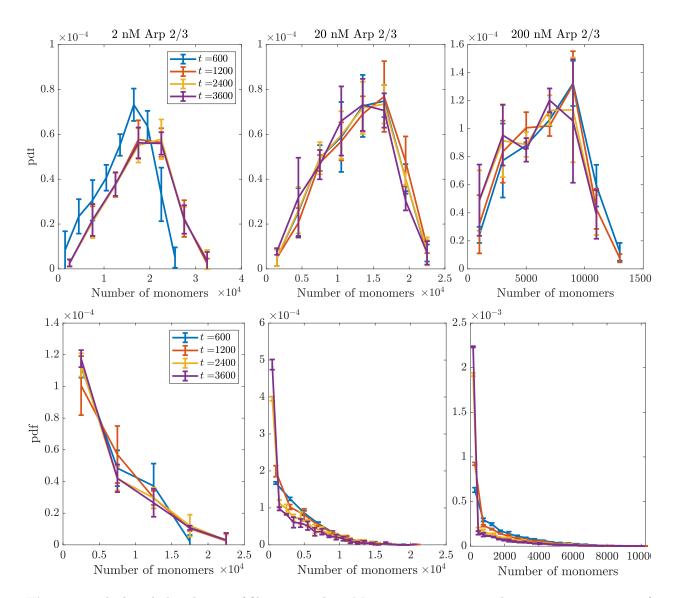


Figure 9: The length distribution of filaments with 5 μ M actin concentration and various concentrations of arp 2/3 (2 nM at left, 20 nM in the middle, and 200 nM at right). We separate the filaments into linear and mother filaments, and show their length distribution at the top, and branches, with their length distribution at the bottom.

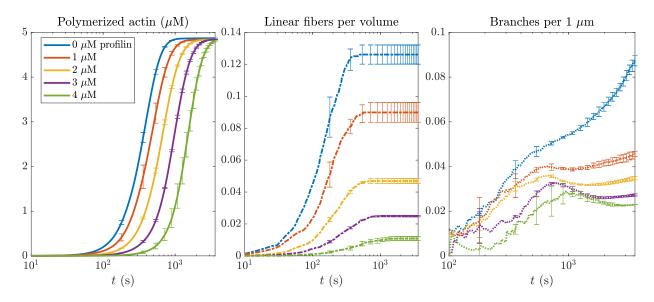


Figure 10: Course of polymerization and assembly with 5 μ M actin concentration, and varying profilin concentration with 20 nM arp 2/3 complex. From left to right, these plots show the total amount of polymerized actin, the number of linear fibers (including mother filaments) per volume, and the number of branches per length. With profilin, the number of branches does not grow in the equilibrium phase.

3.2 Incorporating profilin

Because profilin competes with VCA for binding actin monomers, and because VCA is (in most systems) required for arp 2/3 complex to bind an actin monomer and form a branch, we assume that branching with profilin-bound monomers is impossible. With this assumption, Figure 10 shows the effect of adding profilin to a system of polymerizing actin and arp 2/3 complex.

Our observations for the branching process with profilin are similar to what we saw with linear filaments (Fig. 5). Increasing the profilin concentration slows down the nucleation and growth phase, since there are substantially less fibers nucleated. The branching process with profilin is significantly more interesting. When the profilin concentration is slow $(1-2 \mu M)$, the branching rate initially $(t \le 400 \text{ s})$ is comparable to what it is without profilin. As more of the actin gets polymerized, the branching rate slows down substantially, because more of the actin is bound to profilin and so new branches cannot get nucleated. In systems without profilin, the branch density grows in the equilibrium phase, since free monomers can form new branches (especially when a lot of arp 2/3 is available). But with profilin, the branch rate is constant in the equilibrium phase; because the free actin concentration is low, almost all of it is bound to profilin, and no new branches can form.

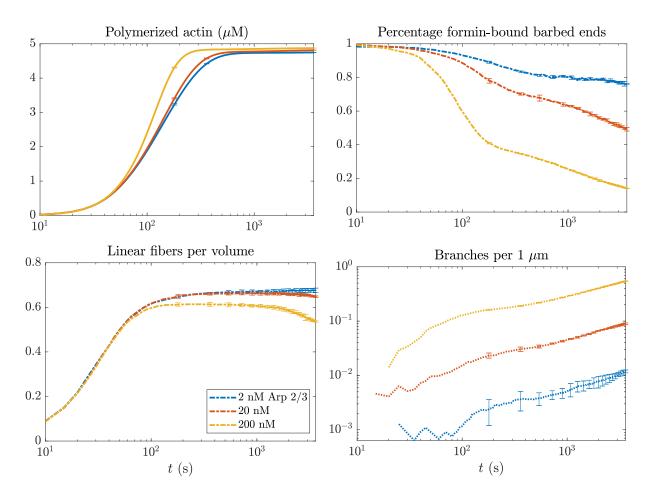


Figure 11: Course of polymerization and assembly with 5 μ M actin concentration, 1 nM formin concentration, and three arp 2/3 concentrations. The top plots show the concentration of polymerized actin and the percentage of formin-bound barbed ends. The bottom plots show the dynamics of fiber and branch growth.

