

¹ SEARCH FOR EXOTIC HIGGS DECAYS TO LIGHT
² NEUTRAL SCALARS IN FINAL STATES WITH
³ BOTTOM QUARKS AND TAU LEPTONS

⁴ KA YU STEPHANIE KWAN

⁵ A DISSERTATION
⁶ PRESENTED TO THE FACULTY
⁷ OF PRINCETON UNIVERSITY
⁸ IN CANDIDACY FOR THE DEGREE
⁹ OF DOCTOR OF PHILOSOPHY

¹⁰ NOT YET RECOMMENDED FOR ACCEPTANCE
¹¹ BY THE DEPARTMENT OF PHYSICS
¹² ADVISER: ISOBEL OJALVO

¹³ MAY 2024

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Abstract

Open questions in particle physics may be addressed by the existence of an extended Higgs sector beyond the Standard Model Higgs boson with mass 125 GeV, which was discovered in 2012 at the Large Hadron Collider (LHC) by the CMS and ATLAS experiments. Many properties of a potential extended Higgs sector remain unconstrained by current measurements, making direct searches of exotic Higgs decays a powerful probe of new physics. The decay of the 125 GeV Higgs boson into two light neutral scalar particles ($h \rightarrow aa$) is allowed in extensions of the Standard Model, such as Two Higgs Doublet Models extended with a scalar singlet (2HDM+S). We present a search at CMS for exotic decays of the 125 GeV Higgs boson to two light neutral scalars, which decay to two bottom quarks and two tau leptons ($h \rightarrow aa \rightarrow bb\tau\tau$). This analysis is combined with a different search where the light scalars decay to two bottom quarks and two muons. The results from the $bb\tau\tau$ analysis and the combined analyses are interpreted in 2HDM+S scenarios. In a different extension of the Standard Model, the Two Real Singlet Model (TRSM), the 125 GeV Higgs boson can decay to two light scalars with unequal mass ($h \rightarrow a_1a_2$). This decay has not been searched for to date at CMS. We present ongoing work on a search for $h \rightarrow a_1a_2$, where the a_2 decays into two a_1 , resulting in four bottom quarks and two tau leptons in the final state, in the $\mu\tau_h$ channel of the $\tau\tau$ decay. Such searches for rare processes will directly benefit from the increased datasets that will be generated by the High-Luminosity LHC (HL-LHC), which is scheduled to increase the LHC’s number of simultaneous proton-proton collisions by a factor of five to seven. To contribute to the performance of the CMS Level-1 Trigger in selecting collisions with interesting physics, this thesis presents an upgraded algorithm for reconstructing electrons and photons in the barrel calorimeter, which will use information with higher spatial granularity to distinguish genuine electrons and photons from background.

⁴²

Acknowledgements

⁴³ Placeholder acknowledgements.

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422	errors and only several of the full set of systematic errors (only those	
423	associated with the lepton energy scales and τ_h identification efficiency).	135

⁴²⁴ Chapter 1

⁴²⁵ Introduction

⁴²⁶ The Standard Model is the current prevailing theoretical framework that encompasses
⁴²⁷ all known elementary particles to date and describes their interactions, yet falls short
⁴²⁸ of describing open problems in physics. Here, we describe the history of the Standard
⁴²⁹ Model and its particle content (Section 1.1), and provide a mathematical motivation of
⁴³⁰ the SM as a gauge theory (Section 1.2). We introduce the Higgs mechanism (Section
⁴³¹ 1.3), and outline two groups of theoretical extensions to the Standard Model that
⁴³² feature extended Higgs sectors (Sections 1.4 and 1.5).

⁴³³ 1.1 History of the Standard Model

⁴³⁴ The building blocks of our modern-day understanding of particle physics were estab-
⁴³⁵ lished over the course of many decades by experimental discoveries and theoretical
⁴³⁶ advances, culminating in the development of a theoretical framework known as the
⁴³⁷ Standard Model (SM). In the 1880s, the electron was the first subatomic particle to
⁴³⁸ be identified, through measurements of particles produced by ionizing gas. By the
⁴³⁹ 1930s, atoms were known to consist mostly of empty space, with protons and neutrons
⁴⁴⁰ concentrated at the center and orbited by electrons. Spurred by advances in parti-
⁴⁴¹ cle accelerator technology, the experimental discoveries of the positron, the muon,

and the pion, painted an increasingly complicated picture of particle physics that could not be described solely with atomic physics [1]. Quantum field theory (QFT) began to be developed in the early 20th century as an extension of the conceptual framework of quantum mechanics to electromagnetic fields [2]. In 1927, Dirac coined the name quantum electrodynamics (QED), which was the first part of QFT that was developed. QED quantized the electromagnetic field and supplied a relativistic theory of the electron, and could be applied to concrete physical processes such as the scattering of high-frequency photons by free electrons (Compton scattering), and the production of electron-positron pairs by photons [2]. In the 1940s the QED-only picture was challenged by the realization that the four-fermion theory of weak interactions had infinities at higher orders of perturbation theory which could not be removed via the technique of renormalization [3], i.e. shifting divergences into parts of the theory that do not influence empirical measurements [2].

In the 1950s and 1960s, QFT was extended to describe not only the electromagnetic force, but also the strong and weak force, with the final picture forming the Standard Model. This took place in the development and maturation of three principles: the quark model, the idea of gauge (or local) symmetry, and spontaneously broken symmetry [3]. In the fully fledged QFT, Lagrangians had to be formed that contained new classes of quantum fields, or particles [2].

The particle content of the Standard Model is summarized in Fig. 1.1. Particles are grouped into fermions, which comprise all known matter, and bosons, which mediate the interactions between particles. Fermions consist of quarks and leptons, and are grouped into three generations. For example, the electron belongs to the first generation of leptons. The second and third generation counterparts of the electron are the muon and the tau lepton, and are over 200 and 30,000 times heavier than the electron respectively. The quarks are also organized into three generations (top and bottom quarks, charm and strange quarks, and up and down quarks), and

469 carry fractional electric charge. Bosons are force carriers; the interaction of fermions
 470 with bosons corresponds to fundamental forces. The Standard Model describes the
 471 electromagnetic force, the strong nuclear force, and the weak nuclear force. Through
 472 the strong force, quarks can form composite particles known as hadrons. Familiar
 473 examples of hadrons are the protons and neutrons in the nucleus of an atom.

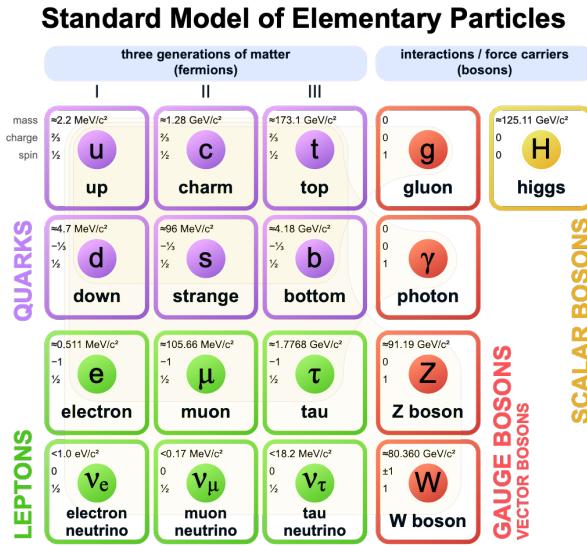


Figure 1.1: Table of Standard Model particles showing the grouping of the fermions into three generations of matter and the bosons, responsible for carrying the three fundamental forces in the Standard Model. The masses, charges, and spins of the particles are shown. The antimatter counterparts of the fermions are not shown. The possible interactions between the fermions and gauge bosons are highlighted.

1.2 The Standard Model as a gauge theory

474 In this section we lay the theoretical foundations of the Standard Model as a gauge
 475 theory, starting from the principle of gauge invariance (gauge symmetry), with local
 476 gauge symmetries giving rise to interactions between particles.
 477

478 Gauge theories of elementary particle interactions originate from a freedom of
 479 choice in the mathematical description of particle fields which has no effect on the
 480 particles' physical states [4]. The existence and form of the particles' interactions,

481 can be deduced from the existence of physically indeterminate, gaugable quantities.

