

# Recent Advances in Quantum Computing: From Annealing Dynamics to Fundamental Tests of Nonlocality

Claude Opus 4.6  
Scientific Research Assistant  
Anthropic

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## Abstract

Quantum computing research spans from practical optimization to fundamental tests of quantum mechanics, with recent advances revealing how quantum-classical boundaries govern system behavior. We review three contemporary studies that exemplify this breadth. First, comparative investigations of simulated versus quantum reverse annealing on D-Wave systems demonstrate that classical relaxation dominates at long timescales, with systematic ground state suppression distinguishing reverse from forward annealing protocols. Second, studies of quantum chaos in periodically driven macrospin ensembles establish that Lyapunov exponents unify classical and quantum dynamical characterization, with quantum-classical correspondence holding only within the Lyapunov time  $t_L \sim 1/\lambda_{\max}$ . Third, advances in Bell nonlocality tests achieve detection efficiency thresholds as low as 50% through optimized one-detector steering protocols and loss-tolerant designs, enabling loophole-free demonstrations with simplified experimental setups. Across these domains, temporal and energetic scales determine when quantum effects persist despite environmental coupling. These findings illuminate fundamental limitations and opportunities for quantum technologies while refining our understanding of quantum mechanics itself.

## 1 Introduction

- Quantum computing stands at the intersection of practical optimization and fundamental physics, with recent advances revealing new insights across multiple domains. While quantum annealing systems promise computational advantages for optimization problems, their dynamics involve complex phase transitions that blur the boundaries between quantum and classical behavior. Simultaneously, investigations into quantum chaos through macrospin systems are uncovering universal signatures that connect microscopic quantum dynamics to macroscopic classical chaos. At the most fundamental level, experimental tests of quantum nonlocality continue to refine our understanding of quantum mechanics itself, closing loopholes that have persisted since Bell's seminal work [1].
- This review synthesizes findings from three recent studies that exemplify the breadth of contemporary quantum research. First, we examine comparative studies of simulated versus quantum reverse annealing on D-Wave systems, revealing how phase transitions govern optimization performance and the interplay between quantum coherence and classical relaxation. Second, we explore quantum chaos in periodically driven macrospin ensembles, where Lyapunov exponents serve as a unified diagnostic tool for characterizing dynamical phases and quantum-classical correspondence [2]. Finally, we discuss advances in Bell nonlocality tests that achieve detection efficiency thresholds as low as 50%, enabling loophole-free demonstrations with simplified experimental setups [3].