3.3 Incorporating formins

We now incorporate formin into the branching dynamics. As in linear filaments, formin can bind to the barbed end of branched filaments and modify the growth rate there.

3.3.1 Without profilin

We first consider 5 μ M actin and 1 nM formin with varying concentrations of arp 2/3, plotting the dynamics in Fig. 11. As we have already seen in Section 2.2, formin speeds up the nucleation phase and produces more filaments. In Fig. 11, we see that, regardless of the arp 2/3 concentration, there are initially about 0.6 fibers per μ m³ produced. What happens next depends on the concentration of arp 2/3 complex. When there is a small amount of arp 2/3, there are a small amount of

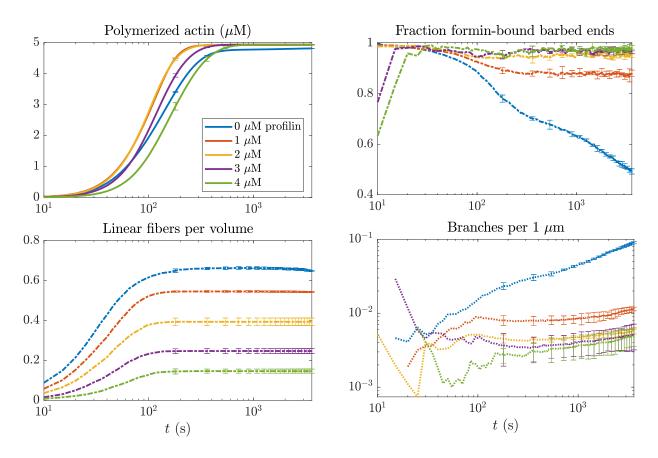


Figure 12: Course of polymerization and assembly with 5 μ M actin concentration, 1 nM formin concentration, 20 nM arp 2/3 concentration, and varying profilin concentration. The top plots show the concentration of polymerized actin and the percentage of formin-bound barbed ends. The bottom plots show the dynamics of fiber and branch growth.

branches that form, and the amount of linear filaments is roughly constant in time. At high arp 2/3 concentrations, a significant amount of branch formation occurs, so that the amount of linear filaments per volume drops. Because there are more branches at higher arp 2/3 concentrations, the percentage of barbed ends with a bound formin drops significantly. Finally, although there are more filaments in the case with formin, the total length of polymerized actin is the same, and so the number of branches per 1 μ m does not change from the case without formin (c.f. Fig. 8).

3.3.2 With profilin

Based on previous results, we now consider adding in profilin to a system with 1 nM formin and 20 nM arp 2/3. Without profilin, we saw that adding formin made little difference in the branch density; will this also be true when we add in profilin?

Figure 12 shows the results for a system that includes formin, arp, and profilin (everything). As we have seen previously, adding profilin reduces the number of fibers and number of branches. This means there are less barbed ends available, and so a higher fraction of the formin is bound to the barbed ends. As a result of profilin accelerating polymerization, the polymerization dynamics with profilin concentrations 1–3 μ M are again faster than without profilin, which is the same conclusion we reached for linear fibers (see Fig. 6). But the faster dynamics have strong consequences for branching: for a fixed parameter set, the number of branches without formin (Fig. 10) is substantially *higher* than with formin. This is because the balance of rates has shifted; with profilin, polymerization of linear filaments is faster, and so more of the actin goes there than into the branches.

3.4 Summary of results

Our results on branching can be summarized as follows

- Branching occurs after the assembly of linear filaments. When the arp 2/3 concentration is high enough, the formation of new branches can compete with linear filaments, and there are less linear filaments at the end of the nucleation phase.
- Without profilin, new branches form in the equilibrium phase, and the number of branches constantly increases. With profilin, the number of branches remains constant in the equilibrium phase, because the free actin is all bound to profilin, and actin-profilin cannot form branches.
- With formin, there are more linear fibers per volume, but the branch density is unchanged from the case without formin (because the amount of polymerized actin is the same). This conclusion changes when we add profilin: because profilin-actin goes into formin-bound fibers at a faster rate, the number of branches drops relative to the case without formin. Thus, if we wanted to use formin to reduce the branch frequency, we need to add profilin as well. (This would be a nice experiment to try.)

Appendix A Algorithm and validation

This appendix contains the details of our algorithm and the validation. For validation, we use ODEs, and close the system by assuming a maximum of five monomers per fiber.

