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# The Singularity Theorems in General Relativity

Damon Binder

ASC Presentation

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- ▶ General Relativity shows that gravity occurs because of the geometry of spacetime

$$R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} = 8\pi T_{\mu\nu}$$

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- ▶ General Relativity shows that gravity occurs because of the geometry of spacetime

$$R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} = 8\pi T_{\mu\nu}$$

- ▶ Predicts singularities - points where matter density becomes infinite and spacetime ceases to exist

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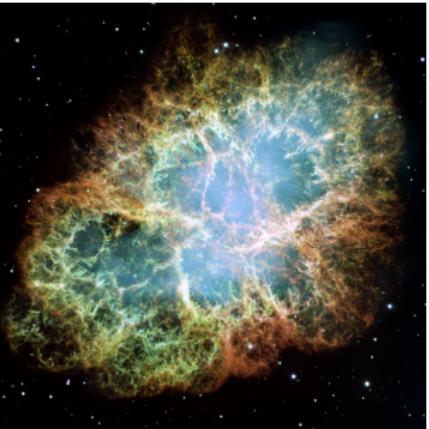
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$$R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} = 8\pi T_{\mu\nu}$$

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  - ▶ Stellar Collapse



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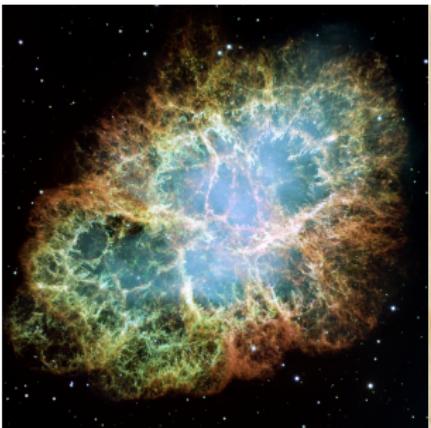
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- ▶ General Relativity shows that gravity occurs because of the geometry of spacetime

$$R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} = 8\pi T_{\mu\nu}$$

- ▶ Predicts singularities - points where matter density becomes infinite and spacetime ceases to exist
  - ▶ Stellar Collapse
  - ▶ Beginning of the Universe



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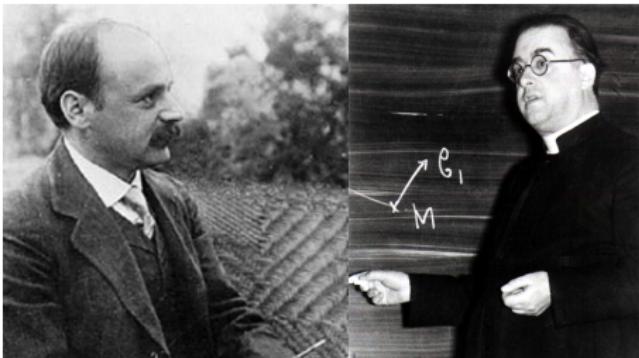
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## Stellar Collapse

- ▶ Schwarzschild in 1916 derives the spherically symmetric solution, which has a singularity

$$ds^2 = - \left(1 - \frac{2M}{r}\right) dt^2 + \left(1 - \frac{2M}{r}\right)^{-1} dr^2 + r^2 d\Omega^2$$

- ▶ Lemaître (1932) and Synge (1950) realise that the  $r = 2M$  surface is not a singularity
- ▶ At  $r = 0$  we have a true, curvature singularity
- ▶ For an observer with  $r < 2M$ , the singularity will be reached in finite time



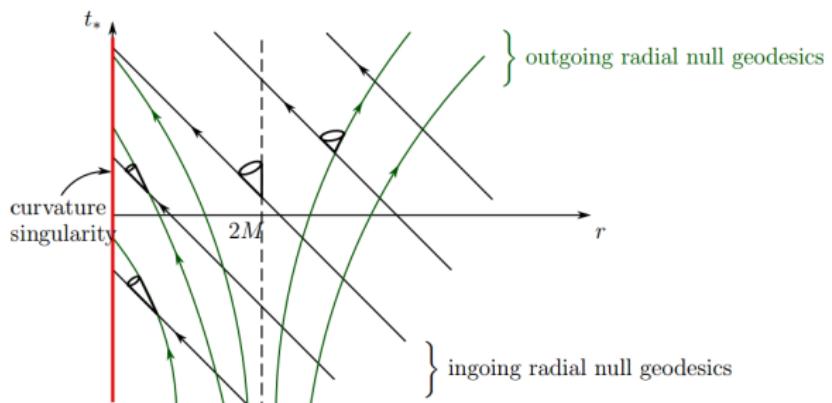
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## Stellar Collapse