482 An example of this gauge invariance is classical physics is the electromagnetic
483 interaction, where the fundamental field is the four-vector potential A^μ [4]. The
484 physical electromagnetic fields and Maxwell's equations arise from the elements of
485 the tensor $F_{\mu\nu}(x) = \partial_\mu A_\nu(x) - \partial_\nu A_\mu(x)$. Any two choices of A^μ that are related by a
486 transformation of the form

$$A_\mu \rightarrow A_\mu + \partial_\mu \alpha \quad (1.1)$$

487 for any real, differentiable function $\alpha(x)$, describe the same physical configuration,
488 and has no effect on Maxwell's equations. This "redundancy" in the choice of gauge
489 in Eqn. 1.1 is called a gauge symmetry.

490 One important consequence of gauge symmetry comes from the application of
491 Noether's theorem, which states that for every global transformation under which the
492 Lagrangian density is invariant, there exists a conserved quantity. If $\mathcal{L}(\Psi(x), \partial_\mu \Psi(x))$
493 is invariant under the transformation of the wave function $\Psi(x) \rightarrow \Psi'(x)$, where
494 $\Psi'(x) = \Psi(x) + \delta\Psi(x)$, then there exists a conserved current

$$\partial_\mu \left(\frac{\partial \mathcal{L}(x)}{\partial(\partial_\mu \Psi(x))} \delta\Psi(x) \right) = 0 \quad (1.2)$$

495 In classical mechanics, the conservation of linear momentum, angular momentum,
496 and energy follows from translational invariance, rotational variance, and invariance
497 under translations in time [4]. Likewise, charge conservation can be shown to arise
498 from the invariance of the Dirac Lagrangian density $\mathcal{L}_{\text{Dirac}} = \bar{\Psi}(i\gamma^\mu \partial_\mu - m)\Psi$ under the
499 particle wavefunction's phase transformation, $\Psi'(x) = \exp(i\epsilon\chi)\Psi(x)$. Thus Noether's
500 theorem establishes a correspondence between a gauge symmetry and a conserved
501 internal property (e.g. charge or momentum).

502 Interactions between particles arise if we modify the wave function with a phase

503 transformation $\Psi'(x) = \exp(ie\chi)\Psi(x)$, and allow the phase χ to be a function of
 504 spacetime [4]. A wave function of the form

$$\Psi'(x) = \exp(ie\chi(x))\Psi(x) \quad (1.3)$$

505 can be verified to *not* be a solution to the Dirac equation for free particles: $(i\gamma^\mu\partial_\mu -$
 506 $m)\Psi(x) = 0$. This necessitates a modified Dirac equation, where the derivative takes
 507 into account that the vector field $V(x)$ needs to be compared at two displaced space-
 508 time points in a curvilinear coordinate system:

$$\mathcal{D}_\mu \equiv \lim_{\Delta x^\mu \rightarrow 0} \frac{V_{||}(x + \Delta x) - V(x)}{\Delta x^\mu} \quad (1.4)$$

509 We define a covariant derivative,

$$D_\mu = \partial_\mu + ieA_\mu \quad (1.5)$$

510 where $A_\mu(x)$ is a 4-vector potential. Thus the modified Dirac equation reads:

$$(i\gamma^\mu\mathcal{D}_\mu - m)\Psi(x) = 0 \quad (1.6)$$

511 The simultaneous gauge transformation $A'_\mu(x) = A_\mu(x) - \partial_\mu\chi(x)$ and wavefunction
 512 transformation $\Psi'(x) = \exp(ie\chi(x))\Psi(x)$ leaves the covariant-derivative form of the
 513 Dirac equation (Eqn 1.1) invariant.

514 The generalization of this result is as follows: if a theory is invariant for unitary
 515 transformations U of the particle states according to

$$\Psi' = U\Psi \quad (1.7)$$

516 One must define a derivative of the form

$$D^\mu = \partial^\mu + igB^\mu \quad (1.8)$$

517 to keep the theory invariant under Eqn. 1.7. The four-potential B^μ represents the
518 interacting four-potential which must be added to keep the theory invariant.

519 In the case of the Standard Model, the theory is built around the gauge trans-
520 formations $G = SU(3) \times SU(2) \times U(1)$. $SU(3)$ is associated to the strong force
521 (subscripted C); $SU(2)$ is associated to the weak force (subscripted L); and $U(1)$ is
522 hypercharge (subscripted Y). The gauge-covariant derivative is

$$\mathcal{D}_\mu = \partial_\mu - ig'B_\mu \frac{Y}{2} - igW_\mu^\alpha \frac{\tau_a}{2} - ig_s G_\mu^k \frac{\lambda_k}{2} \quad (1.9)$$

523 • In the $U(1)_Y$ term, B_μ is the weak hypercharge field.

524 • In the $SU(2)_L$ term, $W_\mu(x) = (W_\mu^1(x), W_\mu^2(x), W_\mu^3(x))$ are a triplet of four-
525 potentials. $\tau/2$ are the Pauli matrices, generators of the $SU(2)$ transformation.

526 • In the $SU(3)_C$ term, the gluon (color) field is G_μ . λ_k are the Gell-Man matrices,
527 generators of the $SU(3)$ transformation.

528 The invariance of the Standard Model under $SU(3)_C \times SU(2)_L \times U(1)_Y$ requires
529 massless fermions and massless force carriers.

530 1.3 The Higgs Mechanism

531 To introduce mass into the theory, i.e. to change the propagation of the gauge par-
532 ticles and all the fermions, the physical vacuum cannot have all the symmetries of
533 the Standard Model Lagrangian [4]. The symmetries of the physical vacuum must
534 be spontaneously broken, without affecting gauge invariance in the Lagrangian. The

535 Higgs mechanism proposes the existence of a scalar field, or fields, with nonzero vac-
 536 um expectation values, which reduce the gauge symmetries of the physical vacuum
 537 from $SU(3)_C \times SU(2)_L \times U(1)_Y$ down to $SU(3)_C \times U(1)_{EM}$.

538 The Higgs field interacts with the gauge bosons and fermions throughout space,
 539 impeding their free propagation. The resulting broken symmetry correctly predicts
 540 the mass ratio of the neutral (Z) and charged (W) massive electroweak bosons, and
 541 predicts that at least one physical degree of freedom in the Higgs field is a particle
 542 degree of freedom, called the Higgs boson. The location of the minimum of the Higgs
 543 potential can be constrained from previously measured Standard Model parameters,
 544 but the shape of the mass distribution of the Higgs boson must be experimentally
 545 measured.

546 The minimal choice of Higgs field comes from the breaking of $SU(2)_L \times U(1)_Y$
 547 down to $U(1)_{EM}$. The smallest $SU(2)$ multiplet is the doublet. The existence of three
 548 massive electroweak bosons leads the Higgs sector to have at least three degrees of
 549 freedom. The minimal single-doublet complex scalar Higgs field is

$$\Phi(x) = \begin{pmatrix} \phi^+(x) \\ \phi^0(x) \end{pmatrix} = \frac{1}{\sqrt{2}} \begin{pmatrix} \phi_1^+(x) + i\phi_2^+(x) \\ \phi_1^0(x) + i\phi_2^0(x) \end{pmatrix} \quad (1.10)$$

550 where ϕ_1^+ , ϕ_2^+ , ϕ_1^0 , and ϕ_2^0 are real (four degrees of freedom). By convention, the
 551 nonzero vacuum expectation value is assigned to ϕ_1^0 .