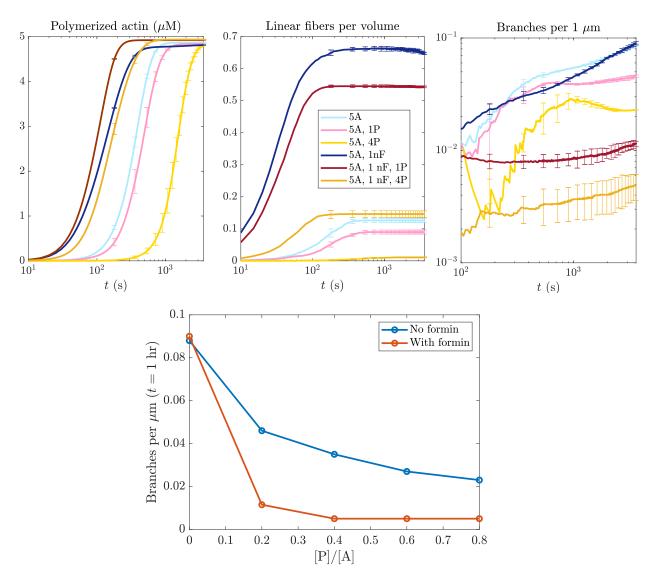


Figure 13: Summary of the role of formin and profilin for branched fiber growth (20 nM arp 2/3 in all cases). Top plots: We show the base 5 μ M actin in light blue, with 1 μ M profilin in pink and 4 μ M profilin in yellow. The darker colors then show the same concentrations, but with 1 nM formin. Bottom plot: the number of branches per μ m after one hour as a function of profilin concentration without (blue) and with (red) formin.

A.1 Stochastic simulation algorithm for linear filaments

To keep our algorithm general, let us suppose that we have P_M proteins which bind to monomers in solution, and are kicked off of the monomers once they bind to fibers. We also assume that there are P_B proteins which bind barbed ends of actin filaments. Letting M be the total number of free actin monomers, and $M^{(0)}$ the amount of free actin monomers attached to no protein. The equilibrium reactions.

$$M^{(0)} + P^{(i)} \stackrel{\kappa^{(i)}}{\longleftrightarrow} M^{(i)} \tag{3a}$$

$$M^{(i)} = \kappa^{(i)} M^{(0)} \left(P^{(i)} - M^{(i)} \right) \to M^{(i)} = \frac{\kappa^{(i)} M^{(0)} P^{(i)}}{1 + \kappa^{(i)} M^{(0)}}$$
(3b)

dictate the number of monomers attached to protein $i \geq 1$. In the case when there is only one monomer, the equation

$$M^{(0)} + M^{(1)} = M^{(0)} + \frac{\kappa^{(1)} M^{(0)} P^{(1)}}{1 + \kappa^{(1)} M^{(0)}} = M$$
(3c)

can be solved to obtain a closed-form solution for the amount of free monomers

$$M^{(0)} = \frac{1}{2\kappa^{(1)}} \left(-\left(1 + \kappa^{(1)} \left(P^{(1)} - M\right)\right) + \sqrt{\left(1 + \kappa^{(1)} \left(P^{(1)} - M\right)\right)^2 + 4\kappa^{(1)}M} \right).$$
 (3d)

Thus, prior to calculating any reaction rates, the first step of our reaction algorithm is to compute the number of monomers attached to each protein using (3).

For our reactions, we let $k_{ij}^{\rm d+}$ be the rate of dimer formation for monomers bound to protein i (starting from 0, to include no bound protein), and barbed end bound to protein j (starting from 0, again to include no protein). Likewise, $k_{ij}^{\rm t+}$ is the rate of trimer formation, and $k_{ij}^{\rm f+}$ is the rate of addition to fibers (length 3 or more). The rate of addition to fibers is broken into $k_{ij}^{\rm f+} = k_i^{\rm p+} + k_{ij}^{\rm b+}$, since the pointed end rate is assumed to depend only on the protein bound to the monomer, while the barbed end addition rate depends on both bound proteins. The rates of unbinding $k_j^{\rm f-} = k^{\rm p-} + k_j^{\rm b-}$ are a function of the barbed-bound protein only. Finally, we let $k_j^{\rm on}$ and $k_j^{\rm off}$ be the rates of binding and unbinding of the barbed end proteins.

Given this notation, the next reaction method proceeds as follows: at the beginning of each step, we compute the total number of monomers bound to each protein via (3). Then, we sample a time for each of the following reactions to occur:

