

RECONSTRUCTION D'IMAGE AVEC LES MACHINES À INFÉRENCES RÉCURRENTIELLES

par

Alexandre Adam

Mémoire présenté en vue de l'obtention du grade de
Maître ès sciences (M.Sc.)

Département de physique
Université de Montréal

Résumé

Abstract

Table des matières

Résumé	ii
Abstract	iii
Liste des tableaux	vi
Liste des figures	vii
Remerciements	x
1 Introduction	1
1.1 Lentilles gravitationnelles de type Galaxie-Galaxie	1
1.1.1 Contexte historique	1
1.1.2 Les angles de déflections	1
1.1.3 Contexte scientifique moderne	5
1.2 Interférométrie par masque non-régulier	5
1.2.1 Contexte historique	5
1.2.2 Les angles de fermeture	5
1.2.3 Contexte scientifique moderne	5
1.3 Auto-encodeur variationnel	5
1.3.1 Description du modèle	5
1.3.2 Le truc de reparamétrisation	7
1.4 Machines à inférence récurrentielles	8
1.4.1 Formalisme des problèmes inverses	8
1.4.2 La relation de récurrence	8

2 Pixelated Reconstruction of Gravitational Lenses using Recurrent Inference Machine	9
Abstract	9
2.1 Introduction	11
2.2 Methods	13
2.2.1 Maximum a posteriori	13
2.2.2 Recurrent Inference Machine	14
2.2.3 The Gradient Model	16
2.2.4 The Forward Model	17
2.2.5 Fine-Tuning	18
2.3 Data	20
2.3.1 COSMOS	20
2.3.2 IllustrisTNG	21
2.3.3 Simulated Observations	23
2.4 Training	24
2.4.1 VAE	24
2.4.2 RIM	25
2.5 Results	26
2.5.1 Goodness of Fit	27
2.5.2 Quality of the Reconstructions	31
2.6 Conclusion	31
Bibliographie	33
A Elastic Weight Consolidation	43
B VAE Architecture and optimisation	45
C RIM architecture and optimisation	48

Liste des tableaux

2.1	Physical model parameters.	24
2.2	SIE parameters.	24
2.3	Hyperparameters for fine-tuning the RIM.	27
2.4	\log_{10} -normal moments of the loss on the test set	28
B.1	Hyperparameters for the background source VAE.	46
B.2	Hyperparameters for the convergence VAE.	47
C.1	Hyperparameters for the RIM.	50

Liste des figures

1.1	Schéma d'une lentille gravitationnelle.	5
1.2	Modèle graphique d'un VAE. Les flèches pleines indiquent le processus génératif, alors que les flèches pointillées indiquent le processus d'inférence.	6
2.1	Example of a simulated lensed image in the test set that exhibits a large deflection in its eastern arc which indicates the presence of a massive object — in this case a dark matter subhalo. The fine-tuning procedure is able to recover this subhalo because of its strong signal in the lensed image and reduces the residuals to noise level.	15
2.2	Rolled computational graph of the RIM. Dashed arrows represent operations not recorded for BPTT.	16
2.3	A single time step of the unrolled computation graph of the RIM. GRU units are placed in the skip connections to guide the reconstruction of the source and convergence. A schematic of the steps to compute the likelihood gradients is shown in the bottom right of the figure, including the Adam processing step of the likelihood gradient.	17
2.4	Examples similar to the test task, also shown in Figure 2.7. The first column shows the ground truth used to simulate the lensed image. The second column shows the baseline prediction that is then encoded in the latent space of the VAE in order to sample the next 4 columns.	20
2.5	Examples of COSMOS galaxy images (top row) and VAE generated samples (bottom row) used as labels in \mathcal{D}	21
2.6	Examples of smoothed Illustris TNG100 convergence map (top row) and VAE generated samples (bottom row) used as labels in \mathcal{D}	22
2.7	Cherry-picked sample of the fine-tuned RIM reconstructions on a test set of 3000 examples. Examples are ordered from the best χ^2 (top) to the worst (bottom). The percentile rank of each example is in the leftmost column. The last example shown has SNR above the threshold defined in Figure 2.9.	26

2.8	Distribution of the goodness of fit for the baseline and fine-tuned network (right panel), as well as log-loss difference between the two network for a given example in the test set (left panel).	27
2.9	Goodness of fit as a function of SNR shows a threshold behavior where our method reaches its limit.	29
2.10	Comparison between baseline (RIM) and fine-tuned (RIM+FT) reconstructions for VAE generated gravitational lensing systems. From top to bottom, we increase SNR. The first 2 rows have noise level reconstruction, while the last 3 row show significant improvement over the baseline. The intensity color scale is chosen to show the reconstruction down to the third decimal place, where the baseline prediction breaks down.	30
2.11	Statistics of the coherence spectrum on the test set. The solid line is the average coherence. The transparent region is the 68% confidence interval. The fine-tuning procedure yields a noticeable improvement on the coherence of the source at all frequencies.	30
C.1	Unrolled computational graph of the RIM. Operations along solid arrows are being recorded for BPTT, while operations along dashed arrows are not. The blue arrows are only used for optimisation during training. During fine-tuning or testing, the loss is computed only as an oracle metric to validate that our methods can recover the ground truth.	49
C.2	30 reconstructions taken at random from the test set of 3000 examples simulated from COSMOS and IllustrisTNG data at high SNR. The colorscale are the same as in Figure 2.7.	51

À Maman et Julia

Remerciements

Chapitre 1

Introduction

1.1 Lentilles gravitationnelles de type Galaxie-Galaxie

1.1.1 Contexte historique

L'idée des lentilles gravitationnelles est attribuée à Fritz Zwicky (1937) qui, suivant les calculs publié par Einstein (1936), est le premier à postuler correctement que l'anneau d'Einstein, produit par la déflection de la lumière d'une source lointaine par le champ gravitationnel d'une galaxie (appelée nébuleuse extragalactique à l'époque) en avant-plan de cette source selon le point de vue d'un observateur sur Terre, pourrait être observé. Dans le même article, Zwicky articule précisément les idées qui nous motivent encore aujourd'hui (presque 100 ans plus tard) à étudier ces objets, c'est-à-dire que les lentilles gravitationnelles permettraient

1. d'imager des galaxies trop lointaine pour que l'on puisse les résoudre avec nos télescopes ;
2. de mesurer directement la masse gravitationnelle de ces galaxies.

Il est à noté qu'Einstein considérait la possibilité d'observer ce phénomène extrêmement improbable. La résolution des télescopes optiques durant la majorité du XX^e siècle étaient limités à environ 0.5 arcsecondes par la turbulence de l'air. Avec une telle résolution, un anneau d'Einstein d'une taille caractéristique de 1 arcseconde apparaîtrait comme un point lumineux étalé, et ne serait donc pas distinguable d'une étoile.

1.1.2 Les angles de déflections

Dans les paragraphes qui suivent, je dérive les équations centrales qui nous permettent d'étudier les lentilles gravitationnelles de type galaxie-galaxie. Des traitements similaires peuvent être trouvé dans les manuels de références de Meneghetti (2013) et Carroll (2003).

Supposons qu'un photon est sur une trajectoire parallèle à l'axe de visée \mathbf{e}_{\parallel} d'un observateur sur Terre. Supposons de plus que la source d'un champ gravitationnel Φ est situé sur l'axe de visée,

ce qui a pour effet de courber la trajectoire de ce photon entre son point d'origine A et son point d'arrivée B . On définit l'angle de déviation comme la déviation totale de cette trajectoire dans la direction perpendiculaire à l'axe de visée de l'observateur. De façon générale, cette déviation s'écrit

$$\boldsymbol{\alpha} = - \int_{\lambda_A}^{\lambda_B} \ddot{\mathbf{x}} \times \mathbf{e}_{\parallel} d\lambda, \quad (1.1)$$

où λ paramétrise la trajectoire du photon $\mathbf{x}(\lambda)$. Le signe négatif nous indique qu'on prend la perspective de l'observateur.

La trajectoire d'un photon est sujette au principe de Fermat, qui stipule que la lumière suit une trajectoire qui extrémise la durée du parcours entre deux points. Dans le langage du calcul des variations, la variation de la durée s'écrit

$$\delta T = \delta \int_A^B n(\mathbf{x}(\ell)) \frac{d\ell}{c} = 0, \quad (1.2)$$

où ℓ est un élément de longueur sur la trajectoire et n est un indice de réfraction. Pour déterminer l'indice de réfraction du champ gravitationnel d'une galaxie, on doit utiliser le formalisme de la relativité générale. Selon le principe d'équivalence (fort), l'effet d'un champ gravitationnel est localement indistinguables d'une accélération causée par la courbure d'un espace-temps décrite par une métrique $g_{\mu\nu}$. La trajectoire d'un photon se trouve alors en cherchant les géodésiques de cet espace-temps. On fait l'approximation que le potentiel Φ d'une galaxie est celui d'un gaz parfait, c'est-à-dire qu'il satisfait une équation de Poisson

$$\nabla^2 \Phi = 4\pi G \rho. \quad (1.3)$$

Dans la limite où ce potentiel est faible $\frac{2\Phi}{c^2} \ll 1$, la métrique $g_{\mu\nu}$ est décrite par une expansion au premier ordre autour de la métrique de Minkowsky

$$ds^2 = g_{\mu\nu} dx^\mu dx^\nu \approx \left(1 + \frac{2\Phi}{c^2}\right) c^2 dt^2 - \left(1 - \frac{2\Phi}{c^2}\right) d\mathbf{x}^2. \quad (1.4)$$

Puisqu'un photon suit une géodésique de l'espace-temps $ds^2 = 0$, on peut déterminer l'indice de réfraction en réarrangeant l'équation (1.4)

$$n \equiv c \left(\frac{\|d\mathbf{x}\|}{dt} \right)^{-1} \approx 1 - \frac{2\Phi}{c^2}. \quad (1.5)$$

En réécrivant l'élément de longueur $d\ell$ en terme du paramètre de la trajectoire $d\ell = \|\frac{d\mathbf{x}}{d\lambda}\| d\lambda$, on peut réécrire l'équation (1.2) sous la forme

$$\delta \int_{\lambda_A}^{\lambda_B} n(\mathbf{x}) \|\dot{\mathbf{x}}\| d\lambda = 0. \quad (1.6)$$

Par correspondance avec la fonctionnelle de l'action $J(x) = \int_{\lambda_0}^{\lambda_1} \mathcal{L}(\lambda, x, \dot{x}) d\lambda$ on trouve que le lagrangien de la trajectoire s'écrit $\mathcal{L} = n(\mathbf{x})\sqrt{\dot{x}^2}$. La trajectoire qui satisfait (1.2) est une solution des équations d'Euler-Lagrange

$$\frac{d}{d\lambda} \frac{\partial \mathcal{L}}{\partial \dot{\mathbf{x}}} - \frac{\partial \mathcal{L}}{\partial \mathbf{x}} = 0. \quad (1.7)$$

On a donc

$$\frac{d}{d\lambda} n \frac{\dot{\mathbf{x}}}{\|\dot{\mathbf{x}}\|} - \|\dot{\mathbf{x}}\| \nabla n = 0, \quad (1.8)$$

Puisque le choix du paramètres λ est libre, on peut le choisir tel que $\|\dot{\mathbf{x}}\| = 1$ en tout point de la trajectoire. Ainsi,

$$\begin{aligned} \frac{d}{d\lambda} n \dot{\mathbf{x}} - \nabla n &= 0 \\ \implies n \ddot{\mathbf{x}} + (\nabla n \cdot \dot{\mathbf{x}}) \dot{\mathbf{x}} - \nabla n &= 0 \end{aligned} \quad (1.9)$$

À ce point de la dérivation, on utilise l'approximation de Born. C'est-à-dire qu'on approxime la trajectoire du photon comme une ligne droite sur l'axe de visée \mathbf{e}_{\parallel} . Cette approximation est justifiée dans le contexte des lentilles gravitationnelles de type galaxie-galaxie, puisque les angles de déviation sont généralement de l'ordre de l'arcseconde ou plus petit. Comme, le vecteur $\dot{\mathbf{x}}$ est tangent à la trajectoire du photon, on obtient

$$\ddot{\mathbf{x}} \times \mathbf{e}_{\parallel} = \frac{1}{n} \nabla_{\perp} n = \nabla_{\perp} \log n \approx -\frac{2}{c^2} \nabla_{\perp} \Phi, \quad (1.10)$$

où ∇_{\perp} est un gradient selon les coordonnées perpendiculaires à \mathbf{e}_{\parallel} . On note que le facteur 2 qui apparaît dans l'équation (1.10) est un effet qui vient de la relativité générale. Ce facteur corrige la solution que l'on aurait obtenu avec une dérivation classique (newtonienne).

On est maintenant en mesure de calculer l'angle de déviation. J'introduit le paramètre d'impact ξ qui est la distance perpendiculaire entre la position d'origine du photon sur le plan de la lentille et l'axe de visé (voir Figure 1.1). Dans le cas où le potentiel est généré par une masse M ponctuelle, c.-à-d. qu'on suppose $\rho = M\delta^3(\mathbf{x})$, où δ est la fonction delta de Dirac, alors le potentiel qui satisfait l'équation de Poisson (1.3) est la fonction de Green $\Phi = -\frac{GM}{\sqrt{\xi^2 + z^2}}$, où z est la coordonné sur l'axe de visée. L'équation (1.1) se réécrit finalement comme

$$\begin{aligned} \alpha(\xi) &= -\frac{2GM}{c^2} \int_{-\infty}^{\infty} \frac{\partial}{\partial \xi} \frac{1}{(\xi^2 + z^2)^{1/2}} dz \\ \implies \alpha(\xi) &= \frac{4GM}{c^2 \xi^2} \xi \end{aligned} \quad (1.11)$$

Cette solution se généralise naturellement à un profil de masse quelconque en assumant qu'il s'exprime comme une somme d'élément de masses $dm = \Sigma d^2 \xi'$, où $\Sigma = \int \rho dz$ est un densité surfacique de masse. L'angle de déviation total mesuré à un point ξ est alors une convolution sur tout le plan

de la lentille (mince) puisque l'équation (1.11) dépend linairement de la masse M :

$$\boldsymbol{\alpha}(\boldsymbol{\xi}) = \frac{4G}{c^2} \int_{\mathbb{R}^2} \Sigma(\boldsymbol{\xi}') \frac{\boldsymbol{\xi} - \boldsymbol{\xi}'}{\|\boldsymbol{\xi} - \boldsymbol{\xi}'\|^2} d^2 \boldsymbol{\xi}' \quad (1.12)$$

L'angle de déviation est une quantité cruciale pour résoudre une lentille gravitationnelle puisqu'il décrit une transformation des coordonnées angulaires du plan de la lentille ($\boldsymbol{\theta}$) vers les coordonnées angulaires du plan de la source ($\boldsymbol{\beta}$). On assume que les distances entre l'observateur et la lentille D_ℓ , entre l'observateur et la source D_s et entre la lentille et la source $D_{\ell s}$, sont beaucoup plus grandes que les distances perpendiculaires à l'axe de visée $\boldsymbol{\xi}$ ou $\boldsymbol{\eta}$ (voir figure 1.1). Cette approximation est justifiée pour les objets qui nous intéressent, pour lesquels les distances parallèles à l'axe de visée sont généralement de l'ordre du Gpc, alors que les distances perpendiculaire sont généralement de l'ordre du kpc ; soit 6 ordres de grandeurs de différences. Ainsi, on peut faire un argument géométrique (euclidien)

$$\begin{aligned} D_s \boldsymbol{\theta} &= \boldsymbol{\eta}' \\ D_s \boldsymbol{\beta} &= \boldsymbol{\eta} \\ D_{\ell s} \boldsymbol{\alpha} &= \boldsymbol{\eta}' - \boldsymbol{\eta} \\ \implies D_s \boldsymbol{\beta} &= D_s \boldsymbol{\theta} - D_{\ell s} \boldsymbol{\alpha} \end{aligned} \quad (1.13)$$

La dernière relation est l'équation maîtresse qui nous permet de tracer les rayons lumineux d'une source vers un détecteur fictif dans nos simulations. On notera que cette relation reste valide pour un univers courbe et/ou en expansion (ç.-à-d. décrit par une géométrie non-euclidienne), à condition qu'on utilise une notion de distance qui satisfait, par définition, la relation trigonométrique euclidienne

$$D \equiv \frac{\xi}{\theta} \quad (1.14)$$

Il est généralement pratique de travailler avec la forme adimensionnelle de l'équation (1.13). On introduit la densité critique

$$\Sigma_c = \frac{c^2}{4\pi G} \frac{D_s}{D_{\ell s} D_\ell}, \quad (1.15)$$

qui nous permet de définir la quantité qu'on nomme convergence $\kappa(\boldsymbol{\theta}) \equiv \frac{\Sigma(\boldsymbol{\theta})}{\Sigma_c}$. On définit ainsi l'angle réduit

$$\hat{\boldsymbol{\alpha}}(\boldsymbol{\theta}) = \frac{1}{\pi} \int_{\mathbb{R}^2} \kappa(\boldsymbol{\theta}') \frac{\boldsymbol{\theta} - \boldsymbol{\theta}'}{\|\boldsymbol{\theta} - \boldsymbol{\theta}'\|} d^2 \boldsymbol{\theta}', \quad (1.16)$$

qui satisfait l'équation de la lentille adimensionnelle

$$\boldsymbol{\beta} = \boldsymbol{\theta} - \hat{\boldsymbol{\alpha}}(\boldsymbol{\theta}). \quad (1.17)$$

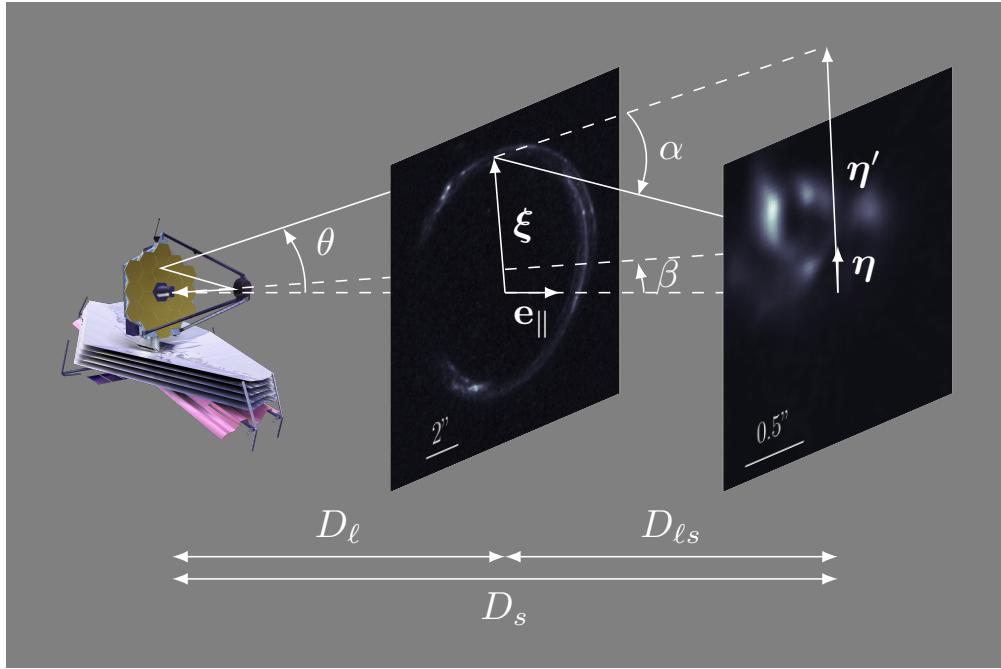


FIGURE 1.1 – Schéma d'une lentille gravitationnelle.

1.1.3 Contexte scientifique moderne

1.2 Interférométrie par masque non-régulier

1.2.1 Contexte historique

1.2.2 Les angles de fermeture

1.2.3 Contexte scientifique moderne

1.3 Auto-encodeur variationnel

1.3.1 Description du modèle

Les auto-encodeurs variationnels (VAE) ont été introduits par [Kingma and Welling \(2013\)](#) comme une approche pour inférer approximativement les variables latentes (ou cachées) qui modélisent une distribution a posteriori définie implicitement via un échantillon de données. Dans cette section, j'introduis les concepts principaux relié à ce type de modélisation. Le lecteur peut aussi se référer au livre blanc de [Kingma and Welling \(2019\)](#).

