

# Chem3208 Lectures 10–14

## Electronic Correlation Methods

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## 1 Introduction

2 Configuration Interaction Methods

3 Møller-Plesset Perturbation Theory

4 Configuration Cluster Methods

5 Quantum Monte Carlo Methods

# The correlation energy

- HF replaces the e-e interaction by an averaged interaction
- The error in the HF method is called the correlation energy

$$E_c = E - E_{\text{HF}}$$

- The correlation energy is small but cannot be neglected!
- HF energy roughly 99% of total but chemistry very sensitive to remaining 1%
- The correlation energy is always negative
- Computing  $E_c$  is one of the central problem of quantum chemistry

# Correlation methods

- ① Configuration Interaction (CID, CIS, CISD, QCISD)
- ② Coupled Cluster (CCD, CCSD, CCSD[T], CCSD(T), CCSDT)
- ③ Møller-Plesset perturbation theory (MP2, MP3, MP4)
- ④ Multireference methods (MCSCF, CASSCF, RASSCF, MRCC, CASPT2)
- ⑤ Explicitly correlated F12 methods (MP2-F12, CCSD-F12, CAS-F12)
- ⑥ Density-functional theory (DFT, TDDFT)
- ⑦ Stochastic Quantum Monte Carlo methods (VMC, DMC, FCIQMC)

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# Configuration Interaction (CI)

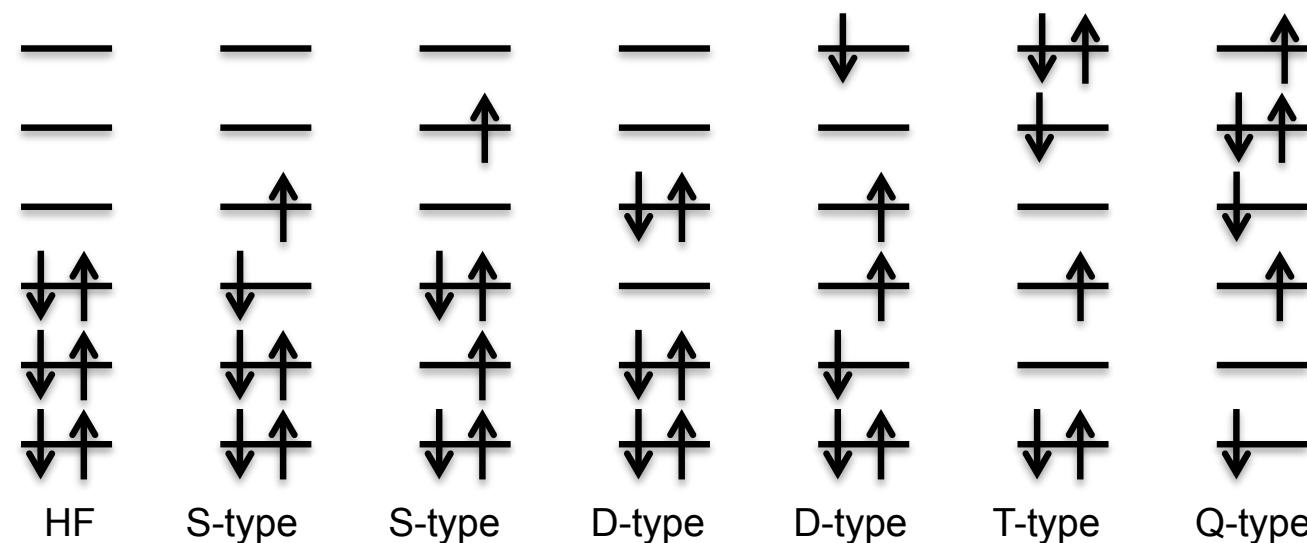
- This is the **oldest** and perhaps the **easiest** method to understand
- CI is based on the **variational principle** (like HF)
- The CI wave function is a **linear combination** of determinants
- CI methods use **excited determinants** to “improve” the HF wave function