552 The minimal self-interacting Higgs potential that is invariant under $SU(2)_L \times$
 553 $U(1)_Y$ is given by

$$V(\Phi^\dagger \Phi) = -\mu^2 \Phi^\dagger \Phi + \lambda (\Phi^\dagger \Phi)^2, \quad \mu^2 > 0, \lambda > 0 \quad (1.11)$$

554 where λ is the coupling strength of the four-point Higgs interaction. The potential

555 energy is minimized at

$$\Phi_{\min} = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v \end{pmatrix}, \text{ where } v = \sqrt{\mu^2/\lambda} \quad (1.12)$$

556 Choosing a fixed orientation of $\langle \Phi \rangle$ out of a continuous set of possible ground states
557 spontaneously breaks the symmetry of the physical vacuum, as illustrated in Fig 1.2.

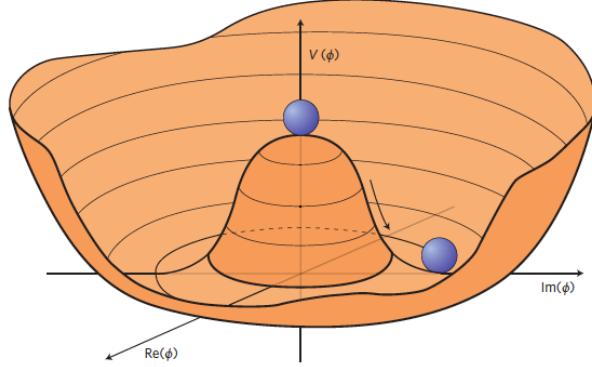


Figure 1.2: An illustration of the Higgs potential [5]. Choosing any of the points at the bottom of the potential breaks spontaneously the rotational $U(1)$ symmetry.

558 The excitations of the Higgs field with respect to the minimum Φ_{\min} are parame-
559 terized by

$$\Phi(x) = \exp(i\xi(x) \cdot \tau) \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v + H(x) \end{pmatrix} \quad (1.13)$$

560 Three degrees of freedom are coupled directly to the electroweak gauge bosons; this
561 is often referred to as the gauge bosons “eating” the Goldstone bosons to form the
562 longitudinal polarizations of the massive spin-1 boson states. The $H(x)$ excitation is
563 in the radial direction and corresponds to the free particle state of the Higgs boson.

564 1.4 Two-Higgs Doublet Models

565 One of the simplest possible extensions to the Standard Model is adding a doublet
566 to the minimal Higgs sector of the Standard Model, which is a $SU(2)_L$ doublet H

567 with hypercharge $Y = +\frac{1}{2}$, denoted here as $H \sim 2_{+1/2}$. These extensions are found
 568 in several theories such as supersymmetry. A general 2HDM can be extended with a
 569 light scalar (2HDM+S) to obtain a rich set of exotic Higgs decays [6].

The charges of the Higgs fields are chosen to be $H_1 \sim 2_{-1/2}$ and $H_2 \sim 2_{+1/2}$, which
 acquire vacuum expectation values $v_{1,2}$ which are assumed to be real and aligned [6].
 Expanding about the minima yields two complex and four real degrees of freedom:

$$H_1 = \frac{1}{\sqrt{2}} \begin{pmatrix} v_1 + H_{1,R}^0 + iH_{1,I}^0 \\ H_{1,R}^- + iH_{1,I}^- \end{pmatrix} \quad (1.14)$$

$$H_2 = \frac{1}{\sqrt{2}} \begin{pmatrix} H_{2,R}^+ + iH_{2,I}^+ \\ v_2 + H_{2,R}^0 + iH_{2,I}^0 \end{pmatrix} \quad (1.15)$$

570 The charged scalar and pseudoscalar mass matrices are diagonalized by a rotation
 571 angle β , defined as $\tan \beta = v_2/v_1$. One charged (complex) field and one neutral
 572 pseudoscalar combination of $H_{1,2,I}^0$ are eaten by the SM gauge bosons after electroweak
 573 symmetry breaking [6]. The other complex field yields two charged mass eigenstates
 574 H^\pm , which are assumed to be heavy. The remaining three degrees of freedom yield
 575 one neutral pseudoscalar mass eigenstate

$$A = H_{1,I}^0 \sin \beta - H_{2,I}^0 \cos \beta \quad (1.16)$$

576 and two neutral scalar mass eigenstates (where $-\pi/2 \leq \alpha \leq \pi/2$)

$$\begin{pmatrix} h \\ H^0 \end{pmatrix} = \begin{pmatrix} -\sin \alpha & \cos \alpha \\ \cos \alpha & \sin \alpha \end{pmatrix} \begin{pmatrix} H_{1,R}^0 \\ H_{2,R}^0 \end{pmatrix} \quad (1.17)$$

577 We assume that the 2HDM is near or in the decoupling limit: $\alpha \rightarrow \pi/2 - \beta$, where the
 578 lightest state in the 2HDM is h , which we identify as the 125 GeV Higgs particle [6].
 579 In this limit, the fermion couplings of h become identical to the Standard Model

580 Higgs, while the gauge boson couplings are very close to Standard Model-like for
581 $\tan \beta \gtrsim 5$. All of the properties of h can be determined by just two parameters: $\tan \beta$
582 and α , and the fermion couplings to the two Higgs doublets.

583 2HDM can be extended by a scalar singlet (2HDM+S) [6]:

$$S = \frac{1}{\sqrt{2}}(S_R + iS_I) \quad (1.18)$$

584 If this singlet only couples to the Higgs doublets $H_{1,2}$ and has no direct Yukawa
585 couplings, all of its couplings to SM fermions result from mixing with $H_{1,2}$. Under
586 these simple assumptions, exotic Higgs decays $h \rightarrow ss \rightarrow X\bar{X}Y\bar{Y}$ or $h \rightarrow aa \rightarrow$
587 $X\bar{X}Y\bar{Y}$, and $h \rightarrow aZ \rightarrow X\bar{X}Y\bar{Y}$ are permitted, where $s(a)$ is a (pseudo)scalar mass
588 eigenstate mostly composed of $S_R(S_I)$, and X, Y are Standard Model fermions or
589 gauge bosons. There are two pseudoscalars in the 2HDM+S, and the mostly singlet-
590 like pseudoscalar can be chosen to be the one lighter than the SM-like Higgs. For
591 $m_a < m_h - m_Z \sim 35$ GeV, the exotic Higgs decay $h \rightarrow Za$ is possible, and for
592 $m_a < m_h/2 \approx 63$ GeV, the exotic Higgs decay $h \rightarrow aa$ is possible.

593 In 2HDM, and by extension 2HDM+S, there are four types of fermion couplings
594 commonly discussed in the literature that forbid flavor-changing neutral currents at
595 tree level [6]. These are referred to as Type I (all fermions couple to H_2), Type II
596 (MSSM-like, d_R and e_R couple to H_1 , u_R to H_2), Type III (lepton-specific, leptons
597 and quarks couple to H_1 and H_2 respectively) and Type IV (flipped, with u_R , e_R
598 coupling to H_2 and d_R to H_1). The exact branching ratios of the pseudoscalars to
599 Standard Model particles vary depending on the 2HDM+S model and the value of
600 $\tan \beta$ (e.g. Fig. 1.3).

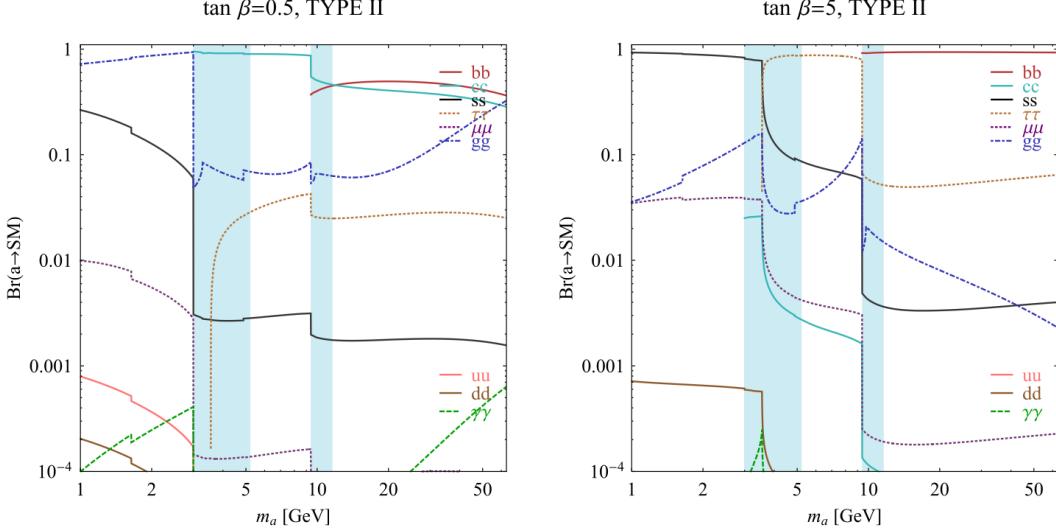


FIG. 7 (color online). Branching ratios of a singletlike pseudoscalar in the 2HDM + S for type-II Yukawa couplings. Decays to quarkonia likely invalidate our simple calculations in the shaded regions.