1. Spontaneous dimer formation. For all i and j=0, dimer formation occurs with rate $(k_{i0}^{d+}/V)M^{(i)}M^{(i)}$.

- 2. Barbed-bound-protein-induced nucleation. For all i and j > 0, dimer formation occurs with rate $(k_{ij}^{d+}/V^2)M^{(i)}M^{(i)}B^{(j)}$.
- 3. Dimer break-up. For all j, each individual dimer breaks up with rate $k_j^{\text{d-}}$.
- 4. Trimer formation. For all i and j, each individual dimer becomes a trimer with rate $(k_{ij}^{t+}/V)M^{(i)}$.
- 5. Trimer breakup. For all j, each individual trimer breaks up with rate $k_j^{\text{t-}}$.
- 6. Fiber growth. For each individual fiber of length 3 or more, fiber growth occurs with rate $(k^{f+}/V)M^{(i)}$, where $k^{f+} = k_i^{p+} + k_{ij}^{b+}$.
- 7. Fiber shrinkage. For each individual fiber of length 4 or more with barbed bound protein j, removal of monomers occurs at the barbed end with rate $k_j^{\text{b-}}$ and at the pointed end with rate $k_j^{\text{p-}}$ (this is actually two reactions).
- 8. Barbed protein binding. Barbed protein j binds to each fiber of length 2 or more with rate $k_j^{\text{on}} B^{(j)}$.
- 9. Barbed protein unbinding. Barbed protein j unbinds from each fiber that it is attached to with rate k_j^{off} .

We use an exact Gillespie simulation algorithm [11] to simulate these reactions. If a reaction has rate r_i , we sample a time for it to occur according to an exponential distribution, $\Delta t_i = -\log(u)/r_i$, where $u \in (0, 1]$. The reaction chosen is then the one that occurs soonest. Once a reaction is chosen, we increment time by this minimum amount and process the reation by changing the state of the system to reflect its outcome (e.g., unbinding from the barbed end removes one monomer from a filament and adds it to the free monomer count). We then repeat this entire process (including the calculation of bound monomers in (3)) until the time step is complete. Because this is an exact simulation algorithm, the time step Δt is arbitrary and is not a restriction on accuracy.

To efficiently implement these reactions, we take an approach where we track a number of fiber nucleates, which we denote as $B_m^{(j)}$ for m = 2, 3. Once fibers reach a size of four, we create an object in space to track each of them individually.

A.1.1 Algorithm validation

We validate our stochastic simulation algorithm by comparing to ODEs. To facilitate this, we must cap the length of fibers at a maximum size M, which we set to M = 5. To formulate the ODEs, we begin by defining the rate at which monomers become dimers with the barbed end bound with protein j. Given the rates and reactions defined in Section A.1, this rate is given by

$$R_2^{(0)} = \sum_{i=0}^{P_M} \left(k_{i0}^{d+} M^{(i)} M^{(i)} \right) - k_0^{d-} B_2^{(0)} \qquad R_2^{(j>0)} = \sum_{i=0}^{P_M} \left(k_{ij}^{d+} M^{(i)} M^{(i)} B^{(j)} \right) - k_j^{d-} B_2^{(j)}$$
(4a)

The rate at which dimers become trimers is given by

$$R_3^{(j)} = \sum_{i=0}^{P_M} \left(k_{ij}^{t+} M^{(i)} B_2^{(j)} \right) - k_j^{t-} B_3^{(j)}, \tag{4b}$$

and the rates at which fibers of larger sizes are formed are

$$R_m^{(j)} = \sum_{i=0}^{P_M} \left(k_{ij}^{f+} M^{(i)} B_{m-1}^{(j)} \right) - k^{f-} B_m^{(j)}, \qquad 4 \le m \le M.$$
 (4c)

Finally, the rate at which barbed binding proteins bind and unbind from the bound end of a fiber of size m is

$$S_m^{(j)} = k_j^{\text{on}} B^{(j)} B_m^{(0)} - k_j^{\text{off}} B_m^{(j)}. \tag{4d}$$

With these intermediates defined, the ODEs that govern the dynamics are

$$\frac{dM}{dt} = \sum_{j=0}^{P_B} \left(-2R_2^{(j)} - \sum_{m=3}^{M} R_m^{(j)} \right)$$
 (4e)

$$\frac{dB^{(j)}}{dt} = -R_2^{(j)} - S_m^{(j)} \tag{4f}$$

$$\frac{dE_{m}^{(j)}}{dt} = \begin{cases}
R_{m}^{(j)} - R_{m+1}^{(j)} - S_{m}^{(j)} & j = 0, \quad 2 \le m < M \\
R_{m}^{(j)} - R_{m+1}^{(j)} + S_{m}^{(j)} & j > 0, \quad 2 \le m < M \\
R_{m}^{(j)} - S_{m}^{(j)} & j = 0, \quad m = M \\
R_{m}^{(j)} + S_{m}^{(j)} & j > 0, \quad m = M
\end{cases} \tag{4g}$$

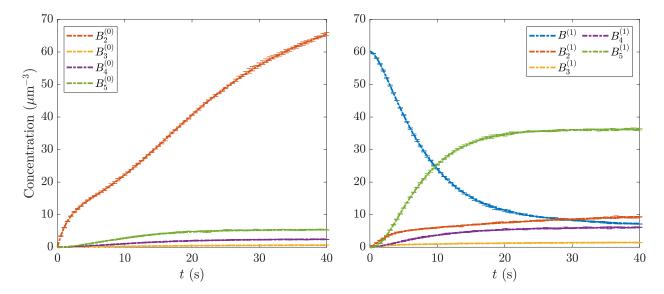


Figure 14: Validating the stochastic simulation algorithm with one monomer-binding and one barbed-binding protein. We consider an actin concentration of 2 μ M in a box of size 3 μ m on each side, and the rates given in (5). The dashed-dotted lines show the concentration (in units μ m⁻³) of each particular species $(B_m^{(j)})$ denotes a polymer of size m with barbed protein j bound to it) predicted by the ODEs (4). The solid lines with error bars show the output from the stochastic simulations. These data points show the mean over 3 sets of 10 simulations, with the error bars the standard error in the mean.