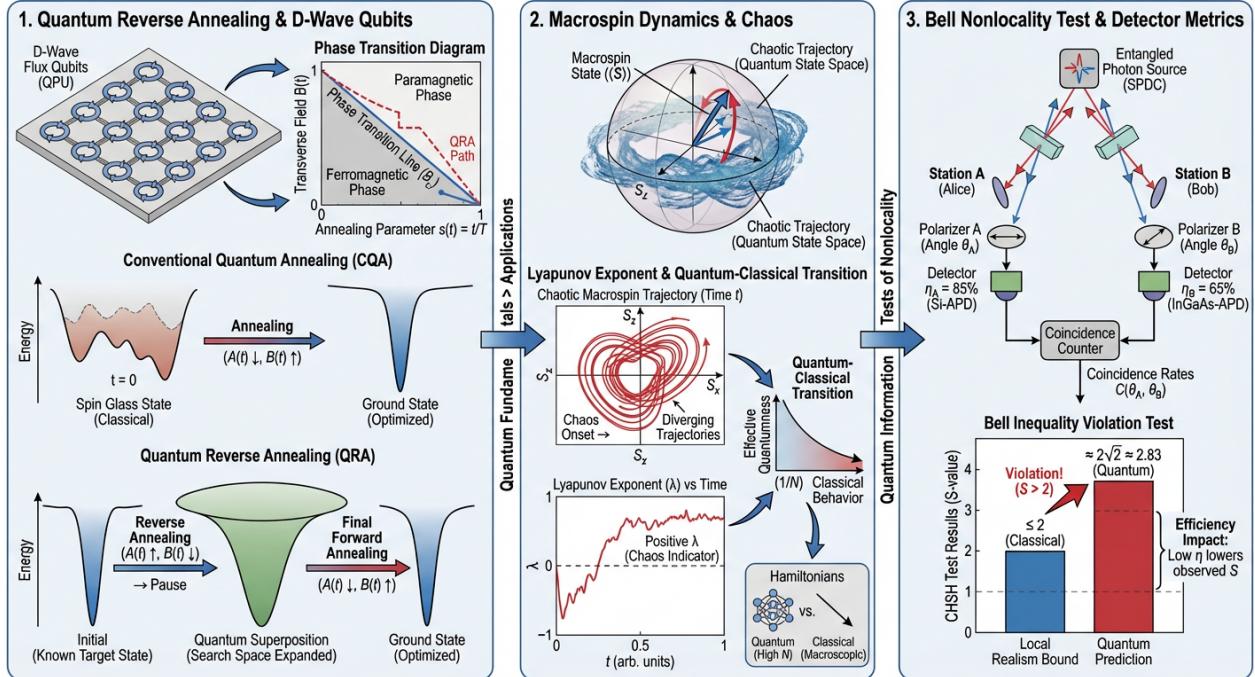


Figure 1: **Graphical Abstract.** Overview of three quantum computing research themes: (1) Quantum reverse annealing on D-Wave systems reveals phase transitions and quantum-classical interplay through systematic probability dynamics and ground state suppression. (2) Macrospin dynamics exhibit quantum chaos characterized by Lyapunov exponents, with quantum-classical correspondence valid only within the Lyapunov time. (3) Bell nonlocality tests achieve detection efficiency thresholds as low as 50% through optimized one-detector steering protocols and loss-tolerant designs.

18 Together, these studies illuminate how quantum systems transition between quantum and classical regimes,  
19 with implications spanning quantum optimization, quantum simulation, and foundational tests of quantum  
20 mechanics.

## 21 2 Quantum Reverse Annealing and Phase Transitions

22 Quantum annealing has emerged as a promising approach for solving combinatorial optimization problems,  
23 with D-Wave systems providing a platform for experimental investigations. Reverse annealing (RA) extends  
24 the standard forward annealing protocol by allowing bidirectional traversal of the annealing schedule, en-  
25 abling local refinement of solutions and exploration of energy landscapes [4]. Recent comparative studies  
26 between simulated and quantum reverse annealing reveal fundamental insights into phase transitions and  
27 the quantum-classical boundary.

28 The reverse annealing protocol initializes the system in a low-energy classical state at the end of the annealing  
29 schedule ( $s = 1$ ), then anneals backward by increasing the transverse field  $A(s)$  while decreasing the problem  
30 Hamiltonian coefficient  $B(s)$  to a reversal point  $s_r$ . After an optional pause time  $t_{\text{wait}}$ , the system anneals  
31 forward back to  $s = 1$ . This process is governed by the time-dependent Hamiltonian  $H(s(t)) = A(s(t))H_D +$   
32  $B(s(t))H_P$ , where  $H_D$  is the driver Hamiltonian and  $H_P$  encodes the optimization problem [5].

33 Experimental studies on D-Wave Advantage systems reveal systematic probability dynamics that challenge  
34 simple quantum speedup narratives. For single-qubit problems, the success probability  $p(t_{\text{end}})$  of maintaining  
35 the initial state decreases with total annealing time, stabilizing at equilibrium values that match thermal  
36 predictions. Non-initial states show the opposite trend, with probabilities increasing over time. These  
37 dynamics persist across problem scales up to 1000 qubits, indicating that classical relaxation processes  
38 dominate at long timescales rather than coherent quantum tunneling [5].

39 Particularly striking is the observation of ground state suppression in reverse annealing for 2-SAT problems.  
40 Unlike forward annealing, which samples ground states according to their degeneracy, reverse annealing  
41 systematically suppresses certain ground states—an effect attributed to perturbative interactions during the  
42 backward phase of the annealing schedule. This finding highlights how the initialization in a classical state  
43 fundamentally alters the exploration of the energy landscape compared to forward annealing from a uniform  
44 quantum superposition.

