

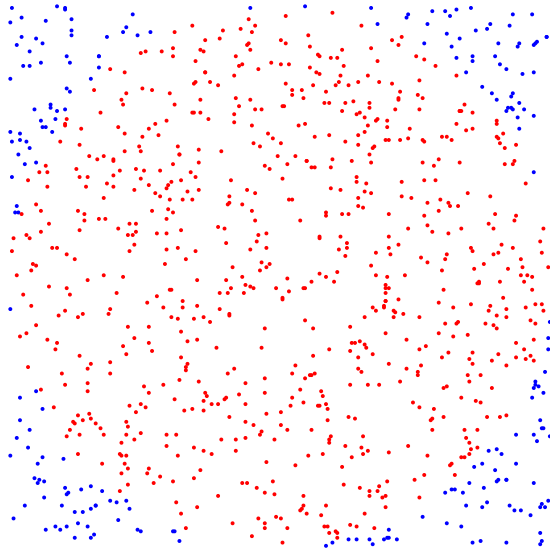
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Monte Carlo (MC) method

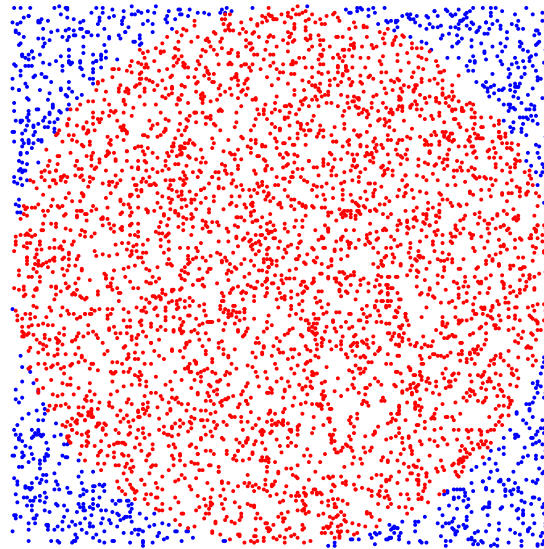
- Monte Carlo is a **numerical integration method**
- It is used in problems where it is **too difficult or impossible to obtain analytical expressions or the dimensionality of the integral is large**
- The method consists in **repeating random sampling many times** to obtain numerical results:
⇒ this is a **non-deterministic** or **stochastic** method.
- **MC converges as $N^{-1/2}$** where N is the number of MC step
- In 1946, **Stanislaw Ulam** was the first mathematician to dignify this approach with a name, in honor of his uncle having a little issue with gambling
- **Nicolas Metropolis** also made important contributions (**Metropolis algorithm**)

Monte Carlo computation of π

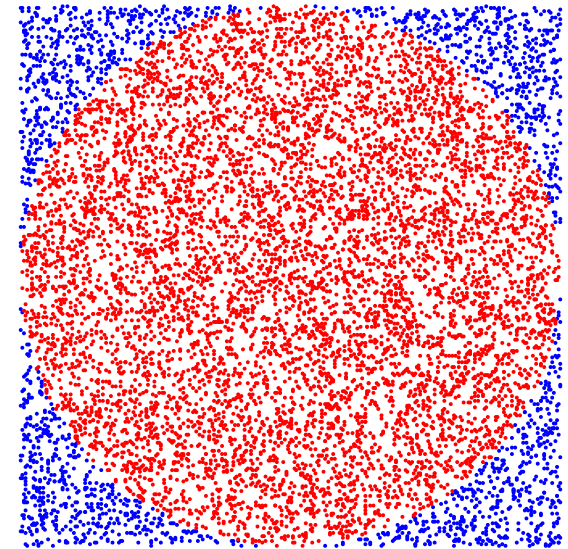
$$\int_{-1}^1 \int_{-1}^1 f(x, y) dx dy = \pi \quad \text{with} \quad f(x, y) = \begin{cases} 1, & x^2 + y^2 \leq 1, \\ 0, & \text{otherwise.} \end{cases}$$



$$\pi \approx \frac{774}{1000} = 3.096$$



$$\pi \approx \frac{3962}{5000} = 3.1696$$



$$\pi \approx \frac{7948}{10000} = 3.1792$$

```
program pi
c  Declare variables
  implicit none
  integer seed,i,n
  double precision one,four,x,y,r,res
c  Input data
  one=1.0d0
  four=4.0d0
  seed=19082016
  res=0.0d0
  n=100000
c  Initialize random number generator
  call srand(seed)
c  Start loop
  do i=1,n
c    Get two random numbers
    x=rand(0)
    y=rand(0)
c    Calculate the radius
    r=x**2+y**2
c    Increment result if in circle
    if(r.le.1) res=res+one
  enddo
c  Print the result
  print*, 'MC estimate of pi = ', four*res/dble(n)
  print*, 'Exact value of pi = ', four*atan(one)
c  End of the program
  return
end
```

Variational Monte Carlo (VMC)

- Within quantum chemistry, **VMC** is used to obtain expectation values (mainly energies)
- In VMC, the expectation value of the Hamiltonian with respect to a **trial wave function** Ψ_T is obtained using a **stochastic integration technique**
- The VMC energy is an upper bound to the exact ground state energy

$$E_{\text{VMC}} = \frac{\int \Psi_T(\mathbf{r}) \mathbf{H} \Psi_T(\mathbf{r}) d\mathbf{r}}{\int \Psi_T(\mathbf{r})^2 d\mathbf{r}} = \frac{\int \frac{\mathbf{H} \Psi_T(\mathbf{r})}{\Psi_T(\mathbf{r})} \Psi_T(\mathbf{r})^2 d\mathbf{r}}{\int \Psi_T(\mathbf{r})^2 d\mathbf{r}} = \frac{\int E_L(\mathbf{r}) \Psi_T(\mathbf{r})^2 d\mathbf{r}}{\int \Psi_T(\mathbf{r})^2 d\mathbf{r}}$$

where

$$E_L(\mathbf{r}) = \frac{\hat{H} \Psi_T(\mathbf{r})}{\Psi_T(\mathbf{r})} \text{ is the local energy}$$

Diffusion Monte Carlo (DMC)

Time-dependent Schrödinger equation **written in imaginary time**:

$$\frac{\partial \Phi(\mathbf{r}, \tau)}{\partial \tau} = (\mathbf{H} - S)\Phi(\mathbf{r}, \tau)$$

- For $\tau \rightarrow \infty$, the solution is the **exact ground state wave function** $\Phi(\mathbf{r})$
- DMC generates **configurations** (or **walkers**) distributed according to the **density** $\rho(\mathbf{r}, \tau) = \Psi_T(\mathbf{r}) \Phi(\mathbf{r}, \tau)$

$$\frac{\partial \rho(\mathbf{r}, \tau)}{\partial \tau} = \underbrace{\frac{1}{2} \nabla^2 \rho(\mathbf{r}, \tau)}_{\text{diffusion}} + \underbrace{\nabla \cdot [F(\mathbf{r}) \rho(\mathbf{r}, \tau)]}_{\text{drift}} - \underbrace{[E_L(\mathbf{r}) - E_T] \rho(\mathbf{r}, \tau)}_{\text{branching}}$$

where

$$F(\mathbf{r}) = \frac{\nabla \Psi_T(\mathbf{r})}{\Psi_T(\mathbf{r})} \quad \text{is the quantum force}$$

If $\Psi_T(\mathbf{r})$ has **exact nodes**, DMC energy = **exact energy** (**fixed-node error**)

Congratulations!

You have survived Part III of Chem3208!

Next time, you will talk about

Density-functional theory

with Andrew Gilbert