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- ▶ For an observer with  $r < 2M$ , the singularity will be reached in finite time



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## Stellar Collapse

- ▶ Chandrasekhar (1931), and Oppenheimer and Volkoff (1939) studied stellar structure, white dwarfs and neutron stars



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## Stellar Collapse

- ▶ Chandrasekhar (1931), and Oppenheimer and Volkoff (1939) studied stellar structure, white dwarfs and neutron stars
- ▶ Found that stellar collapse was inevitable for sufficiently heavy objects



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# Introduction

## Stellar Collapse

- ▶ Chandrasekhar (1931), and Oppenheimer and Volkoff (1939) studied stellar structure, white dwarfs and neutron stars
- ▶ Found that stellar collapse was inevitable for sufficiently heavy objects
- ▶ Oppenheimer and Snyder (1939) found that a ball of dust will collapse into a Schwarzschild black hole



# Introduction

## The Beginning of the Universe

- ▶ Friedmann (1924) and Lemaitre (1927) study isotropic and homogenous universes

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## The Beginning of the Universe

- ▶ Friedmann (1924) and Lemaitre (1927) study isotropic and homogenous universes
- ▶ Lemaitre (1932) demonstrates general existence of singularities in Bianchi Type I class

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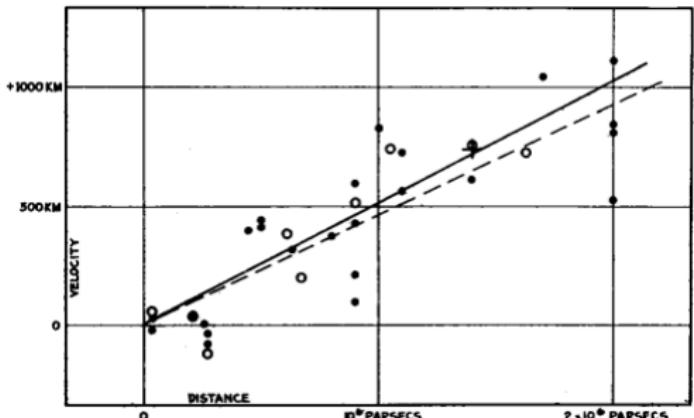
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# Introduction

## The Beginning of the Universe

- ▶ Friedmann (1924) and Lemaitre (1927) study isotropic and homogenous universes
- ▶ Lemaitre (1932) demonstrates general existence of singularities in Bianchi Type I class
- ▶ Since Hubble showed the universe was expanding (1929); these models then suggested that the universe had a beginning of infinite density



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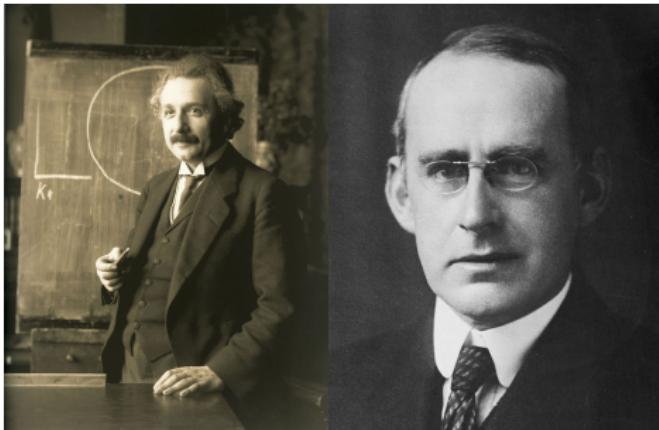
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# Introduction

What did these singularities mean?

- ▶ Early relativists, including Einstein, blamed symmetry assumptions and idealizations about the matter content
- ▶ Eddington infamously opposed Chandrasekhar's work
- ▶ Singularities were rarely studied and therefore poorly understood
- ▶ General study of singularities would have to wait for another generation of relativists



# Introduction

What did these singularities mean?

- ▶ Raychaudhuri (1955) finds first singularity theorem and develops the Raychaudhuri equation



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What did these singularities mean?

