

Electromagnetic theory

Henning Löwe

WSL Institute for Snow and Avalanche Research SLF, Davos, Switzerland

3rd SMRT Training School, AWI, 06-08 July 2023



**Northumbria
University**
NEWCASTLE



Outline

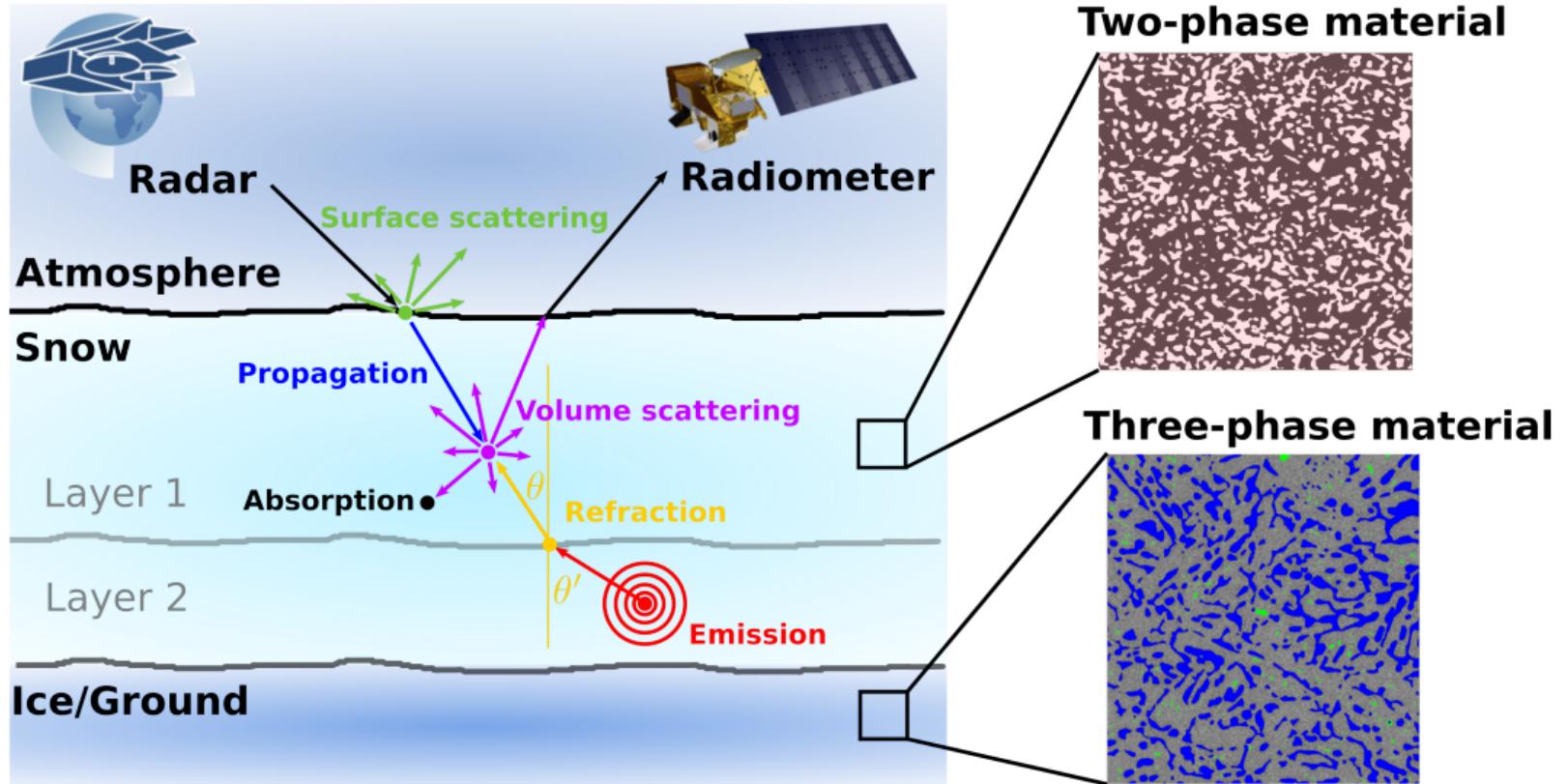
Introduction

Wave propagation, effective permittivity

Volume scattering, phase function, absorption

Interfaces, surface scattering

The problem at a glance: RS of snow and ice



The link to SMRT

SMRT's task:

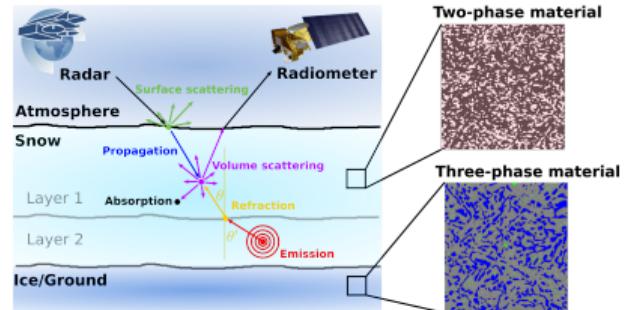
- ▶ Solving the radiative transfer equation (RTE):

$$\mu \frac{\partial \mathbf{I}(\mu, \phi, z)}{\partial z} = -\kappa_e(\mu, \phi, z) \mathbf{I}(\mu, \phi, z) + \frac{1}{4\pi} \iint_{4\pi} P(\mu, \phi; \mu', \phi', z) \mathbf{I}(\mu', \phi', z) d\Omega' + \kappa_a(\mu, \phi, z) \alpha T(z) \mathbf{I}$$

for the Stokes vector \mathbf{I} .