$$\begin{aligned}
 |\Phi_0\rangle = & c_0|\Psi_0\rangle + \sum_{ia} c_i^a |\Psi_i^a\rangle + \sum_{\substack{i < j \\ a < b}} c_{ij}^{ab} |\Psi_{ij}^{ab}\rangle \\
 & + \sum_{\substack{i < j < k \\ a < b < c}} c_{ijk}^{abc} |\Psi_{ijk}^{abc}\rangle + \sum_{\substack{i < j < k < l \\ a < b < c < d}} c_{ijkl}^{abcd} |\Psi_{ijkl}^{abcd}\rangle + \dots
 \end{aligned}$$

- In  $\Psi_{ij}^{ab}$ , the electrons in **occupied** spinorbital  $i$  and  $j$  have been promoted to the **virtual** spinorbitals  $a$  and  $b$

# CI method and Excited determinants

## Excited determinants



## CI wave function

$$|\Phi_0\rangle = c_0|\text{HF}\rangle + c_S|\text{S}\rangle + c_D|\text{D}\rangle + c_T|\text{T}\rangle + c_Q|\text{Q}\rangle + \dots$$

# Truncated CI

- When  $|S\rangle$  (**singlets**) are taken into account: **CIS**

$$|\Phi_{\text{CIS}}\rangle = c_0|\text{HF}\rangle + c_S|S\rangle$$

**NB:** CIS is an **excited state method** (cf Part IV)

- When  $|D\rangle$  (**doubles**) are taken into account: **CID**

$$|\Phi_{\text{CID}}\rangle = c_0|\text{HF}\rangle + c_D|D\rangle$$

**NB:** CID is the **cheapest CI method**

- When  $|S\rangle$  and  $|D\rangle$  are taken into account: **CISD**

$$|\Phi_{\text{CISD}}\rangle = c_0|\text{HF}\rangle + c_S|S\rangle + c_D|D\rangle + c_T|T\rangle$$

**NB:** CISD is the **most commonly-used CI method**

- When  $|S\rangle$ ,  $|D\rangle$  and  $|T\rangle$  (**triples**) are taken into account: **CISDT**

$$|\Phi_{\text{CISDT}}\rangle = c_0|\text{HF}\rangle + c_S|S\rangle + c_D|D\rangle + c_T|T\rangle$$

- CISDTQ**, etc.

# Full CI

- When all possible excitations are taken into account  
this is called a **Full CI calculation (FCI)**

$$|\Phi_{\text{FCI}}\rangle = c_0|\text{HF}\rangle + c_S|\text{S}\rangle + c_D|\text{D}\rangle + c_T|\text{T}\rangle + c_Q|\text{Q}\rangle + \dots$$

- FCI gives the **exact solution of the Schrödinger equation within a given basis**
- For a two-electron system, FCI = CISD
- For a three-electron system, FCI = CISDT
- FCI is becoming more and more fashionable these days (e.g. **FCIQMC method**)
- So, why do we care about other methods?

# Full CI

- When all possible excitations are taken into account this is called a Full CI calculation (**FCI**)

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- FCI is becoming more and more fashionable these days (e.g. **FCIQMC method**)
- So, why do we care about other methods?
- Because FCI is bloody computationally expensive!

# Size of CI Matrix

*“Assume we have 10 electrons in 38 spin MOs: 10 are occupied and 28 are empty”*

- There is  $C_{10}^k$  possible ways of selecting  $k$  electrons out of the 10 occupied orbitals

$$C_n^k = \frac{n!}{k!(n-k)!}$$

- There is  $C_{28}^k$  ways of distributing them out in the 28 virtual orbitals
- For a given excitation level  $k$ , there is  $C_{10}^k C_{28}^k$  excited determinants
- The total number of possible excited determinant is