We validate our code by comparing the ODEs to stochastic simulations with the following rates

$$k_{ij}^{d+} = 3.5 \times 10^{-3} \,\mu\text{M}^{-1} \cdot \text{s}^{-1} \begin{pmatrix} 1 & 0.57 \,\mu\text{M}^{-1} \\ 1.7 & 0.29 \,\mu\text{M}^{-1} \end{pmatrix} \qquad k_{ij}^{t+} = 0.13 \,\mu\text{M}^{-1} \cdot \text{s}^{-1} \begin{pmatrix} 1 & 10 \\ 0.4 & 2 \end{pmatrix}$$

$$k_{ij}^{b+} = 1.6 \,\mu\text{M}^{-1} \cdot \text{s}^{-1} \begin{pmatrix} 1 & 2 \\ 0.2 & 2.5 \end{pmatrix} \qquad k_{i}^{p+} = 1.3 \,\mu\text{M}^{-1} \cdot \text{s}^{-1} \begin{pmatrix} 1 & 0.25 \end{pmatrix} \qquad (5)$$

$$\kappa_{1} = 5 \,\mu\text{M} \qquad P^{(1)} = 0.05 \,\mu\text{M} \qquad B^{(1)}(t=0) = 0.1 \,\mu\text{M} \qquad k_{1}^{\text{on}} = 5 \,\mu\text{M}^{-1} \cdot \text{s}^{-1} \qquad k_{1}^{\text{off}} = 8.1 \times 10^{-2} \,\text{s}^{-1}$$

$$k_{j}^{d-} = 0.041 \,\text{s}^{-1} \begin{pmatrix} 1 & 0.2 \end{pmatrix} \qquad k_{j}^{t-} = 22 \,\text{s}^{-1} \begin{pmatrix} 1 & 0.4 \end{pmatrix} \qquad k_{j}^{p-} = 0.8 \,\text{s}^{-1} \qquad k_{j}^{b-} = 1.4 \,\text{s}^{-1} \begin{pmatrix} 1 & 0.1 \end{pmatrix}$$

Figure 14 shows that the solution of the ODEs (dashed-dotted lines) matches the output from stochastic simulations, thus validating all parts of our algorithm up to the branching of filaments.

A.1.2 Rates for formin and profilin

In the language of our simulation platform, the rates with formin and profilin are

$$k_{1}^{\text{on}} = 29.1 \,\mu\text{M}^{-1} \cdot \text{s}^{-1} \qquad k_{1}^{\text{off}} = 5 \times 10^{-4} \,\text{s}^{-1}$$

$$k_{j}^{\text{d-}} = k_{0}^{\text{d-}} \left(1 \quad 0\right) \qquad k_{j}^{\text{t-}} = k_{0}^{\text{t-}} \left(1 \quad 0\right) \qquad k^{\text{b-}} = k_{0}^{\text{b-}} \left(1 \quad 1\right),$$

$$k_{ij}^{\text{d+}} = k_{0}^{\text{d+}} \begin{pmatrix} 1 \quad 0.057 \,\text{nM}^{-1} \\ 0 \quad 0 \,\text{nM}^{-1} \end{pmatrix} \qquad k_{ij}^{\text{t+}} = k_{0}^{\text{t+}} \begin{pmatrix} 1 \quad 5.5 \times 10^{4} \\ 0 \quad 2.8 \times 10^{5} \end{pmatrix}$$

$$k_{ij}^{\text{b+}} = k_{0}^{\text{b+}} \begin{pmatrix} 1 \quad 0.5 \\ 0.8 \quad 3 \end{pmatrix} \qquad k_{i}^{\text{p+}} = k_{0}^{\text{p+}} \left(1 \quad 0.1\right).$$

where the base rates are given in Table 1. Note the large rate for the trimers; we assume that formin-bound dimers become trimers with rate $\alpha_{\text{for}}k_0^{\text{b+}} + k_0^{\text{p+}}$, which is divided by $k_0^{\text{t+}}$ to obtain the proportion (α_{for} is 0.5 without profilin and 3 with profilin). We assume that formin bound dimers and trimers do not depolymerize.