- ▶ Raychaudhuri (1955) finds first singularity theorem and develops the Raychaudhuri equation
- ▶ Penrose (1965) develops theorem for stellar collapse



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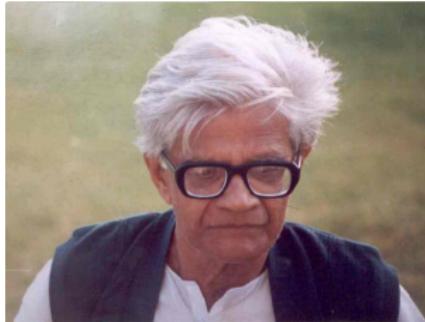
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What did these singularities mean?

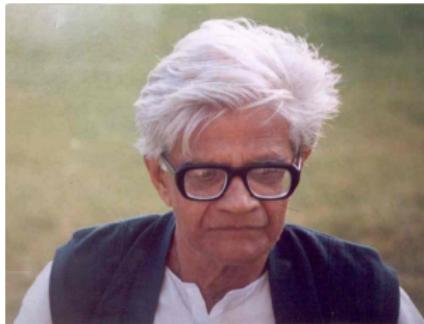
- ▶ Raychaudhuri (1955) finds first singularity theorem and develops the Raychaudhuri equation
- ▶ Penrose (1965) develops theorem for stellar collapse
- ▶ Extended to the early universe by Hawking (1967)



# Introduction

What did these singularities mean?

- ▶ Raychaudhuri (1955) finds first singularity theorem and develops the Raychaudhuri equation
- ▶ Penrose (1965) develops theorem for stellar collapse
- ▶ Extended to the early universe by Hawking (1967)
- ▶ These results were an integral part of the golden age of general relativity



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# Singularities

## Definition of Singularities

- ▶ It is quite difficult to define singularities in GR:

# Singularities

## Definition of Singularities

- ▶ It is quite difficult to define singularities in GR:
  - ▶ Singularities are not part of spacetime

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- ▶ It is quite difficult to define singularities in GR:
  - ▶ Singularities are not part of spacetime
  - ▶ Coordinate system choices makes characterization hard

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# Singularities

## Definition of Singularities

- ▶ It is quite difficult to define singularities in GR:
  - ▶ Singularities are not part of spacetime
  - ▶ Coordinate system choices makes characterization hard
  - ▶ Spaces can have singular behaviour whilst the curvature remains bounded

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# Singularities

## Definition of Singularities

- We need a definition using regular points in spacetime

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# Singularities

## Definition of Singularities

- ▶ We need a definition using regular points in spacetime
- ▶ If a traveller approached a singularity, he would disappear in finite time

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# Singularities

## Definition of Singularities

- ▶ We need a definition using regular points in spacetime
- ▶ If a traveller approached a singularity, he would disappear in finite time
- ▶ If a curve cannot continue backward in time, matter on the curve must have appeared *ab initio*

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# Singularities

## Definition of Singularities

- ▶ We need a definition using regular points in spacetime
- ▶ If a traveller approached a singularity, he would disappear in finite time
- ▶ If a curve cannot continue backward in time, matter on the curve must have appeared *ab initio*
- ▶ So we say that a singularity in a Lorentzian manifold is a non-spacelike curve which cannot be extended in a regular manner yet only takes part of its canonical parameter
- ▶ This is called (non-spacelike) *b-incompleteness*

# Singularities

## Structure of the Theorems

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Singularity theorems all have a basic form:

If a spacetime  $\mathcal{M}$  is sufficiently differentiable, then if

- ▶ Condition on curvature
- ▶ Causality condition
- ▶ Initial/boundary conditions

are satisfied, the spacetime has incomplete causal geodesics

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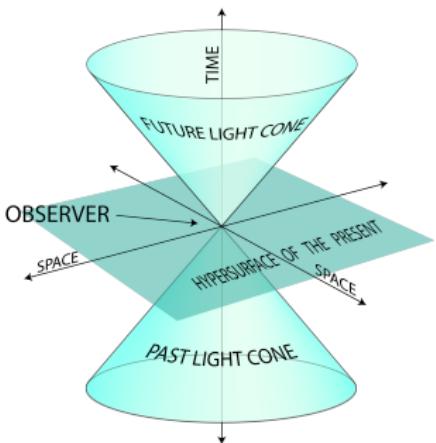
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# Causality