On définit $\mathbf{z} \sim q(\mathbf{z})$ comme une variable latente et \mathbf{x} comme un exemple d'un échantillon de donnée $\mathcal{D} = \{\mathbf{x}^{(i)}\}_{i=1}^N$. Notre objectif est de modéliser la distribution $p(\mathbf{x})$, implicitement décrite par notre échantillon. On suppose, sans perte de généralité, que la distribution de \mathbf{x} fait partie d'une

famille de distribution , caractérisé par θ , conditionnelle à la variable cachée : $p_\theta(\mathbf{x} \mid \mathbf{z})$. Déterminer p_θ est généralement difficile, voir intractable, si la dimensionnalité de \mathbf{x} est grande, ce qui est le cas pour des images pour lesquelles on trouve facilement $\text{dim}(\mathbf{x}) > 10^4$. Pour résoudre cette difficulté, on introduit un modèle paramétrique d'inférence $q_\phi(\mathbf{z} \mid \mathbf{x})$ dont le rôle est de modéliser la distribution a posteriori de la variable latente pour la distribution qui nous intéresse

$$q_\phi(\mathbf{z} \mid \mathbf{x}) \approx p_\theta(\mathbf{z} \mid \mathbf{x}). \quad (1.18)$$

La notion de distance entre ces deux distributions est mesurée par la divergence de Kullback-Leibler $D_{\text{KL}}(\cdot \parallel \cdot) \geq 0$:

$$\begin{aligned} D_{\text{KL}}(q_\phi(\mathbf{z} \mid \mathbf{x}) \parallel p_\theta(\mathbf{z} \mid \mathbf{x})) &= \mathbb{E}_{q_\phi(\mathbf{z} \mid \mathbf{x})} \left[\log q_\phi(\mathbf{z} \mid \mathbf{x}) - \log p_\theta(\mathbf{z} \mid \mathbf{x}) \right] \\ &= \mathbb{E}_{q_\phi(\mathbf{z} \mid \mathbf{x})} \left[\log q_\phi(\mathbf{z} \mid \mathbf{x}) - \log \frac{p_\theta(\mathbf{z}, \mathbf{x})}{p_\theta(\mathbf{x})} \right] \\ &= \underbrace{\log p_\theta(\mathbf{x}) - \mathbb{E}_{q_\phi(\mathbf{z} \mid \mathbf{x})} \left[\log p_\theta(\mathbf{z}, \mathbf{x}) - \log q_\phi(\mathbf{z} \mid \mathbf{x}) \right]}_{\equiv \mathcal{L}_{\phi, \theta}(\mathbf{x})}. \end{aligned} \quad (1.19)$$

On remarque par cette manipulation que la distance D_{KL} , en plus de mesurer la distance entre les deux distributions a posteriori (par définition), mesure aussi la différence entre le terme $\mathcal{L}_{\phi, \theta}(\mathbf{x})$, qu'on nomme limite inférieure sur l'évidence (de l'anglais *evidence lower bound* : ELBO), et la distribution qui nous intéresse $p_\theta(\mathbf{x})$. L'objectif d'un modèle VAE est de minimiser la distance D_{KL} . En observant l'équation (1.19), on réalise que ceci peut être accompli en maximisant la ELBO en relation aux paramètres ϕ et θ , ce qui accomplit deux objectifs simultanément :

1. Améliorer le processus génératif $p_\theta(\mathbf{x})$ puisque $\log p_\theta(\mathbf{x}) \geq \mathcal{L}_{\phi, \theta}(\mathbf{x})$;
2. Améliorer le processus d'inférence puisque $D_{\text{KL}}(q_\phi(\mathbf{z} \mid \mathbf{x}) \parallel p_\theta(\mathbf{z} \mid \mathbf{x})) = \log p_\theta(\mathbf{x}) - \mathcal{L}_{\phi, \theta}(\mathbf{x})$ est simultanément minimisé.

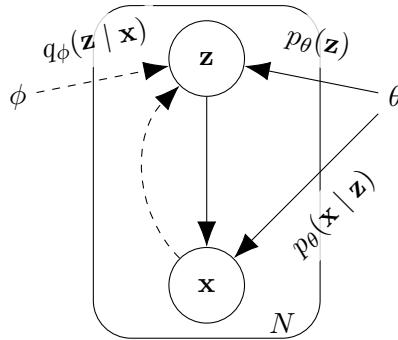


FIGURE 1.2 – Modèle graphique d'un VAE. Les flèches pleines indiquent le processus génératif, alors que les flèches pointillées indiquent le processus d'inférence.

1.3.2 Le truc de reparamétrisation

Le gradient de la ELBO par rapport aux paramètres variationnels, $\nabla_{\phi,\theta}\mathcal{L}_{\phi,\theta}(\mathbf{x})$, est une quantité que l'on aimerait calculer pour faire usage d'algorithmes comme la grimpe de gradient stochastique pour maximiser la ELBO en terme de ϕ et θ . Or, la liste de paramètres ϕ apparaît dans la distribution de prélevement pour calculer l'espérance mathématique $\mathbb{E}_{q_\phi(\mathbf{z}|\mathbf{x})}$ dans la ELBO (1.19). Cette opération n'a pas de dérivée formelle en terme de ϕ .

Pour résoudre ce problème, on utilise le truc de reparamétrisation (Kingma and Welling, 2013), qui consiste à restreindre la forme fonctionnelle de $q_\phi(\mathbf{z} \mid \mathbf{x})$ à une famille paramétrique qui s'exprime comme la transformation différentiable d'une variable aléatoire auxiliaire ϵ . On considère le cas où $q_\phi(\mathbf{z} \mid \mathbf{x})$ et $p(\epsilon)$ font partie de la famille gaussienne isotropique :

$$p(\epsilon) \equiv \mathcal{N}(0, \mathbb{1}); \quad (1.20)$$

$$q_\phi(\mathbf{z} \mid \mathbf{x}) = \mathcal{N}(\boldsymbol{\mu}_\phi(\mathbf{x}), \mathbb{1} e^{\log \boldsymbol{\sigma}_\phi^2(\mathbf{x})}); \quad (1.21)$$

$$\mathbf{z} = \boldsymbol{\mu}_\phi + \boldsymbol{\sigma}_\phi \odot \epsilon. \quad (1.22)$$

\odot symbolise le produit d'Hadamard, ou encore le produit élément-par-élément de vecteurs. Ce choix fait en sorte que les paramètres variationnelles ne participent plus au processus de prélevement, maintenant pris en charge par la variable aléatoire ϵ . Cette propriété est cruciale dans le but de prendre le gradient de la ELBO (1.19). En effet, on peut maintenant échanger les opérateurs $\nabla_{\phi,\theta}$ et $\mathbb{E}_{q_\phi(\mathbf{z}|\mathbf{x})} = \mathbb{E}_{p(\epsilon)}$, ce qui nous permet d'appliquer le gradient à l'intérieur de l'espérance mathématique. De plus, ϕ décrit maintenant une fonction générique dont le rôle est d'inférer les paramètres d'une distribution gaussienne isotropique (1.21), $f_\phi(\mathbf{x}) = (\boldsymbol{\mu}, \log \boldsymbol{\sigma}^2)$, étant donné la valeur d'un échantillon \mathbf{x} . En pratique, on peut construire une approximation de cette fonction avec un réseau de neurones convolutionnelles.

Pour déterminer la forme fonctionnelle de la ELBO, on stipule a priori que la distribution marginale des variables latentes devrait satisfaire

$$p_\theta(\mathbf{z}) = \mathcal{N}(0, \mathbb{1}) \quad (1.23)$$

On est libre de faire ce choix sans pour autant limiter les formes possibles de la distribution qui nous intéresse $p_\theta(\mathbf{x})$. On peut alors exprimer la ELBO comme

$$\mathcal{L}_{\phi,\theta}(\mathbf{x}) = \mathbb{E}_{q_\phi(\mathbf{z}|\mathbf{x})} \left[\log p_\theta(\mathbf{z}, \mathbf{x}) - \log q_\phi(\mathbf{z} \mid \mathbf{x}) \right]; \quad (1.24)$$

$$\implies \mathcal{L}_{\phi,\theta}(\mathbf{x}) = \underbrace{\mathbb{E}_{q_\phi(\mathbf{z}|\mathbf{x})} \left[\log p_\theta(\mathbf{x} \mid \mathbf{z}) \right]}_{\text{terme de reconstruction}} + \underbrace{\mathbb{E}_{q_\phi(\mathbf{z}|\mathbf{x})} \left[\log p_\theta(\mathbf{z}) - \log q_\phi(\mathbf{z} \mid \mathbf{x}) \right]}_{\equiv -D_{\text{KL}}(q_\phi(\mathbf{z}|\mathbf{x}) \parallel p_\theta(\mathbf{z}))}. \quad (1.25)$$

La divergence de KL obtenue au second terme du membre droit de l'équation (1.25) admet une

solution fermée étant donné les familles paramétriques stipulées pour $p_\theta(\mathbf{z})$ (1.23) et $q_\phi(\mathbf{z} \mid \mathbf{x})$ (1.21)

$$-D_{\text{KL}}(q_\phi(\mathbf{z} \mid \mathbf{x}) \parallel p_\theta(\mathbf{z})) = \frac{1}{2} \sum_{j=1}^{\dim(\mathbf{z})} (1 + [\log \boldsymbol{\sigma}_\phi^2]_j - [\boldsymbol{\mu}_\phi]_j - [\boldsymbol{\sigma}_\phi^2]_j) \quad (1.26)$$

Une dérivation de ce terme est donnée dans l'appendice B de [Kingma and Welling \(2013\)](#). Le premier terme du membre droit de l'équation (1.25) est nommé *terme de reconstruction* puisqu'il connecte avec l'objectif des fonctions de type auto-encodeurs d'apprendre une représentation latente d'un échantillon de données. La fondation théorique de ce type d'analyse repose sur le principe du goulot d'information (?), un sujet qui n'est pas abordé dans cette thèse. Le lecteur est invité à se référer à la revue sur le sujet par [Goldfeld and Polyanskiy \(2020\)](#). La reconstruction s'accomplit en utilisant d'abord le modèle d'inférence $\mathbf{z}^{(1:L)} \stackrel{\text{i.i.d}}{\sim} q_\phi(\mathbf{z} \mid \mathbf{x})$ ¹ pour obtenir un échantillon de représentations latentes à partir des équations (1.20) à (1.22), puis en utilisant le modèle génératif $\hat{\mathbf{x}}^{(i)} \sim p_\theta(\mathbf{x} \mid \mathbf{z}^{(i)})$ pour obtenir un échantillon de reconstructions $\hat{\mathbf{x}}^{(1:L)}$ similaire à l'exemple originel \mathbf{x} . Comme on a déjà une variable auxiliaire ϵ qui se charge de l'aspect génératif du modèle, on peut construire une approximation du modèle génératif avec une fonction générique des variables latentes $g_\theta(\mathbf{z})$. Encore une fois, un réseau de neurones convolutionnelles est un choix pratique pour modéliser cette fonction dans le cas où \mathbf{x} est une image. En général, on choisit une erreur quadratique moyenne pour modéliser le terme de reconstruction, de sorte que

$$\mathbb{E}_{q_\phi(\mathbf{z} \mid \mathbf{x})} \left[\log p_\theta(\mathbf{x} \mid \mathbf{z}) \right] \simeq -\frac{1}{L} \sum_{i=1}^L \|\mathbf{x} - (\hat{\mathbf{x}}^{(i)})^2\|_2^2 \quad (1.27)$$

1.4 Machines à inférence récurrentielles

1.4.1 Formalisme des problèmes inverses

1.4.2 La relation de récurrence

1. i.i.d : identiquement et indépendamment distribué.

Chapitre 2

Pixelated Reconstruction of Gravitational Lenses using Recurrent Inference Machine

Abstract

Modeling strong gravitational lenses in order to quantify the distortions in the images of background sources and to reconstruct the mass density in the foreground lenses has traditionally been a difficult computational challenge. As the quality of gravitational lens images increases, the task of fully exploiting the information they contain becomes computationally and algorithmically more difficult. In this work, we use a neural network based on the Recurrent Inference Machine (RIM) to simultaneously reconstruct an undistorted image of the background source and the lens mass density distribution as pixelated maps. The method we present iteratively reconstructs the model parameters (the source and density map pixels) by learning the process of optimization of their likelihood given the data using the physical model (a ray-tracing simulation), regularized by a prior implicitly learned by the neural network through its training data. When compared to more traditional parametric models, the proposed method is significantly more expressive and can reconstruct complex mass distributions, which we demonstrate by using realistic lensing galaxies taken from the cosmological hydrodynamic simulation IllustrisTNG.

2.1 Introduction

Strong gravitational lensing is a natural phenomenon through which multiple distorted images of luminous background objects, i.e. early-type star-forming galaxies, are formed by massive foreground objects along the line of sight (e.g., Vieira et al., 2013; Marrone et al., 2018; Rizzo et al., 2020; Sun et al., 2021). These distortions are tracers of the distribution of mass in foreground objects, independent of the electromagnetic behaviour of these overdensities. As such, this phenomenon offers a powerful probe of the distribution of dark matter and its properties outside of the Milky Way (e.g., Dalal and Kochanek, 2002; Treu and Koopmans, 2004; Hezaveh et al., 2016; Gilman et al., 2020, 2021).

Lens modeling is the process of inferring the parameters describing both the mass distribution in the foreground lens and undistorted image of the background source. This has traditionally been a time- and resource-consuming procedure. A common practice to model the mass of lensing galaxies is to assume that their density profiles follow simple parametric forms, e.g., a power law $\rho \propto r^{-\gamma'}$. These profiles generally provide a good fit to low-resolution data and are easy to work with due to their small number of parameters (e.g., Koopmans et al., 2006; Barnabè et al., 2009; Auger et al., 2010). However, as high-resolution and high signal-to-noise ratio (SNR) images become available, lens analysis with simple models requires the introduction of additional parameters representing the true complexity of the mass distribution in lensing galaxies and their immediate environments (e.g., Sluse et al., 2017; Wong et al., 2017; Birrer et al., 2019; Rusu et al., 2020, 2017; Li et al., 2021). This approach becomes intractable as the quality of images increases. For example, no simple parametric model of the Hubble Space Telescope (HST) Wide Field Camera 3 (WFC3) images of the Cosmic Horseshoe (J1148+1930) — initially discovered by Belokurov et al. (2007) — has been able to model the fine features of the extended arc (e.g., Bellagamba et al., 2016; James et al., 2018; Cheng et al., 2019; Schuldt et al., 2019).

Free-form methods — also misleadingly called nonparametric methods — attempt to relax the assumptions about the smoothness and symmetries of these parametric profiles by changing their parametric support to more expressive families like regular (or adaptive) grid representations and meshfree representations. But, this added flexibility comes at a price, which is (often) a high-dimensional inference problem that is under-constrained, meaning that imposing a prior on the reconstructed parameters becomes essential to penalize unphysical solutions and avoid overfitting the data. Free-form methods have a rich history in cluster-scale lensing (Bartelmann et al., 1996; Seitz et al., 1998; Abdelsalam et al., 1998a,b; Bradač et al., 2005; Diego et al., 2005; Cacciato et al., 2006; Diego et al., 2007; Liesenborgs et al., 2006, 2007; Jee et al., 2007; Coe et al., 2008; Merten et al., 2009; Deb et al., 2012; Merten, 2016; Ghosh et al., 2020; Torres-Ballesteros and Castañeda, 2022) and weak lensing (Kaiser and Squires, 1993; Marshall, 2001; Massey et al., 2007; Deb et al., 2008; Simon et al., 2012; Leonard et al., 2012; Lanusse et al., 2016; Jeffrey et al., 2020; Starck et al., 2021; Remy et al., 2022). They strive to make better use of the information contained in lensing features like resolved arc details, multiple image position, flux ratios, image deformation (i.e. weak

lensing constraints) or even the null space of an image in order to place better constraints on the morphology of the mass density of the lens. On the other hand, comparatively less work has been done in the context of galaxy-galaxy lensing system to tackle free-form mass modelling directly ([Saha and Williams, 1997, 2004](#); [Birrer et al., 2015](#); [Coles et al., 2014](#)). The main reason for this state of affairs is the difficulty of specifying an appropriate prior which regularizes the problem over its non-linear parameters while maintaining computational tractability and sufficient flexibility to model a large variety of systems.

In view of this, the focus of the field has instead been on using free-form methods for the background source reconstruction only. A number of well-established procedures exists for linear inversion of pixellated-source models in the context of traditional maximum likelihood modeling. These methods, originally developed by [Warren and Dye \(2003\)](#); [Suyu et al. \(2006\)](#), are based mainly on imposing a quadratic-log prior to the source pixels to regularize the optimisation. Subsequently, there have been multiple attempts at building a bridge toward free-form methods in order to correct simplistic assumptions on the mass density model by using linear corrections of the lensing potential ([Koopmans, 2005](#); [Suyu and Blandford, 2006](#); [Vegetti and Koopmans, 2009](#); [Vegetti et al., 2012](#)), or iteratively specified priors for shapelets corrections ([Birrer and Amara, 2018](#); [Nightingale et al., 2018](#)). But these extensions are generally limited by assumptions regarding the accuracy of the initial parametric models, and thus we are left wanting for a more flexible approach.

Over the recent years, deep learning methods have proven extremely successful at modeling quickly and accurately strong lensing systems ([Hezaveh et al., 2017](#); [Perreault Levasseur et al., 2017](#); [Morningstar et al., 2018](#); [Coogan et al., 2020](#); [Park et al., 2021](#); [Legin et al., 2021](#); [Wagner-Carena et al., 2021](#); [Schuldt et al., 2022](#); [Wagner-Carena et al., 2022](#); [Karchev et al., 2022](#); [Anau Montel et al., 2022](#); [Mishra-Sharma and Yang, 2022](#)). More specifically, [Morningstar et al. \(2019\)](#) demonstrated that recurrent convolutional neural networks can learn implicitly complex prior distributions from their training data to successfully reconstruct pixelated undistorted images of strongly lensed sources, circumventing the need to specify explicitly a prior distribution over those parameters. Motivated by this success, we propose a method that extends this framework to solve the full lensing problem and simultaneously reconstruct a pixelated lensing mass map and a pixelated undistorted background source.

The method we propose here is based on the Recurrent Inference Machine (thereafter referred to as RIM), originally developed by [Putzky and Welling \(2017\)](#). In this framework, we aim to learn an iterative inference algorithm, moving away from hand-chosen inference algorithms and hand-crafted priors. Instead, the prior is learned implicitly through the dataset used to train the neural network that update the solution parameters at each iteration.

In this paper, we present a new architecture for the neural network based on a U-net architecture ([Ronneberger et al., 2015](#)), tailored to our highly non-linear inverse problem. We also introduce a fine-tuning prescription which allows us to exploit directly the prior encoded in the neural network parameters in order to perform statistically significant — noise-level — reconstructions of high-SNR

galaxy-galaxy lensing systems simulated using IllustrisTNG (Nelson et al., 2019) projected density maps and background galaxy images collected from the COSMOS survey (Koekemoer et al., 2007; Scoville et al., 2007).

The paper is organised as follows. Section 2.2 details the inference pipeline. In Section 2.3, we present the data creation and preprocessing for training the RIM and the generative models used in this paper. In Section 2.4, we report the training strategy for the models used in this work. In Section 2.5, we report and discuss our results on a held-out test set of gravitational lenses. Section 2.6 concludes and situates our finding within the larger context of studying gravitational lensing.

2.2 Methods

In this section, we present the steps to build a free-form inference pipeline with a RIM, beginning with a general introduction about MAP inference with a Gaussian likelihood in Section 2.2.1. In Section 2.2.2, we motivate the use of a Recurrent Inference Machine to solve this problem and describe the computational graph of the RIM as well as the optimisation problem of learning the gradient model. The architecture of the gradient model is described in Section 2.2.3. We describe the raytracing simulation in Section 2.2.4. Finally, we describe the fine-tuning procedure and transfer learning technique applied to reach noise level reconstructions in Section 2.2.5.

2.2.1 Maximum a posteriori

The task of reconstructing a signal vector $\mathbf{x} \in \mathcal{X}$ given observed data $\mathbf{y} \in \mathcal{Y}$ is formulated as an ill-posed inverse problem with a known forward model F and additive noise distribution. We assume a Gaussian distribution with known covariance matrix C , such that

$$\begin{aligned} \mathbf{y} &= F(\mathbf{x}) + \boldsymbol{\eta}; \\ \boldsymbol{\eta} &\sim \mathcal{N}(0, C). \end{aligned} \tag{2.1}$$

In our case study, F is a many-to-one non-linear mapping between the model space \mathcal{X} and the data space \mathcal{Y} . Finding a unique solution for this ill-posed inverse problem requires strong inductive biases to be introduced in the inference procedure in order to favour certain hypotheses over others. In the maximum a posteriori (MAP) optimization framework, the inductive biases are encoded in a prior $p(\mathbf{x})$ that is introduced as a probability distribution over \mathcal{X} which reduces the relevant space to explore during inference. The MAP solution is the hypothesis that maximizes the product of the likelihood $p(\mathbf{y} | \mathbf{x})$ and the prior :

$$\hat{\mathbf{x}}_{\text{MAP}} = \underset{\mathbf{x} \in \mathcal{X}}{\operatorname{argmax}} \log p(\mathbf{y} | \mathbf{x}) + \log p(\mathbf{x}). \tag{2.2}$$

By restricting ourselves to Gaussian noise models, the likelihood can be calculated directly and takes the form

$$\log p(\mathbf{y} \mid \mathbf{x}) \propto -(\mathbf{y} - F(\mathbf{x}))^T C^{-1} (\mathbf{y} - F(\mathbf{x})) \quad (2.3)$$

However, the prior distribution is harder to define. It is problem-dependent and requires expert knowledge of the model domain.