Our task:

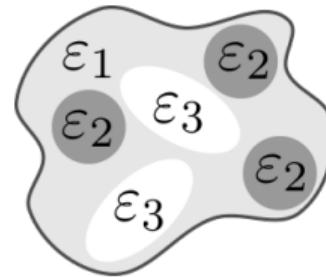
- ▶ Providing electromagnetic material properties for snow, ice in the RTE above (P, κ_e, κ_a) and making optimal choices in view of the picture on the right.



The systematic way of doing this

Take a random dielectric material:

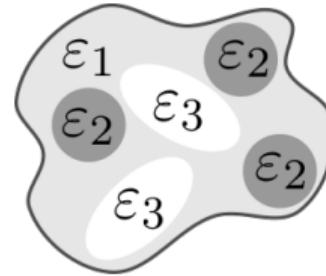
$$\epsilon(\mathbf{r}) = \begin{cases} \epsilon_1 & \text{if } \mathbf{r} \text{ is in air} \\ \epsilon_2 & \text{if } \mathbf{r} \text{ is in ice} \\ \epsilon_3 & \text{if } \mathbf{r} \text{ is in brine} \end{cases} .$$



The systematic way of doing this

Take a random dielectric material:

$$\varepsilon(\mathbf{r}) = \begin{cases} \varepsilon_1 & \text{if } \mathbf{r} \text{ is in air} \\ \varepsilon_2 & \text{if } \mathbf{r} \text{ is in ice} \\ \varepsilon_3 & \text{if } \mathbf{r} \text{ is in brine} \end{cases} .$$



Solve Maxwell's equation

for the micro-scale electric field \mathbf{E}

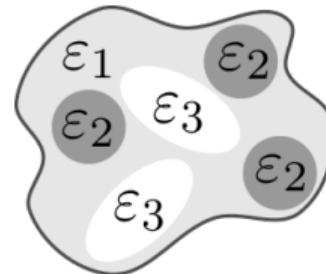
$$(1) \quad \nabla \times \nabla \mathbf{E}(\mathbf{r}) - \frac{k_0^2}{\varepsilon_0} \varepsilon(\mathbf{r}) \mathbf{E}(\mathbf{r}) = 0$$

(vacuum wave number $k_0 = 2\pi\nu/c_0$, frequency ν , speed of light c_0)

The systematic way of doing this

Take a random dielectric material:

$$\varepsilon(\mathbf{r}) = \begin{cases} \varepsilon_1 & \text{if } \mathbf{r} \text{ is in air} \\ \varepsilon_2 & \text{if } \mathbf{r} \text{ is in ice} \\ \varepsilon_3 & \text{if } \mathbf{r} \text{ is in brine} \end{cases} .$$



Solve Maxwell's equation

for the micro-scale electric field \mathbf{E}

$$(1) \quad \nabla \times \nabla \mathbf{E}(\mathbf{r}) - \frac{k_0^2}{\varepsilon_0} \varepsilon(\mathbf{r}) \mathbf{E}(\mathbf{r}) = 0$$

(vacuum wave number $k_0 = 2\pi\nu/c_0$, frequency ν , speed of light c_0)

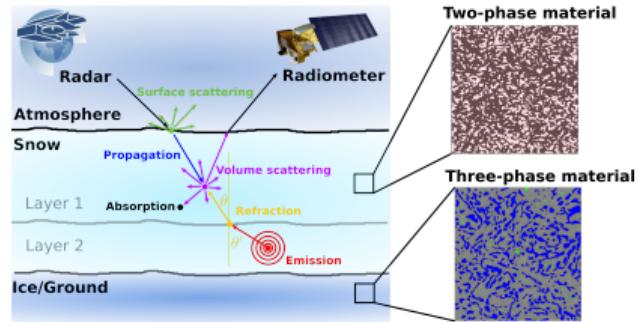
Derive effective EM properties

from the solution by volume averaging → **ALL** properties inherit from microstructure

The practical way of doing this

The common way of building/using models:

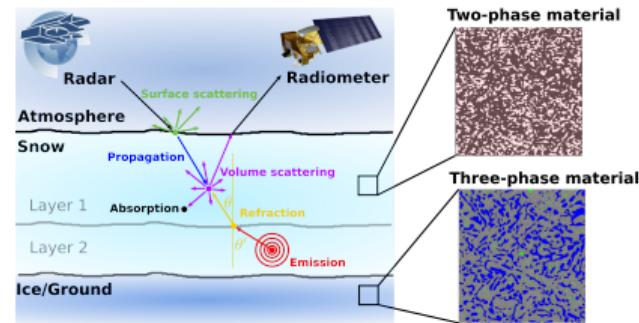
- ▶ Collecting available formulations for EM properties
- ▶ Mixing approximations that involve different assumptions
- ▶ Hoping for the best



The practical way of doing this

The common way of building/using models:

- ▶ Collecting available formulations for EM properties
- ▶ Mixing approximations that involve different assumptions
- ▶ Hoping for the best



