

We have coordinate:  $(r, \theta, \phi)$

$$r > 0, \quad 0 \leq \theta < \pi, \quad 0 \leq \phi < 2\pi$$



lets define  $(u, v, w)$  to be components in  $r, \theta, \phi$ .

Each grid  $(r_i, \theta_i, \phi_i)$  is the center of the cell.

A cell is given by  $r_{i-1/2} \leq r < r_{i+1/2}$ .

$$\theta_{j-1/2} \leq \theta < \theta_{j+1/2}$$

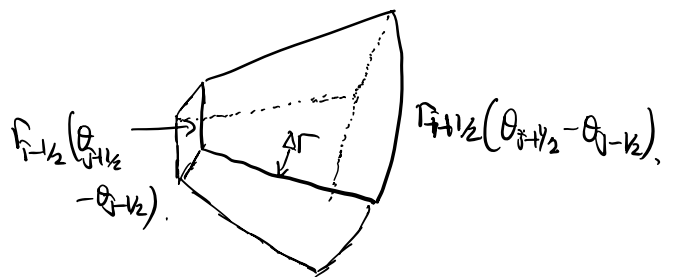
$$\phi_{k-1/2} \leq \phi < \phi_{k+1/2}$$



The cell walls in  $r$  and  $\theta$  direction are curved,  
but straight in  $\phi$  direction

Given a scalar quantity:  $q$ :

$$\frac{\partial q}{\partial t} + \vec{\nabla}(q\vec{u}) = 0$$



$$\hookrightarrow = \frac{\partial q}{\partial t} + \frac{1}{r^2} \frac{\partial(r^2 q u)}{\partial r} + \frac{1}{r \sin \theta} \frac{\partial(\sin \theta q v)}{\partial \theta} + \frac{1}{r \sin \theta} \frac{\partial(q w)}{\partial \phi} = 0$$

We see there are geometric factors?

Given cell  $(i, j, k)$ :  $r_{i-1/2} \leq r < r_{i+1/2}$ ,  $\theta_{j-1/2} \leq \theta < \theta_{j+1/2}$

$$\phi_{k-1/2} \leq \phi < \phi_{k+1/2}$$

Volume:  $V_{ijk} = \frac{1}{3}(\bar{r}_{i+1/2}^3 - \bar{r}_{i-1/2}^3)(\cos\theta_{j-1/2} - \cos\theta_{j+1/2})(\phi_{k+1/2} - \phi_{k-1/2})$

The  $i+1/2$  surface of the cell in  $r$ -direction (i.e. when  $r = \text{constant}$ )

$$S_{(r), i+1/2, j, k} = \bar{r}_{i+1/2}^2 (\cos\theta_{j-1/2} - \cos\theta_{j+1/2})(\phi_{k+1/2} - \phi_{k-1/2})$$

Similarly: The  $\theta+1/2$  surface of the cell in  $\theta$ -direction (i.e.  $\theta = \text{constant}$ )

$$S_{(\theta), i, j+1/2, k} = \frac{1}{2}(\bar{r}_{i+1/2}^2 - \bar{r}_{i-1/2}^2) \sin\theta_{j+1/2} (\phi_{k+1/2} - \phi_{k-1/2})$$

For  $k+1/2$ :

$$S_{(\phi), i, j, k+1/2} = \frac{1}{2}(\bar{r}_{i+1/2}^2 - \bar{r}_{i-1/2}^2)(\theta_{j+1/2} - \theta_{j-1/2})$$

If we difference equations in conservative form:

$$\begin{aligned} \frac{\partial(\rho_{ijk} V_{ijk})}{\partial t} &= F_{i-1/2, j, k} S_{(r), i-1/2, j, k} - F_{i+1/2, j, k} S_{(r), i+1/2, j, k} \\ &+ F_{i, j-1/2, k} S_{(\theta), i, j-1/2, k} - F_{i, j+1/2, k} S_{(\theta), i, j+1/2, k} \\ &+ F_{i, j, k-1/2} S_{(\phi), i, j, k-1/2} - F_{i, j, k+1/2} S_{(\phi), i, j, k+1/2} \end{aligned}$$

$$\frac{\partial(\rho_{ijk} V_{ijk})}{\partial t} + (\Delta_r F_r S_r)_{i, j, k} + (\Delta_\theta F_\theta S_\theta)_{i, j, k} + (\Delta_\phi F_\phi S_\phi)_{i, j, k} = 0$$

↑  
Above works for scalar quantity like  $p$  and  $e_{bt}$ .

## ★ Euler Equation in Polar Coordinates:

Density:

$$\frac{\partial \rho}{\partial t} + \vec{\nabla} \cdot (\rho \vec{u}) = 0$$

$$\hookrightarrow \frac{\partial \rho}{\partial t} + \frac{1}{r^2} \frac{\partial (r^2 \rho u)}{\partial r} + \frac{1}{r \sin \theta} \frac{\partial (\sin \theta \rho v)}{\partial \theta} + \frac{1}{r \sin \theta} \frac{\partial (\rho w)}{\partial \phi} = 0$$

Energy:

$$\frac{\partial (\rho e)}{\partial t} + \vec{\nabla} \cdot (\rho e \vec{u} + p \vec{u}) = 0$$

$$\hookrightarrow \frac{\partial (\rho e)}{\partial t} + \frac{1}{r^2} \frac{\partial (r^2 (\rho e + p) u)}{\partial r} + \frac{1}{r \sin \theta} \frac{\partial (\sin \theta (\rho e + p) v)}{\partial \theta} + \frac{1}{r \sin \theta} \frac{\partial ((\rho e + p) w)}{\partial \phi} = 0$$