2.2.2 Recurrent Inference Machine

Instead of handcrafting such a distribution, we attempt to build an inference machine with an implicit prior built in the training set \mathcal{D} and encoded in a deep neural network architecture (Bengio, 2009). The RIM (Putzky and Welling, 2017) is a form of learned gradient-based inference algorithm intended to solve inverse problems of the form of equation (2.1). This framework has mainly been applied in the context of linear and under-constrained inverse problems — i.e. where the function F can be represented in a matrix form and where a prior on the solution space is required to produce a unique solution — for which the prior on the parameters \mathbf{x} , $p(\mathbf{x})$, is either intractable or hard to compute (Morningstar et al., 2018, 2019; Lønning et al., 2019). The use of the RIM to solve non-linear inverse problems was first investigated in (Modi et al., 2021). In our case, the inverse problem is non-linear, as it is given by equation (2.7).

The governing equation for the RIM is a recurrent relation that takes the general form

$$\hat{\mathbf{x}}^{(t+1)} = \hat{\mathbf{x}}^{(t)} + g_\varphi(\hat{\mathbf{x}}^{(t)}, \mathbf{y}, \nabla_{\hat{\mathbf{x}}^{(t)}} \log p(\mathbf{y} \mid \hat{\mathbf{x}}^{(t)})). \quad (2.4)$$

In the text, we will often use the shorthand notation $\nabla_{\mathbf{y} \mid \mathbf{x}} \equiv \nabla_{\mathbf{x}} \log p(\mathbf{y} \mid \mathbf{x})$ to refer to the gradient of the likelihood. By minimizing a weighted mean squared loss backpropagated through time,

$$\mathcal{L}_\varphi(\mathbf{x}, \mathbf{y}) = \frac{1}{T} \sum_{t=1}^T \sum_{i=1}^M \mathbf{w}_i (\hat{\mathbf{x}}_i^{(t)} - \mathbf{x}_i)^2, \quad (2.5)$$

the neural network g_φ learns to optimize the parameters \mathbf{x} given a likelihood function. The converged parameters of the neural network given the training set \mathcal{D} , $\varphi_{\mathcal{D}}^*$, are those that minimize the cost — or empirical risk — which is defined as the expectation of the loss over \mathcal{D}

$$\varphi_{\mathcal{D}}^* = \operatorname{argmin}_{\varphi} \mathbb{E}_{(\mathbf{x}, \mathbf{y}) \sim \mathcal{D}} [\mathcal{L}_\varphi(\mathbf{x}, \mathbf{y})]. \quad (2.6)$$

Unlike previous works (Andrychowicz et al., 2016; Putzky and Welling, 2017; Morningstar et al., 2018, 2019; Lønning et al., 2019), the data vector \mathbf{y} — or observation — is fed to the neural network in order to learn the initialization of the parameters ($\mathbf{x}^{(0)} = g_\varphi(0, \mathbf{y}, 0)$) as well as their optimization. We found empirically that this significantly improves the performance of the model for our problem and it avoids situations where the model would get stuck in local minima at test time due to poor initialization.

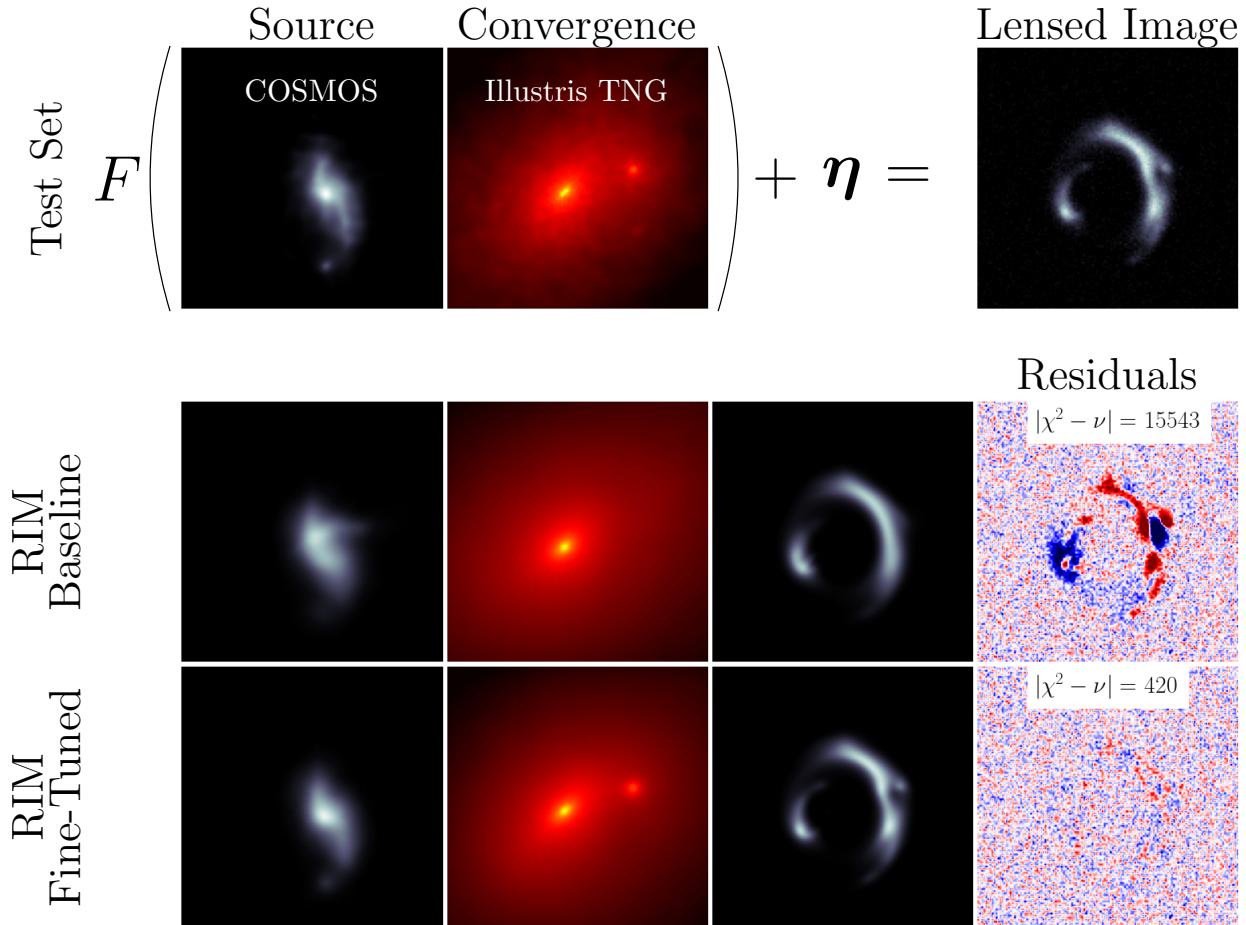


FIGURE 2.1 – Example of a simulated lensed image in the test set that exhibits a large deflection in its eastern arc which indicates the presence of a massive object — in this case a dark matter subhalo. The fine-tuning procedure is able to recover this subhalo because of its strong signal in the lensed image and reduces the residuals to noise level.

We follow previous works in setting a uniform weight over the time steps ($\mathbf{w}^{(t)} = \frac{\mathbf{w}}{T}$). The choice of the pixel weights \mathbf{w}_i is informed by our empirical observations when training the network. Details are reported in appendix C.

In Figure 2.2, we show the rolled computational graph of the RIM. During training of the gradient model g_φ , operations along the solid arrows are being recorded for backpropagation through time. The recording is stopped along the dashed arrow since these operations are part of the forward modelling process. By avoiding the computation of these gradients, training time is reduced and knowledge about the inner workings of a specific likelihood (and forward model) is insulated from the optimization algorithm. This is analogous to a common RNN use-case like text generation, where the process responsible for producing the next element in a time series is a black box to the optimization algorithm.

The gradient of the likelihood is computed using automatic differentiation. Following (Modi

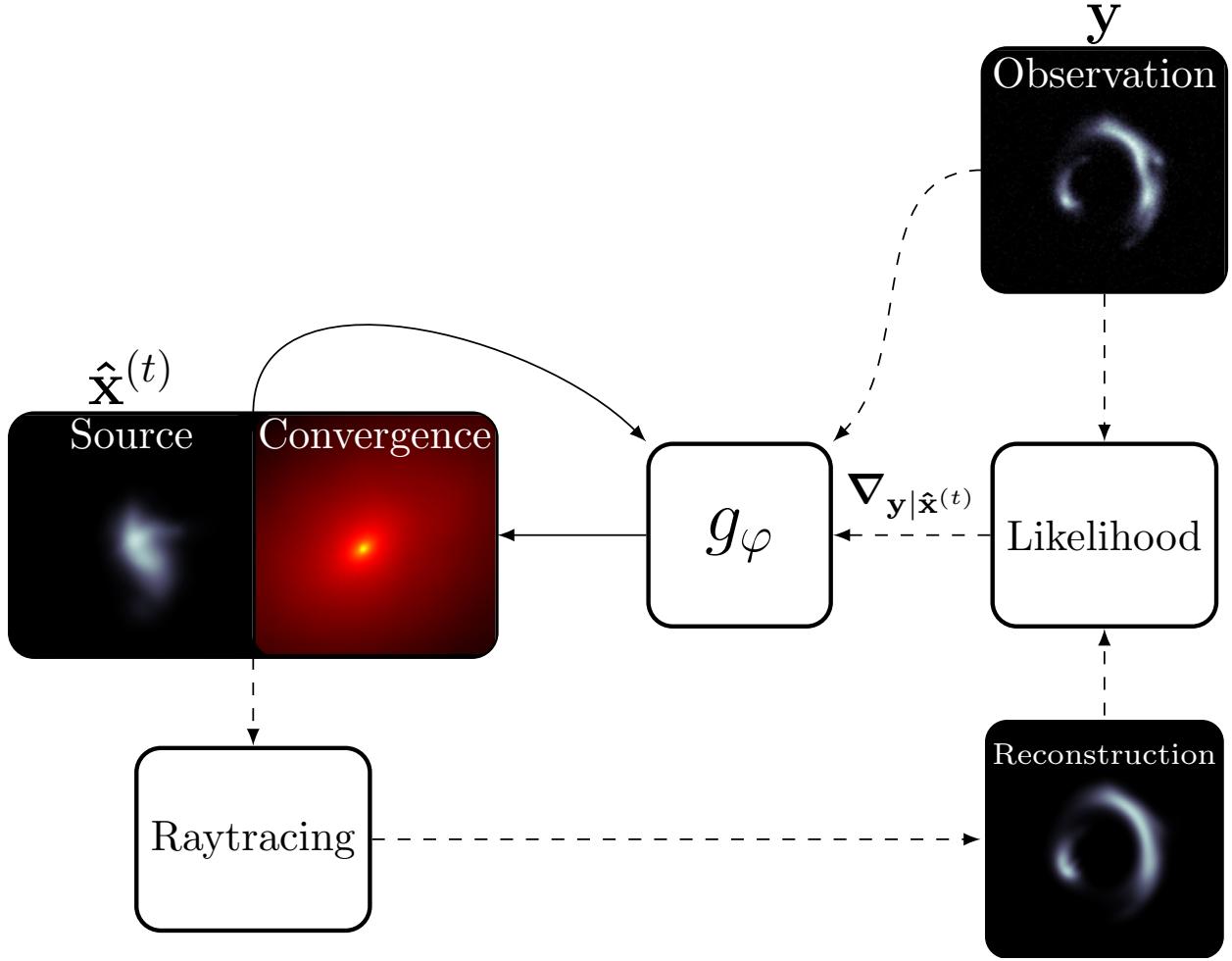


FIGURE 2.2 – Rolled computational graph of the RIM. Dashed arrows represent operations not recorded for BPTT.

et al., 2021), we preprocess the gradients using the Adam algorithm (Kingma and Welling, 2013). For clarity, we only illustrated this step in Figure 2.3.

2.2.3 The Gradient Model

The neural network architecture is illustrated in Figure 2.3, which shows a single time step of the unrolled computation graph of the RIM. We use a U-net (Ronneberger et al., 2015) architecture with Gated Recurrent Units (GRU : Cho et al., 2014) placed in each skip connections.

Each GRU cell has its own memory tensor that is updated through time at each iteration of equation 2.4. The shape of a memory tensor is set to match the feature tensor fed into it from the parent layer in the network graph. Instead of learning a compressed representation like in the hourglass architecture (i.e. autoencoder), the U-net architecture naturally separates the spatial

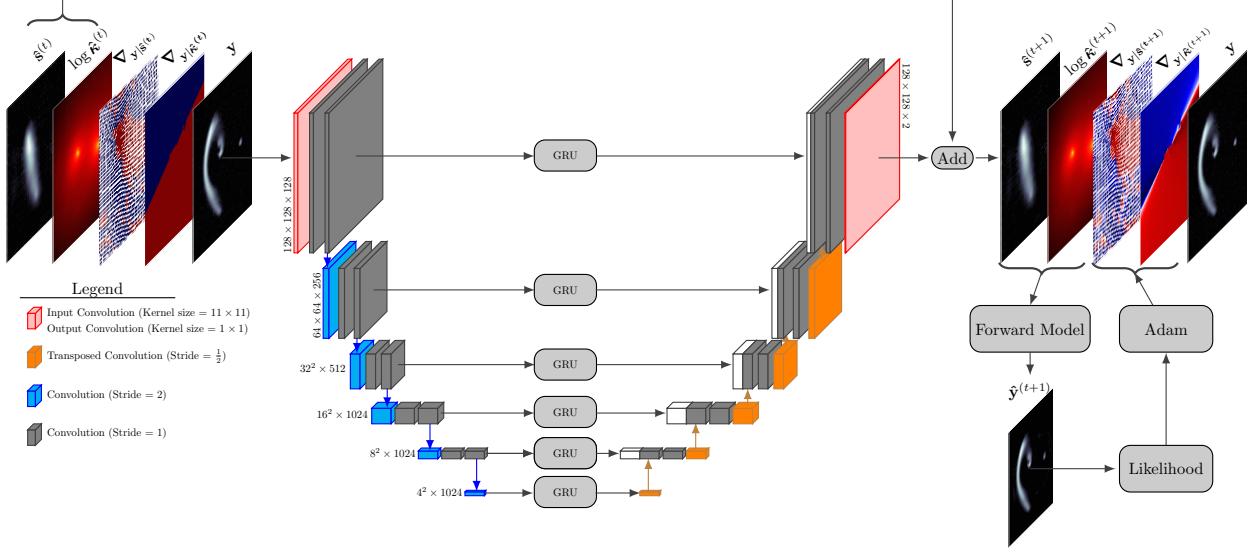


FIGURE 2.3 – A single time step of the unrolled computation graph of the RIM. GRU units are placed in the skip connections to guide the reconstruction of the source and convergence. A schematic of the steps to compute the likelihood gradients is shown in the bottom right of the figure, including the Adam processing step of the likelihood gradient.

frequency components of the signal into its vertical levels. The first level generally encodes high frequency features while the lower level encodes low frequency features (due to downsampling of the feature maps). Adding an independent memory unit at each level preserve this property.

Convolutional layers with a stride of 2 are used for downsampling and stride of $\frac{1}{2}$ for upsampling of the feature maps (identified in blue and orange respectively in figure 2.3). Half-stride convolutions are implemented in practice with the transposed convolution layers from Tensorflow (Abadi et al., 2015). Most layers use a kernel size of 3×3 , except the first and last layer. The first layer has larger receptive field (11×11) in order to capture more details in the input tensor. The last layer has kernels of size 1×1 . A tanh activation function is used for each convolutional layer, including strided convolutions, except for the output layer. The U-net outputs an image tensor with two channels, one dedicated for the update of the source and the other to the update of the convergence (see figure 2.3).

2.2.4 The Forward Model

An observation is simulated by ray tracing the brightness distribution of the background source to the foreground coordinate system. In our case, the coordinate systems have discretized representations. Each pixel of an image is labeled with a subscript index i , which we distinguish from a parenthesized superscript index (i) that refers to the member of a set or list of tensors. For clarity, we omit the superscript index in what follows.

Each pixel is associated with an intensity value and a coordinate vector. The foreground pixel

coordinates $\boldsymbol{\theta}_i$ and the source pixel coordinates $\boldsymbol{\beta}_i$ are related by the lens equation

$$\boldsymbol{\beta}_i = \boldsymbol{\theta}_i - \boldsymbol{\alpha}(\boldsymbol{\theta}_i), \quad (2.7)$$

where $\boldsymbol{\alpha}$ is a deflection angle. It is obtained from the projected surface density field κ — also referred to as convergence — by the integral

$$\boldsymbol{\alpha}(\boldsymbol{\theta}_i) = \frac{1}{\pi} \int_{\mathbb{R}^2} \kappa(\boldsymbol{\theta}') \frac{\boldsymbol{\theta}_i - \boldsymbol{\theta}'}{\|\boldsymbol{\theta}_i - \boldsymbol{\theta}'\|^2} d^2\boldsymbol{\theta}'. \quad (2.8)$$

The intensity of a pixel in a simulated observation is obtained by bilinear interpolation of the source brightness distribution at the coordinate $\boldsymbol{\beta}_i$. In this work, the convergence also has a discrete representation. Thus, we approximate this integral by a discrete global convolution. Taking advantage of the convolution theorem, this operation can be computed in near-linear time using the Fast Fourier Transform (FFT).

Assuming the observation has M^2 pixels, the convolution kernel would have $(2M + 1)^2$ pixels. Both the convergence tensor and the kernel tensor are zero-padded to a size of $(4M + 1)^2$ pixels in order to approximate a linear convolution and significantly reduce aliasing.

A blurring operator — convolution by a point spread function (PSF) — is then applied to the lensed image to replicate the response of an imaging system. This operator is implemented as a GPU-accelerated matrix operation since the blurring kernels used in this paper have a significant proportion of their energy distribution encircled inside a small pixel radius.

2.2.5 Fine-Tuning

Objective function

Once the gradient model is trained, the RIM is a baseline estimator of the parameters \mathbf{x} given a noisy observation \mathbf{y} , a PSF and a noise covariance matrix. We now concern ourselves with a strategy to improve this estimator. This is important for observations with high SNR, for which the estimator must be extremely accurate to model all the fine features present in the arcs. The metric for the goodness of fit is the reduced chi squared $\chi_{\nu}^2 = \frac{\chi^2}{\nu}$, where ν is the total number of degrees of freedom which corresponds to the total amount of pixels in \mathbf{y} in this work. Generally, our goal will be to reach $\chi_{\nu}^2 = 1$, or equivalently $|\chi^2 - \nu| = 0$, which indicates that the RIM hypothesis has reconstructed all the signal to be recovered from the observation. We note that such a problem is exceedingly difficult at high SNR. In this regime, our hypothesis has to be both precise and accurate in order to satisfy the chi squared test. This is unlike noisy observations, where strong assumptions about the model parameters are required and often sufficient to reconstruct the remaining information left in the observed vector \mathbf{y} .

We observe that we can optimize the network parameters w.r.t to the chi squared directly

$$\hat{\varphi}_{\text{MAP}} = \underset{\varphi}{\operatorname{argmax}} \frac{1}{T} \sum_{t=1}^T \log p(\mathbf{y} \mid \hat{\mathbf{x}}^{(t)}) + \log p(\varphi). \quad (2.9)$$

Unlike the loss in equation (2.5), this objective function makes no use of labels (\mathbf{x}). This allows us to use equation (2.9) at test time in order to fine-tune the RIM's weights to a specific test example.

Transfer Learning

We now address the issue of transferring knowledge from the training task, problem (2.6), to a test task specific to an observation, problem (2.9). The reader might refer to reviews on transfer learning (Pan and Yang, 2010; Zhuang et al., 2019) for a broad overview of the field. The strategy we outline fall into the category of inductive transfer learning.

Since the data likelihood $p(\mathbf{y} \mid \mathbf{x})$ does not contain *a priori* information about the solution $\hat{\varphi}_{\text{MAP}}$, inductive biases must be introduced to make the problem (2.9) well-posed. Thus, we

(\mathcal{H}_1) initialize the network parameters with $\varphi_{\mathcal{D}}^*$;

(\mathcal{H}_2) apply early stopping when a maximum number of steps is reached or $\chi_{\nu}^2 \leq 1$;

(\mathcal{H}_3) use a small learning rate.

(\mathcal{H}_2) and (\mathcal{H}_3) encode the assumption that the optimal estimator is to be found *near* the initialization.

As it turns out, (\mathcal{H}_1) is not strong enough to preserve the knowledge learned from the training task. This has long been observed in the literature and was coined as the catastrophic interference phenomenon in connectionist networks (McCloskey and Cohen, 1989; Ratcliff, 1990). In summary, a sequential learning problem exhibits catastrophic forgetting of old knowledge when confronted with new examples (possibly from a different distribution or process), in a manner

(CF₁) proportional to the amount of learning;

(CF₂) strongly dependant to the disruption of the parameters involved in representing the old knowledge.

While (\mathcal{H}_2) and (\mathcal{H}_3) can potentially alleviate (CF₁), (CF₂) is not trivially addressed by the inductive biases introduced so far.