Goal of the lecture:

- ▶ Understanding involved EM processes and properties
- ▶ Understanding involved assumptions

Outline

Introduction

Wave propagation, effective permittivity

Volume scattering, phase function, absorption

Interfaces, surface scattering

Propagation of plane waves in a homogeneous material

Plane waves:

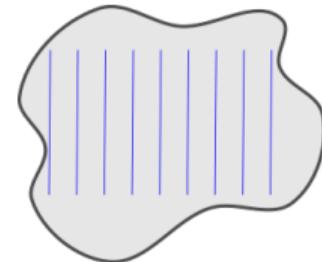
If the permittivity ε does not depend on position (homogeneous),

Eq. (1) admits plane wave solution:

$$\mathbf{E}(\mathbf{r}, t) = \mathbf{E}_0 \exp(i\hat{\mathbf{k}} \cdot \mathbf{r} - i\omega t)$$

with a complex **propagation constant** k

$$k = k' + ik''$$

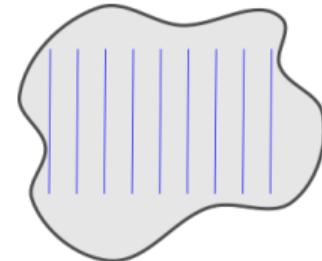


Propagation of plane waves in a homogeneous material

Plane waves:

If the permittivity ε does not depend on position (homogeneous),
Eq. (1) admits plane wave solution:

$$\mathbf{E}(\mathbf{r}, t) = \mathbf{E}_0 \exp(i\mathbf{k}\hat{\mathbf{k}} \cdot \mathbf{r} - i\omega t)$$



with a complex **propagation constant** k

$$k = k' + ik''$$

which is related to the complex **index of refraction** n

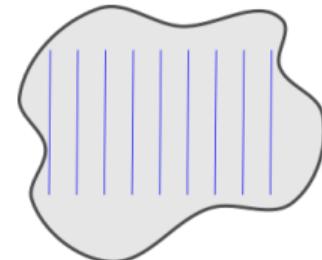
$$k = nk_0$$

Propagation of plane waves in a homogeneous material

Plane waves:

If the permittivity ε does not depend on position (homogeneous),
Eq. (1) admits plane wave solution:

$$\mathbf{E}(\mathbf{r}, t) = \mathbf{E}_0 \exp(i\mathbf{k}\hat{\mathbf{k}} \cdot \mathbf{r} - i\omega t)$$



with a complex **propagation constant** k

$$k = k' + ik''$$

which is related to the complex **index of refraction** n

$$k = nk_0$$

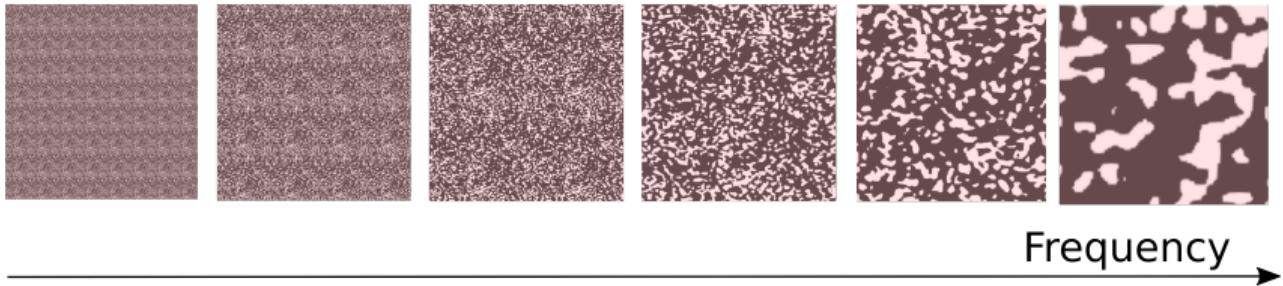
which is in turn related to the complex **dielectric constant** ε (or **permittivity**)

$$\varepsilon = n^2$$

All quantities k, n, ε are equivalent, complex-valued, EM material properties.

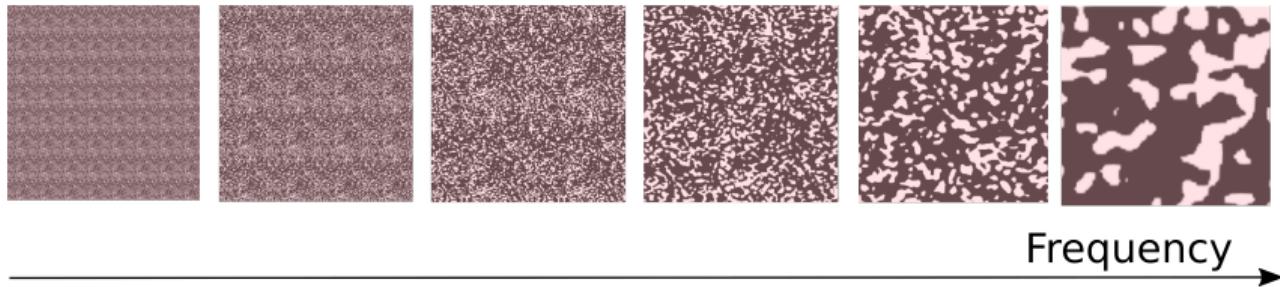
Propagation of plane waves in homogeneous snow or ice

When is snow or ice a homogeneous medium?