Figure 1.3: Branching ratios of a singlet-like pseudoscalar in Type II 2HDM+S for $\tan\beta = 0.5$ (left) and $\tan\beta = 5$ (right) from [6], showing the dependence of the branching ratios on $\tan\beta$, as well as the prominence of the branching ratios to bb and $\tau\tau$, the channels searched for in the analysis presented here.

1.5 Two Real Singlet Model

The two real singlet model (TRSM) adds two real singlet degrees of freedom to the Standard Model. These are written as two real singlet fields S and X . Depending on the vacuum expectation values acquired by the scalars, different phases of the model can be realized [7]. To reduce the number of free parameters, two discrete \mathbb{Z}_2 symmetries are introduced. The fields are decomposed as

$$\Phi = \begin{pmatrix} 0 \\ \frac{\phi_h + v}{\sqrt{2}} \end{pmatrix}, S = \frac{\phi_S + v_S}{\sqrt{2}}, X = \frac{\phi_X + v_X}{\sqrt{2}} \quad (1.19)$$

To achieve electroweak-breaking symmetry, $v = v_{SM} \sim 246$ GeV is necessary. If the vacuum expectation values $v_S, v_X \neq 0$ the \mathbb{Z}_2 are spontaneously broken, and the fields $\phi_{h,S,X}$ mix into three physical scalar states. This is called the broken phase and leads to the most interesting collider phenomenology.

611 The mass eigenstates $h_{1,2,3}$ are related to the fields $\phi_{h,S,X}$ through a 3×3 orthogonal
612 mixing matrix denoted R . The mass eigenstates are assumed to be ordered $M_1 \leq$
613 $M_2 \leq M_3$. R is parameterized by the three mixing angles θ_{hS} , θ_{hX} , θ_{SX} . The nine
614 parameters of the scalar potential can be expressed in terms of the three physical
615 Higgs masses, the three mixing angles, and the three vacuum expectation values.

616 After fixing one of the Higgs masses to the mass of the observed Higgs boson, and
617 fixing the Higgs doublet vacuum expectation value to its Standard Model value, there
618 are seven remaining free parameters of the TRSM [7].

619 In one benchmark scenario of TRSM [7], the heaviest scalar state h_3 is identified
620 with the 125 GeV Higgs, h_{125} , and it can decay asymmetrically $h_{125} \rightarrow h_1 h_2$, which
621 we also denote $h \rightarrow a_1 a_2$ to highlight the similarity with the symmetric decay $h \rightarrow aa$
622 typically interpreted in 2HDM+S as discussed. The parameter values in TRSM are
623 chosen such that the coupling of h_3 to Standard Model particles are nearly identical
624 to the Standard Model predictions.

625 In benchmark scenario 1 (benchmark plane 1, or BP1) (Fig. 1.4) [7], the maximal
626 branching ratios for $h_3 \rightarrow h_1 h_2$ reach up to 7 – 8% which translates into a signal
627 rate of around 3 pb. These maximal branching ratios are reached in the intermediate
628 mass state for h_2 , $M_2 \sim 60 – 80$ GeV. For $M_2 < 40$ GeV, although phase space opens
629 up significantly for light decay products, the branching ratio becomes smaller.

630 If the decay channel $h_2 \rightarrow h_1 h_1$ is kinematically open (i.e. $M_2 > 2M_1$), it is the
631 dominant decay mode leading to a significant rate for the $h_1 h_1 h_1$ final state, in a
632 “cascade” decay. In BP1, $BR(h_2 \rightarrow h_1 h_1) \simeq 100\%$ above the red line in Fig. 1.4. If,
633 in addition, $M_1 \gtrsim 10$ GeV, the h_1 decays dominantly to $b\bar{b}$ leading to a sizable rate
634 for the $b\bar{b}b\bar{b}b\bar{b}$ final state as shown in Fig. 1.4 (*bottom right*).

635 If the $h_2 \rightarrow h_1 h_1$ decay is kinematically closed (i.e. $M_2 < 2M_1$), both scalars decay
636 directly to Standard Model particles, with branching ratios identical to a Standard
637 Model-like Higgs boson, i.e. with the $b\bar{b}b\bar{b}$ final state dominating, as shown in Fig. 1.4

₆₃₈ (*bottom left*), while at smaller masses, combinations with τ leptons and eventually
₆₃₉ final states with charm quarks and muons become relevant [7].