We follow the work of Kirkpatrick et al. (2016) to define a prior distribution over φ that address this issue

$$\log p(\varphi) \propto -\frac{\lambda}{2} \sum_j \operatorname{diag}(\mathcal{I}(\varphi_{\mathcal{D}}^*))_j (\varphi_j - [\varphi_{\mathcal{D}}^*]_j)^2. \quad (2.10)$$

$\operatorname{diag}(\mathcal{I}(\varphi_{\mathcal{D}}^*))$ is the diagonal of the Fisher information matrix encoding the amount of information that some set of gravitational lensing systems from the training set, and similar to the observed test task, carries about the baseline RIM weights $\varphi_{\mathcal{D}}^*$ — the parameters that minimize the empirical risk

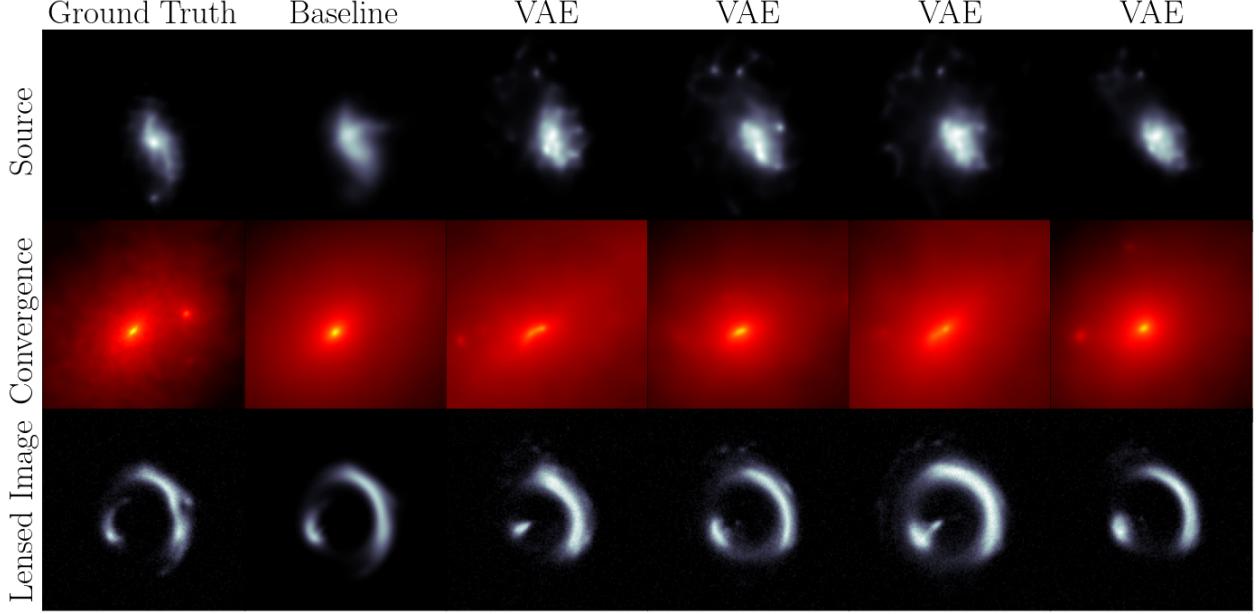


FIGURE 2.4 – Examples similar to the test task, also shown in Figure 2.7. The first column shows the ground truth used to simulate the lensed image. The second column shows the baseline prediction that is then encoded in the latent space of the VAE in order to sample the next 4 columns.

(equation 2.6). We can also understand this prior using the Cramér-Rao lower bound (Rao, 1945; Cramér, 1946). The prior can thus be framed as a multivariate Gaussian distribution characterised by a diagonal covariance matrix with $\text{diag}(\mathcal{I})$ as its inverse and by $\varphi_{\mathcal{D}}^*$ as its first moment. Within this view, the Lagrange multiplier is tuning our estimated uncertainty about the neural network weights for the particular task at hand. We've included a derivation of this term in the appendix A.

Examples are drawn from the set of training examples similar to the test task by sampling the latent space of both the source VAE and the convergence VAE near the baseline prediction of the RIM. Figure 2.4 illustrates what we mean by *similar*.

2.3 Data

2.3.1 COSMOS

The background source brightness distributions are taken from the Hubble Space Telescope (HST) Advanced Camera for Surveys Wide Field Channel COSMOS field (Koekemoer et al., 2007; Scoville et al., 2007), a 1.64 deg^2 contiguous survey acquired in the F814W filter. A dataset of mag limited ($F814W < 23.5$) deblended galaxy postage stamps (Leauthaud et al., 2007) was compiled as part of the GREAT3 challenge (Mandelbaum et al., 2014). The data is publicly available (Mandelbaum et al., 2012), and the preprocessing is done through the open source software **GALSIM** (Rowe et al., 2015).

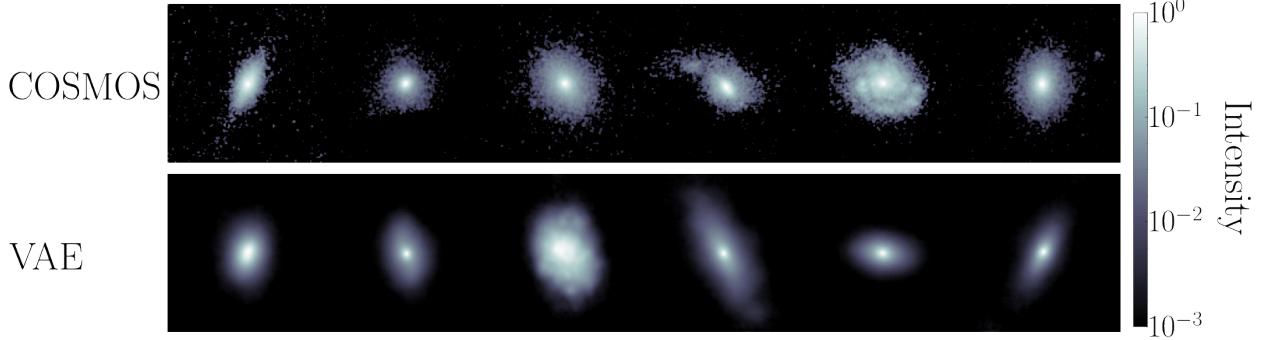


FIGURE 2.5 – Examples of COSMOS galaxy images (top row) and VAE generated samples (bottom row) used as labels in \mathcal{D} .

We applied the `marginal` selection criteria (see the `COSMOSCatalog` class) and imposed a flux per image greater than 50 photons $\text{cm}^{-2} \text{s}^{-1}$. This final set has a total of 13 321 individual images. Each image is convolved with its original PSF and drawn into a postage stamps of 158^2 pixels. Each image is then background subtracted, randomly shifted, rotated by an angle multiple of 90° , cropped down to 128^2 pixels and normalized to pixel intensities in the range $[0, 1]$.

We split each unique galaxies into a training set (90%) and a test set (10%). The training set is used to train a VAE and produce simulated observations to train the RIM.

2.3.2 IllustrisTNG

Smooth Particle Lensing

To compute a convergence map from an N-body simulation, we follow [Aubert et al. \(2007\)](#) in treating each particle as flow tracers instead of describing their density as Dirac delta functions. Smoothing each particle density on a non-singular kernel reduces the particle noise affecting all important lensing quantities — most importantly the convergence. At the same time, the choice of the kernel size is important to preserve substructures in the lens that we might potentially be interested in. Following [Rau et al. \(2013\)](#), we use Gaussian smoothing with an adaptive kernel size determined by the distance of the 64th nearest neighbours of a given particle $D_{64,i}$.

$$\kappa(\mathbf{x}) = \frac{1}{\Sigma_{\text{crit}}} \sum_{i=1}^{N_{\text{part}}} \frac{m_i}{2\pi\hat{\ell}_i^2} \exp\left(-\frac{1}{2} \frac{(\mathbf{x} - \mathbf{x}_i)^2}{\hat{\ell}_i^2}\right) \quad (2.11)$$

$$\hat{\ell}_i = \sqrt{\frac{103}{1024}} D_{64,i}.$$

The nearest neighbours are found by fitting a k-d tree — implemented in `scikit-learn` ([Pedregosa et al., 2011](#)) — to the N_{part} particles in a cylinder centered on the centre of mass of the halo of

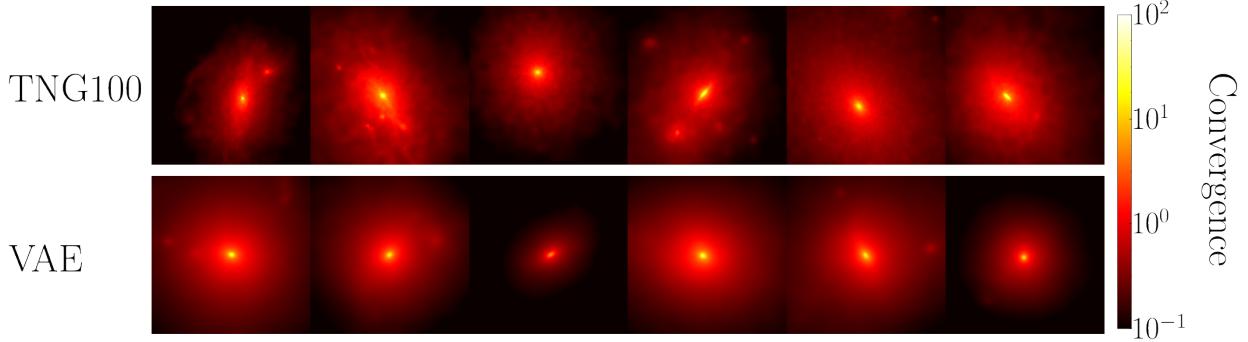


FIGURE 2.6 – Examples of smoothed Illustris TNG100 convergence map (top row) and VAE generated samples (bottom row) used as labels in \mathcal{D} .

interest. The critical surface density is defined as

$$\Sigma_{\text{crit}} = \frac{4\pi G}{c^2} \frac{D_\ell D_{\ell s}}{D_s}, \quad (2.12)$$

where D_ℓ , D_s and $D_{\ell s}$ are angular diameter distance to the lens, source and between the lens and the source respectively, G is the gravitational constant and c the speed of light.

Preprocessing

The projected surface density maps (convergence) of lensing galaxies were made using the redshift $z = 0$ snapshot of the IllustrisTNG-100 simulation (Nelson et al., 2019) in order to produce physically realistic realizations of dark matter and baryonic matter halos. We selected 1604 halos with the criteria that they have a total dark matter mass of at least $9 \times 10^{11} M_\odot$. We then collected all dark matter, gas, stars and black holes particles from the data associated to the galaxy cluster within which the halo resides in to create a smoothed projected surface density maps around the centroid of the halo as prescribed in section 2.3.2.

We adopt the Λ CDM cosmology from Planck Collaboration (2020) with $h = 0.68$ to compute angular diameter distances. We also fix the source redshift to $z_s = 1.5$ and the deflector redshift to $z_\ell = 0.5$. We note that changing the redshifts or the cosmology only amount in a rescaling of the κ map by a global scalar, not the morphology of the profiles. Thus, this choice does not change the generality of our method. The smoothed distributions from equation (2.11) are rendered into a regular grid of 188^2 pixels with a comoving field of view of $105 \text{ kpc}/h$. To avoid edge effects in the pixelated maps, we include particles outside of the field of view in the sum of equation (2.11).

Before applying augmentation or considering different projections, our dataset of halos is split into a training set (90%) and a test set (10%) in order to make sure that the test set consists only of convergence maps unseen by the RIM during training. We take 3 different projections (xy , xz and yz) of each 3D particle distributions, which amounts to a dataset with a total of 4812 individual convergence maps. Random rotations by an angle multiple of 90° and random shifts to the pixel

coordinates are applied to each image. The κ maps are then rescaled by a random factor to change their estimated Einstein radius to the range [0.5, 2.5] seconds of arc. The Einstein radius is defined as

$$\theta_E = \sqrt{\frac{4GM(\theta_E)}{c^2} \frac{D_{\ell s}}{D_\ell D_s}} \quad (2.13)$$

where $M(\theta_E)$ is the mass enclosed inside the Einstein radius. In practice, we estimate this quantity by summing over the mass of pixels with a value greater than the critical density ($\kappa > 1$). For data augmentation purposes, this procedure gives a good enough estimate of the size of the lensed image that will be produced by some κ map. We test multiple scaling factors for each κ map, then uniformly sample between those that produce an estimated Einstein radius within the desired range. This step is used to remove any bias in the Einstein radius that might come from the mass function of the simulation.

The final maps are cropped down to 128^2 pixels. Placed at a redshift $z_\ell = 0.5$, a κ map will thus span an angular field of view of $7.69''$ with a resolution similar to HST. With these augmentation procedures, a total of 50 000 maps are created from the training split and 5 000 from the test split. The training set is used to train a VAE and produce simulated observations to train the RIM.

2.3.3 Simulated Observations

Having defined a source map and a convergence map, we apply the ray tracing simulation prescribed in section 2.2.4 to produce an observation with observational effects that crudely correspond to HST images.

For each observation, a Gaussian point spread function is created with a full width at half maximum (FWHM) randomly generated from a truncated normal distribution. The support of the distribution is truncated below by the angular size of a single pixel and above by the angular size of 4 pixels. White noise with a standard deviation randomly generated from a truncated normal distribution is then added to the convolved observation to simulate SNR conditions between 10 dB and 30 dB. For simplicity, we define $\text{SNR} = \frac{1}{\sigma}$. This definition is equivalent to the peak signal-to-noise ratio.

As a validation criteria for each simulated image, we impose a minimum magnification of 3. Thus, we make sure that most pixel coordinates in the image plane will be mapped inside the source coordinate system through the lens equation (2.7).

400 000 observations are simulated from random pairs of COSMOS sources and IllustrisTNG convergence training splits in order to train the RIM. An additional 200 000 observations are created from pairs of COSMOS source and pixelated SIE convergence map. The parameters for these κ maps are listed in table 2.2. We found this addition to be beneficial to learning since it adds an inductive bias in the learning favoring isothermal profiles. We expect some lensing configurations like large Einstein rings or double images to poorly constrain the inner structure of the mass distribution.

TABLE 2.1 – Physical model parameters.

Parameter	Distribution/Value
Lens redshift z_ℓ	0.5
Source redshift z_s	1.5
Field of view ('')	7.69
Source field of view ('')	3
PSF FWHM ('')	$\mathcal{TN}(0.06, 0.3; 0.08, 0.05)$ ¹
Noise amplitude σ	$\mathcal{TN}(0.001, 0.1; 0.01, 0.03)$

Building an inference pipeline with strong constraints on the slope of the profile, other than the lensed image, goes beyond the scope of this work. As such, imposing an implicit prior for the slope through the dataset is sufficient for our goal. It is also motivated by the *bulge-halo conspiracy* — the observation that most lensing configurations observed in the sky can be explained to first order approximation by an average slope consistent with an isothermal profile (Auger et al., 2010; Dutton and Treu, 2014).

TABLE 2.2 – SIE parameters.

Parameter	Distribution
Radial shift ('')	$\mathcal{U}(0, 0.1)$
Azimutal shift	$\mathcal{U}(0, 2\pi)$
Orientation	$\mathcal{U}(0, \pi)$
θ_E ('')	$\mathcal{U}(0.5, 2.5)$
Ellipticity	$\mathcal{U}(0, 0.6)$

1 600 000 simulated observations are generated from the VAE background sources and convergence maps as part of the training set. In principle, we could continuously generate examples from the VAE. However, having a fixed amount let us apply some validation check to each examples in order to avoid configurations like a single image of the background source or an Einstein ring cropped by the field of view.

2.4 Training

2.4.1 VAE

As mentionned in Kingma and Welling (2019), direct optimisation of the ELBO can prove difficult because the reconstruction term $\log p_\theta(\mathbf{x} | \mathbf{z})$ is relatively weak compared to the Kullback Leibler (KL) divergence term. To alleviate this issue, we follow the work of Bowman et al. (2015) and Kaae Sønderby et al. (2016) in setting a warm-up schedule for the KL term in thstarting from $\beta = 0.1$ up to β_{\max} .

Usually, $\beta_{\max} = 1$ is considered optimal since it matches the original ELBO objective derived

by Kingma and Welling (2013). But, we are more interested in the sharpness of our samples and accurate inference around small regions of the latent space for fine-tuning. Thus, setting $\beta_{\max} < 1$ allows us to increase the size of the information bottleneck (or latent space) of the VAE and improve the reconstruction cost of the model. This is a variant of the β -VAE (Higgins et al., 2017), where $\beta > 1$ was found to improve disentangling of the latent space (Burgess et al., 2018).

The value for β_{\max} and the steepness of the schedule are grid searched alongside the architecture for the VAE. Our criteria for an optimal model is a VAE that achieve the lowest reconstruction error in order to produce sharp images. At the same time, the KL divergence value should be smaller than an empirically defined threshold to respect the latent space prior. This value is found in practice by manually looking at the quality of generated samples for different VAE hyperparameters. A similar method is explored and formalized in the InfoVAE framework (Zhao et al., 2017).

A notable element of the VAE architecture is the use of a fully connected layer to reshape the features of the convolutional layer into the chosen latent space dimension. Following the work of Lanusse et al. (2021), we introduce an ℓ_2 penalty between the input and output of the bottleneck dense layers to encourage an identity mapping. This regularisation term is slowly removed during training.

2.4.2 RIM

The architecture of the gradient model was grid searched on smaller dataset ($\lesssim 10\,000$ examples) in order to quickly identify a small grid of valid hyperparameters. Then, the best hyperparameters were identified using a two-stage training process on the training dataset. In the first stage, we trained 24 different architectures from this small hyperparameter set for approximately 4 days (wall time using a single Nvidia A100 gpu). Different architectures would have a training time much longer than others, and this was factored in the architecture selection process. For example, adding more time steps (T) to the recurrent relation (2.4) would yield better generalisation on the test set, but this would come at great costs to training time until convergence. Following this first stage, 4 architectures were deemed efficient enough to be trained for an additional 6 days. We only report the results for the best architectures out of these 4.

Each reconstruction is performed by fine-tuning the baseline model on a test task composed of an observation vector, a PSF and a noise covariance. In practice, fine-tuning the test of 3 000 examples can be accomplished in parallel so as to be done in at most a few days by spreading the computation on ~ 10 Nvidia A100 GPUs (or 10 hours on ~ 100 GPUs). Each reconstruction uses at most 2000 steps, which turns out to be approximately 20 minutes (wall-time) per reconstruction. Early stopping is applied when the χ^2 reaches noise level. The hyperparameters for this procedure are reported in Table 2.3.

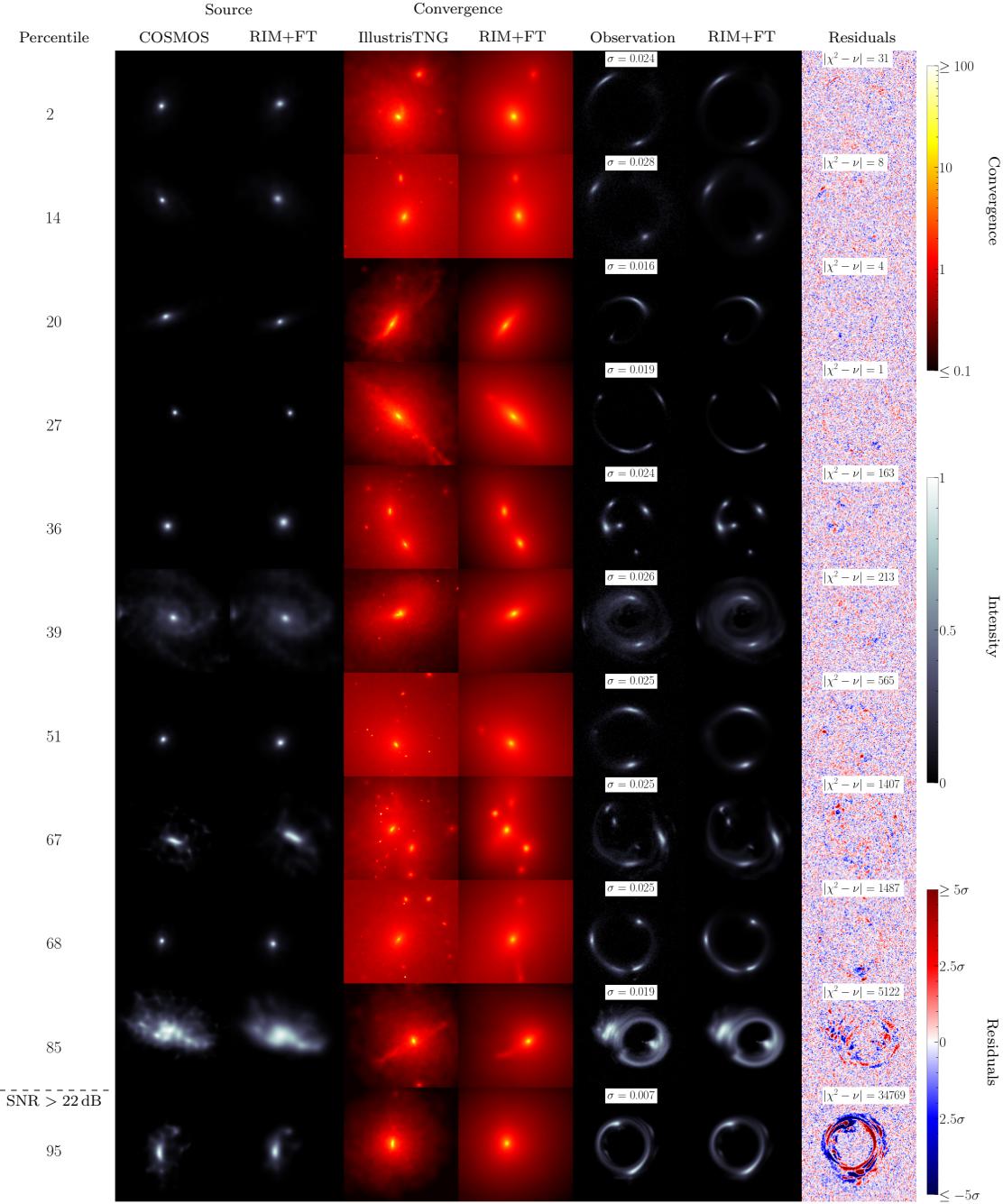


FIGURE 2.7 – Cherry-picked sample of the fine-tuned RIM reconstructions on a test set of 3000 examples. Examples are ordered from the best χ^2 (top) to the worst (bottom). The percentile rank of each example is in the leftmost column. The last example shown has SNR above the threshold defined in Figure 2.9.