## Causal Structure

- At every point on the manifold we have lightcone



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- Gives us timelike, spacelike, and null vectors

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# Causality

## Causal Structure

- ▶  $x \ll y$  means there is a timelike curve from  $x$  to  $y$

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# Causality

## Causal Structure

- ▶  $x \ll y$  means there is a timelike curve from  $x$  to  $y$
- ▶ There is a hierarchy of causal conditions

# Causality

## Causal Structure

- ▶  $x \ll y$  means there is a timelike curve from  $x$  to  $y$
- ▶ There is a hierarchy of causal conditions
  - ▶ Chronology condition forbids closed timelike curves  
 $(x \not\ll x)$

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## Causal Structure

- ▶  $x \ll y$  means there is a timelike curve from  $x$  to  $y$
- ▶ There is a hierarchy of causal conditions
  - ▶ Chronology condition forbids closed timelike curves  
( $x \not\ll x$ )
  - ▶ Strong causality means that for every  $p \in \mathcal{M}$  there is a neighbourhood of  $p$  for which no timelike curve passes through more than once

# Causality

## Causal Structure

- ▶  $x \ll y$  means there is a timelike curve from  $x$  to  $y$
- ▶ There is a hierarchy of causal conditions
  - ▶ Chronology condition forbids closed timelike curves ( $x \not\ll x$ )
  - ▶ Strong causality means that for every  $p \in \mathcal{M}$  there is a neighbourhood of  $p$  for which no timelike curve passes through more than once
  - ▶ Stably causal means that under a small change of the metric, the space will still not violate the chronology condition

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# Causality

## Cauchy Surfaces

- ▶ A Cauchy surface  $\mathcal{S}$  for a set  $\mathcal{N}$  is a surface which every non-spacelike curve in  $\mathcal{N}$  intersects exactly once. We say  $\mathcal{N}$  is globally hyperbolic

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- ▶ Theorem: If  $\mathcal{S}$  is a Cauchy surface of  $\mathcal{N}$ , then  $\mathcal{N}$  is homeomorphic to  $\mathcal{S} \times \mathbb{R}$

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- ▶ Theorem: If  $\mathcal{S}$  is a Cauchy surface of  $\mathcal{N}$ , then  $\mathcal{N}$  is homeomorphic to  $\mathcal{S} \times \mathbb{R}$
- ▶ Theorem: Between any two points  $p, q$  in  $\mathcal{N}$  with  $p \ll q$  there is a non-spacelike geodesic of maximum length

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# Geodesic Focussing

## Geodesic Congruences

- ▶ Take a spacelike hypersurface  $\mathcal{S}$  on our manifold  $\mathcal{M}$

# Geodesic Focussing

## Geodesic Congruences

- ▶ Take a spacelike hypersurface  $\mathcal{S}$  on our manifold  $\mathcal{M}$
- ▶ At every point  $s \in \mathcal{S}$  there is a timelike unit vector  $u^\alpha$ , and so a geodesic  $\gamma_p(\tau)$  with  $\gamma_p(0) = p$  and  $\dot{\gamma}_p(0) = u^\alpha$

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# Geodesic Focussing

## Geodesic Congruences

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- ▶ We have a vector field  $\dot{\gamma}_p(\tau)$  in some region of spacetime, along with a series of surfaces  $\mathcal{S}_\tau$

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# Geodesic Focussing

## Geodesic Congruences

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- ▶ We have a vector field  $\dot{\gamma}_p(\tau)$  in some region of spacetime, along with a series of surfaces  $\mathcal{S}_\tau$
- ▶ The expansion scalar

$$\theta = \nabla^a \gamma_a(\tau),$$

where  $a$  is the derivative on  $\mathcal{S}_\tau$ , is the infinitesimal expansion of  $\mathcal{S}_\tau$

- ▶ If  $\theta > 0$  the surface is expanding, if  $\theta < 0$  the surface is shrinking

# Geodesic Focussing

## Geodesic Congruences

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- ▶ From the Raychaudhuri equation, we can derive

$$\frac{d\theta}{d\tau} \leq -\frac{1}{3}\theta^2 - R_{\alpha\beta}u^\alpha u^\beta$$

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# Geodesic Focussing

## Geodesic Congruences

- ▶ From the Raychaudhuri equation, we can derive

$$\frac{d\theta}{d\tau} \leq -\frac{1}{3}\theta^2 - R_{\alpha\beta}u^\alpha u^\beta$$

- ▶ If  $R_{\alpha\beta}u^\alpha u^\beta \geq 0$ , then

$$\frac{d\theta}{d\tau} \leq -\frac{1}{3}\theta^2 \leq 0$$

so gravity is attractive!