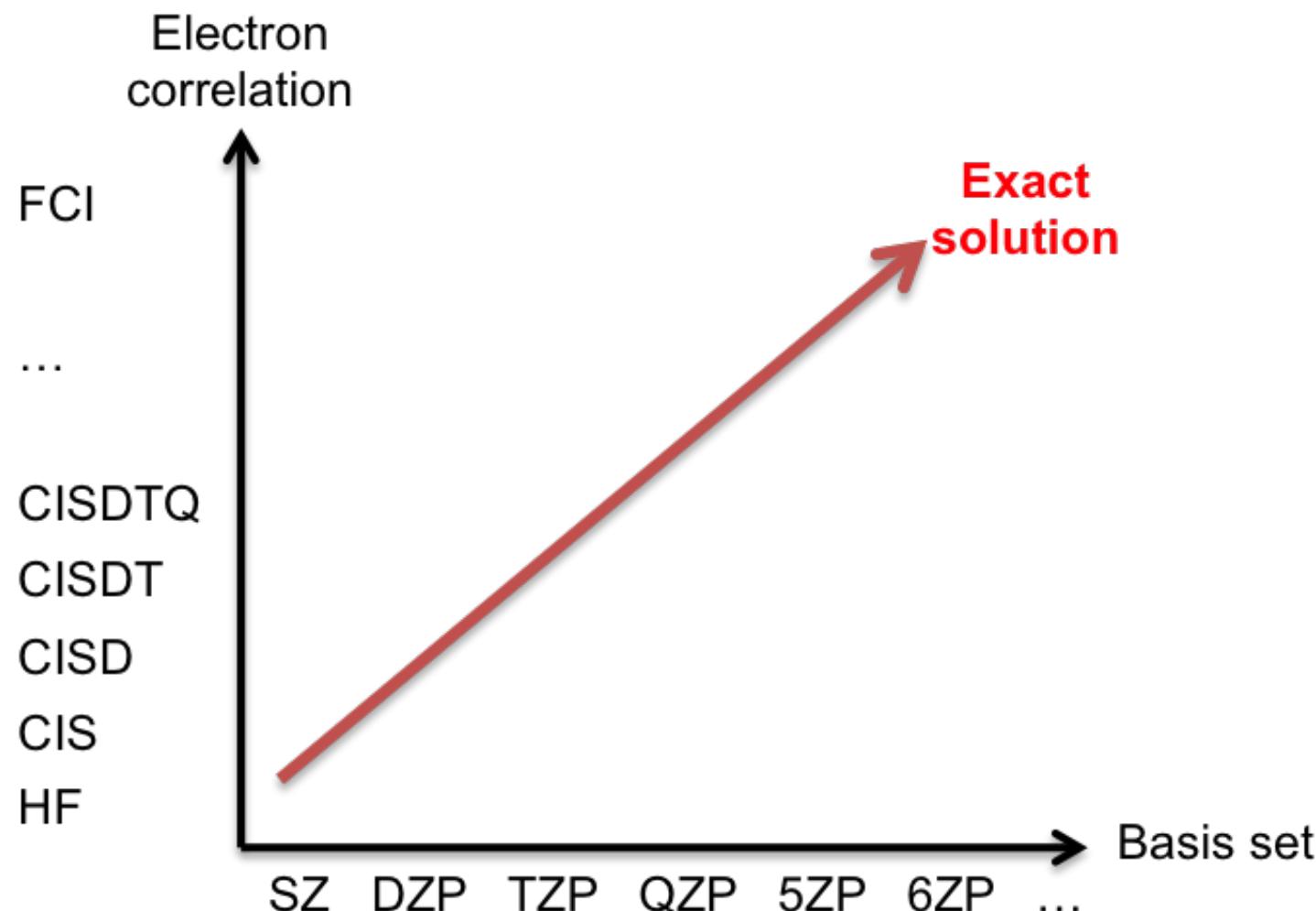
$$\sum_{k=0}^{10} C_{10}^k C_{28}^k = C_{38}^{10} = 472,733,756$$

- This is a lot...