2.5 Results

In this section, we present the performance of our approach on the held out test set. A sample of 3000 reconstruction problems is generated from the held-out HST and IllustrisTNG data with

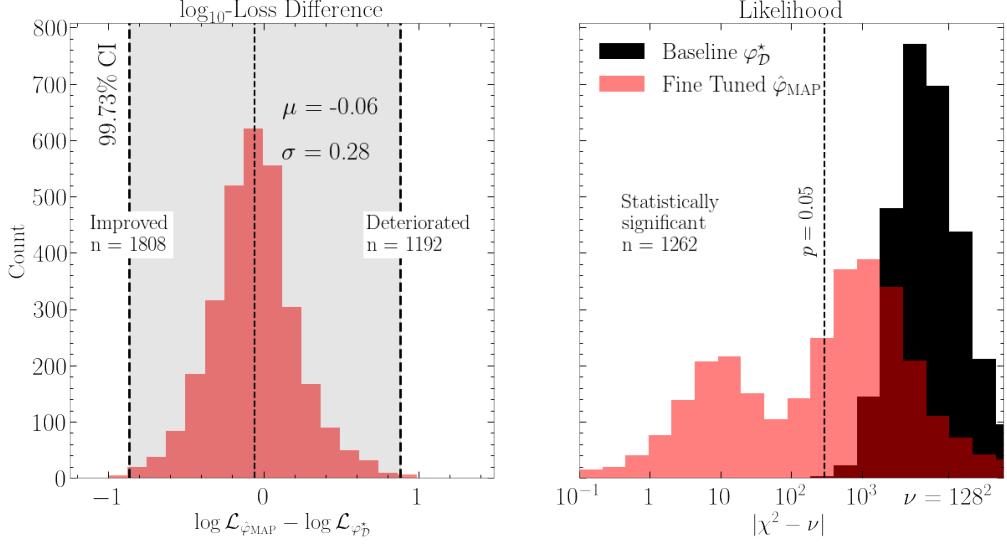


FIGURE 2.8 – Distribution of the goodness of fit for the baseline and fine-tuned network (right panel), as well as log-loss difference between the two network for a given example in the test set (left panel).

TABLE 2.3 – Hyperparameters for fine-tuning the RIM.

Parameter	Value
Optimizer	RMSProp
Learning rate	10^{-6}
Maximum number of steps	2 000
λ	2×10^5
ℓ_2	0
Number of samples from VAE	200
Latent space distribution	$\mathcal{N}(\mathbf{z}^{(T)}, \sigma = 0.3)^2$

noise conditions and PSFs similar to the training set.

2.5.1 Goodness of Fit

Figure 2.7 is a cherry picked sample of the reconstruction from the test set. The samples are selected to showcase the wide range of lensing configuration that our approach can successfully solve at high SNR. We made a point to select mostly examples that have a lot of structure in their convergence map to distinguish our approach from existing analytical methods. We did not make an emphasis in selecting complicated sources since we judge that free-form reconstruction of the source is essentially a solved problem. Indeed, many methods can reconstruct free-form sources once the convergence map is known or well constrained (Warren and Dye, 2003; Suyu et al., 2006; Vegetti and Koopmans, 2009; Birrer and Amara, 2018; Morningstar et al., 2019; Galan et al., 2021; Karchev et al., 2022; Mishra-Sharma and Yang, 2022).

To offset our selection bias, we selected samples in different percentile from the test set rank

TABLE 2.4 – \log_{10} -normal moments of the loss on the test set

Model	$\mu(\log \mathcal{L})$	$\sigma(\log \mathcal{L})$
Baseline (φ_D^*)	-1.96	0.36
Fine-tuned ($\hat{\varphi}_{\text{MAP}}$)	-2.02	0.37

ordered by the χ^2 metric. We also show a randomly selected sample from the test set in Figure C.2.

Figure 2.8 shows a comparison between the goodness of fit of the baseline model and the fine-tuned prediction. Since we empirically observe that the distribution of the loss on the test set (and the training set) follows a log-normal distribution, we find much more informative to look at the log-loss distribution to extract information about the fine-tuning procedure. The left panel of Figure 2.8 shows the distribution of the log-loss difference between the fine-tuned prediction and the baseline model. This distribution shows that the fine-tuning procedure loss is constrained within ~ 1 order of magnitude of the original loss with a 99.73 % probability. We find that the log-loss difference has a scatter of $\sigma = 0.28$, which is smaller than the scatter of the baseline log-loss over the entire test set $\sigma(\log \mathcal{L}(\varphi_D^*)) = 0.36$ reported in Table 2.4. We note that this metric is not optimized during fine-tuning, and is only computed as an oracle metric — still, the fine-tuning procedure does not significantly deteriorate or improve the loss of the baseline prediction in average. We report the first 2 moments of the loss log-normal distribution for the baseline and the fine-tuned reconstructions in Table 2.4 in order to compare explicitly the loss distributions. As can be seen in this table, there is no significant difference between the two distribution. This statement can be proven for the measured mean values — $\mu(\log \mathcal{L}(\hat{\varphi}_{\text{MAP}})) = \mu(\log \mathcal{L}(\varphi_D^*))$ — using the two-sided normal p-value test ([Casella and Berger, 2001](#)), which we find satisfy the null hypothesis with $p = 0.87288$ ($Z = -0.16$). All those observations support our claim that EWC regularisation preserves the prior learned during pretraining, or at least that it preserves the surrogate measures of the prior we reported.

The right panel of Figure 2.8 shows that the χ^2 of the reconstruction is improved substantially compared to the baseline model. The χ^2 is directly optimized by the fine-tuning procedure (equation (2.9)). Thus, this improvement is to be expected. Out of the 3000 reconstructions, 1262 reach $|\chi^2 - \nu| \leq 296$, which is the criteria that satisfy the null hypothesis for the number of degrees of freedom $\nu = 128^2$. For those reconstructions that do not satisfy this criteria, we still observe a significant improvement in the distribution of goodness of fit compared to the baseline (black histogram in Figure 2.8).

To understand why some of the reconstructions do not reach noise level, we report the joint distribution of the reduced chi squared and the SNR in Figure 2.9. Two behaviors can be identified. For SNR below a certain threshold, the goodness of fit of the fine-tuned model is essentially flat, with a certain scatter, around the noise level. This scatter increases as a function of SNR, which reflects the fact that high-SNR reconstructions are much more difficult to perform than low-SNR observations (for any model or algorithm, not just ours) because they require very precise assumptions about the solution. Near the threshold, this scatter is at its peak. This is SNR regime

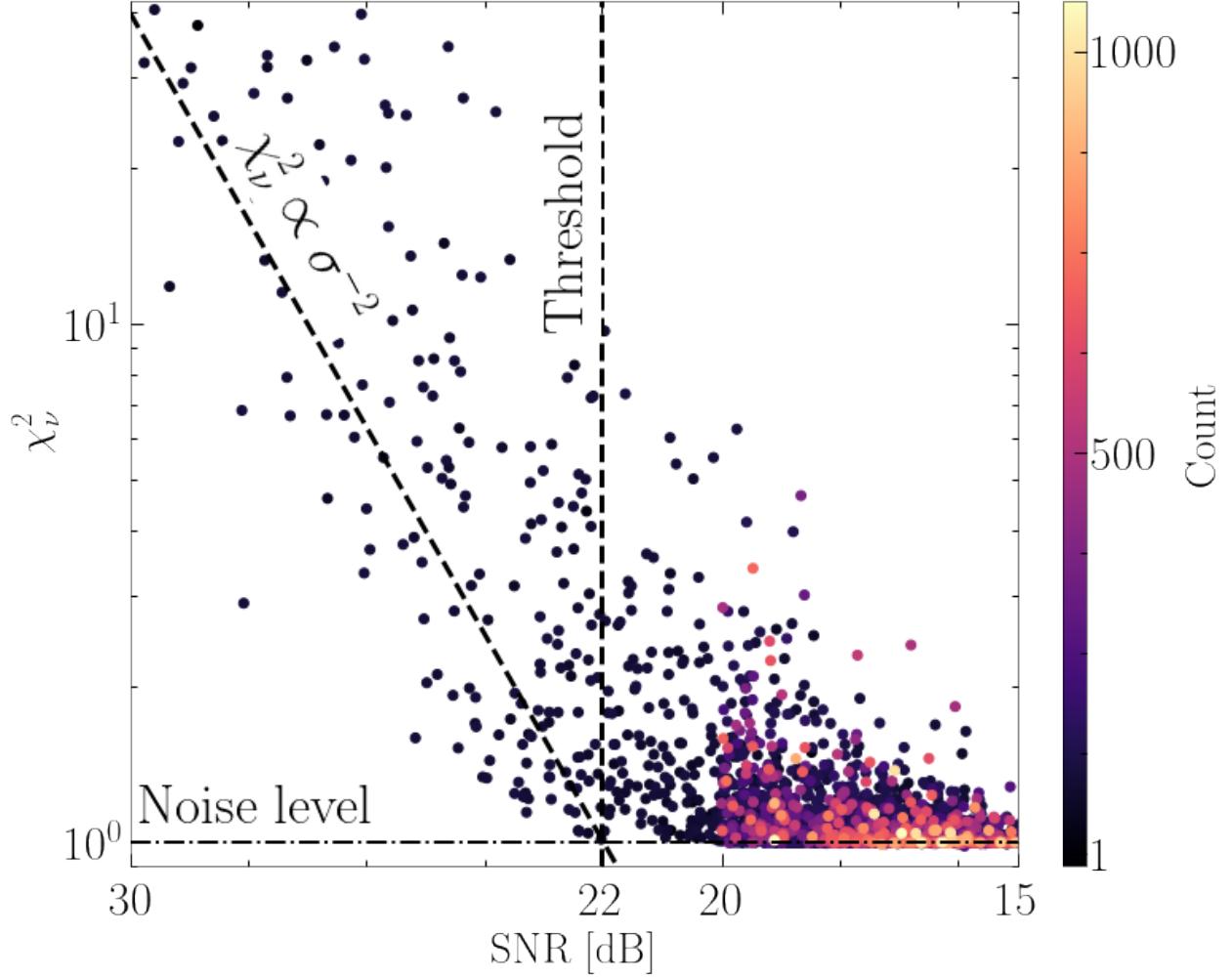


FIGURE 2.9 – Goodness of fit as a function of SNR shows a threshold behavior where our method reaches its limit.

($18 \text{ dB} \lesssim \text{SNR} \leq 22 \text{ dB}$) is where most of the fine-tuned reconstructions which do not reach noise level are found ($|\chi^2 - \nu| > 296$). For SNR above the threshold, the goodness of fit follows the trend $\chi^2 \propto \sigma^{-2}$, which means the reconstructions have stopped improving on par with the noise level. We define an optimistic threshold at the data point with largest SNR value which still have statistically significant residuals ($|\chi^2 - \nu| < 296$). We find this threshold value to be 22 dB based on the test set.

To understand this threshold behavior, we inspected some source reconstructions with intensity plotted in log scale in Figure 2.10. As can be seen in this figure, the source pixel intensities of the RIM prediction in the regime $\text{SNR} \gtrsim 22 \text{ dB}$, meaning the third decimal place of the reconstructed intensity values shown in pale blue, do not have a coherent structure and correspond mostly to numerical errors, even for the fine-tuned network. Reducing the noise level in the observed lensed image does not change this behavior, which indicates that this is a limitation of the neural network itself. Thus, we argue that the threshold can be explained in part because the inductive biases learned by the neural network during training prohibits the fine-tuning procedure to reconstruct

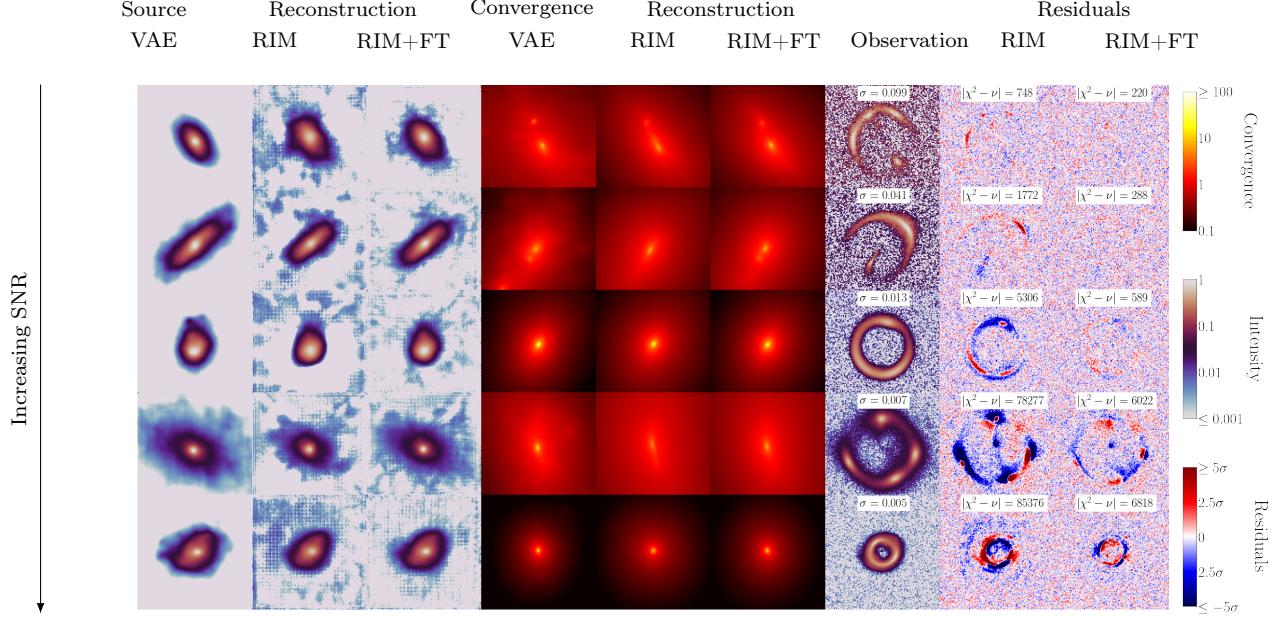


FIGURE 2.10 – Comparison between baseline (RIM) and fine-tuned (RIM+FT) reconstructions for VAE generated gravitational lensing systems. From top to bottom, we increase SNR. The first 2 rows have noise level reconstruction, while the last 3 row show significant improvement over the baseline. The intensity color scale is chosen to show the reconstruction down to the third decimal place, where the baseline prediction breaks down.

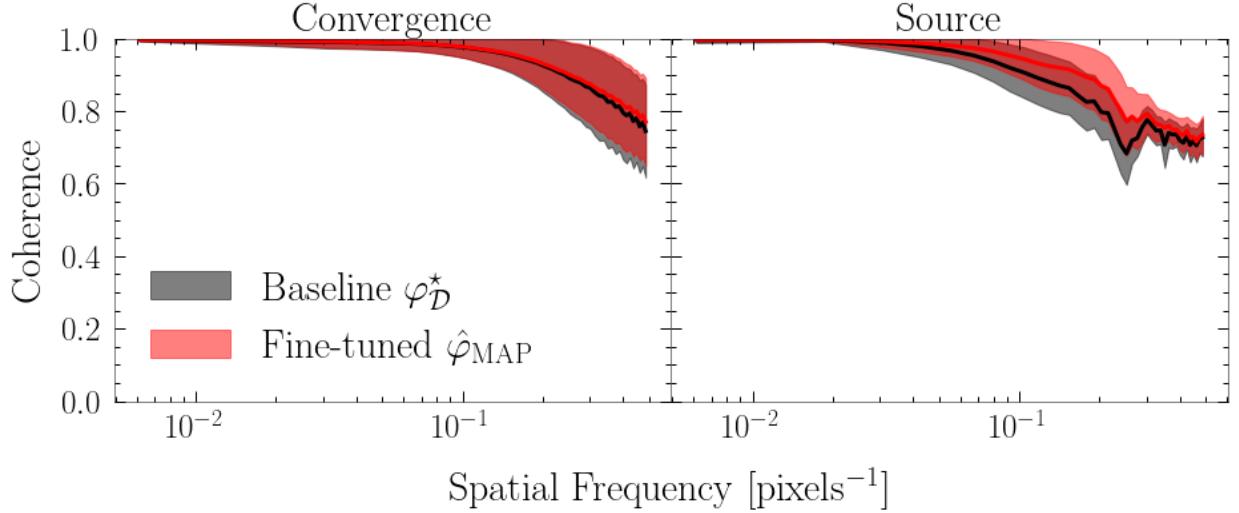


FIGURE 2.11 – Statistics of the coherence spectrum on the test set. The solid line is the average coherence. The transparent region is the 68% confidence interval. The fine-tuning procedure yields a noticeable improvement on the coherence of the source at all frequencies.

the signal past the threshold value via the EWC prior term. Reducing the strength of the prior will not improve the reconstruction since fine-tuning with $\lambda \ll 2 \times 10^5$ will quickly produce noise leakage in the source reconstruction.

Since the threshold value found is large enough to apply our method to most known gravitational lens images (Bolton et al., 2008; Shu et al., 2017), we leave further investigations to future works. We note that we purposefully ignored the low SNR regime since analytical methods perform well in that regime — at low SNR, simple Sersic (Sérsic, 1963) and SIE (Keeton, 2001) models will often give good fit to the data since most of the information about the source and the convergence morphology is lost.

2.5.2 Quality of the Reconstructions

In addition to a visual inspection of the reconstructed sources and convergences, we compute the coherence spectrum to quantitatively assess the quality of the reconstructions

$$\gamma(k) = \frac{P_{12}(k)}{\sqrt{P_{11}(k)P_{22}(k)}}. \quad (2.14)$$

$P_{ij}(k)$ is the cross power spectrum of images i and j at the wavenumber k . Figure 2.11 shows the mean value and the 68% inclusion interval of $\gamma(k)$ for the convergence and source maps in a test set of 3000 examples. The fine-tuning procedure, shown in red, is able to improve significantly the coherence of the baseline background source, shown in black, at all scales. The coherence spectrum of the convergence remains unchanged by the fine-tuning procedure. Still, we note that many examples in the dataset showcase significant improvement which we illustrate in Figure 2.1.

2.6 Conclusion

The results obtained here demonstrate the effectiveness of machine learning methods for inferring pixelated maps of the distribution of mass in lensing galaxies. Since this is a heavily under-constrained problem, stringent priors are needed to avoid overfitting the data, a task that has traditionally been difficult to accomplish (e.g., Saha and Williams, 1997). The model proposed here can implicitly learn these priors from a set of training data.

The flexible and expressive form of the reconstructions means that, in principle, any lensing system (e.g., a single simple galaxy, or a group of complex galaxies) could be analyzed by this model, without any need for pre-determining the model parameterization. This is of high value given the diversity of observed lensing systems, and their relevance for constraining astrophysical and cosmological parameters.

Another significant advantage of the proposed method is that the true physical model (a ray-tracing simulation code here) is used at every iteration to calculate the likelihood of the RIM predictions given the observations, which significantly helps with the interpretability of the model and the obtained results.

Finally, perhaps the most important limitation of the method is the fact that, in its current form,

the model only provides point estimates of the parameters of interest. Quantifying the posteriors of such high-dimensional data may require a generative process. In future work, we will explore the possibility of sampling from latent variables within the RIM to obtain samples from the posteriors.

Software and data

The source code, as well as the various scripts and parameters used to produce the model and results is available as open-source software under the package `Censai`³. The model parameters, as well as convergence maps used to train these models and the test set examples and reconstructions results are also available as open-source datasets hosted by Zenodo⁴. This research made use of `Tensorflow` (Abadi et al., 2015), `Tensorflow-Probability` (Dillon et al., 2017), `Numpy` (Harris et al., 2020), `Scipy` (Virtanen et al., 2020), `Matplotlib` (Hunter, 2007), `Scikit-image` (Van der Walt et al., 2014), `IPython` (Pérez and Granger, 2007), `Pandas` (Wes McKinney, 2010; pandas development team, 2020), `Scikit-learn` (Pedregosa et al., 2011), `Astropy` (Astropy Collaboration et al., 2013, 2018) and `GalSim` (Rowe et al., 2015).