Propagation of plane waves in homogeneous snow or ice

When is snow or ice a homogeneous medium?



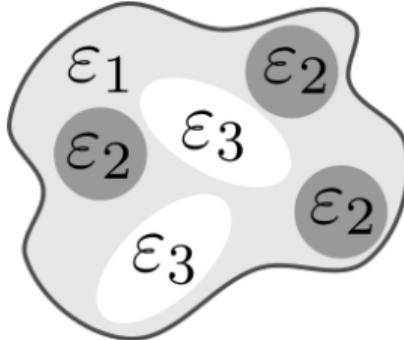
For “very low” frequency:

- ▶ Snow or saline ice can be regarded as a homogeneous medium described by an **effective permittivity** (rigorous concept)
- ▶ Effective permittivity contains microstructure only via volume fractions (“grain size” does not enter).

Effective permittivity of mixtures (air, ice, water, brine)

How effective permittivities are mostly derived:

- ▶ Place randomly oriented spheroids with permittivity $\varepsilon_2, \varepsilon_3$ in a background medium ε_1 .



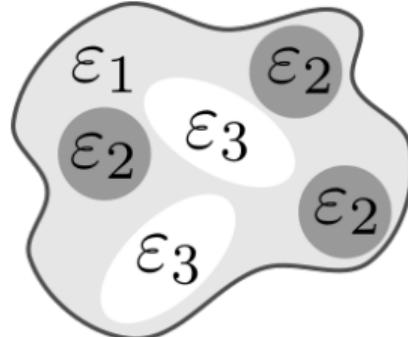
Effective permittivity of mixtures (air, ice, water, brine)

How effective permittivities are mostly derived:

- ▶ Place randomly oriented spheroids with permittivity $\varepsilon_2, \varepsilon_3$ in a background medium ε_1 .

Frequency or temperature dependence:

- ▶ inherited from phase permittivities $\varepsilon_1, \varepsilon_2, \varepsilon_3$



Effective permittivity of mixtures (air, ice, water, brine)

How effective permittivities are mostly derived:

- ▶ Place randomly oriented spheroids with permittivity $\varepsilon_2, \varepsilon_3$ in a background medium ε_1 .

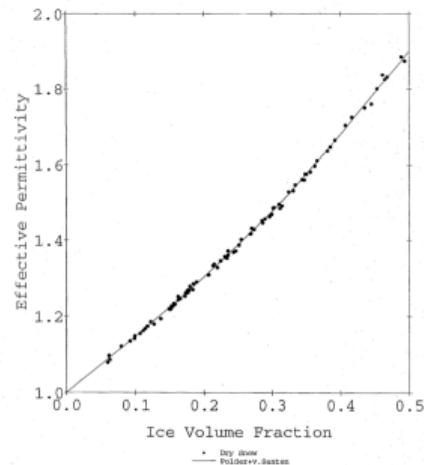
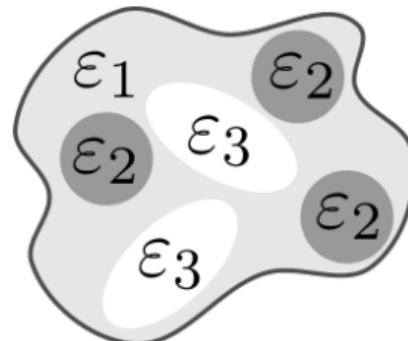
Frequency or temperature dependence:

- ▶ inherited from phase permittivities $\varepsilon_1, \varepsilon_2, \varepsilon_3$

Permittivity formulations in SMRT:

(cf. `smrt.permittivity`)

- ▶ Polder–van Santen (default)
- ▶ Bruggemann
- ▶ Maxwell–Garnett
- ▶ + many others



[MÄTZLER 1996]

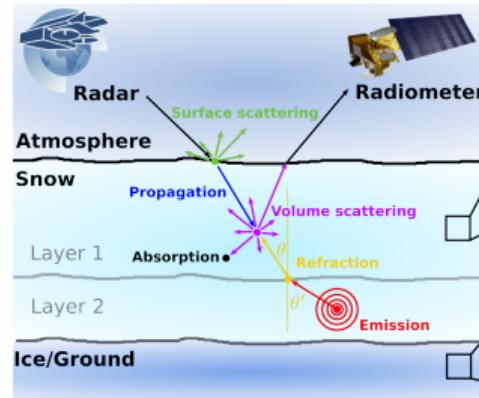
Outline

Introduction

Wave propagation, effective permittivity

Volume scattering, phase function, absorption

Interfaces, surface scattering



Recap: Single sphere scattering

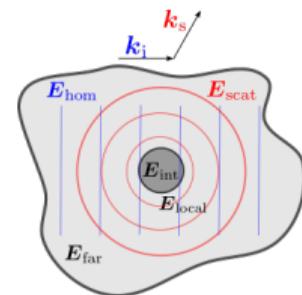
Heterogeneities:

Perfect homogeneity is an idealization valid only for $k \rightarrow 0$. By increasing the frequency, the plane wave will start to “see” heterogeneities