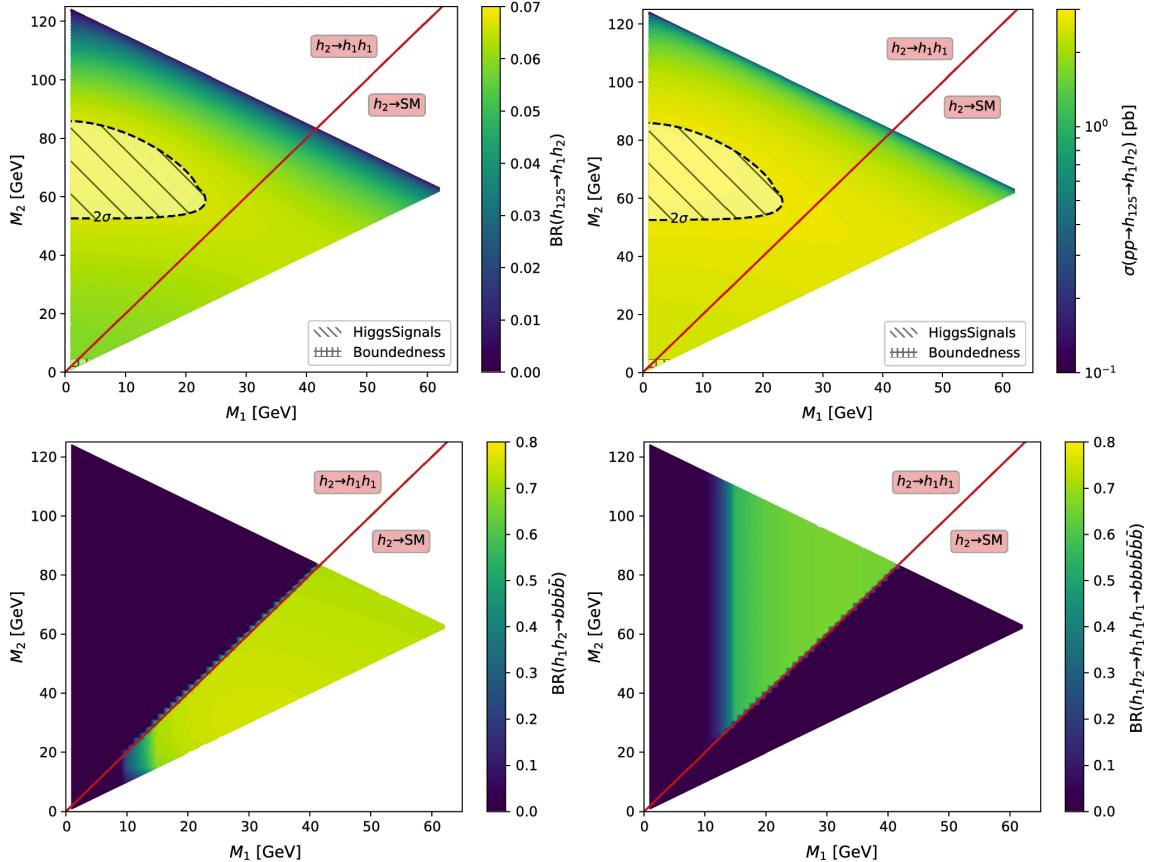


Figure 1.4: Benchmark plane BP1 for benchmark scenario 1 from [7], for the decay signature $h_{125} \rightarrow h_1 h_2$ with $h_{125} \equiv h_3$, defined in the (M_1, M_2) plane. The color code shows $\text{BR}(h_3 \rightarrow h_1 h_2)$ (*top left*) and the 13 TeV LHC signal rate for $pp \rightarrow h_3 \rightarrow h_1 h_2$ (*top right*). The red line separates the region $M_2 > 2M_1$, where $\text{BR}(h_2 \rightarrow h_1 h_1) \sim 100\%$, from the region $M_2 < 2M_1$, where $\text{BR}(h_2 \rightarrow F_{SM}) \sim 100\%$. The *bottom left* and *right* show the branching ratio of the $h_1 h_2$ into (respectively) $b\bar{b}b\bar{b}$, and through a $h_2 \rightarrow h_1 h_1$ cascade to $b\bar{b}b\bar{b}b\bar{b}$. The hatched region indicates where the decay rate slightly exceeds the 2σ upper limit inferred from the LHC Higgs rate measurements, though the region depends on the parameter choices and experimental searches should cover the whole mass range.

640 **Chapter 2**

641 **The Large Hadron Collider and the**
642 **CMS Experiment**

643 This chapter introduces the key aspects of the CERN Large Hadron Collider (LHC)
644 and the Compact Muon Solenoid (CMS) experiment where the work for this thesis was
645 conducted. Section 2.1 describes the history of accelerator developments at CERN
646 that led to the construction of the LHC, the current LHC configuration, and the
647 largest experiments located at the LHC. The concepts of beam luminosity and pile-
648 up, which are critical for understanding and measuring high-energy particle collisions,
649 are described in Section 2.2 and discussed in the context of the High-Luminosity
650 LHC (HL-LHC) upgrade in Section 2.3. Lastly, Section 2.4 describes the design
651 and function of CMS and its subdetectors, and terminates in a description of data
652 processing at CMS, beginning from online event filtering in the Level-1 Trigger, to
653 processing in the High-Level Trigger, to offline particle reconstruction, and finally
654 long-term storage and processing of measured events.

655 2.1 The Large Hadron Collider

656 CERN, the European Organization for Nuclear Research, is an international organiza-
657 tion based in Meyrin, Switzerland which operates the world's largest particle physics
658 laboratory, and is the site of the Large Hadron Collider (LHC) [8]. The very first
659 accelerator built at CERN was the 600 MeV Synchrocyclotron (SC), which initially
660 provided beams for CERN's first experiments. The newer and more powerful Proton
661 Synchrotron (PS), which could accelerate particles to an energy of 28 GeV, began op-
662 erations in 1959 and is still in use today. The first hadron collider at CERN was the
663 Intersecting Storage Rings (ISR), which consisted of two interlaced rings each with a
664 diameter of 200. The ISR collided protons at a center-of-mass energy of 62 GeV and
665 began measuring collisions in 1971. In 1968 CERN began to accelerate heavy ions
666 in the Super Proton Synchrotron (SPS), which is 7 kilometers in circumference and
667 was the first of CERN's giant underground rings to be built. The SPS became the
668 forefront of CERN's particle physics program in 1976, and in 1981 was converted into
669 a proton-antiproton collider. The final and largest underground ring constructed at
670 CERN was the Large Electron-Positron (LEP) collider, which was commissioned in
671 July 1989 and hosted 5176 magnets and 128 accelerating cavities located around a
672 27-kilometer circumference. Over 11 years of research, four detectors, ALEPH, DEL-
673 PHI, L3, and OPAL measured the collisions, with collision energies reaching up to
674 209 GeV in the year 2000. In November 2000, LEP was closed down to make way for
675 the construction of the LHC in the same tunnel.

676 In its current configuration, the LHC accelerator complex at CERN is a suc-
677 cession of machines that accelerate particles in stages until they reach their final energy
678 of 6.5 TeV per beam [9] [10]. In Linear accelerator 4 (Linac4), negative hydrogen
679 ions (hydrogen atoms with an additional electron) are accelerated to 160 MeV, and
680 stripped of their two electrons, leaving only protons, before entering the Proton Syn-
681 chrotron Booster (PSB). These protons are accelerated to 2 GeV, then to 26 GeV in

682 the Proton Synchrotron (PS), and 450 GeV in the Super Proton Synchrotron (SPS).
683 The protons are transferred to the two beam pipes of the Large Hadron Collider
684 (LHC). The LHC is a 27-kilometer ring of superconducting magnets, inside which
685 one beam circulates clockwise and the other counterclockwise. Each LHC ring takes
686 4 minutes and 20 seconds to fill, and it takes about 20 minutes for the protons to
687 reach their maximum energy. During normal operating conditions, beams circulate
688 for many hours inside the LHC ring.

689 The beams of particles in the LHC are made to collide at a center-of-mass energy
690 of up to 14 TeV, at four positions at particle detector experiments located around
691 the ring: ATLAS, CMS, ALICE, and LHCb. An aerial view of the four major
692 experiments' locations is shown in Fig. 2.1 [11]. ATLAS and CMS are the two
693 general-purpose detectors with broad physics programmes spanning Standard Model
694 measurements and searches for signatures of new physics [12] [13]. The two experi-
695 ments use different technical solutions and different magnet system designs. ALICE
696 is a general-purpose detector dedicated to measuring LHC heavy-ion collisions, and
697 is designed to address the physics of strongly interacting matter, and the properties
698 of quark-gluon plasma [14]. The LHCb experiment specializes in investigating CP vi-
699 olation through measuring the differences in matter and antimatter, by using a series
700 of subdetectors to detect mainly forward particles close to the beam direction [15].

701 **2.2 Luminosity and pile-up**

702 In order to search for rare processes, such as those resulting from a Higgs, W, or Z
703 boson, a large number of parton interactions per second are required at the LHC.
704 The number of events generated per second by the LHC collisions is given by

$$N_{event} = \mathcal{L} \cdot \sigma_{event} \quad (2.1)$$

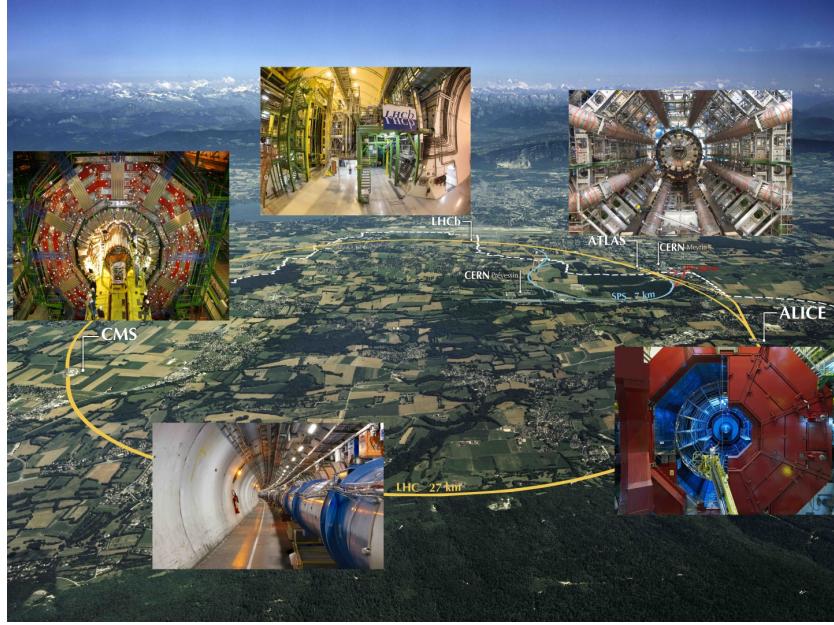


Figure 2.1: Aerial view of the Large Hadron Collider (LHC) spanning the border of France and Switzerland, and the four major experiments located around the ring: CMS (Compact Muon Solenoid), LHCb (LHC beauty), ATLAS (A Toroidal LHC Apparatus), and ALICE (A Large Ion Collider Experiment) [11].