Acknowledgements

We thank Ronan Legin for fruitful discussion and insights about training the neural network and comments about the manuscript. This research was supported by the Schmidt Futures Foundation. The work was also enabled in part by computational resources provided by Calcul Quebec, Compute Canada and the Digital Research Alliance of Canada. Y.H. and L.P. acknowledge support from the National Sciences and Engineering Council of Canada grant RGPIN-2020-05102, the Fonds de recherche du Québec grant 2022-NC-301305, and the Canada Research Chairs Program. A.A. was supported by an IVADO scholarship.

3.  <https://github.com/AlexandreAdam/Censai>

4.  <https://doi.org/10.5281/zenodo.6555463>

Bibliographie

- M. Abadi, A. Agarwal, P. Barham, E. Brevdo, Z. Chen, C. Citro, G. S. Corrado, A. Davis, J. Dean, M. Devin, S. Ghemawat, I. Goodfellow, A. Harp, G. Irving, M. Isard, Y. Jia, R. Jozefowicz, L. Kaiser, M. Kudlur, J. Levenberg, D. Mané, R. Monga, S. Moore, D. Murray, C. Olah, M. Schuster, J. Shlens, B. Steiner, I. Sutskever, K. Talwar, P. Tucker, V. Vanhoucke, V. Vasudevan, F. Viégas, O. Vinyals, P. Warden, M. Wattenberg, M. Wicke, Y. Yu, and X. Zheng. TensorFlow : Large-scale machine learning on heterogeneous systems, 2015. URL <https://www.tensorflow.org/>. Software available from tensorflow.org.
- H. M. Abdelsalam, P. Saha, and L. L. R. Williams. Nonparametric Reconstruction of Abell 2218 from Combined Weak and Strong Lensing. *AJ*, 116(4) :1541–1552, Oct. 1998a. doi : 10.1086/300546.
- H. M. Abdelsalam, P. Saha, and L. L. R. Williams. Non-parametric reconstruction of cluster mass distribution from strong lensing : modelling Abell 370. *MNRAS*, 294 :734–746, Mar. 1998b. doi : 10.1046/j.1365-8711.1998.01356.x.
- N. Anau Montel, A. Coogan, C. Correa, K. Karchev, and C. Weniger. Estimating the warm dark matter mass from strong lensing images with truncated marginal neural ratio estimation. *arXiv e-prints*, art. arXiv :2205.09126, May 2022.
- M. Andrychowicz, M. Denil, S. Gomez, M. W. Hoffman, D. Pfau, T. Schaul, B. Shillingford, and N. de Freitas. Learning to learn by gradient descent by gradient descent. *arXiv e-prints*, art. arXiv :1606.04474, June 2016.
- Astropy Collaboration, T. P. Robitaille, E. J. Tollerud, P. Greenfield, M. Droettboom, E. Bray, T. Aldcroft, M. Davis, A. Ginsburg, A. M. Price-Whelan, W. E. Kerzendorf, A. Conley, N. Crighton, K. Barbary, D. Muna, H. Ferguson, F. Grollier, M. M. Parikh, P. H. Nair, H. M. Unther, C. Deil, J. Woillez, S. Conseil, R. Kramer, J. E. H. Turner, L. Singer, R. Fox, B. A. Weaver, V. Zabalza, Z. I. Edwards, K. Azalee Bostroem, D. J. Burke, A. R. Casey, S. M. Crawford, N. Dencheva, J. Ely, T. Jenness, K. Labrie, P. L. Lim, F. Pierfederici, A. Pontzen, A. Ptak, B. Refsdal, M. Servillat, and O. Streicher. Astropy : A community Python package for astronomy. *A&A*, 558 :A33, Oct. 2013. doi : 10.1051/0004-6361/201322068.
- Astropy Collaboration, A. M. Price-Whelan, B. M. Sipőcz, H. M. Günther, P. L. Lim, S. M. Crawford, S. Conseil, D. L. Shupe, M. W. Craig, N. Dencheva, A. Ginsburg, J. T. Vand erPlas, L. D. Bradley, D. Pérez-Suárez, M. de Val-Borro, T. L. Aldcroft, K. L. Cruz, T. P. Robitaille, E. J. Tollerud, C. Ardelean, T. Babej, Y. P. Bach, M. Bachetti, A. V. Bakanov, S. P. Bamford, G. Barentsen, P. Barmby, A. Baumbach, K. L. Berry, F. Biscani, M. Boquien, K. A. Bostroem, L. G. Bouma, G. B. Brammer, E. M. Bray, H. Breytenbach, H. Buddelmeijer, D. J. Burke, G. Calderone, J. L. Cano Rodríguez, M. Cara, J. V. M. Cardoso, S. Cheedella, Y. Copin, L. Corrales, D. Crichton, D. D'Avella, C. Deil, É. Depagne, J. P. Dietrich, A. Donath, M. Droettboom, N. Earl, T. Erben, S. Fabbro, L. A. Ferreira, T. Finethy, R. T. Fox, L. H. Garrison, S. L. J. Gibbons, D. A. Goldstein, R. Gommers, J. P. Greco, P. Greenfield, A. M. Groener, F. Grollier, A. Hagen, P. Hirst, D. Homeier, A. J. Horton, G. Hosseinzadeh, L. Hu, J. S. Hunkeler, Ž. Ivezić, A. Jain, T. Jenness, G. Kanarek, S. Kendrew, N. S. Kern, W. E. Kerzendorf, A. Khvalko, J. King, D. Kirkby, A. M. Kulkarni, A. Kumar, A. Lee, D. Lenz, S. P. Littlefair, Z. Ma, D. M. Macleod, M. Mastropietro, C. McCully, S. Montagnac, B. M. Morris, M. Mueller, S. J. Mumford, D. Muna, N. A. Murphy, S. Nelson, G. H.

- Nguyen, J. P. Ninan, M. Nöthe, S. Ogaz, S. Oh, J. K. Parejko, N. Parley, S. Pascual, R. Patil, A. A. Patil, A. L. Plunkett, J. X. Prochaska, T. Rastogi, V. Reddy Janga, J. Sabater, P. Sakurikar, M. Seifert, L. E. Sherbert, H. Sherwood-Taylor, A. Y. Shih, J. Sick, M. T. Silbiger, S. Singanamalla, L. P. Singer, P. H. Sladen, K. A. Sooley, S. Sornarajah, O. Streicher, P. Teuben, S. W. Thomas, G. R. Tremblay, J. E. H. Turner, V. Terrón, M. H. van Kerkwijk, A. de la Vega, L. L. Watkins, B. A. Weaver, J. B. Whitmore, J. Woillez, V. Zabalza, and Astropy Contributors. The Astropy Project : Building an Open-science Project and Status of the v2.0 Core Package. AJ, 156(3) :123, Sept. 2018. doi : 10.3847/1538-3881/aabc4f.
- D. Aubert, A. Amara, and R. B. Metcalf. Smooth Particle Lensing. MNRAS, 376(1) :113–124, Mar. 2007. doi : 10.1111/j.1365-2966.2006.11296.x.
- M. W. Auger, T. Treu, A. S. Bolton, R. Gavazzi, L. V. E. Koopmans, P. J. Marshall, L. A. Moustakas, and S. Burles. The Sloan Lens ACS Survey. X. Stellar, Dynamical, and Total Mass Correlations of Massive Early-type Galaxies. ApJ, 724(1) :511–525, Nov. 2010. doi : 10.1088/0004-637X/724/1/511.
- M. Barnabè, O. Czoske, L. V. E. Koopmans, T. Treu, A. S. Bolton, and R. Gavazzi. Two-dimensional kinematics of SLACS lenses - II. Combined lensing and dynamics analysis of early-type galaxies at $z = 0.08\text{--}0.33$. MNRAS, 399(1) :21–36, Oct. 2009. doi : 10.1111/j.1365-2966.2009.14941.x.
- M. Bartelmann, R. Narayan, S. Seitz, and P. Schneider. Maximum-likelihood Cluster Reconstruction. ApJ, 464 :L115, June 1996. doi : 10.1086/310114.
- F. Bellagamba, N. Tessore, and R. B. Metcalf. Zooming into the Cosmic Horseshoe : new insights on the lens profile and the source shape. *Monthly Notices of the Royal Astronomical Society*, 464(4) :4823–4834, 10 2016. ISSN 0035-8711. doi : 10.1093/mnras/stw2726. URL <https://doi.org/10.1093/mnras/stw2726>.
- V. Belokurov, N. W. Evans, A. Moiseev, L. J. King, P. C. Hewett, M. Pettini, L. Wyrzykowski, R. G. McMahon, M. C. Smith, G. Gilmore, S. F. Sanchez, A. Udalski, S. Koposov, D. B. Zucker, and C. J. Walcher. The Cosmic Horseshoe : Discovery of an Einstein Ring around a Giant Luminous Red Galaxy. ApJ, 671(1) :L9–L12, Dec. 2007. doi : 10.1086/524948.
- Y. Bengio. Learning deep architectures for ai. *Found. Trends Mach. Learn.*, 2(1) :1–127, jan 2009. ISSN 1935-8237. doi : 10.1561/2200000006. URL <https://doi.org/10.1561/2200000006>.
- S. Birrer and A. Amara. lenstronomy : Multi-purpose gravitational lens modelling software package. *Physics of the Dark Universe*, 22 :189–201, Dec. 2018. doi : 10.1016/j.dark.2018.11.002.
- S. Birrer, A. Amara, and A. Refregier. Gravitational Lens Modeling with Basis Sets. ApJ, 813(2) :102, Nov. 2015. doi : 10.1088/0004-637X/813/2/102.
- S. Birrer, T. Treu, C. E. Rusu, V. Bonvin, C. D. Fassnacht, J. H. H. Chan, A. Agnello, A. J. Shajib, G. C. F. Chen, M. Auger, F. Courbin, S. Hilbert, D. Sluse, S. H. Suyu, K. C. Wong, P. Marshall, B. C. Lemaux, and G. Meylan. H0LiCOW - IX. Cosmographic analysis of the doubly imaged quasar SDSS 1206+4332 and a new measurement of the Hubble constant. MNRAS, 484(4) :4726–4753, Apr. 2019. doi : 10.1093/mnras/stz200.
- A. S. Bolton, S. Burles, L. V. E. Koopmans, T. Treu, R. Gavazzi, L. A. Moustakas, R. Wayth, and D. J. Schlegel. The Sloan Lens ACS Survey. V. The Full ACS Strong-Lens Sample. ApJ, 682(2) :964–984, Aug. 2008. doi : 10.1086/589327.
- S. R. Bowman, L. Vilnis, O. Vinyals, A. M. Dai, R. Jozefowicz, and S. Bengio. Generating Sentences from a Continuous Space. *arXiv e-prints*, art. arXiv :1511.06349, Nov. 2015.

- M. Bradač, P. Schneider, M. Lombardi, and T. Erben. Strong and weak lensing united. I. The combined strong and weak lensing cluster mass reconstruction method. *A&A*, 437(1) :39–48, July 2005. doi : 10.1051/0004-6361:20042233.
- C. P. Burgess, I. Higgins, A. Pal, L. Matthey, N. Watters, G. Desjardins, and A. Lerchner. Understanding disentangling in β -VAE. *arXiv e-prints*, art. arXiv :1804.03599, Apr. 2018.
- M. Cacciato, M. Bartelmann, M. Meneghetti, and L. Moscardini. Combining weak and strong lensing in cluster potential reconstruction. *A&A*, 458(2) :349–356, Nov. 2006. doi : 10.1051/0004-6361:20054582.
- S. Carroll. *Spacetime and Geometry : An Introduction to General Relativity*. Benjamin Cummings, 2003. ISBN 0805387323.
- G. Casella and R. Berger. *Statistical Inference*. Duxbury Resource Center, June 2001. ISBN 0534243126.
- J. Cheng, M. P. Wiesner, E.-H. Peng, W. Cui, J. R. Peterson, and G. Li. Adaptive Grid Lens Modeling of the Cosmic Horseshoe Using Hubble Space Telescope Imaging. *ApJ*, 872(2) :185, Feb. 2019. doi : 10.3847/1538-4357/ab0029.
- K. Cho, B. van Merriënboer, C. Gulcehre, D. Bahdanau, F. Bougares, H. Schwenk, and Y. Bengio. Learning Phrase Representations using RNN Encoder-Decoder for Statistical Machine Translation. *arXiv e-prints*, art. arXiv :1406.1078, June 2014.
- D. Coe, E. Fuselier, N. Benítez, T. Broadhurst, B. Frye, and H. Ford. LensPerfect : Gravitational Lens Mass Map Reconstructions Yielding Exact Reproduction of All Multiple Images. *ApJ*, 681(2) :814–830, July 2008. doi : 10.1086/588250.
- J. P. Coles, J. I. Read, and P. Saha. Gravitational lens recovery with GLASS : measuring the mass profile and shape of a lens. *MNRAS*, 445(3) :2181–2197, Dec. 2014. doi : 10.1093/mnras/stu1781.
- A. Coogan, K. Karchev, and C. Weniger. Targeted Likelihood-Free Inference of Dark Matter Substructure in Strongly-Lensed Galaxies. *arXiv e-prints*, art. arXiv :2010.07032, Oct. 2020.
- H. Cramér. *Mathematical methods of statistics*, volume 9. Princeton University Press, Princeton, NJ, 1946.
- N. Dalal and C. S. Kochanek. Direct Detection of Cold Dark Matter Substructure. *ApJ*, 572(1) :25–33, June 2002. doi : 10.1086/340303.
- S. Deb, D. M. Goldberg, and V. J. Ramdass. Reconstruction of Cluster Masses Using Particle Based Lensing. I. Application to Weak Lensing. *ApJ*, 687(1) :39–49, Nov. 2008. doi : 10.1086/590544.
- S. Deb, A. Morandi, K. Pedersen, S. Riemer-Sorensen, D. M. Goldberg, and H. Dahle. Mass Reconstruction using Particle Based Lensing II : Quantifying substructure with Strong+Weak lensing and X-rays. *arXiv e-prints*, art. arXiv :1201.3636, Jan. 2012.
- J. M. Diego, P. Protopapas, H. B. Sandvik, and M. Tegmark. Non-parametric inversion of strong lensing systems. *MNRAS*, 360(2) :477–491, June 2005. doi : 10.1111/j.1365-2966.2005.09021.x.
- J. M. Diego, M. Tegmark, P. Protopapas, and H. B. Sandvik. Combined reconstruction of weak and strong lensing data with WSLAP. *MNRAS*, 375(3) :958–970, Mar. 2007. doi : 10.1111/j.1365-2966.2007.11380.x.
- J. V. Dillon, I. Langmore, D. Tran, E. Brevdo, S. Vasudevan, D. Moore, B. Patton, A. Alemi, M. Hoffman, and R. A. Saurous. TensorFlow Distributions. *arXiv e-prints*, art. arXiv :1711.10604, Nov. 2017.

- A. A. Dutton and T. Treu. The bulge–halo conspiracy in massive elliptical galaxies : implications for the stellar initial mass function and halo response to baryonic processes. *Monthly Notices of the Royal Astronomical Society*, 438(4) :3594–3602, 01 2014. ISSN 0035-8711. doi : 10.1093/mnras/stt2489. URL <https://doi.org/10.1093/mnras/stt2489>.
- A. Einstein. Lens-Like Action of a Star by the Deviation of Light in the Gravitational Field. *Science*, 84(2188) :506–507, 1936. doi : 10.1126/science.84.2188.506. URL <https://www.science.org/doi/abs/10.1126/science.84.2188.506>.
- A. Galan, A. Peel, R. Joseph, F. Courbin, and J. L. Starck. SLITRONOMY : Towards a fully wavelet-based strong lensing inversion technique. *A&A*, 647 :A176, Mar. 2021. doi : 10.1051/0004-6361/202039363.
- A. Ghosh, L. L. R. Williams, and J. Liesenborgs. Free-form gralre lens inversion of galaxy clusters with up to 1000 multiple images. *MNRAS*, 494(3) :3998–4014, May 2020. doi : 10.1093/mnras/staa962.
- D. Gilman, S. Birrer, A. Nierenberg, T. Treu, X. Du, and A. Benson. Warm dark matter chills out : constraints on the halo mass function and the free-streaming length of dark matter with eight quadruple-image strong gravitational lenses. *MNRAS*, 491(4) :6077–6101, Feb. 2020. doi : 10.1093/mnras/stz3480.
- D. Gilman, J. Bovy, T. Treu, A. Nierenberg, S. Birrer, A. Benson, and O. Sameie. Strong lensing signatures of self-interacting dark matter in low-mass haloes. *MNRAS*, 507(2) :2432–2447, Oct. 2021. doi : 10.1093/mnras/stab2335.
- Z. Goldfeld and Y. Polyanskiy. The Information Bottleneck Problem and Its Applications in Machine Learning. *arXiv e-prints*, art. arXiv :2004.14941, Apr. 2020.
- C. R. Harris, K. J. Millman, S. J. van der Walt, R. Gommers, P. Virtanen, D. Cournapeau, E. Wieser, J. Taylor, S. Berg, N. J. Smith, R. Kern, M. Picus, S. Hoyer, M. H. van Kerkwijk, M. Brett, A. Haldane, J. F. del Río, M. Wiebe, P. Peterson, P. Gérard-Marchant, K. Sheppard, T. Reddy, W. Weckesser, H. Abbasi, C. Gohlke, and T. E. Oliphant. Array programming with NumPy. *Nature*, 585(7825) :357–362, Sept. 2020. doi : 10.1038/s41586-020-2649-2. URL <https://doi.org/10.1038/s41586-020-2649-2>.
- Y. D. Hezaveh, N. Dalal, D. P. Marrone, Y.-Y. Mao, W. Morningstar, D. Wen, R. D. Blandford, J. E. Carlstrom, C. D. Fassnacht, G. P. Holder, A. Kemball, P. J. Marshall, N. Murray, L. Perreault Levasseur, J. D. Vieira, and R. H. Wechsler. Detection of Lensing Substructure Using ALMA Observations of the Dusty Galaxy SDP.81. *ApJ*, 823(1) :37, May 2016. doi : 10.3847/0004-637X/823/1/37.
- Y. D. Hezaveh, L. Perreault Levasseur, and P. J. Marshall. Fast automated analysis of strong gravitational lenses with convolutional neural networks. *Nature*, 548(7669) :555–557, Aug. 2017. doi : 10.1038/nature23463.
- I. Higgins, L. Matthey, A. Pal, C. P. Burgess, X. Glorot, M. M. Botvinick, S. Mohamed, and A. Lerchner. beta-vae : Learning basic visual concepts with a constrained variational framework. In *ICLR*, 2017.
- J. D. Hunter. Matplotlib : A 2d graphics environment. *Computing in Science & Engineering*, 9(3) :90–95, 2007. doi : 10.1109/MCSE.2007.55.
- B. L. James, M. Auger, M. Pettini, D. P. Stark, V. Belokurov, and S. Carniani. Mapping UV properties throughout the Cosmic Horseshoe : lessons from VLT-MUSE. *MNRAS*, 476(2) :1726–1740, May 2018. doi : 10.1093/mnras/sty315.
- M. J. Jee, H. C. Ford, G. D. Illingworth, R. L. White, T. J. Broadhurst, D. A. Coe, G. R. Meurer, A. van der Wel, N. Benítez, J. P. Blakeslee, R. J. Bouwens, L. D. Bradley, R. Demarco, N. L. Homeier, A. R. Martel, and S. Mei. Discovery of a Ringlike Dark Matter Structure in the Core of the Galaxy Cluster Cl 0024+17. *ApJ*, 661(2) :728–749, June 2007. doi : 10.1086/517498.