A.2 Branched filaments

Given the set of rules formulated in Section 3, we can now formulate the reactions involving arp 2/3 complex. Letting D be the number of free arp 2/3 complexes, the reactions are

- 1. Branch formation. Each fiber (or branch) of length 4 or more can form a branch with rate $r = \left(k_i^{\text{br}+}/V^2\right) DM^{(i)}m$, where m is the number of monomers on the fiber. If this reaction is selected, we form a new branch with one monomer at a 70 degree angle from the mother filament (chosen at random from the branches once a particular structure is selected).
- 2. Branch unbinding. Branches of length one can come off the mother filament with rate k^{br} .
- 3. Branch polymerization. For each i, branches can polymerize from their barbed ends only with rate $\left(k_{ij}^{\rm b+}/V\right)M^{(i)}$.
- 4. Branch depolymerization. Branches can depolymerize from their barbed ends with rate k^{b} , as long as the geometry permits it (see rules above).
- 5. Barbed protein binding and unbinding. These are the same reactions as for linear fibers.

Because of the spatial complexity of branched networks, it is difficult to validate depolymerization using ODEs, as it becomes complicated to work out the probability that an unbinding event, if selected, will actually be successful. Because of this, we divide our validation work into two steps: all events *other than* depolymerization, and depolymerization alone. The first case allows us to work with ODEs, as we have done previously, while the second requires us to synthesize a specific structure and look at its depolymerization.

A.2.1 ODE-based validation (excludes depolymerization)

In the presence of branching, additional terms are required for the ODEs. We let D be the number of arp 2/3 complexes, and $D_m^{(j)}$ be a branch of length m with barbed-bound protein j. If $k^{\text{br}+}$ is a rate of branching in units of concentration⁻² × s⁻¹ per monomer of mother filament, the total rate of branch formation is

$$R_b = D\left(\sum_{i=0}^{P_M} k_i^{\text{br}+} M^{(i)}\right) \left(\sum_{j=0}^{P_B} \sum_{m=4}^M m D_m^{(j)}\right) - k^{\text{br}-} \sum_j D_1^{(j)}.$$
 (6a)

In addition to forming and being removed, branches exchange barbed proteins with rates given by (4d), and the flux between branches of size m and size m-1 is given by

$$C_m^{(j)} = \sum_{i=0}^{P_M} \left(k_{ij}^{b+} M^{(i)} D_{m-1}^{(j)} \right) - k^{b-} D_m^{(j)}, \qquad 2 \le m \le M, \tag{6b}$$

which is (4c) with polymerization from the barbed end only. It is straightforward to combine these rates into equations for D and $D_m^{(j)}$. The only complication becomes if a depolymerization event is successful. For branching, we combine the rates in (5) with the additional branching rates,

$$k_i^{\text{br}+} = 10^{-2} \,\mu\text{M}^{-2} \cdot \text{s}^{-1} \, \left(1 \quad 1.3\right) \qquad k^{\text{br}-} = 0.5 \,\text{s}^{-1} D(t=0) = 0.1 \,\mu\text{M}$$
 (7)

to validate our code without depolymerization of barbed and pointed ends (i.e., $k^{b-} = k^{p-} = 0$). Figure 15 shows that the trajectories we obtain from simulations (mean over 10 trajectories repeated 3 times for error bars) match those predicted by the ODEs.

A.2.2 Structure-based validation for depolymerization

Let us consider the structure shown in Fig. 16. It is composed of a mother filament with 4 monomers, and then three branches with 6, 8, and 10 monomers. Because each branch is attached to the previous one, our rules state that this filament can only depolymerize linearly. The average time to depolymerize the branches (remove all the blue monomers) should be $21/k^{\rm b-} + 3/k^{\rm br}$, since there are 21 monomers not attached to arp 2/3 and 3 monomers attached to arp 2/3. If $k^{\rm b-} = 1.4/{\rm s}$

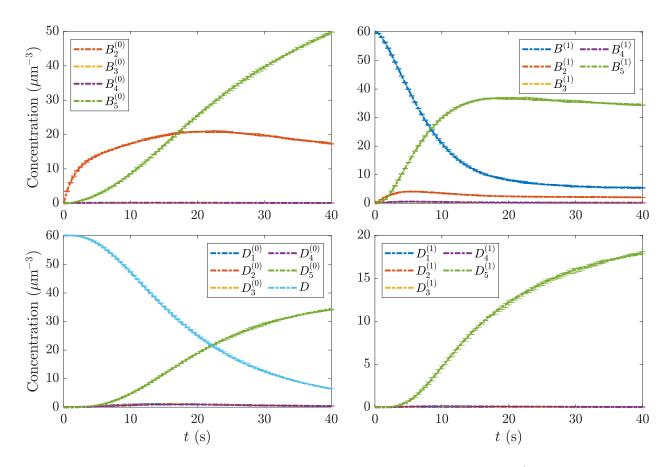


Figure 15: Validating the stochastic simulation algorithm with formins and arp 2/3 complex, where we only allow growth of fibers. We compare our stochastic simulations to ODEs (based on (4) and (6) with rates given in (5) and (7)).