where σ_{event} is the cross-section for the event under study, and \mathcal{L} the instantaneous luminosity. The instantaneous luminosity is measured in units of $\text{cm}^{-2} \text{ s}^{-1}$, and depends only on the beam parameters, and can be written for a Gaussian beam distribution as:

$$\mathcal{L} = \frac{N_b^2 n_b f_{rev} \gamma_r}{4\pi \epsilon_n \beta^*} F \quad (2.2)$$

where the parameters are as defined, along with some example typical nominal values in Phase-1 of the LHC [16] [17]:

- N_b is the number of particles per bunch ($N_b \approx 1.15 \times 10^{11}$ protons per bunch)
- n_b is the number of bunches per beam (maximum 2808),
- f_{rev} is the revolution frequency ($\approx 11 \text{ kHz}$),
- γ_r is the relativistic gamma factor,

- ϵ_n is the normalized transverse beam emittance (area in a transverse plane occupied by the beam particles),
 - β^* is the beta function at the collision point ($\beta^* = 0.55$ m),
 - and F is the geometric luminosity reduction factor due to the crossing angle at the interaction points ($F \approx 0.84$ for Phase-1. Note that complete overlap would give $F = 1$).

721 Peak luminosity at interaction points 1 and 5 reach values of $\sim 1.0 \times 10^{34} \text{ cm}^{-2} \text{ s}^{-1}$,
 722 with peak luminosity per bunch crossing reaching $\sim 3.56 \times 10^{34} \text{ cm}^{-2} \text{ s}^{-1}$.

Per Eqn. 2.1, the integrated luminosity over time is proportional to the number of events produced, and the size of LHC datasets is commonly presented in terms of integrated luminosity. Collider operation aims to optimize the integrated luminosity. Thus the exploration of rare events in the LHC collisions requires both high beam energies and high beam intensities.

The interaction vertex corresponding to the hard scattering of the protons is called the primary interaction vertex, or primary vertex (PV). The LHC’s nominal beam luminosities are sufficiently large for multiple proton-proton collisions to occur in the same time window of 25 nanoseconds in which proton bunches collide [18]. To measure a proton-proton collision, the primary vertices must be separated from overlapping collisions, called “pile-up” collisions.

The pile-up is defined as the average number of pp collisions per bunch crossing, and can be estimated from the inelastic pp cross section of $\sigma_{\text{inel}} = 68.6$ millibarns at a center-of-mass energy of $\sqrt{s} = 13$ TeV [19]:

$$\text{Pile-up} = \frac{\mathcal{L} \times \sigma_{\text{inel}}}{n_b \cdot f} \sim 22 \quad (2.3)$$

⁷³⁷ A distribution of pile-up in the data-taking years 2016-2018 is shown in Fig. 2.2.

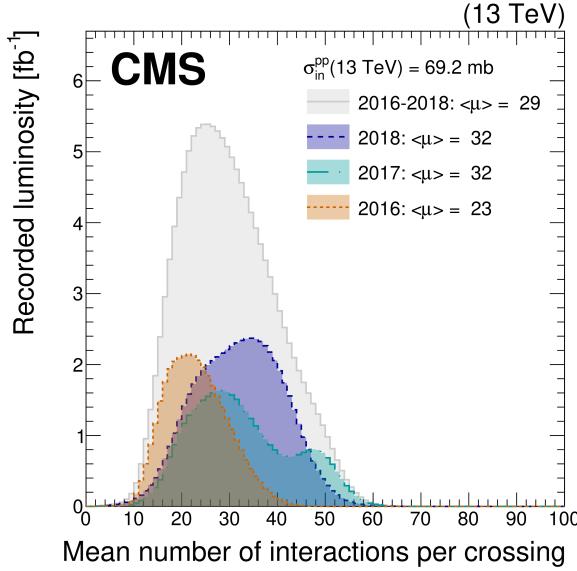


Figure 2.2: Distribution of the mean number of inelastic collisions per bunch crossing (pile-up) in data [18], for proton-proton collisions in 2016 (*dotted orange*), 2017 (*dotted light blue*), 2018 (*dotted dark blue*), and integrated over 2016-2018 (*solid grey*). A cross-section of inelastic proton-proton collisions of 69.2 mbarns is assumed. In the running conditions of the High-Luminosity LHC, pile-up will reach unprecedented levels of up to 200 per bunch crossing [20].

These multiple collisions will lead to higher occupancies in the detector, and particles originating from the pile-up interactions can be confused with those originating from the primary vertex. Thus, higher luminosities create more intense pile-up conditions, posing a greater challenge to detector performance and particle reconstruction and identification.

2.3 The High-Luminosity LHC

The High-Luminosity LHC (HL-LHC) is a major upgrade of the LHC scheduled to take place in the late 2020s, that will increase the instantaneous luminosity by a factor of five beyond the original design value, and the integrated luminosity by a factor of ten [20]. This will be accomplished through accelerator technological advances: for instance, reduction of the interaction point β^* from 0.55 m down to 0.15

749 m by installation of new final-focusing magnets, and improvements in the geometric
750 luminosity loss factor $F \approx 1$ through the installation of crab cavities that optimize
751 the orientation of colliding bunches. A further discussion of the HL-LHC upgrades
752 for the CMS detector follows in Chapter 3.

753 2.4 The CMS Detector

754 We give a brief overview of the Compact Muon Solenoid (CMS) experiment here
755 and discuss each of the subdetectors in more detail in the following sections. The
756 CMS experiment was conceived to study proton-proton and lead-lead collisions at
757 a center-of-mass energy of 14 TeV (5.5 TeV nucleon-nucleon) and at luminosities up
758 to $10^{34} \text{ cm}^{-2} \text{ s}^{-1}$ ($10^{27} \text{ cm}^{-2} \text{ s}^{-1}$) [21] [22]. Starting from the beam interaction region
759 at the center of the CMS detector, particles first pass through a silicon pixel and
760 strip tracker, in which charged-particle trajectories (tracks) and origins (vertices)
761 are reconstructed from signals (hits) in the sensitive layers. The tracker, electro-
762 magnetic calorimeter (ECAL), and hadronic calorimeter (HCAL) are immersed in a
763 high-magnetic-field superconducting solenoid that bends the trajectories of charged
764 particles. After passing through the tracker, electrons and photons are then absorbed
765 in the electromagnetic calorimeter (ECAL) comprised of lead-tungstate scintillating-
766 crystals. The corresponding electromagnetic showers are detected as clusters of energy
767 recording in neighboring cells, from which the direction and energy of the particles can
768 be determined. Charged and neutral hadrons may initiate a hadronic shower in the
769 ECAL as well, which is then fully absorbed in the hadron calorimeter (HCAL). The
770 resulting clusters are used to estimate their direction and energies. Muons and neu-
771 trinos pass through the calorimeters with little to no interactions. Neutrinos escaped
772 undetected; muons produce hits in additional gas-ionization chamber muon detectors
773 housed in the iron yoke of the flux-return. A sketch of example particle interactions

in a transverse slice of the CMS detector is shown in Fig. 2.3. The collision data is recorded with the use of the Level-1 (L1) trigger (discussed in greater detail in 2.5.5), the High-Level Trigger (HLT), and data acquisition systems ensuring high efficiency in selecting physics events of interest.