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# Geodesic Focussing

## Geodesic Congruences

- ▶ From the Raychaudhuri equation, we can derive

$$\frac{d\theta}{d\tau} \leq -\frac{1}{3}\theta^2 - R_{\alpha\beta}u^\alpha u^\beta$$

- ▶ If  $R_{\alpha\beta}u^\alpha u^\beta \geq 0$ , then

$$\frac{d\theta}{d\tau} \leq -\frac{1}{3}\theta^2 \leq 0$$

so gravity is attractive!

- ▶ If  $\theta < 0$  at some  $\tau_0$ , then  $\theta$  will reach negative infinity at finite affine parameter

# Geodesic Focussing

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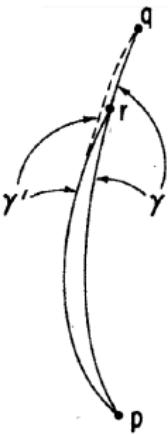
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Source: *General Relativity*, Wald (1984)

# Geodesic Focussing

## Energy Conditions

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- ▶ If  $R_{\alpha\beta}u^\alpha u^\beta \geq 0$ , then by the Einstein field equation

$$T_{\alpha\beta}u^\alpha u^\beta \geq \frac{1}{2}T$$

- ▶ This is known as the strong energy condition
- ▶ Seems to hold in most physically relevant scenarios
- ▶ For a perfect fluid with density of water, pressure would need to be less than  $-10^{15}$  atmospheres to violate the condition

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# Cosmological Theorems

The story so far

## Geodesic Focussing

- ▶ If strong energy condition holds, converging geodesics will focus at finite parameter
- ▶ If a geodesic has a focus point between two points  $p, q$ , then there is a longer path from  $p$  to  $q$

## Causality

- ▶ Between two points  $p, q$  in a globally hyperbolic set  $\mathcal{N}$ , there is a geodesic of maximum length.

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# Cosmological Theorems

## First Theorem

### Theorem

Let  $\mathcal{M}$  be a globally hyperbolic manifold satisfying the strong energy condition. Suppose  $\mathcal{S}$  is a Cauchy surface for which the geodesic expansion  $\theta \leq C \leq 0$  everywhere. Then no geodesic passing through  $\mathcal{S}$  can be extended further into the past than  $3/|C|$ .

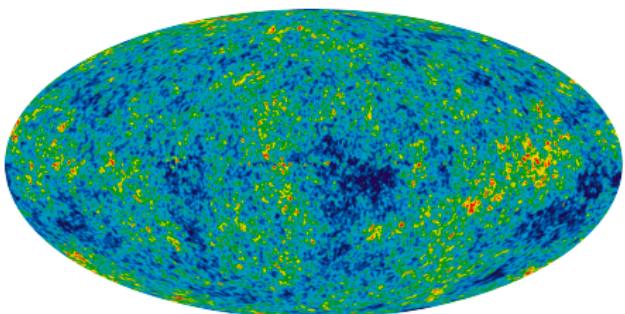
### Proof Sketch

Assume  $\gamma$  is a geodesic which can be extended further than  $3/|C|$  into the past, and let  $p$  be a point on this geodesic lying further than  $3/|C|$  in the past. Then there is a maximum length path from  $p$  to  $\mathcal{S}$  with length above  $3/|C|$ . But Raychaudhuri equations shows any curve of this length must contain a focal point.

# Cosmological Theorems

## Initial Conditions for Cosmology

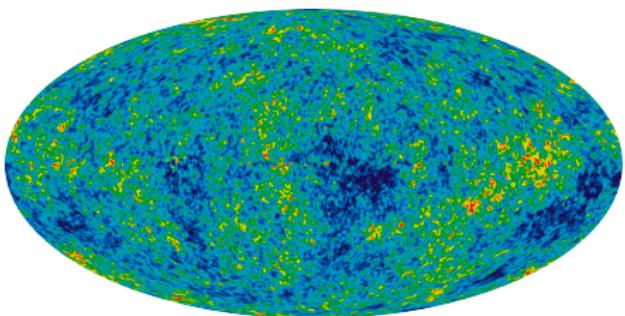
- ▶ Cosmological observations show that the universe is expanding
- ▶ In the  $\Lambda$ CMD model, the universe was fairly uniformly expanding at the time of nucleosynthesis,  $z \approx 10^8$ .