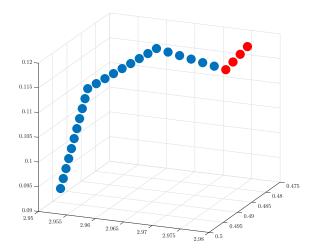


Figure 16: The branched geometry for the depolymerization test. Red points show the mother filament, blue the branches.

and $k^{\rm br-}=0.5/{\rm s}$, this comes out to 6+15=21 s. To verify that our code reproduces this depolymerization time, we perform a set of 1000 simulations with depolymerization only and record the time it takes for the filament to collapse down to the mother filament. This time is 20.9 ± 0.2 (error bar generated by repeating 20 times), which validates our implementation.

In the case when the mother filament has length 8 instead of 4, the time does not change significantly (21.1 \pm 0.5), since the three monomers at the pointed end can unbind rapidly ($k^{\rm p-}$ = 0.8/s). The time does change significantly when we put all the branches onto the same barbed end of the mother filament. In this case, all of the branches can depolymerize simultaneously, and the measured time for depolymerization is 9.9 ± 0.1 s, which is slightly longer than the theoretical time for the longest branch ($9/k^{\rm b-} + 1/k^{\rm br-} = 8.4$ s) to unbind, in accordance with the theoretical expectation that the expected value maximum is longer than the expected value of the longest event.

References

- [1] Bruce Alberts. Molecular biology of the cell. WW Norton & Company, 2017.
- [2] Deb Sankar Banerjee and Shiladitya Banerjee. Emergence and maintenance of variable-length actin filaments in a limiting pool of building blocks. *Biophysical Journal*, 121(12):2436–2448, 2022.
- [3] Deb Sankar Banerjee and Shiladitya Banerjee. Size regulation of multiple organelles competing for a limiting subunit pool. *PLOS Computational Biology*, 18(6):e1010253, 2022.
- [4] Christopher C Beltzner and Thomas D Pollard. Pathway of actin filament branch formation by arp2/3 complex. *Journal of Biological Chemistry*, 283(11):7135–7144, 2008.
- [5] Thomas A Burke, Jenna R Christensen, Elisabeth Barone, Cristian Suarez, Vladimir Sirotkin, and David R Kovar. Homeostatic actin cytoskeleton networks are regulated by assembly factor competition for monomers. *Current Biology*, 24(5):579–585, 2014.
- [6] Marie-France Carlier and Shashank Shekhar. Global treadmilling coordinates actin turnover and controls the size of actin networks. Nature Reviews Molecular Cell Biology, 18(6):389–401, 2017.

- [7] Fung-Yi Chan, Ana M Silva, Joana Saramago, Joana Pereira-Sousa, Hailey E Brighton, Marisa Pereira, Karen Oegema, Reto Gassmann, and Ana Xavier Carvalho. The arp2/3 complex prevents excessive formin activity during cytokinesis. *Molecular biology of the cell*, 30(1):96– 107, 2019.
- [8] Aravind Chandrasekaran, Edward Giniger, and Garegin A Papoian. Nucleation causes an actin network to fragment into multiple high-density domains. *Biophysical Journal*, 121(17):3200–3212, 2022.
- [9] Leah Edelstein-Keshet and G Bard Ermentrout. Models for the length distributions of actin filaments: I. simple polymerization and fragmentation. Bulletin of mathematical biology, 60(3):449–475, 1998.
- [10] Joseph Fass, Chi Pak, James Bamburg, and Alex Mogilner. Stochastic simulation of actin dynamics reveals the role of annealing and fragmentation. *Journal of theoretical biology*, 252(1):173–183, 2008.
- [11] Daniel T Gillespie. Stochastic simulation of chemical kinetics. *Annu. Rev. Phys. Chem.*, 58:35–55, 2007.
- [12] David Harbage and Jané Kondev. Exact length distribution of filamentous structures assembled from a finite pool of subunits. The Journal of Physical Chemistry B, 120(26):6225–6230, 2016.
- [13] Elizabeth S Harris, Fang Li, and Henry N Higgs. The mouse formin, frlα, slows actin filament barbed end elongation, competes with capping protein, accelerates polymerization from monomers, and severs filaments. Journal of Biological Chemistry, 279(19):20076–20087, 2004.
- [14] Jifeng Hu, Anastasios Matzavinos, and Hans G Othmer. A theoretical approach to actin filament dynamics. *Journal of Statistical Physics*, 128(1-2):111–138, 2007.
- [15] David R Kovar, Elizabeth S Harris, Rachel Mahaffy, Henry N Higgs, and Thomas D Pollard. Control of the assembly of atp-and adp-actin by formins and profilin. Cell, 124(2):423–435, 2006.
- [16] Fang Li and Henry N Higgs. The mouse formin mdia1 is a potent actin nucleation factor regulated by autoinhibition. *Current biology*, 13(15):1335–1340, 2003.