- N. Jeffrey, F. Lanusse, O. Lahav, and J.-L. Starck. Deep learning dark matter map reconstructions from DES SV weak lensing data. *MNRAS*, 492(4) :5023–5029, Mar. 2020. doi : 10.1093/mnras/staa127.
- C. Kaae Sønderby, T. Raiko, L. Maaløe, S. Kaae Sønderby, and O. Winther. Ladder Variational Autoencoders. *arXiv e-prints*, art. arXiv :1602.02282, Feb. 2016.
- N. Kaiser and G. Squires. Mapping the Dark Matter with Weak Gravitational Lensing. *ApJ*, 404 :441, Feb. 1993. doi : 10.1086/172297.
- K. Karchev, A. Coogan, and C. Weniger. Strong-lensing source reconstruction with variationally optimized Gaussian processes. *MNRAS*, 512(1) :661–685, May 2022. doi : 10.1093/mnras/stac311.
- C. R. Keeton. A Catalog of Mass Models for Gravitational Lensing. *arXiv e-prints*, art. astro-ph/0102341, Feb. 2001.
- D. P. Kingma and M. Welling. Auto-Encoding Variational Bayes. *arXiv e-prints*, art. arXiv :1312.6114, Dec. 2013.
- D. P. Kingma and M. Welling. An Introduction to Variational Autoencoders. *arXiv e-prints*, art. arXiv :1906.02691, June 2019.
- J. Kirkpatrick, R. Pascanu, N. Rabinowitz, J. Veness, G. Desjardins, A. A. Rusu, K. Milan, J. Quan, T. Ramalho, A. Grabska-Barwinska, D. Hassabis, C. Clopath, D. Kumaran, and R. Hadsell. Overcoming catastrophic forgetting in neural networks. *arXiv e-prints*, art. arXiv :1612.00796, Dec. 2016.
- A. M. Koekemoer, H. Aussel, D. Calzetti, P. Capak, M. Giavalisco, J.-P. Kneib, A. Leauthaud, O. Le Fevre, H. J. McCracken, R. Massey, B. Mobasher, J. Rhodes, N. Scoville, and P. L. Shopbell. The COSMOS Survey : Hubble Space Telescope Advanced Camera for Surveys Observations and Data Processing. *The Astrophysical Journal Supplement Series*, 172(1) :196–202, sep 2007. ISSN 0067-0049. doi : 10.1086/520086.
- L. V. E. Koopmans. Gravitational imaging of cold dark matter substructures. *MNRAS*, 363(4) :1136–1144, Nov. 2005. doi : 10.1111/j.1365-2966.2005.09523.x.
- L. V. E. Koopmans, T. Treu, A. S. Bolton, S. Burles, and L. A. Moustakas. The Sloan Lens ACS Survey. III. The Structure and Formation of Early-Type Galaxies and Their Evolution since $z \sim 1$. *ApJ*, 649(2) :599–615, Oct. 2006. doi : 10.1086/505696.
- F. Lanusse, J. L. Starck, A. Leonard, and S. Pires. High resolution weak lensing mass mapping combining shear and flexion. *A&A*, 591 :A2, June 2016. doi : 10.1051/0004-6361/201628278.
- F. Lanusse, R. Mandelbaum, S. Ravanbakhsh, C.-L. Li, P. Freeman, and B. Póczos. Deep generative models for galaxy image simulations. *MNRAS*, 504(4) :5543–5555, July 2021. doi : 10.1093/mnras/stab1214.
- A. Leauthaud, R. Massey, J.-P. Kneib, J. Rhodes, D. E. Johnston, P. Capak, C. Heymans, R. S. Ellis, A. M. Koekemoer, O. L. Fèvre, Y. Mellier, A. Réfrégier, A. C. Robin, N. Scoville, L. Tasca, J. E. Taylor, and L. V. Waerbeke. Weak Gravitational Lensing with COSMOS : Galaxy Selection and Shape Measurements. *The Astrophysical Journal Supplement Series*, 172(1) :219, sep 2007. ISSN 0067-0049. doi : 10.1086/516598.
- R. Legin, Y. Hezaveh, L. Perreault Levasseur, and B. Wandelt. Simulation-Based Inference of Strong Gravitational Lensing Parameters. *arXiv e-prints*, art. arXiv :2112.05278, Dec. 2021.
- A. Leonard, F. X. Dupé, and J. L. Starck. A compressed sensing approach to 3D weak lensing. *A&A*, 539 :A85, Mar. 2012. doi : 10.1051/0004-6361/201117642.
- N. Li, C. Becker, and S. Dye. The impact of line-of-sight structures on measuring H_0 with strong lensing time delays. *MNRAS*, 504(2) :2224–2234, June 2021. doi : 10.1093/mnras/stab984.

- J. Liesenborgs, S. De Rijcke, and H. Dejonghe. A genetic algorithm for the non-parametric inversion of strong lensing systems. *MNRAS*, 367(3) :1209–1216, Apr. 2006. doi : 10.1111/j.1365-2966.2006.10040.x.
- J. Liesenborgs, S. de Rijcke, H. Dejonghe, and P. Bekaert. Non-parametric inversion of gravitational lensing systems with few images using a multi-objective genetic algorithm. *MNRAS*, 380(4) :1729–1736, Oct. 2007. doi : 10.1111/j.1365-2966.2007.12236.x.
- K. Lønning, P. Putzky, J. J. Sonke, L. Reneman, M. W. Caan, and M. Welling. Recurrent inference machines for reconstructing heterogeneous MRI data. *Medical Image Analysis*, 53 :64–78, apr 2019. ISSN 13618423. doi : 10.1016/j.media.2019.01.005.
- R. Mandelbaum, C. Lackner, A. Leauthaud, and B. Rowe. COSMOS real galaxy dataset. *Zenodo*, jan 2012. URL <https://zenodo.org/record/3242143>.
- R. Mandelbaum, B. Rowe, J. Bosch, C. Chang, F. Courbin, M. Gill, M. Jarvis, A. Kannawadi, T. Kacprzak, C. Lackner, A. Leauthaud, H. Miyatake, R. Nakajima, J. Rhodes, M. Simet, J. Zuntz, B. Armstrong, S. Bridle, J. Coupon, J. P. Dietrich, M. Gentile, C. Heymans, A. S. Jurling, S. M. Kent, D. Kirkby, D. Margala, R. Massey, P. Melchior, J. Peterson, A. Roodman, and T. Schrabback. The Third Gravitational Lensing Accuracy Testing (GREAT3) Challenge Handbook. *The Astrophysical Journal Supplement Series*, 212(1) :5, apr 2014. ISSN 0067-0049. doi : 10.1088/0067-0049/212/1/5. URL <https://iopscience.iop.org/article/10.1088/0067-0049/212/1/5>
<https://iopscience.iop.org/article/10.1088/0067-0049/212/1/5/meta>.
- D. P. Marrone, J. S. Spilker, C. C. Hayward, J. D. Vieira, M. Aravena, M. L. N. Ashby, M. B. Bayliss, M. Béthermin, M. Brodwin, M. S. Bothwell, J. E. Carlstrom, S. C. Chapman, C.-C. Chen, T. M. Crawford, D. J. M. Cunningham, C. De Breuck, C. D. Fassnacht, A. H. Gonzalez, T. R. Greve, Y. D. Hezaveh, K. Lacaille, K. C. Litke, S. Lower, J. Ma, M. Malkan, T. B. Miller, W. R. Morningstar, E. J. Murphy, D. Narayanan, K. A. Phadke, K. M. Rotermund, J. Sreevani, B. Stalder, A. A. Stark, M. L. Strandet, M. Tang, and A. Weiß. Galaxy growth in a massive halo in the first billion years of cosmic history. *Nature*, 553(7686) :51–54, Jan. 2018. doi : 10.1038/nature24629.
- P. J. Marshall. Maximum-entropy reconstruction of the distribution of mass in cluster MS1054-03 from weak lensing data. In D. M. Neumann and J. T. V. Tran, editors, *Clusters of Galaxies and the High Redshift Universe Observed in X-rays*, page 47, Jan. 2001.
- R. Massey, J. Rhodes, R. Ellis, N. Scoville, A. Leauthaud, A. Finoguenov, P. Capak, D. Bacon, H. Aussel, J.-P. Kneib, A. Koekemoer, H. McCracken, B. Mobasher, S. Pires, A. Refregier, S. Sasaki, J.-L. Starck, Y. Taniguchi, A. Taylor, and J. Taylor. Dark matter maps reveal cosmic scaffolding. *Nature*, 445(7125) :286–290, Jan. 2007. doi : 10.1038/nature05497.
- M. McCloskey and N. J. Cohen. Catastrophic Interference in Connectionist Networks : The Sequential Learning Problem. In G. H. Bower, editor, *Psychology of Learning and Motivation - Advances in Research and Theory*, volume 24 of *Psychology of Learning and Motivation*, pages 109–165. Academic Press, 1989. doi : 10.1016/S0079-7421(08)60536-8. URL <https://www.sciencedirect.com/science/article/pii/S0079742108605368>.
- M. Meneghetti. *Introduction to Gravitational Lensing*. Springer Cham, 2013. doi : 10.1007/978-3-030-73582-1.
- J. Merten. Mesh-free free-form lensing - I. Methodology and application to mass reconstruction. *MNRAS*, 461(3) :2328–2345, Sept. 2016. doi : 10.1093/mnras/stw1413.
- J. Merten, M. Cacciato, M. Meneghetti, C. Mignone, and M. Bartelmann. Combining weak and strong cluster lensing : applications to simulations and MS 2137. *A&A*, 500(2) :681–691, June 2009. doi : 10.1051/0004-6361/200810372.
- S. Mishra-Sharma and G. Yang. Strong lensing source reconstruction using continuous neural fields, 2022. URL <https://arxiv.org/abs/2206.14820>.

- C. Modi, F. Lanusse, U. Seljak, D. N. Spergel, and L. Perreault-Levasseur. CosmicRIM : Reconstructing Early Universe by Combining Differentiable Simulations with Recurrent Inference Machines. *arXiv e-prints*, art. arXiv :2104.12864, Apr. 2021.
- W. R. Morningstar, Y. D. Hezaveh, L. P. Levasseur, R. D. Blandford, P. J. Marshall, P. Putzky, and R. H. Wechsler. Analyzing Interferometric Observations of Strong Gravitational Lenses with Recurrent and Convolutional Neural Networks. *arXiv e-prints*, 2018.
- W. R. Morningstar, L. P. Levasseur, Y. D. Hezaveh, R. Blandford, P. Marshall, P. Putzky, T. D. Rueter, R. Wechsler, and M. Welling. Data-driven Reconstruction of Gravitationally Lensed Galaxies Using Recurrent Inference Machines. *The Astrophysical Journal*, 883(1) :14, 2019. ISSN 1538-4357. doi : 10.3847/1538-4357/ab35d7.
- D. Nelson, V. Springel, A. Pillepich, V. Rodriguez-Gomez, P. Torrey, S. Genel, M. Vogelsberger, R. Pakmor, F. Marinacci, R. Weinberger, L. Kelley, M. Lovell, B. Diemer, and L. Hernquist. The IllustrisTNG simulations : public data release. *MNRAS*, 6(1), 2019. ISSN 2197-7909. doi : 10.1186/s40668-019-0028-x. URL www.tng-project.org/data.
- J. W. Nightingale, S. Dye, and R. J. Massey. AutoLens : automated modeling of a strong lens's light, mass, and source. *MNRAS*, 478(4) :4738–4784, Aug. 2018. doi : 10.1093/mnras/sty1264.
- S. J. Pan and Q. Yang. A survey on transfer learning. *IEEE Transactions on Knowledge and Data Engineering*, 22 : 1345–1359, 2010.
- T. pandas development team. pandas-dev/pandas : Pandas, Feb. 2020. URL <https://doi.org/10.5281/zenodo.3509134>.
- J. W. Park, S. Wagner-Carena, S. Birrer, P. J. Marshall, J. Y.-Y. Lin, A. Roodman, and LSST Dark Energy Science Collaboration. Large-scale Gravitational Lens Modeling with Bayesian Neural Networks for Accurate and Precise Inference of the Hubble Constant. *ApJ*, 910(1) :39, Mar. 2021. doi : 10.3847/1538-4357/abdfc4.
- F. Pedregosa, G. Varoquaux, A. Gramfort, V. Michel, B. Thirion, O. Grisel, M. Blondel, P. Prettenhofer, R. Weiss, V. Dubourg, J. Vanderplas, A. Passos, D. Cournapeau, M. Brucher, M. Perrot, and E. Duchesnay. Scikit-learn : Machine learning in Python. *Journal of Machine Learning Research*, 12 :2825–2830, 2011.
- F. Pérez and B. E. Granger. IPython : a system for interactive scientific computing. *Computing in Science and Engineering*, 9(3) :21–29, May 2007. ISSN 1521-9615. doi : 10.1109/MCSE.2007.53. URL <https://ipython.org>.
- L. Perreault Levasseur, Y. D. Hezaveh, and R. H. Wechsler. Uncertainties in Parameters Estimated with Neural Networks : Application to Strong Gravitational Lensing. *ApJ*, 850(1) :L7, Nov. 2017. doi : 10.3847/2041-8213/aa9704.
- Planck Collaboration. Planck 2018 results. VI. Cosmological parameters. *A&A*, 641 :A6, Sept. 2020. doi : 10.1051/0004-6361/201833910.
- P. Putzky and M. Welling. Recurrent Inference Machines for Solving Inverse Problems. *arXiv e-prints*, 2017.
- C. Rao. Information and the accuracy attainable in the estimation of statistical parameters. *Bulletin fo the Calcutta Mathematical Society*, 1945.
- R. Ratcliff. Connectionist models of recognition memory : Constraints imposed by learning and forgetting functions. *Psychological Review*, 97(2) :285–308, 1990. doi : 10.1037/0033-295X.97.2.285.
- S. Rau, S. Vegetti, and S. D. White. The effect of particle noise in N-body simulations of gravitational lensing. *Monthly Notices of the Royal Astronomical Society*, 430(3) :2232–2248, apr 2013. ISSN 13652966. doi : 10.1093/mnras/stt043.

- B. Remy, F. Lanusse, N. Jeffrey, J. Liu, J.-L. Starck, K. Osato, and T. Schrabback. Probabilistic Mass Mapping with Neural Score Estimation. *arXiv e-prints*, art. arXiv :2201.05561, Jan. 2022.
- F. Rizzo, S. Vegetti, D. Powell, F. Fraternali, J. P. McKean, H. R. Stacey, and S. D. M. White. A dynamically cold disk galaxy in the early Universe. *Nature*, 584(7820) :201–204, Aug. 2020. doi : 10.1038/s41586-020-2572-6.
- O. Ronneberger, P. Fischer, and T. Brox. U-Net : Convolutional Networks for Biomedical Image Segmentation. *arXiv e-prints*, art. arXiv :1505.04597, May 2015.
- B. T. Rowe, M. Jarvis, R. Mandelbaum, G. M. Bernstein, J. Bosch, M. Simet, J. E. Meyers, T. Kacprzak, R. Nakajima, J. Zuntz, H. Miyatake, J. P. Dietrich, R. Armstrong, P. Melchior, and M. S. Gill. GalSim : The modular galaxy image simulation toolkit. *Astronomy and Computing*, 10 :121–150, apr 2015. ISSN 22131337. doi : 10.1016/j.ascom.2015.02.002.
- B. T. P. Rowe, M. Jarvis, R. Mandelbaum, G. M. Bernstein, J. Bosch, M. Simet, J. E. Meyers, T. Kacprzak, R. Nakajima, J. Zuntz, H. Miyatake, J. P. Dietrich, R. Armstrong, P. Melchior, and M. S. S. Gill. GALSIM : The modular galaxy image simulation toolkit. *Astronomy and Computing*, 10 :121–150, Apr. 2015. doi : 10.1016/j.ascom.2015.02.002.
- C. E. Rusu, C. D. Fassnacht, D. Sluse, S. Hilbert, K. C. Wong, K.-H. Huang, S. H. Suyu, T. E. Collett, P. J. Marshall, T. Treu, and L. V. E. Koopmans. H0LiCOW - III. Quantifying the effect of mass along the line of sight to the gravitational lens HE 0435-1223 through weighted galaxy counts. *MNRAS*, 467(4) :4220–4242, June 2017. doi : 10.1093/mnras/stx285.
- C. E. Rusu, K. C. Wong, V. Bonvin, D. Sluse, S. H. Suyu, C. D. Fassnacht, J. H. H. Chan, S. Hilbert, M. W. Auger, A. Sonnenfeld, S. Birrer, F. Courbin, T. Treu, G. C. F. Chen, A. Halkola, L. V. E. Koopmans, P. J. Marshall, and A. J. Shajib. H0LiCOW XII. Lens mass model of WFI2033-4723 and blind measurement of its time-delay distance and H_0 . *MNRAS*, 498(1) :1440–1468, Oct. 2020. doi : 10.1093/mnras/stz3451.
- P. Saha and L. L. R. Williams. Non-parametric reconstruction of the galaxy lens in PG 1115+080. *MNRAS*, 292(1) :148–156, Nov. 1997. doi : 10.1093/mnras/292.1.148.
- P. Saha and L. L. R. Williams. A Portable Modeler of Lensed Quasars. *AJ*, 127(5) :2604–2616, May 2004. doi : 10.1086/383544.
- S. Schuldt, G. Chirivi, S. H. Suyu, A. Yıldırım, A. Sonnenfeld, A. Halkola, and G. F. Lewis. Inner dark matter distribution of the Cosmic Horseshoe (J1148+1930) with gravitational lensing and dynamics. *A&A*, 631 :A40, Nov. 2019. doi : 10.1051/0004-6361/201935042.
- S. Schuldt, R. Cañameras, Y. Shu, S. H. Suyu, S. Taubenberger, T. Meinhardt, and L. Leal-Taixé. HOLISMOKEs – IX. Neural network inference of strong-lens parameters and uncertainties from ground-based images. *arXiv e-prints*, art. arXiv :2206.11279, June 2022.
- N. Scoville, H. Aussel, M. Brusa, P. Capak, C. M. Carollo, M. Elvis, M. Giavalisco, L. Guzzo, G. Hasinger, C. Impey, J.-P. Kneib, O. LeFevre, S. J. Lilly, B. Mobasher, A. Renzini, R. M. Rich, D. B. Sanders, E. Schinnerer, D. Schminovich, P. Shopbell, Y. Taniguchi, and N. D. Tyson. The Cosmic Evolution Survey (COSMOS) : Overview. *The Astrophysical Journal Supplement Series*, 172(1) :1–8, sep 2007. ISSN 0067-0049. doi : 10.1086/516585.
- S. Seitz, P. Schneider, and M. Bartelmann. Entropy-regularized maximum-likelihood cluster mass reconstruction. *A&A*, 337 :325–337, Sept. 1998.
- J. L. Sérsic. Influence of the atmospheric and instrumental dispersion on the brightness distribution in a galaxy. *Boletin de la Asociacion Argentina de Astronomia La Plata Argentina*, 6 :41–43, Feb. 1963.

- Y. Shu, J. R. Brownstein, A. S. Bolton, L. V. E. Koopmans, T. Treu, A. D. Montero-Dorta, M. W. Auger, O. Czoske, R. Gavazzi, P. J. Marshall, and L. A. Moustakas. The Sloan Lens ACS Survey. XIII. Discovery of 40 New Galaxy-scale Strong Lenses. *ApJ*, 851(1) :48, Dec. 2017. doi : 10.3847/1538-4357/aa9794.
- P. Simon, C. Heymans, T. Schrabback, A. N. Taylor, M. E. Gray, L. van Waerbeke, C. Wolf, D. Bacon, M. Barden, A. Böhm, B. Häuftler, K. Jahnke, S. Jogee, E. van Kampen, K. Meisenheimer, and C. Y. Peng. Spatial matter density mapping of the STAGES Abell A901/2 supercluster field with 3D lensing. *MNRAS*, 419(2) :998–1016, Jan. 2012. doi : 10.1111/j.1365-2966.2011.19760.x.
- D. Sluse, A. Sonnenfeld, N. Rumbaugh, C. E. Rusu, C. D. Fassnacht, T. Treu, S. H. Suyu, K. C. Wong, M. W. Auger, V. Bonvin, T. Collett, F. Courbin, S. Hilbert, L. V. E. Koopmans, P. J. Marshall, G. Meylan, C. Spinello, and M. Tewes. H0LiCOW - II. Spectroscopic survey and galaxy-group identification of the strong gravitational lens system HE 0435-1223. *MNRAS*, 470(4) :4838–4857, Oct. 2017. doi : 10.1093/mnras/stx1484.
- J. L. Starck, K. E. Themelis, N. Jeffrey, A. Peel, and F. Lanusse. Weak-lensing mass reconstruction using sparsity and a Gaussian random field. *A&A*, 649 :A99, May 2021. doi : 10.1051/0004-6361/202039451.
- F. Sun, E. Egami, P. G. Pérez-González, I. Smail, K. I. Caputi, F. E. Bauer, T. D. Rawle, S. Fujimoto, K. Kohno, U. Dudzevičiūtė, H. Atek, M. Bianconi, S. C. Chapman, F. Combes, M. Jauzac, J.-B. Jolly, A. M. Koekemoer, G. E. Magdis, G. Rodighiero, W. Rujopakarn, D. Schaerer, C. L. Steinhardt, P. Van der Werf, G. L. Walth, and J. R. Weaver. Extensive Lensing Survey of Optical and Near-infrared Dark Objects (El Sonido) : HST H-faint Galaxies behind 101 Lensing Clusters. *ApJ*, 922(2) :114, Dec. 2021. doi : 10.3847/1538-4357/ac2578. URL <https://ui.adsabs.harvard.edu/abs/2021ApJ...922..114S>.
- S. H. Suyu and R. D. Blandford. The anatomy of a quadruply imaged gravitational lens system. *MNRAS*, 366(1) :39–48, Feb. 2006. doi : 10.1111/j.1365-2966.2005.09854.x.
- S. H. Suyu, P. J. Marshall, M. P. Hobson, and R. D. Blandford. A Bayesian analysis of regularized source inversions in gravitational lensing. *MNRAS*, 371(2) :983–998, Sept. 2006. doi : 10.1111/j.1365-2966.2006.10733.x.
- N. Tishby, F. C. Pereira, and W. Bialek. The information bottleneck method. In *Proc. of the 37-th Annual Allerton Conference on Communication, Control and Computing*, pages 368–377, 1999. URL <https://arxiv.org/abs/physics/0004057>.
- D. A. Torres-Ballesteros and L. Castañeda. relensing : Reconstructing the mass profile of galaxy clusters from gravitational lensing. *arXiv e-prints*, art. arXiv :2201.10076, Jan. 2022.
- T. Treu and L. V. E. Koopmans. Massive Dark Matter Halos and Evolution of Early-Type Galaxies to $z \sim 1$. *ApJ*, 611(2) :739–760, Aug. 2004. doi : 10.1086/422245.
- S. Van der Walt, J. L. Schönberger, J. Nunez-Iglesias, F. Boulogne, J. D. Warner, N. Yager, E. Gouillart, and T. Yu. scikit-image : image processing in python. *PeerJ*, 2 :e453, 2014.
- S. Vegetti and L. V. E. Koopmans. Bayesian strong gravitational-lens modelling on adaptive grids : objective detection of mass substructure in Galaxies. *MNRAS*, 392(3) :945–963, Jan. 2009. doi : 10.1111/j.1365-2966.2008.14005.x.
- S. Vegetti, D. J. Lagattuta, J. P. McKean, M. W. Auger, C. D. Fassnacht, and L. V. E. Koopmans. Gravitational detection of a low-mass dark satellite galaxy at cosmological distance. *Nature*, 481(7381) :341–343, Jan. 2012. doi : 10.1038/nature10669.
- J. D. Vieira, D. P. Marrone, S. C. Chapman, C. De Breuck, Y. D. Hezaveh, A. Weiβ, J. E. Aguirre, K. A. Aird, M. Aravena, M. L. N. Ashby, M. Bayliss, B. A. Benson, A. D. Biggs, L. E. Bleem, J. J. Bock, M. Bothwell, C. M. Bradford, M. Brodwin, J. E. Carlstrom, C. L. Chang, T. M. Crawford, A. T. Crites, T. de Haan, M. A. Dobbs,