For  $N = 10$  and  $K = 38$ :

$k$	Num. of excitations
0	1
1	280
2	17,010
3	393,120
4	4,299,750
5	24,766,560
6	79,115,400
7	142,084,800
8	139,864,725
9	69,069,000
10	13,123,110
Tot.	472,733,756

# Pople diagram



# How to do a CI calculation?

- Do a HF calculation to obtain occupied and virtual MOs
- Compute the CI matrix  $\mathbf{H}_{\text{CI}}$  (cf next slide)
- Diagonalize  $\mathbf{H}_{\text{CI}}$ :  
$$\mathbf{H}_{\text{CI}} \mathbf{C}_{\text{CI}} = \mathbf{E}_{\text{CI}} \mathbf{H}_{\text{CI}}$$
- The lowest eigenvalue corresponds to the ground state energy.
- The others correspond to excited states
- In reality, one never builds the entire  $\mathbf{H}_{\text{CI}}$  because it is too big  
⇒ Davidson diagonalization

# The FCI matrix $\mathbf{H}_{\text{FCI}}$

$$|\Phi_0\rangle = c_0|\text{HF}\rangle + c_S|\text{S}\rangle + c_D|\text{D}\rangle + c_T|\text{T}\rangle + c_Q|\text{Q}\rangle + \dots$$

Before pruning:

	$ \text{HF}\rangle$	$ \text{S}\rangle$	$ \text{D}\rangle$	$ \text{T}\rangle$	$ \text{Q}\rangle$	$\dots$
$\langle \text{HF} $	$\langle \text{HF} \text{H} \text{HF}\rangle$	$\langle \text{HF} \text{H} \text{S}\rangle$	$\langle \text{HF} \text{H} \text{D}\rangle$	$\langle \text{HF} \text{H} \text{T}\rangle$	$\langle \text{HF} \text{H} \text{Q}\rangle$	$\dots$
$\langle \text{S} $	$\langle \text{S} \text{H} \text{HF}\rangle$	$\langle \text{S} \text{H} \text{S}\rangle$	$\langle \text{S} \text{H} \text{D}\rangle$	$\langle \text{S} \text{H} \text{T}\rangle$	$\langle \text{S} \text{H} \text{Q}\rangle$	$\dots$
$\langle \text{D} $	$\langle \text{D} \text{H} \text{HF}\rangle$	$\langle \text{D} \text{H} \text{S}\rangle$	$\langle \text{D} \text{H} \text{D}\rangle$	$\langle \text{D} \text{H} \text{T}\rangle$	$\langle \text{D} \text{H} \text{Q}\rangle$	$\dots$
$\langle \text{T} $	$\langle \text{T} \text{H} \text{HF}\rangle$	$\langle \text{T} \text{H} \text{S}\rangle$	$\langle \text{T} \text{H} \text{D}\rangle$	$\langle \text{T} \text{H} \text{T}\rangle$	$\langle \text{T} \text{H} \text{Q}\rangle$	$\dots$
$\langle \text{Q} $	$\langle \text{Q} \text{H} \text{HF}\rangle$	$\langle \text{Q} \text{H} \text{S}\rangle$	$\langle \text{Q} \text{H} \text{D}\rangle$	$\langle \text{Q} \text{H} \text{T}\rangle$	$\langle \text{Q} \text{H} \text{Q}\rangle$	$\dots$
$\vdots$	$\vdots$	$\vdots$	$\vdots$	$\vdots$	$\vdots$	$\vdots$

After pruning:

