from IPython.display import Image import matplotlib.pyplot as plt import numpy as np

## Background

In []: # Import packages

The problem that I chose to work on was a variation of the classic sphereical shell problem. This problem set. The problem set. The problem is of a sphere with a thick spherical shell that has a charge density that varies with radius. This shell is from radius a to b and has a charge density of  $ho=k/r^2$ . The important consideration for this problem is that the electric field is not a single continous function. This forces us to solve inside the sphere, inside the shell, and outside of the sphere. The variables a,b, and k are all constants that we can choose to be whatever we want. For this problem I chose a = 2, b = 4, and k = 1.

## Solution

Gauss's law seemly highly applicable here from the spherical symmetry of the problem. This allows me to define Gaussian surface to simplify the solution.

First, we calculate the E-field inside of the sphere (r < a). We can do this by using Gauss's law and a Gaussian surface of a sphere with radius r. This gives us the following equation:

$$\oint ec{E} \cdot dec{A} = rac{Q_{enc}}{\epsilon_0}$$

Knowing we have a spherical surface of radius r and a  $Q_{enc}$  of 0 (since there is no charge inside the sphere) we can simplify the equation to the following:

$$E(4\pi r^2)=0$$

Next, we calculate the E-field inside of the shell (a < r < b). We can do this again by using Gauss's law and a Gaussian surface of a sphere with radius r. This gives us the following equation:

 $E(4\pi r^2)=rac{Q_{enc}}{\epsilon_0}$ 

Here, 
$$Q_{enc}$$
 is not as trivial to calculate. We know that the charge density is  $\rho=k/r^2$  and that the integration factor is  $d\tau=r^2sin(\theta)drd\theta d\phi$ . This allows us to calculate the charge of the shell as follows: 
$$Q_{enc}=\int \rho d\tau=\int_a^r \frac{k}{r^2}r^2sin(\theta)drd\theta d\phi$$

By spherical symmetry we can integrate out the  $\theta$  and  $\phi$  terms to get the following:

$$Q_{enc}=4\pi k\int_{a}^{r}dr=4\pi k(r-a)$$

This means that the E-field inside the shell is:

This gives us an E-field of 0 inside the sphere.

$$ec{E} = rac{k}{\epsilon_0} rac{(r-a)}{r^2} \hat{r}$$

Finally, we calculate the E-field outside of the sphere (r > b). This is done in the same way as the previous calculation except the upper bound on the integral is b. This gives us the following:

$$ec{E} = rac{k}{\epsilon_0}rac{(b-a)}{r^2}\hat{r}$$

## Code

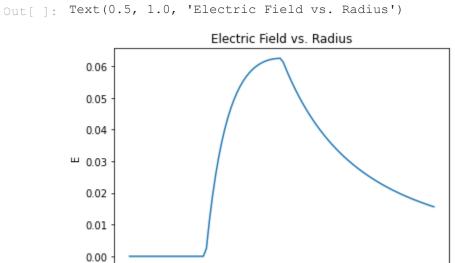
```
# Define our electric field functions for each region
E1 = lambda r: 0
E2 = lambda r,k,epp,a: k/epp * ((r - a)/r**2)
def Efield(r):
    k = 0.5
    epp = 1
    a = 2
   b = 4
    if r < a:
        return E1(r)
    elif r < b:</pre>
        return E2(r,k,epp,a)
    else:
        return E3(r,k,epp,a,b)
```

To get a feeling for what the electric field looks like from the equations that we solved for above, I decided to plot the E-field as a function of radius below. The result makes sense as the E-field is 0 inside the sphere, increases as we get a proportion of  $\frac{1}{r^2}$  more charge from the charge density inside the shell, and decreases  $\frac{1}{m^2}$  as we are outside. This  $\frac{1}{m^2}$  is especially expected as we think about the behavior of the E-field from a point charge as we move away from the sphere.

plt.title('Electric Field vs. Radius')