- E. B. Fomalont, C. D. Fassnacht, E. M. George, M. D. Gladders, A. H. Gonzalez, T. R. Greve, B. Gullberg, N. W. Halverson, F. W. High, G. P. Holder, W. L. Holzapfel, S. Hoover, J. D. Hrubes, T. R. Hunter, R. Keisler, A. T. Lee, E. M. Leitch, M. Lueker, D. Luong-van, M. Malkan, V. McIntyre, J. J. McMahon, J. Mehl, K. M. Menten, S. S. Meyer, L. M. Mocanu, E. J. Murphy, T. Natoli, S. Padin, T. Plagge, C. L. Reichardt, A. Rest, J. Ruel, J. E. Ruhl, K. Sharon, K. K. Schaffer, L. Shaw, E. Shirokoff, J. S. Spilker, B. Stalder, Z. Staniszewski, A. A. Stark, K. Story, K. Vanderlinde, N. Welikala, and R. Williamson. Dusty starburst galaxies in the early Universe as revealed by gravitational lensing. *Nature*, 495(7441) :344–347, Mar. 2013. doi : 10.1038/nature12001.
- P. Virtanen, R. Gommers, T. E. Oliphant, M. Haberland, T. Reddy, D. Cournapeau, E. Burovski, P. Peterson, W. Weckesser, J. Bright, S. J. van der Walt, M. Brett, J. Wilson, K. J. Millman, N. Mayorov, A. R. J. Nelson, E. Jones, R. Kern, E. Larson, C. J. Carey, İ. Polat, Y. Feng, E. W. Moore, J. VanderPlas, D. Laxalde, J. Perktold, R. Cimrman, I. Henriksen, E. A. Quintero, C. R. Harris, A. M. Archibald, A. H. Ribeiro, F. Pedregosa, P. van Mulbregt, and SciPy 1.0 Contributors. SciPy 1.0 : Fundamental Algorithms for Scientific Computing in Python. *Nature Methods*, 17 :261–272, 2020. doi : 10.1038/s41592-019-0686-2.
- S. Wagner-Carena, J. W. Park, S. Birrer, P. J. Marshall, A. Roodman, R. H. Wechsler, and LSST Dark Energy Science Collaboration. Hierarchical Inference with Bayesian Neural Networks : An Application to Strong Gravitational Lensing. *ApJ*, 909(2) :187, Mar. 2021. doi : 10.3847/1538-4357/abdf59.
- S. Wagner-Carena, J. Aalbers, S. Birrer, E. O. Nadler, E. Darragh-Ford, P. J. Marshall, and R. H. Wechsler. From Images to Dark Matter : End-To-End Inference of Substructure From Hundreds of Strong Gravitational Lenses. *arXiv e-prints*, art. arXiv :2203.00690, Mar. 2022.
- S. J. Warren and S. Dye. Semilinear Gravitational Lens Inversion. *ApJ*, 590(2) :673–682, June 2003. doi : 10.1086/375132.
- Wes McKinney. Data Structures for Statistical Computing in Python. In Stéfan van der Walt and Jarrod Millman, editors, *Proceedings of the 9th Python in Science Conference*, pages 56 – 61, 2010. doi : 10.25080/Majora-92bf1922-00a.
- K. C. Wong, S. H. Suyu, M. W. Auger, V. Bonvin, F. Courbin, C. D. Fassnacht, A. Halkola, C. E. Rusu, D. Sluse, A. Sonnenfeld, T. Treu, T. E. Collett, S. Hilbert, L. V. E. Koopmans, P. J. Marshall, and N. Rumbaugh. H0LiCOW - IV. Lens mass model of HE 0435-1223 and blind measurement of its time-delay distance for cosmology. *MNRAS*, 465(4) :4895–4913, Mar. 2017. doi : 10.1093/mnras/stw3077.
- S. Zhao, J. Song, and S. Ermon. InfoVAE : Information Maximizing Variational Autoencoders. *arXiv e-prints*, art. arXiv :1706.02262, June 2017.
- F. Zhuang, Z. Qi, K. Duan, D. Xi, Y. Zhu, H. Zhu, H. Xiong, and Q. He. A Comprehensive Survey on Transfer Learning. *arXiv e-prints*, art. arXiv :1911.02685, Nov. 2019.
- F. Zwicky. Nebulae as gravitational lenses. *Phys. Rev.*, 51 :290–290, Feb 1937. doi : 10.1103/PhysRev.51.290. URL <https://link.aps.org/doi/10.1103/PhysRev.51.290>.

Annexe A

Elastic Weight Consolidation

Suppose we are given a training set \mathcal{D} and a test task \mathcal{T} . The posterior of the RIM parameters φ can be rewritten using the Bayes rule as

$$p(\varphi \mid \mathcal{D}, \mathcal{T}) = \frac{p(\mathcal{T} \mid \mathcal{D}, \varphi)p(\varphi \mid \mathcal{D})}{p(\mathcal{T} \mid \mathcal{D})}. \quad (\text{A.1})$$

We suppose that φ encode information about \mathcal{D} , while \mathcal{T} was unseen by φ . It follows that \mathcal{T} and \mathcal{D} are conditionally independent when given φ . We do not make the stronger assumption that \mathcal{D} and \mathcal{T} are completely independent. In fact, such an assumption would contradict the premiss of our work that building a dataset \mathcal{D} can inform a machine (RIM) about task \mathcal{T} — or that, more broadly, \mathcal{D} contains information about \mathcal{T} .

We rewrite the marginal $p(\mathcal{T} \mid \mathcal{D})$ using the Bayes rule in order to extract $p(\mathcal{D} \mid \mathcal{T})$, the sampling distribution used to compute the Fisher diagonal elements

$$p(\varphi \mid \mathcal{D}, \mathcal{T}) = \frac{p(\mathcal{T} \mid \varphi)p(\varphi \mid \mathcal{D})}{p(\mathcal{D} \mid \mathcal{T})} \frac{p(\mathcal{D})}{p(\mathcal{T})}. \quad (\text{A.2})$$

The log-likelihood $\log p(\mathcal{T} \mid \varphi)$ is equivalent to the negative of the loss function for the particular task at hand. In this work, we assign a uniform probability density to $p(\mathcal{T})$ and $p(\mathcal{D})$ in order to ignore them.

We now turn to the prior $p(\varphi \mid \mathcal{D})$, which appears as a conditional relative to the training dataset. We use the Laplace approximation around the maxima $\varphi_{\mathcal{D}}^*$ to evaluate the prior, where $\varphi_{\mathcal{D}}^*$ are the trained parameters of the RIM that minimize the empirical risk (equation (2.6)). The Taylor expansion of the prior around this maxima yields

$$\log p(\varphi \mid \mathcal{D}) \approx \log p(\varphi_{\mathcal{D}}^* \mid \mathcal{D}) + \underbrace{\frac{1}{2}(\varphi - \varphi_{\mathcal{D}}^*)^T \left(\frac{\partial^2 \log p(\varphi \mid \mathcal{D})}{\partial^2 \varphi} \Big|_{\varphi_{\mathcal{D}}^*} \right)}_{\mathbf{H}(\varphi_{\mathcal{D}}^*)} (\varphi - \varphi_{\mathcal{D}}^*). \quad (\text{A.3})$$

Since $\varphi_{\mathcal{D}}^*$ is an extrema of the prior, the linear term vanishes. The empirical estimate of the negative hessian matrix is the observed Fisher information matrix which can be written as

$$\mathcal{I}(\varphi_{\mathcal{D}}^*) = -\mathbb{E}_{\mathcal{D}|\mathcal{T}}[\mathbf{H}(\varphi_{\mathcal{D}}^*)] = \mathbb{E}_{\mathcal{D}|\mathcal{T}} \left[\left(\left(\frac{\partial \log p(\varphi | \mathcal{D})}{\partial \varphi} \right) \left(\frac{\partial \log p(\varphi | \mathcal{D})}{\partial \varphi} \right)^T \right) \Big|_{\varphi_{\mathcal{D}}^*} \right]. \quad (\text{A.4})$$

The expectation is taken over the sample space $p(\mathcal{D} | \mathcal{T})$ since the network parameters are held fixed during sampling. In order to compute the Fisher score, we apply the Bayes rule to the prior to extract a loss function, which we take to be proportional to the training loss (equation (2.5)) and the χ^2 :

$$\log p(\varphi | (\mathbf{x}, \mathbf{y}) = \mathcal{D}) \propto -\mathcal{L}_\varphi(\mathbf{x}, \mathbf{y}) + \frac{1}{T} \sum_{t=1}^T \log p(\mathbf{y} | \hat{\mathbf{x}}^{(t)}) - \frac{\ell_2}{2} \|\varphi\|_2^2 \quad (\text{A.5})$$

We find in practice the the ℓ_2 term has little effect on the Fisher diagonal and our results. Thus, we set $\ell_2 = 0$.

Since the full Fisher matrix is intractable for a neural network, we approximate the quadratic term of the prior with the diagonal of the Fisher matrix following [Kirkpatrick et al. \(2016\)](#). For an optimisation problem, the first term of (A.3) is constant. Thus, the posterior becomes proportional to

$$\log p(\varphi | \mathcal{D}, \mathcal{T}) \propto \log p(\mathcal{T} | \varphi) - \frac{\lambda}{2} \sum_j \text{diag}(\mathcal{I}(\varphi_{\mathcal{D}}^*))_j (\varphi_j - [\varphi_{\mathcal{D}}^*]_j)^2. \quad (\text{A.6})$$

The Lagrange multiplier λ is introduced to tune our uncertainty about the network parameters during fine-tuning.

Annexe B

VAE Architecture and optimisation

For the following architectures, we employ the notion of *level* to mean layers in the encoder and the decoder with the same resolution. In each level, we place a block of convolutional layers before downsampling (encoder) or after upsampling (decoder). These operations are done with strided convolutions like in the U-net architecture of the RIM.

TABLE B.1 – Hyperparameters for the background source VAE.

Parameter	Value
Input preprocessing	1
<i>Architecture</i>	
Levels (encoder and decoder)	3
Convolutional layer per level	2
Latent space dimension	32
Hidden Activations	Leaky ReLU
Output Activation	Sigmoid
Filters (first level)	16
Filters scaling factor (per level)	2
Number of parameters	3 567 361
<i>Optimization</i>	
Optimizer	Adam
Initial learning rate	10^{-4}
Learning rate schedule	Exponential Decay
Decay rate	0.5
Decay steps	30 000
Number of steps	500 000
β_{\max}	0.1
Batch size	20

TABLE B.2 – Hyperparameters for the convergence VAE.

Parameter	Value
Input preprocessing	\log_{10}
<i>Architecture</i>	
Levels (encoder and decoder)	4
Convolutional layer per level	1
Latent space dimension	16
Hidden Activations	Leaky ReLU
Output Activation	$\mathbb{1}$
Filters (first level)	16
Filters scaling factor (per level)	2
Number of parameters	1 980 033
<i>Optimization</i>	
Optimizer	Adam
Initial learning rate	10^{-4}
Learning rate schedule	Exponential Decay
Decay rate	0.7
Decay steps	20 000
Number of steps	155 000
β_{\max}	0.2
Batch size	32

Annexe C

RIM architecture and optimisation

The notion of link function $\Psi : \Xi \rightarrow \mathcal{X}$, introduced by Putzky and Welling (2017), is an invertible transformation between the network prediction space $\xi \in \Xi$ and the forward modelling space $\mathbf{x} \in \mathcal{X}$. This is a different notion from preprocessing, discussed in section 2.3, because this transformation is applied inside the recurrent relation 2.4 as opposed to before training. In the case where the forward model has some restricted support or it is found that some transformation helps the training, then the link function chosen must be implemented as part of the network architecture as shown in the unrolled computational graph in Figure C.1. Also, the loss \mathcal{L}_φ must be computed in the Ξ space in order to avoid gradient vanishing problems when Ψ is a non-linear mapping, which happens if the non-linear link function is applied in an operation recorded for backpropagation through time (BPTT).

For the convergence, we use an exponential link function with base 10 : $\hat{\kappa} = \Psi(\xi) = 10^\xi$. This Ψ encodes the non-negativity of the convergence. Furthermore, it is a power transformation that leaves the linked pixel values ξ_i normally distributed, thus improving the learning through the non-linearities in the neural network. The pixel weights \mathbf{w}_i in the loss function (2.5) are chosen to encode the fact that the pixel with critical mass density ($\kappa_i > 1$) have a stronger effect on the lensing configuration than other pixels. We find in practice that the weights

$$\mathbf{w}_i = \frac{\sqrt{\kappa_i}}{\sum_i \kappa_i}, \quad (\text{C.1})$$

encode this knowledge in the loss function and improved both the empirical risk and the goodness of fit of the baseline model on early test runs.

For the source, we found that we do not need a link function — its performance is generally better compared to other link function we tried like sigmoid and power transforms — and we found that the pixel weights can be taken to be uniform, i.e. $\mathbf{w}_i = \frac{1}{M}$.

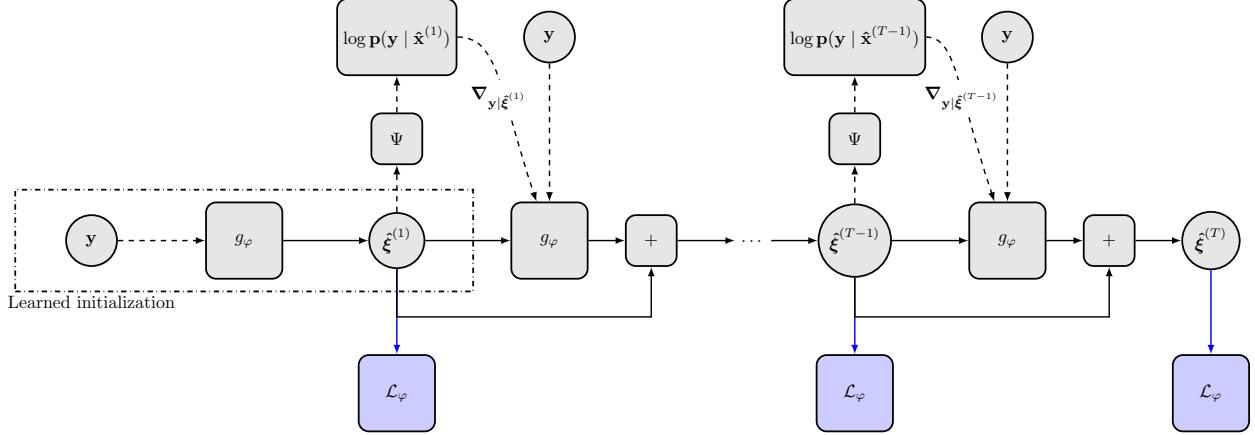


FIGURE C.1 – Unrolled computational graph of the RIM. Operations along solid arrows are being recorded for BPTT, while operations along dashed arrows are not. The blue arrows are only used for optimisation during training. During fine-tuning or testing, the loss is computed only as an oracle metric to validate that our methods can recover the ground truth.

TABLE C.1 – Hyperparameters for the RIM.

Parameter	Value
Source link function	1
κ link function	$10^{\frac{1}{6}}$
<i>Architecture</i>	
Recurrent steps (T)	8
Number of parameters	348 546 818
<i>First Stage Optimisation</i>	
Optimizer	Adamax
Initial learning rate	10^{-4}
Learning rate schedule	Exponential Decay
Decay rate	0.95
Decay steps	100 000
Number of steps	610 000
Batch size	1
<i>Second Stage Optimisation</i>	
Optimizer	Adamax
Initial learning rate	6×10^{-5}
Learning rate schedule	Exponential Decay
Decay rate	0.9
Decay steps	100 000
Number of steps	870 000
Batch size	1

COSMOS RIM+FT IllustrisTNG RIM+FT Observation RIM+FT COSMOS RIM+FT IllustrisTNG RIM+FT Observation RIM+FT

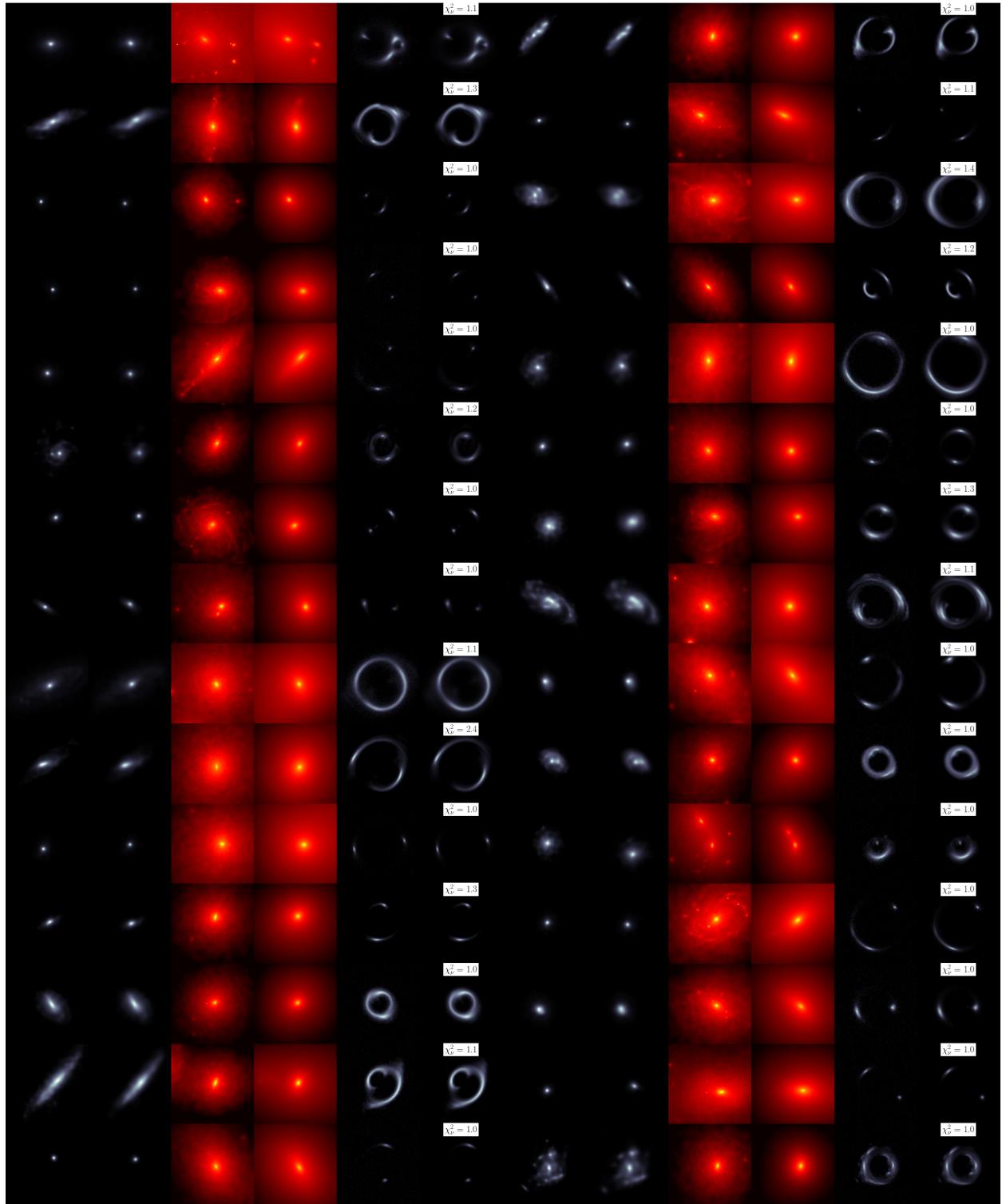


FIGURE C.2 – 30 reconstructions taken at random from the test set of 3000 examples simulated from COSMOS and IllustrisTNG data at high SNR. The colorscale are the same as in Figure 2.7.