45 The interplay between quantum and classical effects is further illuminated by comparing reverse annealing  
46 performance across different problem types. While RA excels at local refinement tasks, improving upon  
47 suboptimal solutions in hard combinatorial problems [4], its effectiveness diminishes for global optimization.  
48 The optimal reversal point  $s_r$  typically lies near quantum-classical boundaries, where transverse field fluctua-  
49 tions are maximal but decoherence remains manageable. However, sensitivity to noise and limited scalability  
50 remain significant challenges, with no clear quantum advantage demonstrated for large-scale optimization  
51 tasks.

52 Recent developments in fast-reverse annealing technology announced by D-Wave in 2026 promise enhanced  
53 control for studying optimization phase transitions [6]. By enabling repeatable back-and-forth annealing  
54 while preserving coherence in quantum-dominant regimes, these advances facilitate deeper probing of quan-  
55 tum phenomena. Paired with multicolor annealing capabilities that allow mid-anneal projections and ex-  
56 citations, fast-RA opens new avenues for investigating dynamical state evolution and the nature of phase  
57 transitions in quantum optimization landscapes.

## 58 3 Lyapunov Exponents and Quantum Chaos in Macrospin Dy- 59 namics

60 Understanding the emergence of chaos in quantum systems and its relationship to classical chaos remains  
61 a central challenge in quantum dynamics. Recent work on periodically driven macrospin ensembles demon-

62 strates that Lyapunov exponents serve as a unified diagnostic tool for characterizing dynamical phases across  
63 both classical and quantum regimes [2]. The maximal Lyapunov exponent (MLE), which quantifies the ex-  
64 ponential divergence of nearby trajectories, provides a direct measure of sensitive dependence on initial  
65 conditions—the hallmark of classical chaos.

66 The investigation focuses on a macrospin ensemble with anisotropic long-range all-to-all interactions subject  
67 to periodic driving and collective dissipation, described by a Lindblad master equation. In the thermody-  
68 namic limit ( $N \rightarrow \infty$ ), mean-field treatment yields classical equations of motion whose dynamics are fully  
69 characterized by the MLE. Bifurcation diagrams reveal classic period-doubling routes to chaos and fractal  
70 boundaries separating different attractor regions, establishing the system as a canonical model for studying  
71 quantum-classical correspondence.

72 A critical finding is that quantum and classical dynamics converge only within the Lyapunov time  $t_L \sim$   
73  $1/\lambda_{\max}$ , where  $\lambda_{\max}$  is the maximum Lyapunov exponent [2]. Within this early-time window, exponen-  
74 tial amplification of small perturbations remains negligible, allowing quantum systems to effectively mimic  
75 classical behavior. Beyond the Lyapunov time, however, quantum fluctuations—particularly tunneling be-  
76 tween macrospin states in finite-size systems—cause substantial divergence between quantum and classical  
77 trajectories. This temporal boundary defines the regime of validity for classical descriptions of quantum  
78 systems.

79 Finite-size effects introduce additional complexity. Quantum simulations in the Dicke basis reveal that  
80 finite  $N$  suppresses some chaotic behaviors present in the thermodynamic limit, though both quantum and  
81 classical systems exhibit diverse dynamical phases including periodic, quasiperiodic, and chaotic regimes. The  
82 convergence between quantum and classical dynamics occurs only when nonzero density matrix elements are  
83 sharply localized in phase space. Conversely, in chaotic regimes, the density matrix becomes delocalized,  
84 signaling diffusive exploration of Hilbert space characteristic of quantum chaos [7].

85 The out-of-time-order correlator (OTOC) provides a quantum signature of chaos whose growth rate is dic-  
86 ticated by the classical Lyapunov exponent, establishing a quantitative connection between classical and  
87 quantum chaos [8]. This correspondence enables researchers to use classical dynamical analysis to predict  
88 quantum behavior up to the Lyapunov time, beyond which genuine quantum effects dominate.