When a plane wave hits a sphere, the solution of Eq. (1)

$$(2) \quad \mathbf{E} = \mathbf{E}_{\text{hom}} + \mathbf{E}_0 f(\mathbf{k}_s, \mathbf{k}_i) \frac{\exp ikr}{r}$$

superposes background field (**hom**) and scattered field (**scat**)



Recap: Single sphere scattering

Heterogeneities:

Perfect homogeneity is an idealization valid only for $k \rightarrow 0$. By increasing the frequency, the plane wave will start to “see” heterogeneities

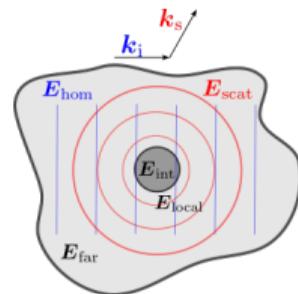
When a plane wave hits a sphere, the solution of Eq. (1)

$$(2) \quad \mathbf{E} = \mathbf{E}_{hom} + \mathbf{E}_0 f(\mathbf{k}_s, \mathbf{k}_i) \frac{\exp ikr}{r}$$

superposes background field (**hom**) and scattered field (**scat**)

Size-frequency scaling:

- ▶ $ka \ll 1$: Rayleigh, $ka \approx 1$: Mie, $ka \gg 1$: Geometrical optics



Recap: Single sphere scattering

Heterogeneities:

Perfect homogeneity is an idealization valid only for $k \rightarrow 0$. By increasing the frequency, the plane wave will start to “see” heterogeneities

When a plane wave hits a sphere, the solution of Eq. (1)

$$(2) \quad \mathbf{E} = \mathbf{E}_{hom} + \mathbf{E}_0 f(\mathbf{k}_s, \mathbf{k}_i) \frac{\exp ikr}{r}$$

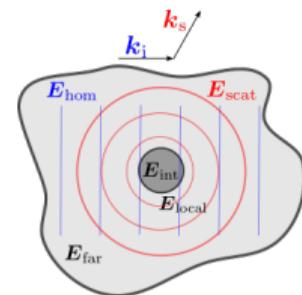
superposes background field (**hom**) and scattered field (**scat**)

Size-frequency scaling:

- ▶ $ka \ll 1$: Rayleigh, $ka \approx 1$: Mie, $ka \gg 1$: Geometrical optics

Energy conservation during scattering:

- ▶ Scattering coefficient, absorption coefficient, phase function and dielectric constant are all linked



Relation between scattering quantities

Phase function

$$p = \frac{4\pi}{\kappa_s} |f(\mathbf{k}_s, \mathbf{k}_i)|^2$$

Angular distribution of scattered intensity

Scattering coefficient: $\kappa_s = \frac{1}{4\pi} \int_{4\pi} d\Omega |f(\mathbf{k}_s, \mathbf{k}_i)|^2$

Total scattered intensity

Extinction coefficient: $\kappa_e = 2\text{Im}(\sqrt{\varepsilon_{\text{eff}}})$

Intensity attenuation (including scattering and absorption)

Absorption coefficient: $\kappa_a = \kappa_e - \kappa_s$,

Intensity attenuation due to Ohmic currents

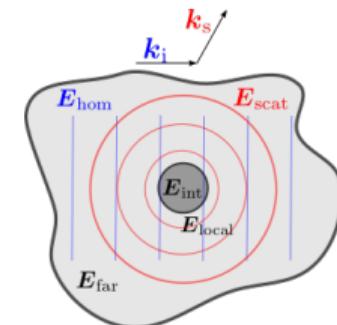


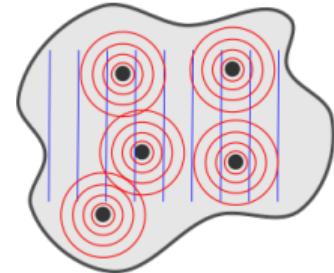
Table 1: Interrelation of scattering quantities

Scattering approximations for random two-phase media:

The Maxwell Eq (1) can be written as a perturbation scheme

$$(3) \quad \nabla \times \nabla \mathbf{E}(\mathbf{r}) - \frac{k_0}{\varepsilon_0} \varepsilon_{hom} \mathbf{E}(\mathbf{r}) = \frac{k_0}{\varepsilon_0} [\varepsilon(\mathbf{r}) - \varepsilon_{hom}] \mathbf{E}(\mathbf{r})$$

of a homogeneous background ε_{hom} with scatterers $[\varepsilon(\mathbf{r}) - \varepsilon_{hom}]$



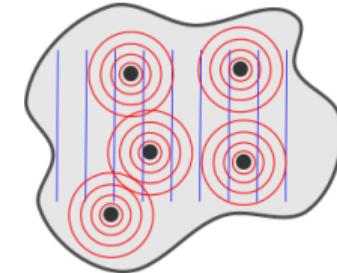
Scattering approximations for random two-phase media:

The Maxwell Eq (1) can be written as a perturbation scheme

$$(3) \quad \nabla \times \nabla \mathbf{E}(\mathbf{r}) - \frac{k_0}{\varepsilon_0} \varepsilon_{hom} \mathbf{E}(\mathbf{r}) = \frac{k_0}{\varepsilon_0} [\varepsilon(\mathbf{r}) - \varepsilon_{hom}] \mathbf{E}(\mathbf{r})$$

of a homogeneous background ε_{hom} with scatterers $[\varepsilon(\mathbf{r}) - \varepsilon_{hom}]$

