

1 Electrons in a metal as an ideal Fermi gas (7 points)

A metal contains conduction electrons in a volume V . The temperature is $T = 0$ so that the chemical potential is $\mu = \varepsilon_F$. The magnetic moment of the electrons is μ_m . A magnetic field B is switched on, hence the energy of particles with spin parallel (+) and anti-parallel (-) to the magnetic field is

$$\varepsilon_{\pm}(p) = \frac{p^2}{2m} \mp \mu_m B \quad (1)$$

respectively. In the following the magnetic field is assumed so weak that only its first order effect should be considered.

a) Calculate the mean number $\langle N_{\pm} \rangle$ of electrons with spin parallel and anti-parallel to the magnetic field, respectively. (3 points)

$$Z = Z_+ + Z_- \quad (2)$$

$$Z_{\pm} = \int_0^{\infty} d\varepsilon \cdot \left(1 + \exp(-\beta\varepsilon) \right) \cdot n_{\pm}(\varepsilon) \quad (3)$$

$$\langle N \rangle = -\frac{1}{\beta} \partial_{\mu} \ln(Z) \quad (4)$$

b) Calculate the mean magnetization M and the magnetic susceptibility χ . (2 points)

$$M = \mu_m \cdot (\langle N_+ \rangle - \langle N_- \rangle) \quad (5)$$

$$\chi = \left(\frac{\partial M}{\partial B} \right) \quad (6)$$

c) Express the chemical potential μ in terms of the mean total number of electrons $\langle N \rangle$ and eliminate the chemical potential from the formula for the susceptibility. (2 points)

$$\langle N \rangle = \langle N_+ \rangle + \langle N_- \rangle \quad (7)$$

Q
1 0/7
2 0/4
3 4/4
Σ 4/15

0/7

2 Particle-hole symmetry for fermions (4 points)

Consider a system of spin 1/2 fermions at temperature T , that occupy a finite number M of energy levels ε_i with $i = 1, 2, \dots, M$.

a) Determine the grand canonical partition sum Z_G . (1 point)

HINT: Pay attention to the degeneracy.

$$Z_G = \sum_{i=1}^M \left[1 + 2e^{-\beta(\varepsilon_i - \mu)} + e^{-\beta(2\varepsilon_i - 2\mu)} \right] \quad (8)$$

b) Use the relation $F = -k_B T \ln Z_G + \mu \langle N \rangle$ between the free energy, the grand canonical partition sum, the chemical potential and the mean particle number to show: the thermodynamic properties of the system stay the same if one, instead of considering the $\langle N \rangle$ particles, distributes $2M - \langle N \rangle$ "holes" with chemical potential μ on the energy levels $-\varepsilon_i$.

→ for each state ε_i , have

spin states \emptyset (no particle)

\uparrow, \downarrow (2x single particle state)

$\uparrow\downarrow$ (2 particle state)

$$Z_i = 1 + 2e^{-\beta(\varepsilon_i - \mu)} + e^{-\beta(2\varepsilon_i - 2\mu)}$$

$$= (1 + e^{-\beta(\varepsilon_i - \mu)})^2$$

0/4

3 Ultra-relativistic ideal Fermi fluid (4 points)

An ultra-relativistic ideal Fermi fluid is contained in a volume V . The chemical potential is μ . Consider only the special case $T = 0$ in the following. Calculate the mean particle number $\langle N \rangle$ and the mean total energy $\langle E \rangle$. Express $\langle E \rangle$ in terms of $\langle N \rangle$ and μ .

HINT: Ultra-relativistic means that you can neglect the mass term in the relativistic formula for the energy of the particles:

$$\varepsilon(\vec{p}) = \sqrt{m^2 c^4 + \vec{p}^2 c^2} \approx pc. \quad (9)$$

We use the following identity from the lecture script:

$$\sum_{\vec{k}, m_s} (...) = 2 \cdot \frac{V}{h^3} \int d\vec{p} (...) = \frac{2V}{h^3} \int_0^\infty dp 4\pi p^2 (...). \quad (10)$$

Thus, the mean particle number is given by

$$\langle N \rangle = \sum_{\vec{k}, m_s} n_{\vec{k}, m_s} = \frac{8\pi V}{c^3 h^3} \int_0^\infty d\varepsilon \varepsilon^2 n(\varepsilon), \quad (11)$$

where we used the linear dispersion relation $\varepsilon(p) = pc$ for ultra-relativistic particles. In the limit $T = 0$, the Fermi function becomes $n(\varepsilon) = \theta(\mu - \varepsilon)$:

$$\langle N \rangle = \frac{8\pi V}{c^3 h^3} \int_0^\mu d\varepsilon \varepsilon^2 = \frac{8\pi V}{3c^3 h^3} \mu^3. \quad (12)$$

A similar calculation is done for the mean energy:

$$\langle E \rangle = \frac{8\pi V}{c^3 h^3} \int_0^\infty d\varepsilon \varepsilon^2 n(\varepsilon) \varepsilon = \frac{8\pi V}{c^3 h^3} \int_0^\mu d\varepsilon \varepsilon^3 = \frac{2\pi V}{c^3 h^3} \mu^4 = \frac{3}{4} \langle N \rangle \mu. \quad (13)$$

4/4

$2 \mu^4$