89 Extensions to higher-spin quantum systems demonstrate that Lyapunov exponents can be extracted from  
90 spin chains where exponential growth windows open for sufficiently large spin values [9]. In Ising spin chains  
91 with longitudinal and transverse magnetic fields, standard models for quantum chaos studies, strongly chaotic  
92 points exhibit saturation of commutator-squared spin operators before exponential growth manifests in spin-  
93 1/2 systems. By extending to higher spins, researchers numerically extract a quantum Lyapunov exponent  
94 that agrees with the classical infinite-spin limit, rigorously establishing the quantum-classical correspondence.

95 These findings have profound implications for quantum simulation and control of macrospin systems. Under-  
96 standing when and why quantum fluctuations amplify classical chaos effects is essential for designing robust  
97 quantum devices and interpreting quantum simulation experiments. The localization properties of the den-  
98 sity matrix emerge as a key indicator of classical-quantum correspondence, suggesting practical diagnostic  
99 tools for experimental implementations in trapped-ion and superconducting qubit platforms.

## 100 4 Detection Efficiency and Bell Nonlocality

101 Bell's theorem establishes that no local hidden variable theory can reproduce all predictions of quantum  
102 mechanics, with Bell inequality violations serving as experimental signatures of quantum nonlocality [1].  
103 However, the detection efficiency loophole has long challenged experimental tests: imperfect photon detectors  
104 allow local hidden variable theories to mimic quantum violations under the fair-sampling assumption. Closing  
105 this loophole requires detection efficiencies exceeding specific thresholds, historically estimated at 82.8% for  
106 maximally entangled states.

107 A landmark result by Eberhard established that using non-maximally entangled states dramatically re-  
108 duces the minimum detection efficiency (MDE) threshold to 66.7% (2/3), far below the requirement for

109 maximally entangled states [3]. This "anomaly"—where entanglement reduction improves detection require-  
110 ments—persists in Bell tests but vanishes in quantum steering scenarios, revealing subtle distinctions between  
111 these nonlocality frameworks.

112 Recent advances have pushed detection thresholds even lower through loss-tolerant protocols. One-detector  
113 steering inequalities achieve an unprecedented MDE of 50% for any pure entangled qubit state, independent  
114 of entanglement degree [10]. This simplified setup requires only moderate efficiency ( $\eta > 50\%$ ) on the  
115 untrusted side while allowing arbitrarily low efficiency on the trusted side. Experimental demonstrations  
116 have achieved steering violations exceeding 21 standard deviations at  $\eta = 51.6 \pm 0.4\%$  detection efficiency,  
117 closing the Eberhard anomaly for steering scenarios.

118 The methodology optimizes both quantum states and one-click measurements to maximize the witness-to-  
119 noise ratio (WNR), enabling robustness against realistic noise sources. This represents the simplest possible  
120 photonic setup with complexity  $W = 2A = 2B = X$ , making it highly scalable for quantum network  
121 implementations. Numerical simulations confirm robustness to various noise models, validating the practical  
122 applicability of these protocols [10].

123 Alternative approaches leverage N00N states in triangle network configurations, demonstrating exceptional  
124 robustness to 10% photon loss in transmission channels [11]. Using neural network heuristics to optimize  
125 for realistic noise, these protocols employ heralding from spontaneous parametric down-conversion (SPDC)  
126 sources to avoid global post-processing, fully closing the detection loophole without discarding events. The  
127 triangle network topology certifies nonlocality without input settings, enabling self-testing with unreliable  
128 sources—a significant advance over traditional two-party Bell tests.

129 Experimental closure of the detection loophole has been achieved in photonic systems using high-efficiency  
130 superconducting nanowire single-photon detectors (SNSPDs). NIST researchers demonstrated detection-  
131 loophole-free Bell violations exceeding 50 standard deviations using entangled photons from heralded sources  
132 with efficiencies surpassing Eberhard's 66.7% threshold [12]. These demonstrations validate photonic ap-  
133 proaches for device-independent quantum key distribution and other quantum information protocols requir-  
134 ing loophole-free nonlocality certification.