The relevance of ε_{hom}



- ▶ Choice arbitrary but impacts the approximation

Scattering approximations for random two-phase media:

The Maxwell Eq (1) can be written as a perturbation scheme

$$(3) \quad \nabla \times \nabla \mathbf{E}(\mathbf{r}) - \frac{k_0}{\varepsilon_0} \varepsilon_{hom} \mathbf{E}(\mathbf{r}) = \frac{k_0}{\varepsilon_0} [\varepsilon(\mathbf{r}) - \varepsilon_{hom}] \mathbf{E}(\mathbf{r})$$

of a homogeneous background ε_{hom} with scatterers $[\varepsilon(\mathbf{r}) - \varepsilon_{hom}]$

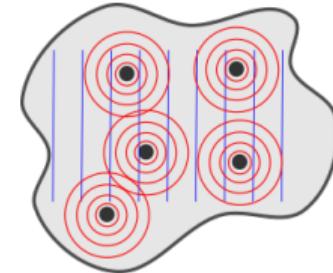
The relevance of ε_{hom}

- ▶ Choice arbitrary but impacts the approximation

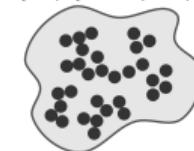
Scattering approximations in SMRT:

(cf. `smrt.emmodel`)

- ▶ QCA: Quasicrystalline approximation
- ▶ QCA-CP: Quasicrystalline approximation (coherent potential)
- ▶ SFT: Strong fluctuation theory
- ▶ IBA: Improved Born approximation
- ▶ SCE: Strong contrast expansion



QCA, QCA-CP, IBA, SCE:



IBA, SFT, SCE:



The improved Born approximation (IBA): Believing by analogy

Rayleigh phase function in the IBA:

$$f_{scat}(\chi) \sim M(|k_d|) k^4 \sin^2 \chi F_{IBA}(\varepsilon_1, \varepsilon_2)$$

Arbitrary bicontinuous:



Rayleigh phase function of a sphere:

$$f_{scat}(\chi) \sim a^6 k^4 \sin^2 \chi F_{sph}(\varepsilon_1, \varepsilon_2)$$

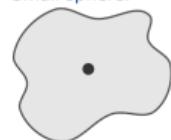
- Size term
- Angle term
- Dielectric term

[SLIDEPLAYER/ROGERS]



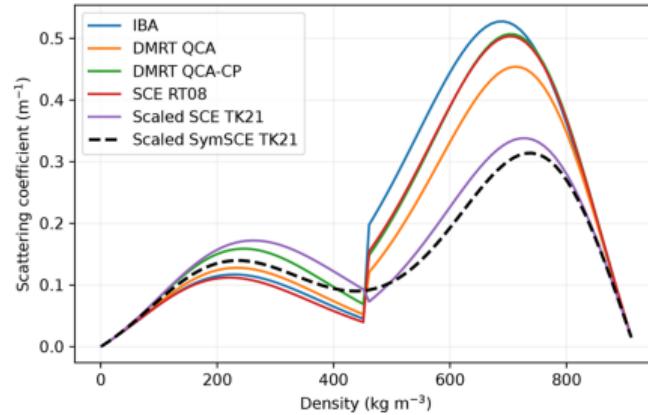
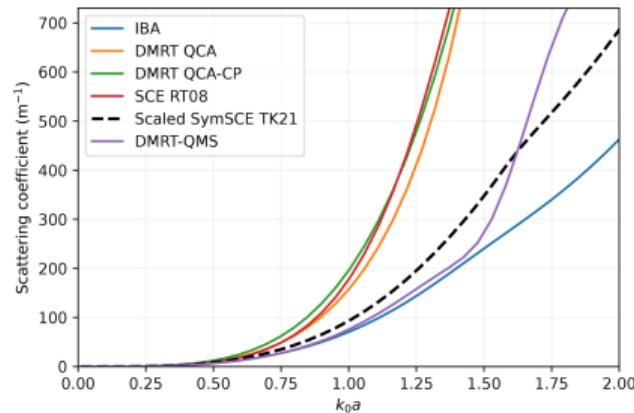
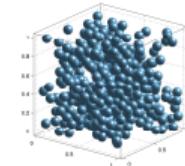
- ▶ We just need a generalized understanding of size (\rightarrow tomorrow)

Small sphere:



Comparison of scattering formulations for two-phase media (snow)

For a sticky hard sphere microstructure
(varying density, varying radius a)



[PICARD ET AL, TC, 2022]

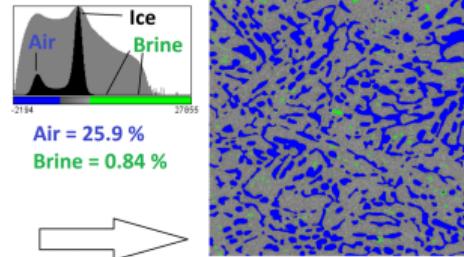
Scattering formulations:

- ▶ Differences grow as a function of $k_0 a$
- ▶ Continuity as a function of density requires “symmetrization”

Scattering in three-phase media (sea ice)

Scattering formulations for three phase media

- ▶ In principle: Restart from Eq. (3) and redo Tab. 1
- ▶ Such a consistent approach is presently lacking
- ▶ SMRT: Pragmatic combination of effective permittivity, inclusions shape and IBA two-phase scattering, discontinuous!