- [17] James Liman, Carlos Bueno, Yossi Eliaz, Nicholas P Schafer, M Neal Waxham, Peter G Wolynes, Herbert Levine, and Margaret S Cheung. The role of the arp2/3 complex in shaping the dynamics and structures of branched actomyosin networks. Proceedings of the National Academy of Sciences, 117(20):10825–10831, 2020.
- [18] Sophie G Martin, W Hayes McDonald, John R Yates, and Fred Chang. Tea4p links microtubule plus ends with the formin for3p in the establishment of cell polarity. *Developmental cell*, 8(4):479–491, 2005.
- [19] Shane G McInally, Jane Kondev, and Bruce L Goode. Scaling of subcellular actin structures with cell length through decelerated growth. *Elife*, 10:e68424, 2021.
- [20] Shane G McInally, Alexander JB Reading, Aldric Rosario, Predrag R Jelenkovic, Bruce Goode, and Jane Kondev. Length control emerges from cytoskeletal network geometry. bioRxiv, pages 2023–11, 2023.
- [21] Lishibanya Mohapatra, Thibaut J Lagny, David Harbage, Predrag R Jelenkovic, and Jane Kondev. The limiting-pool mechanism fails to control the size of multiple organelles. *Cell systems*, 4(5):559–567, 2017.
- [22] Qin Ni and Garegin A Papoian. Turnover versus treadmilling in actin network assembly and remodeling. *Cytoskeleton*, 76(11-12):562–570, 2019.
- [23] Aditya Paul and Thomas Pollard. The role of the fh1 domain and profilin in formin-mediated actin-filament elongation and nucleation. *Current Biology*, 18(1):9–19, 2008.
- [24] Aditya S Paul and Thomas D Pollard. Review of the mechanism of processive actin filament elongation by formins. *Cell motility and the cytoskeleton*, 66(8):606–617, 2009.
- [25] Thomas D Pollard. Regulation of actin filament assembly by arp2/3 complex and formins. Annu. Rev. Biophys. Biomol. Struct., 36:451–477, 2007.
- [26] Thomas D Pollard and John A Cooper. Quantitative analysis of the effect of acanthamoeba profilin on actin filament nucleation and elongation. *Biochemistry*, 23(26):6631–6641, 1984.
- [27] Aldric Rosario, Shane G McInally, Predrag R Jelenkovic, Bruce L Goode, and Jane Kondev. Universal length fluctuations of actin structures found in cells. bioRxiv, pages 2023–07, 2023.

- [28] Aaron D Rosenbloom, Elizabeth W Kovar, David R Kovar, Leslie M Loew, and Thomas D Pollard. Mechanism of actin filament nucleation. *Biophysical Journal*, 120(20):4399–4417, 2021.
- [29] Clarence E Schutt, James C Myslik, Michael D Rozycki, Nalin CW Goonesekere, and Uno Lindberg. The structure of crystalline profilin–β-actin. *Nature*, 365(6449):810–816, 1993.
- [30] Shashank Shekhar, Mikael Kerleau, Sonja Kühn, Julien Pernier, Guillaume Romet-Lemonne, Antoine Jégou, and Marie-France Carlier. Formin and capping protein together embrace the actin filament in a ménage à trois. *Nature communications*, 6(1):8730, 2015.
- [31] Benjamin A Smith, Karen Daugherty-Clarke, Bruce L Goode, and Jeff Gelles. Pathway of actin filament branch formation by arp2/3 complex revealed by single-molecule imaging. *Proceedings* of the National Academy of Sciences, 110(4):1285–1290, 2013.
- [32] Cristian Suarez, Robert T Carroll, Thomas A Burke, Jenna R Christensen, Andrew J Bestul, Jennifer A Sees, Michael L James, Vladimir Sirotkin, and David R Kovar. Profilin regulates f-actin network homeostasis by favoring formin over arp2/3 complex. *Developmental cell*, 32(1):43–53, 2015.
- [33] Cristian Suarez and David R Kovar. Internetwork competition for monomers governs actin cytoskeleton organization. *Nature reviews Molecular cell biology*, 17(12):799–810, 2016.
- [34] Cristian Suarez, Patrick M McCall, Margaret L Gardel, and David R Kovar. When is "enough" enough? *Cell systems*, 4(5):480–482, 2017.
- [35] Yi Xie, Rashmi Budhathoki, and J Todd Blankenship. Combinatorial deployment of f-actin regulators to build complex 3d actin structures in vivo. *Elife*, 10:e63046, 2021.
- [36] Liyuan Xu, Lingyan Cao, Jiejie Li, and Christopher J Staiger. Cooperative actin filament nucleation by the arp2/3 complex and formins maintains the homeostatic cortical array in arabidopsis epidermal cells. *The Plant Cell*, page koad301, 2023.
- [37] Mark E Zweifel, Laura A Sherer, Biswaprakash Mahanta, and Naomi Courtemanche. Nucleation limits the lengths of actin filaments assembled by formin. *Biophysical journal*, 120(20):4442–4456, 2021.