# Cosmological Theorems

## Initial Conditions for Cosmology

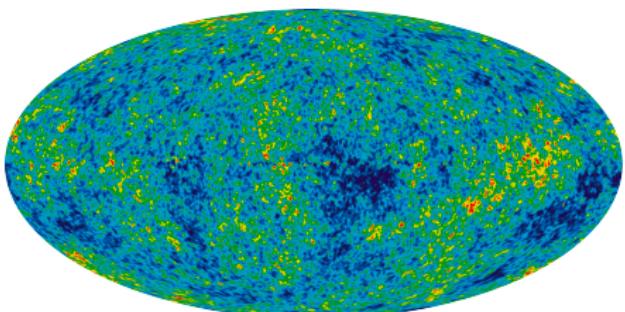
- ▶ Cosmological observations show that the universe is expanding
- ▶ In the  $\Lambda$ CMD model, the universe was fairly uniformly expanding at the time of nucleosynthesis,  $z \approx 10^8$ .
- ▶ Our theorem suggests that either



# Cosmological Theorems

## Initial Conditions for Cosmology

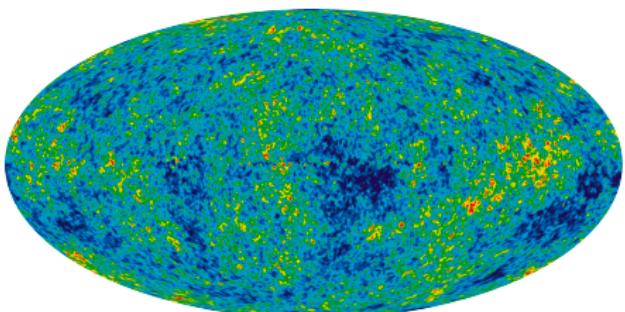
- ▶ Cosmological observations show that the universe is expanding
- ▶ In the  $\Lambda$ CMD model, the universe was fairly uniformly expanding at the time of nucleosynthesis,  $z \approx 10^8$ .
- ▶ Our theorem suggests that either
  - ▶ The universe had a singular beginning



# Cosmological Theorems

## Initial Conditions for Cosmology

- ▶ Cosmological observations show that the universe is expanding
- ▶ In the  $\Lambda$ CMD model, the universe was fairly uniformly expanding at the time of nucleosynthesis,  $z \approx 10^8$ .
- ▶ Our theorem suggests that either
  - ▶ The universe had a singular beginning
  - ▶ The strong energy condition is incorrect



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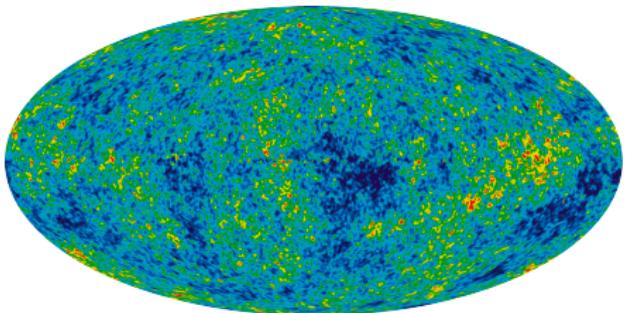
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# Cosmological Theorems

## Initial Conditions for Cosmology

- ▶ Cosmological observations show that the universe is expanding
- ▶ In the  $\Lambda$ CMD model, the universe was fairly uniformly expanding at the time of nucleosynthesis,  $z \approx 10^8$ .
- ▶ Our theorem suggests that either
  - ▶ The universe had a singular beginning
  - ▶ The strong energy condition is incorrect
  - ▶ The universe is not globally hyperbolic



# Cosmological Theorems

## Generalizations

- More sophisticated versions exist

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# Cosmological Theorems

## Generalizations

- ▶ More sophisticated versions exist
- ▶ Hawking (1967) replaced globally hyperbolic with strongly causal and removed the expanding everywhere assumption

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- ▶ More sophisticated versions exist
- ▶ Hawking (1967) replaced globally hyperbolic with strongly causal and removed the expanding everywhere assumption
- ▶ Hawking (1967) removed the globally hyperbolic conditions, assumes existence of compact spacelike three-surface. Weakens conclusion to the existence of an incomplete geodesic.