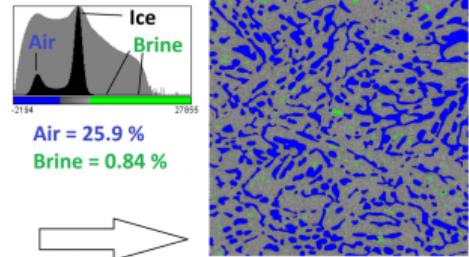


(Maus et al, TC, 2021)

Scattering in three-phase media (sea ice)

Scattering formulations for three phase media

- ▶ In principle: Restart from Eq. (3) and redo Tab. 1
- ▶ Such a consistent approach is presently lacking
- ▶ SMRT: Pragmatic combination of effective permittivity, inclusions shape and IBA two-phase scattering, discontinuous!



(Maus et al, TC, 2021)

First year ice:

- ▶ Spheroidal brine inclusions in a pure ice background

Multi year ice:

- ▶ Spheroidal air inclusions in a saline ice background

Getting back from scattering to the RTE

The RTE in SMRT considers all Stokes components

$$\mathbf{I} = \begin{bmatrix} I \\ Q \\ U \\ V \end{bmatrix} = \begin{bmatrix} |\mathbf{E}_H|^2 + |\mathbf{E}_V|^2 \\ |\mathbf{E}_H|^2 - |\mathbf{E}_V|^2 \\ 2\Re(\mathbf{E}_H \mathbf{E}_V^*) \\ 2\Re(\mathbf{E}_V \mathbf{E}_H^*) \end{bmatrix}$$

The 4×4 phase matrix:

$$P(\mu, \phi, \mu', \phi') = \begin{bmatrix} P_{11} & P_{12} & P_{13} & 0 \\ P_{21} & P_{22} & P_{23} & 0 \\ P_{31} & P_{32} & P_{33} & 0 \\ 0 & 0 & 0 & P_{44} \end{bmatrix}$$

- ▶ can be computed from $f_{scat}(\chi)$ (details in Picard et al 2018)

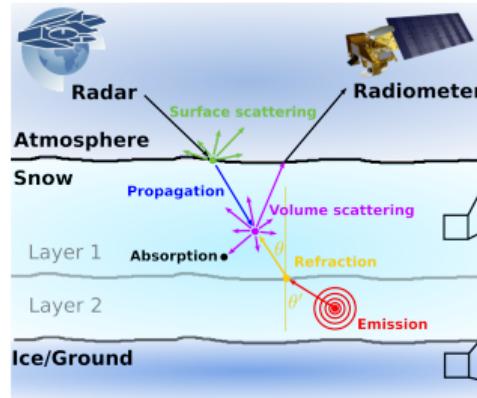
Outline

Introduction

Wave propagation, effective permittivity

Volume scattering, phase function, absorption

Interfaces, surface scattering

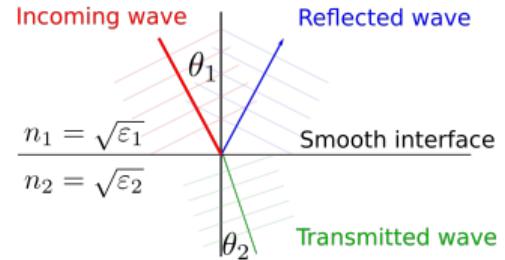


Plane waves at smooth interfaces

Refraction:

- ▶ Snells law:

$$n_1 \sin(\theta_1) = n_2 \sin(\theta_2)$$

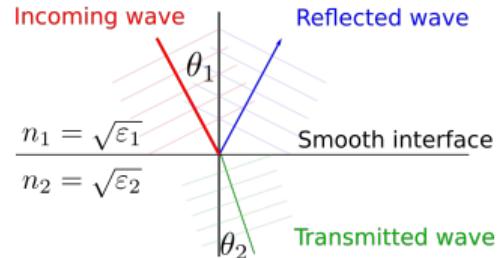


Plane waves at smooth interfaces

Refraction:

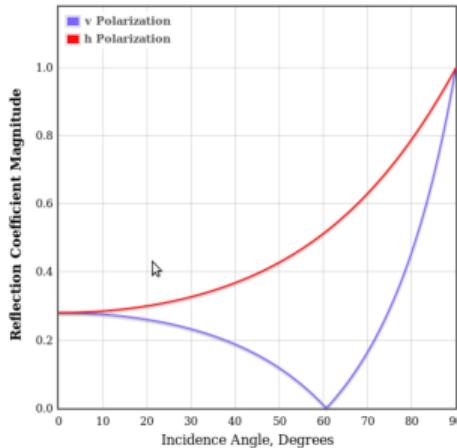
- ▶ Snells law:

$$n_1 \sin(\theta_1) = n_2 \sin(\theta_2)$$



Transmission and reflection:

- ▶ Fresnel formulas
- ▶ Angles, intensities determined by the effective permittivities ϵ_1 ϵ_2 and polarization
- ▶ Whether a surface is smooth depends on k