Momentum Equation:  $\frac{\partial (\rho \vec{u})}{\partial t} + \vec{\nabla} \cdot (\rho \vec{u} \otimes \vec{u}) + \nabla p = 0$

$r$ -radial:

$$\frac{\partial (\rho u)}{\partial t} + \frac{1}{r^2} \frac{\partial (r^2 \rho u^2)}{\partial r} + \frac{\partial p}{\partial r} + \frac{1}{r \sin \theta} \frac{\partial (\sin \theta \rho u v)}{\partial \theta}$$

$$+ \frac{1}{r \sin \theta} \frac{\partial(\rho u w)}{\partial \phi} - \frac{\rho(v^2 + w^2)}{r} = 0$$

In discrete form for Riemann solver:

$$\begin{aligned} \frac{\partial(\rho u V)}{\partial t} + \Delta_r(F_{rr} S_r) + \Delta_\theta(F_{r\theta} S_\theta) + \Delta_\phi(F_{r\phi} S_\phi) \\ = \left[ \frac{2P}{r} + \frac{\rho(v^2 + w^2)}{r} \right] V \end{aligned}$$

where 
$$\left. \begin{aligned} F_{rr} &= \rho u^2 + P \\ F_{r\theta} &= \rho u v \\ F_{r\phi} &= \rho u w \end{aligned} \right\} \text{Momentum Fluxes.}$$

$\theta$ -Equation:

$$\begin{aligned} \frac{\partial(\rho v)}{\partial t} + \frac{1}{r^2} \frac{\partial(r^2 \rho w)}{\partial r} + \frac{1}{r \sin \theta} \frac{\partial(\sin \theta \rho v^2)}{\partial \theta} \\ + \frac{1}{r} \frac{\partial P}{\partial \theta} + \frac{1}{r \sin \theta} \frac{\partial(\rho w v)}{\partial \phi} + \rho \frac{w v}{r} \\ - \frac{\cos \theta}{\sin \theta} \rho \frac{w^2}{r} = 0 \end{aligned}$$

In discrete Riemann solver:

$$\frac{\partial(\rho v V)}{\partial t} + \Delta_r(F_{\theta r} S_r) + \Delta_\theta(F_{\phi\theta} S_\theta) + \Delta_\phi(F_{\phi\phi} S_\phi)$$

$$= - \left[ \rho \frac{vu}{r} - \frac{\cos\theta}{\sin\theta} \frac{\rho w^2 + p}{r} \right] V$$

for  $F_{\theta r} = \rho v u$      $F_{\phi\theta} = \rho v^2 + p$  ,  $F_{\phi\phi} = \rho v w$

$\phi$ -momentum equation:

$$\frac{\partial(\rho w)}{\partial t} + \frac{1}{r^2} \frac{\partial(r^2 \rho u w)}{\partial r} + \frac{1}{r \sin\theta} \frac{\partial(\sin\theta \rho v w)}{\partial \theta}$$

$$+ \frac{1}{r \sin\theta} \frac{\partial(\rho w^2 + p)}{\partial \phi} + \frac{\rho u w}{r} + \frac{\cos\theta}{\sin\theta} \frac{\rho v w}{r} = 0$$

In discrete Riemann solver:

$$\frac{\partial(\rho w V)}{\partial t} + \Delta_r(F_{\phi r} S_r) + \Delta_\theta(F_{\phi\theta} S_\theta) + \Delta_\phi(F_{\phi\phi} S_\phi)$$

$$= - \left[ \frac{\rho u w}{r} + \frac{\cos\theta}{\sin\theta} \frac{\rho v w}{r} \right] V$$

$F_{\phi r} = \rho w u$      $F_{\phi\theta} = \rho w v$      $F_{\phi\phi} = \rho w^2 + p$

procedure in applying Riemann solver to polar coordinates:

1) precompute cell volumes, cell surfaces and cosine and sine values.

2) Do strong-splitting.

3) Compute left and right interface states, and feed them to Riemann solvers.

4) Add source terms if needed?

Replacing  $\phi$ -momentum with angular momentum in  $r, \theta$  advection

$$\text{let } l = w \sin \theta r$$

If we replace  $w$  by  $l$

$$\frac{\partial(lv)}{\partial t} + \frac{\partial(lvl)}{\partial r} + \frac{1}{r} \frac{\partial(lvl)}{\partial \theta} + \frac{\partial(lw^2 + p)}{\partial \phi}$$

$$+ 2 \sin \theta l w + \cos \theta p v w = 0$$

then rewrite:

$$\frac{\partial(lv)}{\partial t} + \frac{1}{r^2} \frac{\partial(r^2 lvl)}{\partial r} + \frac{1}{r \sin \theta} \frac{\partial(\sin \theta lvl)}{\partial \theta} + \frac{\partial(lw^2 + p)}{\partial \phi} = 0$$

then discrete Riemann solver has form:

$$\frac{\partial(lvl)}{\partial t} + \Delta_r(r \sin \theta F_{\phi r} S_r) + \Delta_{\theta}(r \sin \theta F_{\phi \theta} S_{\theta}) + \Delta_{\phi}(r \sin \theta F_{\phi \phi} S_{\phi}) = 0$$

↑

Then

No more

$$\tilde{F}_{\phi r} = r \sin \theta F_{\phi r} = l u l$$

Source  
terms.

$$\tilde{F}_{\phi \theta} = r \sin \theta F_{\phi \theta} = l v l$$

$$\tilde{F}_{\phi \phi} = r \sin \theta F_{\phi \phi} = l u l + P r \sin \theta$$