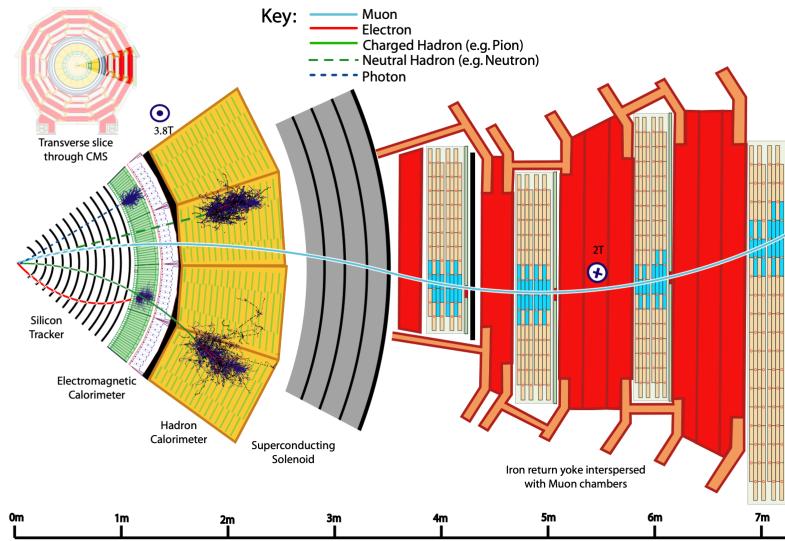


Figure 2.3: Sketch of particle trajectories of muons, electrons, charged and neutral hadrons, and photons in a transverse cross-section of the CMS detector [22].

CMS uses a right-handed coordinate system [21]. The origin is centered at the nominal collision point inside the experiment. The x axis points towards the center of the LHC, and the y axis points vertically upwards. The z axis points along the beam direction. The azimuthal angle, ϕ , is measured from the x axis in the x - y plane, and the radial coordinate in this plane is denoted by r . The polar angle, θ , is measured from the z axis. The pseudorapidity, η , is defined as $\eta = -\ln \tan(\theta/2)$. The momentum and energy transverse to the beam direction, denoted by p_T and E_T respectively, are computed from the x and y components. The momentum imbalance in the transverse plane is called the missing transverse momentum, and its magnitude is denoted by E_T^{miss} .

2.5 Sub-detectors of CMS

This section details the sub-detectors of CMS that operate to identify and precisely measure muons, electrons, photons, and jets over a large energy range.

2.5.1 Inner tracking system

The CMS Tracker performs robust tracking and detailed vertex reconstruction in the 4 T magnetic field of the superconducting solenoidal magnet. The primary sensors used in the tracker are p^+ on n -bulk devices, which allow high voltage operation and are radiation-resistant [23] [24]. The active envelope of the CMS Tracker extends to a radius of 115 cm, over a length of approximately 270 cm on each side of the interaction point [23]. Charged particles in the region $|\eta| \lesssim 1.6$ benefit from the full momentum measurement precision. In this region, a charged particle with p_T of 1000 GeV has a sagitta of $\sim 195 \mu\text{m}$. The Tracker acceptance extends further to $|\eta| = 2.5$, with a reduced radius of approximately 50 cm.

The high magnetic field of CMS causes low p_T charged particles to travel in helical trajectories with small radii. The majority of events contain particles with a steeply falling p_T spectrum, resulting in a track density which rapidly decreases at higher radii.

A schematic view of the current Phase-1 CMS tracker [25], including the pixel detector, is shown in Fig. 2.4. The Phase-1 pixel detector consists of three barrel layers (BPIX) at radii of 4.4 cm, 7.3 cm, and 10.2 cm, and two forward/backward disks (FPIX) at longitudinal positions of ± 34.5 cm and ± 46.5 cm, and extending in radius from about 6 cm to 15 cm. These pixelated detectors produce 3D measurements along the paths of charged particles with single hit resolutions between 10-20 μm .

After the pixel and on their way out of the tracker, particles pass through the silicon strip tracker which reaches out to a radius of 130 cm (Fig. 2.4). The sensor el-

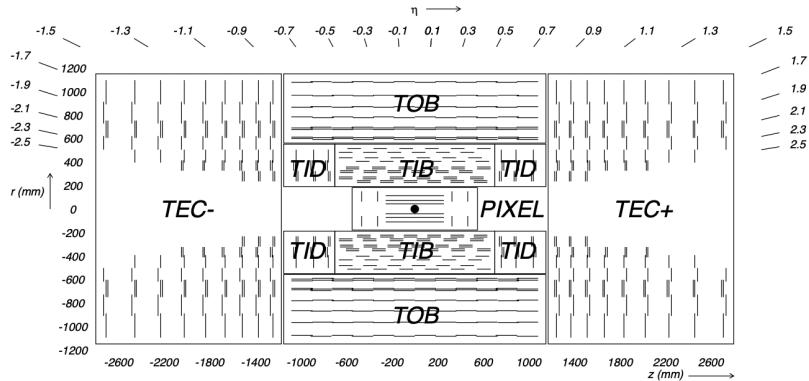


Figure 2.4: Cross section of the current Phase-1 CMS tracker [25]. Each line represents a detector module. Double lines indicate back-to-back modules which deliver two-dimensional (stereo) hits in the strip tracker.

ements in the strip tracker are single-sided p -on- n type silicon micro-strip sensors [21].
 The silicon strip detector consists of four inner barrel (TIB) layers assembled in shells,
 with two inner endcaps (TID), each composed of three small discs. The outer barrel
 (TOB) consists of six concentric layers. Two endcaps (TEC) close off the tracker on
 either end.

2.5.2 ECAL

The electromagnetic calorimeter (ECAL) of CMS measures electromagnetic energy deposits with high granularity. One of the driving criteria in the design was the capability of detecting the Standard Model Higgs boson decay to two photons (in fact, the channel in which the 125 GeV Higgs boson was discovered at CMS). ECAL is a hermetic homogeneous calorimeter comprised of 61,200 lead tungstate (PbWO_4) crystals mounted in the central barrel, with 7,324 crystals in each of the two endcaps [21]. A preshower detector is located in front of the endcap crystals. Avalanche photodiodes (APDs) are used as photodetectors in the barrel and vacuum phototriodes (VPTs) in the endcaps.

The design of the ECAL is driven by the behaviour of high-energy electrons, which

predominantly lose energy in matter via bremsstrahlung, and high-energy photons by e^+e^- pair production. The characteristic amount of matter traversed for these interactions is the radiation length X^0 , usually measured in units of g cm $^{-2}$. The radiation length is also the mean distance over which a high-energy electron loses all but $1/e$ of its energy via bremsstrahlung [26]. Thus high granularity in η and ϕ , and the length of the ECAL crystals, is designed to capture the shower of e/γ produced by electrons and photons.

The barrel part of the ECAL (EB) covers the pseudorapidity range $|\eta| < 1.479$ [21]. The barrel granularity is 360-fold in ϕ and (2×85) -fold in η . The crystal cross-section corresponds to approximately 0.0174×0.0174 in $\eta - \phi$ or 22×22 mm 2 at the front face of the crystal, and 26×26 mm 2 at the rear face. The crystal length is 230 mm, corresponding to $25.8 X_0$.

The ECAL read-out acquires the signals of the photodetectors [21]. At each bunch crossing, digital sums representing the energy deposit in a trigger tower, comprising 5×5 crystals in $\eta \times \phi$, are generated and sent to the Level-1 trigger system (detailed in Section 2.5.5).

2.5.3 HCAL

The hadronic calorimeter (HCAL) of CMS measures hadronic energy, which is key to characterizing the presence of apparent missing transverse energy which could arise from hadron jets and neutrinos or exotic particles [21]. A schematic of the components of HCAL are shown in Fig. 2.5. The HCAL barrel (HB) and endcaps (HE) are located outside of the tracker and the ECAL, spanning a radius of 1.77 m (outer extent of ECAL) up to 2.95 m (inner extent of the magnet coil). An outer hadron calorimeter (HO) is placed outside the solenoid to complement the barrel calorimeter. Beyond $|\eta| = 3$, the forward hadron calorimeter (HF) at 11.2 m from the interaction point extend the pseudorapidity coverage to $|\eta| = 5.2$.