# Cosmological Theorems

## Generalizations

- ▶ More sophisticated versions exist
- ▶ Hawking (1967) replaced globally hyperbolic with strongly causal and removed the expanding everywhere assumption
- ▶ Hawking (1967) removed the globally hyperbolic conditions, assumes existence of compact spacelike three-surface. Weakens conclusion to the existence of an incomplete geodesic.
- ▶ Weakening causality condition cannot save us?

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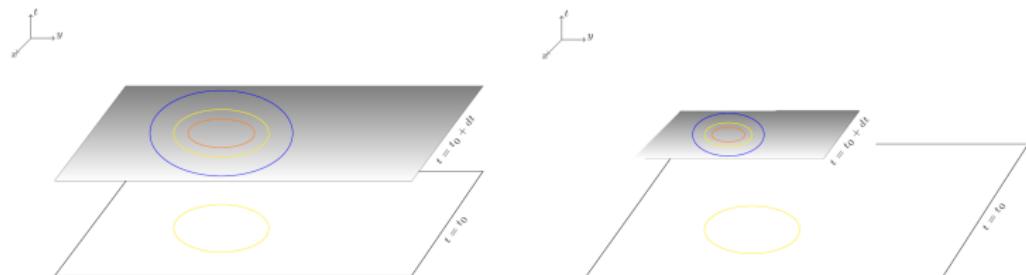
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# Stellar Collapse

## Trapped Surfaces

- ▶ A closed trapped surface is intuitively a surface so strongly bound that light cannot escape
- ▶ Because they occur in the Schwarzschild solution, they will occur in any stellar collapse which is sufficiently spherical



Source: <http://arxiv.org/abs/1410.5226>

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# Stellar Collapse

## The Theorems

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### Theorem (Penrose 1965):

If spacetime has a non-compact Cauchy surface and a closed future-trapped surface, and if the strong energy condition holds, then there are future incomplete null geodesics

### Theorem (Hawking and Penrose 1970):

If the strong, chronology and generic conditions hold and if there is a closed trapped surface, then spacetime is causal geodesically incomplete

- ▶ The generic condition requires that

$$K_{[a}R_{b]cd[e}K_{f]}K^cK^d \neq 0$$

for some tangent vector  $K^a$  to every geodesic.

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# Stellar Collapse

## Black Holes

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- ▶ What happens to these stars?

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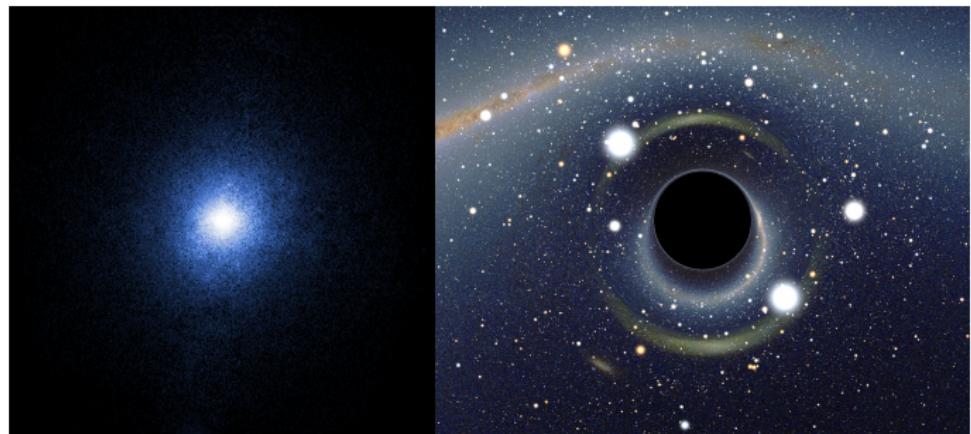
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# Stellar Collapse