	$ \text{HF}\rangle$	$ \text{S}\rangle$	$ \text{D}\rangle$	$ \text{T}\rangle$	$ \text{Q}\rangle$	$\dots$
$\langle \text{HF} $	$\langle \text{HF} \text{H} \text{HF}\rangle$	0	$\langle \text{HF} \text{H} \text{D}\rangle$	0	0	$\dots$
$\langle \text{S} $	0	$\langle \text{S} \text{H} \text{S}\rangle$	$\langle \text{S} \text{H} \text{D}\rangle$	$\langle \text{S} \text{H} \text{T}\rangle$	0	$\dots$
$\langle \text{D} $	$\langle \text{D} \text{H} \text{HF}\rangle$	$\langle \text{D} \text{H} \text{S}\rangle$	$\langle \text{D} \text{H} \text{D}\rangle$	$\langle \text{D} \text{H} \text{T}\rangle$	$\langle \text{D} \text{H} \text{Q}\rangle$	$\dots$
$\langle \text{T} $	0	$\langle \text{T} \text{H} \text{S}\rangle$	$\langle \text{T} \text{H} \text{D}\rangle$	$\langle \text{T} \text{H} \text{T}\rangle$	$\langle \text{T} \text{H} \text{Q}\rangle$	$\dots$
$\langle \text{Q} $	0	0	$\langle \text{Q} \text{H} \text{D}\rangle$	$\langle \text{Q} \text{H} \text{T}\rangle$	$\langle \text{Q} \text{H} \text{Q}\rangle$	$\dots$

# Rules & Observations

- ① No coupling between HF ground state  $|HF\rangle$  and single excitations  $|S\rangle$   
 $\Rightarrow$  Brillouin theorem

$$\langle HF | H | S \rangle = 0$$

- ② No coupling between  $|HF\rangle$  and triples  $|T\rangle$ , quadruples  $|Q\rangle$ , etc.  
 $\Rightarrow$  Slater-Condon rules

$$\langle HF | H | T \rangle = \langle HF | H | Q \rangle = \dots = 0$$

$$\langle S | H | Q \rangle = \dots = 0$$

- ③  $|S\rangle$  have small effect but mix indirectly with  $|D\rangle$   
 $\Rightarrow$  CID  $\neq$  CISD

$$\langle HF | H | S \rangle = 0 \text{ but } \langle S | H | D \rangle \neq 0$$

- ④  $|D\rangle$  have large effect and  $|Q\rangle$  more important than  $|T\rangle$   
 $\Rightarrow$  CID gives most of the correlation energy

$$\langle HF | H | D \rangle \gg \langle HF | H | Q \rangle \gg \langle HF | H | T \rangle$$

# Example

Weights of excited configurations for Ne

Excit. level	Weight
0	$9.6 \times 10^{-1}$
1	$9.8 \times 10^{-4}$
2	$3.4 \times 10^{-2}$
3	$3.7 \times 10^{-4}$
4	$4.5 \times 10^{-4}$
5	$1.9 \times 10^{-5}$
6	$1.7 \times 10^{-6}$
7	$1.4 \times 10^{-7}$
8	$1.1 \times 10^{-9}$

Correlation energy of Be and Method scaling

Method	$\Delta E_c$	%	Scaling
HF	0	0	$N^4$
CIS	0	0	$N^5$
CISD	0.075277	96.05	$N^6$
CISDT	0.075465	96.29	$N^8$
CISDTQ	0.078372	100	$N^{10}$
FCI	0.078372	100	$e^N$

# Size consistency and size extensivity

- Truncated CI methods are **size inconsistent** i.e.

$$2E_c(H_2) \neq E_c(H_2 \cdots H_2)$$

- Size consistent defines for **non-interacting fragment**
- **Size extensivity** refers to the scaling of  $E_c$  with the number of electrons
- **NB:** FCI is size consistent and size extensive

# Problem: CID for H<sub>2</sub>

*"We consider the ground state of the H<sub>2</sub> molecule in a minimal basis set with two spatial MOs  $\psi_1$  and  $\psi_2$ ." We give the following integrals:*

$$h_1 = -1.2529 \quad h_2 = -0.4756$$

$$J_{11} = 0.6746 \quad J_{22} = 0.6975 \quad J_{12} = 0.6636 \quad K_{12} = 0.1813$$

- ① Calculate the HF energy
- ② Knowing that the CID matrix is

$$\mathbf{H}_{\text{CID}} = \begin{pmatrix} 2h_1 + (11|11) & (12|12) \\ (12|12) & 2h_2 + (22|22) \end{pmatrix}$$

Calculate the ground state CID energy

- ③ Combining the results of the previous two questions, calculate the correlation energy