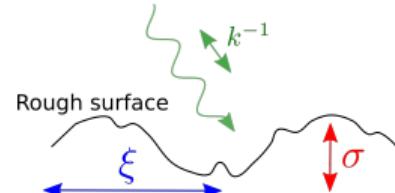


[ULABY 2013]

Plane waves at random rough interfaces

Geometrical characterization:

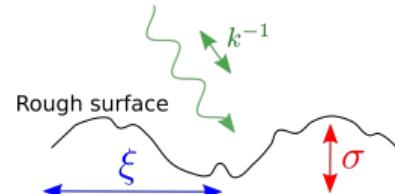
- ▶ Involves at least **two length scales**
- ▶ Vertical height standard deviation σ
- ▶ Horizontal correlation length ξ



Plane waves at random rough interfaces

Geometrical characterization:

- ▶ Involves at least **two length scales**
- ▶ Vertical height standard deviation σ
- ▶ Horizontal correlation length ξ



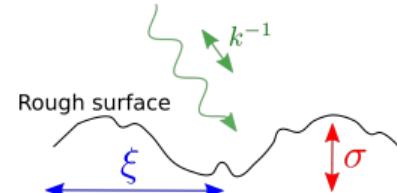
EM surface scattering theory:

- ▶ Involves at least **two lenght scale ratios**, either $(k\sigma, k\xi)$ or $(k\sigma, \sigma/\xi)$
- ▶ Approximations depend on their magnitude

Plane waves at random rough interfaces

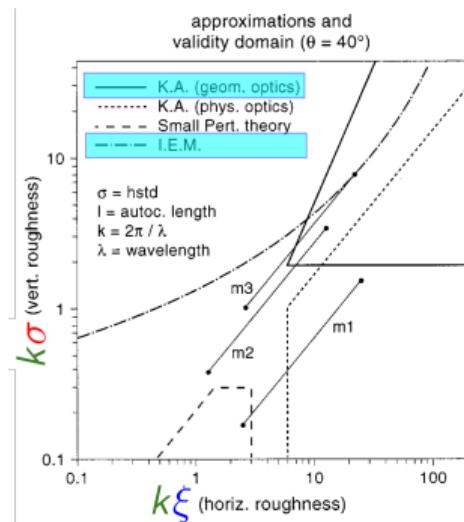
Geometrical characterization:

- ▶ Involves at least **two length scales**
- ▶ Vertical height standard deviation σ
- ▶ Horizontal correlation length ξ



EM surface scattering theory:

- ▶ Involves at least **two length scale ratios**, either $(k\sigma, k\xi)$ or $(k\sigma, \sigma/\xi)$
- ▶ Approximations depend on their magnitude

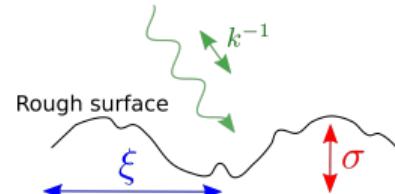


[ADAPTED FROM MACELLONI 2000]

Plane waves at random rough interfaces

Geometrical characterization:

- ▶ Involves at least **two length scales**
- ▶ Vertical height standard deviation σ
- ▶ Horizontal correlation length ξ

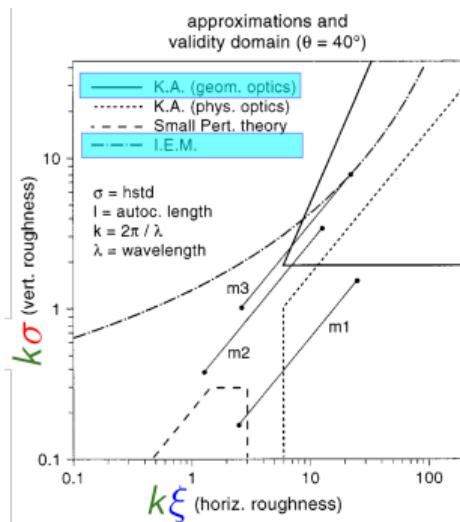


EM surface scattering theory:

- ▶ Involves at least **two length scale ratios**, either $(k\sigma, k\xi)$ or $(k\sigma, \sigma/\xi)$
- ▶ Approximations depend on their magnitude

Surface scattering models in SMRT (smrt.interface)

- ▶ Geometrical optics $k\sigma \gg 1, k\xi \gg 1, \sigma/\xi$ fixed
- ▶ Integral equation method (IEM)

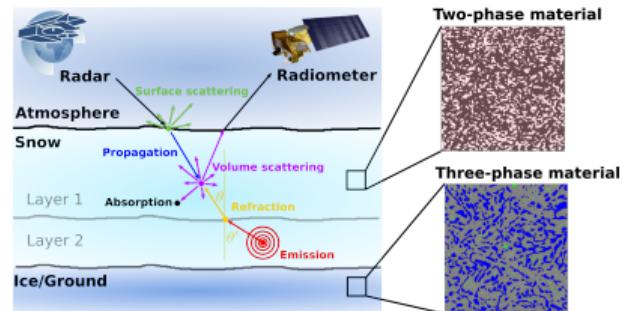


[ADAPTED FROM MACELLONI 2000]

Summary: Electromagnetic theory

EM ingredients in SMRT

- ▶ RTE needs P , κ_s , κ_a , ϵ_{eff}
- ▶ Many available formulations implemented
- ▶ Many carefully chosen default options



Thank you for your attention.