## Black Holes

- ▶ What happens to these stars?
- ▶ They will settle down to become black holes



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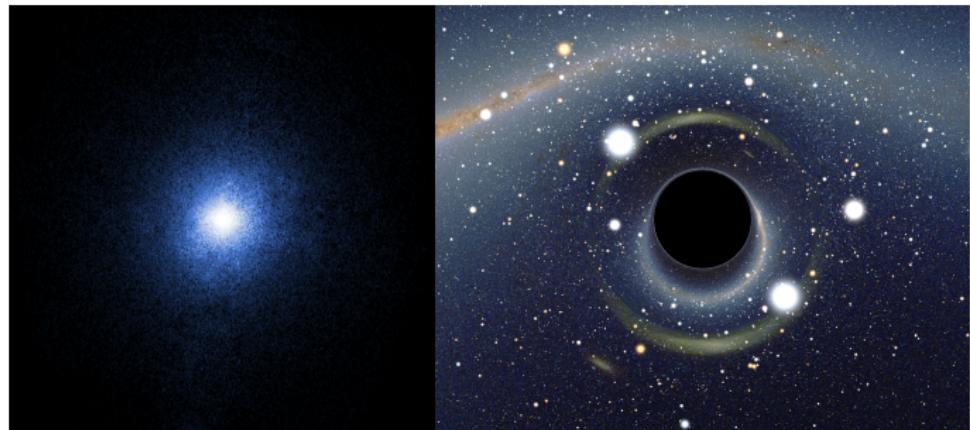
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# Stellar Collapse

## Black Holes

- ▶ What happens to these stars?
- ▶ They will settle down to become black holes
  - ▶ Schwarzschild solution



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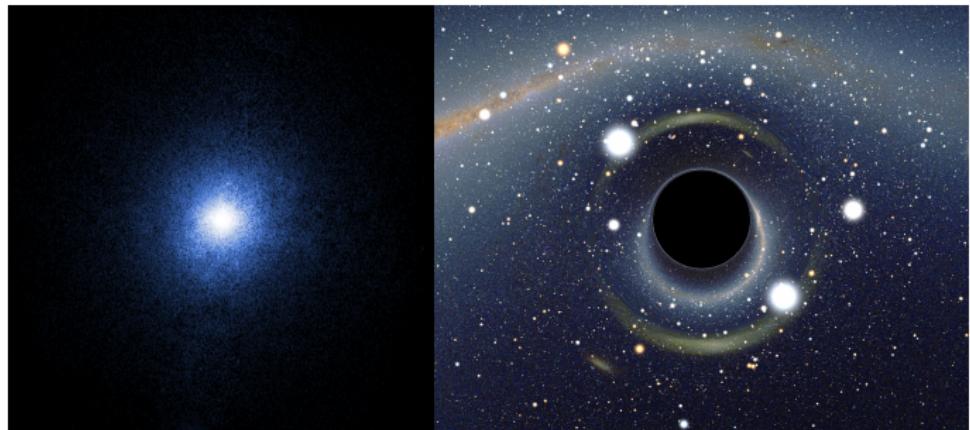
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# Stellar Collapse

## Black Holes

- ▶ What happens to these stars?
- ▶ They will settle down to become black holes
  - ▶ Schwarzschild solution
  - ▶ Kerr solution



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# Conclusion

What next?

- ▶ We seem to have shown there are singularities
  - ▶ Stellar collapse
  - ▶ The beginning of the universe

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# Conclusion

What next?

- ▶ We seem to have shown there are singularities
  - ▶ Stellar collapse
  - ▶ The beginning of the universe
- ▶ What is the nature of these singularities?

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# Conclusion

What next?

## Possibility 1

The singularity is a region of infinite curvature and density.

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# Conclusion

What next?

## Possibility 1

The singularity is a region of infinite curvature and density.

## Possibility 2

A more subtle pathology occurs, perhaps causal conditions are violated.

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# Conclusion

What next?

## Possibility 1

The singularity is a region of infinite curvature and density.

## Possibility 2

A more subtle pathology occurs, perhaps causal conditions are violated.

## Possibility 3

General relativity and/or the strong energy condition breaks down - a more complete theory (quantum gravity?) is needed.

# Conclusion

What next?

- ▶ Singularity theorems are still active

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# Conclusion

What next?

- ▶ Singularity theorems are still active
- ▶ Can we apply the theorems to different physical situations?

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# Conclusion

What next?

- ▶ Singularity theorems are still active
- ▶ Can we apply the theorems to different physical situations?
- ▶ Can we relax assumptions in the theorems?

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# Conclusion

What next?

- ▶ Singularity theorems are still active
- ▶ Can we apply the theorems to different physical situations?
- ▶ Can we relax assumptions in the theorems?
- ▶ Can we better characterize singularities?

# Conclusion

What next?

- ▶ Thank you for listening
- ▶ Any questions?

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