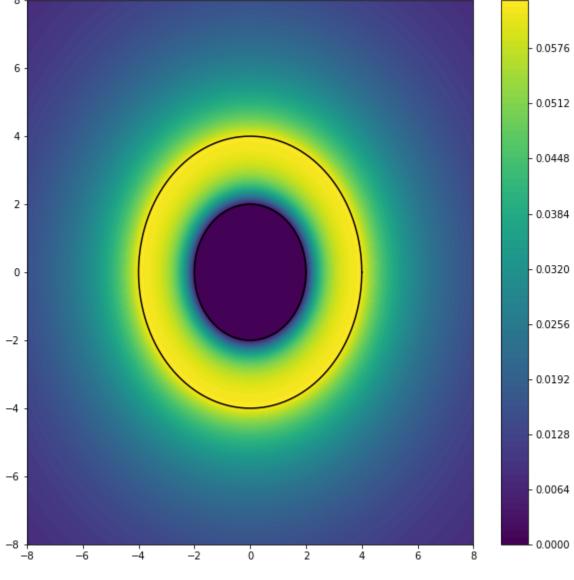
r = np.linspace(0,8,100)E = [Efield(r) for r in r]

plt.plot(r,E) plt.xlabel('r') plt.ylabel('E')



This graph represents a cross section of the E-field as a function of radius to better visualize the E-field. The E-field is in the  $\hat{r}$  direction.

X,Y = np.meshgrid(np.linspace(-8,8,500),np.linspace(-8,8,500))Z = np.zeros(X.shape) for i in range(len(X)): for j in range(len(Y)): Z[i,j] = Efield(np.sqrt(X[i,j]\*\*2 + Y[i,j]\*\*2))plt.figure(figsize=(10,10)) plt.plot(2\*np.cos(np.linspace(0,2\*np.pi,1000)),2\*np.sin(np.linspace(0,2\*np.pi,1000)), color='black') plt.plot(4\*np.cos(np.linspace(0,2\*np.pi,1000)),4\*np.sin(np.linspace(0,2\*np.pi,1000)), color='black') plt.contourf(X,Y,Z,100) plt.colorbar() plt.show()



The plotting code below shows again a plot of the e-field but with the use of quivers on a 3D plot. Additionally the shells are plotted to better visualize the problem.

In []: # adaptation from ICA def electric\_field3D(x\_points, y\_points, z\_points):  $k = 8.99e9 \# Nm^2/C^2$ , Coulomb's constant # Initialize electric field components to zero E x = np.zeros like(x points)E\_y = np.zeros\_like(y\_points) E\_z = np.zeros\_like(z\_points) # Calculate electric field components due to the point charge at each point on the grid for i in range(x\_points.shape[0]): for j in range(y\_points.shape[0]): for k in range(z\_points.shape[0]):  $r_x = x_points[i, j, k]$  $r_y = y_points[i, j, k]$  $r_z = z_points[i, j, k]$  $r_{magnitude} = np.sqrt(r_x**2 + r_y**2 + r_z**2)$ if r\_magnitude != 0: # Avoid division by zero  $r_unit_x = r_x / r_magnitude$ r\_unit\_y = r\_y / r\_magnitude  $r_unit_z = r_z / r_magnitude$  $E_x[i, j, k] = Efield(r_magnitude) * r_unit_x$  $E_y[i, j, k] = Efield(r_magnitude) * r_unit_y$ E\_z[i, j, k] = Efield(r\_magnitude) \* r\_unit\_z return E\_x, E\_y, E\_z

x = np.linspace(-4, 4, 8)y = np.linspace(-4, 4, 8)z = np.linspace(-4, 4, 8)X, Y, Z = np.meshgrid(x, y, z) $E_x$ ,  $E_y$ ,  $E_z$  = electric\_field3D(X, Y, Z) # Create a 3D quiver plot fig = plt.figure(figsize=(10, 10)) ax = fig.add\_subplot(111, projection='3d') ax.quiver(X, Y, Z, E\_x, E\_y, E\_z, length=10, color='b') ax.set xlabel('X-axis') ax.set\_ylabel('Y-axis') ax.set\_zlabel('Z-axis') ax.set\_title('Electric Field of a Charged Sphere') # Plot the shells for r in [2,4]: # Create the mesh in spherical coordinates phi = np.linspace(0, np.pi, 1000) theta = np.linspace(0, 2\*np.pi, 1000)phi, theta = np.meshgrid(phi, theta) # The Cartesian coordinates of the unit sphere x = r\*np.sin(phi)\*np.cos(theta)y = r\*np.sin(phi)\*np.sin(theta)z = r\*np.cos(phi)# Plot the surface ax.plot\_surface(x, y, z, color='b', alpha=0.1) plt.show()

## Electric Field of a Charged Sphere