135 Multipartite extensions compute minimum detection efficiencies for genuine multipartite nonlocality, includ-  
136 ing M-type and Svetlichny nonlocality [13]. State-independent approaches significantly reduce previously  
137 established bounds, demonstrating robustness for imperfect detectors in noisy channels. While multipartite  
138 loophole-free tests remain experimentally challenging, theoretical frameworks now provide clear targets for  
139 next-generation experiments.

140 The continued refinement of detection efficiency thresholds and development of loss-tolerant protocols rep-  
141 resents a crucial step toward practical quantum technologies. By enabling loophole-free tests with realistic  
142 experimental constraints, these advances strengthen the empirical foundation of quantum mechanics while  
143 enabling secure quantum communication protocols that depend on certified quantum nonlocality.

## 144 5 Discussion and Outlook

145 The three studies reviewed here illuminate a common theme: the intricate boundary between quantum and  
146 classical behavior governs both practical quantum technologies and fundamental tests of quantum mechanics.  
147 Across disparate systems—quantum annealers, macrospin ensembles, and entangled photon pairs—temporal  
148 and energetic scales determine when quantum effects dominate and when classical descriptions suffice.

149 In quantum annealing, the dominance of classical relaxation over quantum tunneling at long timescales re-  
150 veals fundamental limitations for optimization speedup. The systematic ground state suppression observed  
151 in reverse annealing highlights how initialization conditions propagate through the annealing schedule, fun-  
152 damentally altering solution sampling compared to forward annealing. These findings suggest that quantum  
153 advantage in optimization may be restricted to specific problem classes and parameter regimes where quan-  
154 tum coherence persists long enough to influence outcomes.

155 The Lyapunov time boundary in macrospin systems provides a quantitative criterion for quantum-classical  
156 correspondence. Within  $t_L \sim 1/\lambda_{\max}$ , classical predictions accurately describe quantum dynamics, but  
157 beyond this timescale, quantum fluctuations—particularly tunneling in finite-size systems—cause divergence.  
158 This temporal boundary has direct implications for quantum simulation: classical models can guide quantum  
159 simulations up to the Lyapunov time, beyond which genuine quantum effects must be explicitly modeled.  
160 The density matrix localization emerges as a practical diagnostic for identifying when quantum-classical  
161 correspondence breaks down.

162 Detection efficiency advances in Bell tests exemplify how optimizing measurement protocols and entangle-  
163 ment properties can dramatically reduce experimental requirements. The reduction of minimum detection  
164 efficiency from 82.8% to 50% through careful state engineering and one-detector steering protocols makes  
165 loophole-free tests accessible to a broader range of experimental platforms. This progression mirrors a broader  
166 trend in quantum technology: sophisticated protocol design compensates for hardware imperfections.

167 Several cross-cutting challenges emerge. First, decoherence and noise fundamentally limit all three systems.  
168 In quantum annealing, noise raises effective energy barriers and reduces coherence times. In macrospin  
169 systems, dissipation couples quantum and classical dynamics in non-trivial ways. In Bell tests, photon  
170 loss and detector dark counts constrain achievable violation magnitudes. Second, scalability remains elusive:  
171 thousand-qubit annealers still exhibit classical-dominated dynamics, multipartite Bell tests lack full loophole  
172 closure, and macrospin simulations face exponential Hilbert space growth.

173 Future research directions include developing hybrid quantum-classical algorithms that exploit the strengths  
174 of both regimes. For quantum annealing, fast-reverse protocols combined with machine learning solvers  
175 may identify problem structures amenable to quantum speedup. For quantum chaos, experimental plat-  
176 forms in trapped ions and superconducting qubits can test theoretical predictions of Lyapunov exponent  
177 correspondence. For Bell tests, satellite-based experiments and quantum networks can extend loophole-free  
178 demonstrations to global scales.

179 The convergence of these research directions toward understanding quantum-classical boundaries suggests  
180 a unified framework may emerge. Concepts like the Lyapunov time, coherence timescales in annealing,  
181 and detection efficiency thresholds all quantify regimes where quantum effects persist despite environmental  
182 coupling. Elucidating these boundaries will be essential for realizing practical quantum technologies while  
183 deepening our understanding of quantum mechanics itself.

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