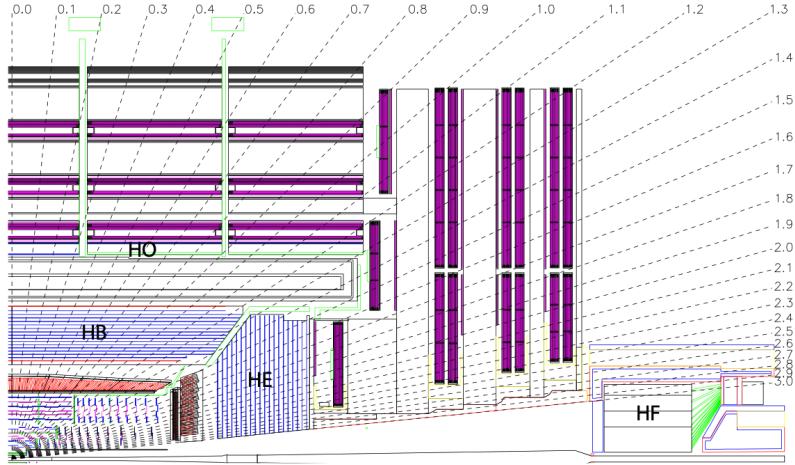


Figure 2.5: Longitudinal view of the CMS detector showing the hadron calorimeter barrel (HB), endcap (HE), outer (HO), and forward (HF) calorimeters from [21].

855 The HB is a sampling calorimeter covering the pseudorapidity range $|\eta| < 1.3$ [21].
 856 It consists of 36 identical azimuthal wedges which form two half-barrels (HB+ and HB-
 857), with a segmentation of $(\Delta\eta, \Delta\phi) = (0.087, 0.087)$. The HE covers pseudorapidity
 858 $1.3 < |\eta| < 3$. The HB and endcap HE calorimeters are sampling calorimeters which
 859 use brass as the absorber and plastic scintillator as the active material. Light from
 860 the plastic scintillator is wavelength-shifted and captured in optic fibers which are
 861 read out by front-end electronics [27].

862 The HF is a Cherenkov calorimeter based on a steel absorber and quartz fibers
 863 which run longitudinally through the absorber and collect Cherenkov light, primarily
 864 from the electromagnetic component of showers developed in the calorimeter [27].
 865 Photomultiplier tubes are used to collect light from the quartz fibers. The HF is
 866 designed to survive in the harsh radiation conditions and high particle flux of the for-
 867 ward region. On average, 760 GeV per proton-proton interaction is deposited into the
 868 two forward calorimeters, compared to only 100 GeV for the rest of the detector [21].
 869 Furthermore, this energy has a pronounced maximum at the highest rapidities.

870 **2.5.4 Muon detectors**

871 The CMS muon system is designed to have the capability of reconstructing the mo-
872 mentum and charge of muons over the kinematic range of the LHC, since muons are a
873 powerful handle on signatures of interesting processes over the high background rate
874 of the LHC [21]. For instance, the decay of the Standard Model Higgs boson into
875 ZZ , which in turn decay to 4 leptons, can be reconstructed with high 4-particle mass
876 resolution if all the leptons are muons, since muons are less affected than electrons
877 by radiative losses in the tracker material.

878 The muon system consists of a cylindrical barrel section and two planar endcap
879 regions [21]. The barrel muon detector consists of drift tube (DT) chambers covering
880 the pseudorapidity region $|\eta| < 1.2$ (Fig. 2.6). The DTs can be used as tracking
881 detectors due to the barrel region's characteristic low neutron-induced backgrounds,
882 low muon rate, and relatively uniform 4T magnetic field contained in the steel yoke.

883 In the two endcap regions, the muon rates and background levels are high and the
884 magnetic field is large and non-uniform [21]. Here, the muon system uses cathode
885 strip chambers (CSCs) to identify muons between $0.9 < |\eta| < 2.4$. The cathode strips
886 of each chamber run radially outwards and provide a precision measurement in the
887 $r - \phi$ bending plane. The anode wires run approximately perpendicular to the strips
888 and are read out in order to measure η and the beam-crossing time of a muon.

889 In addition to the DT and CSC, a dedicated trigger system consisting of resistive
890 plate chambers (RPCs) in the barrel and endcap regions provide a fast, independent,
891 and highly-segmented trigger with a sharp p_T threshold over a large portion of the
892 pseudorapidity range ($|\eta| < 1.6$) of the muon system [21]. RPCs have good time
893 resolution but coarser position resolution compared to the DTs or CSCs. The RPCs
894 also play a role in resolving ambiguities in reconstructing tracks from multiple hits in
895 a chamber.

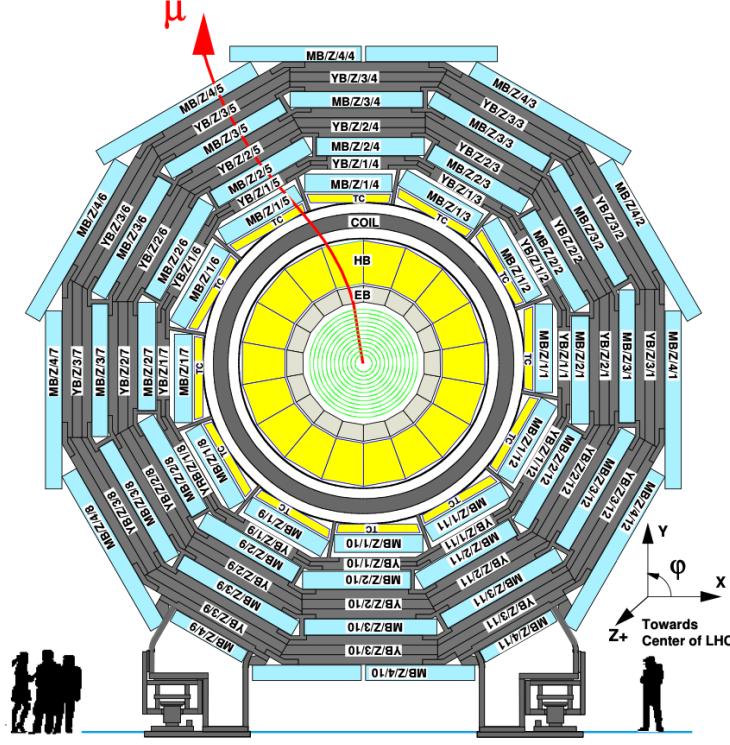


Figure 2.6: Layout of the CMS barrel muon drift tube (DT) chambers in one of the five wheels from [21]. The DTs are organized in 12 sectors of the yoke barrel (YB). In each of the 12 sectors of the yoke, there are 4 muon chambers per wheel (MB1, MB2, MB3, and MB4).

2.5.5 The Level-1 Trigger

The design performance of the LHC corresponds to an instantaneous luminosity of $10^{34} \text{ cm}^{-2} \text{ s}^{-1}$ with a 25 ns bunch crossing rate, giving an average pile-up (number of simultaneous events) of 25 per bunch crossing [28]. However, during Run 2, in 2017 and 2018 the LHC was able to surpass this goal with a mean number of 32 interactions per bunch crossing, and reaching over 50 interactions in short periods (Fig. 2.2). The large number of events from inelastic collisions (minimum bias events) per bunch crossing, combined with the small cross-sections of possible physics discovery signatures, necessitates a sophisticated event selection system for filtering this large event rate, as it is impossible to save all events. This data filtering system is imple-

mented by CMS in two stages. The first stage is the Level-1 (L1) Trigger, which is deployed in custom electronic hardware systems and is responsible for reducing the event rate to around 100 kHz. The second stage is the High-Level Trigger (HLT) which is described in Section 2.5.6. This section describes the Phase-1 configuration of the Level-1 Trigger.

The L1 Trigger data flow of Phase-1 is shown in Fig. 2.7 [28], with organization into the L1 calorimeter trigger, the L1 muon trigger, and the L1 Global Trigger (GT).

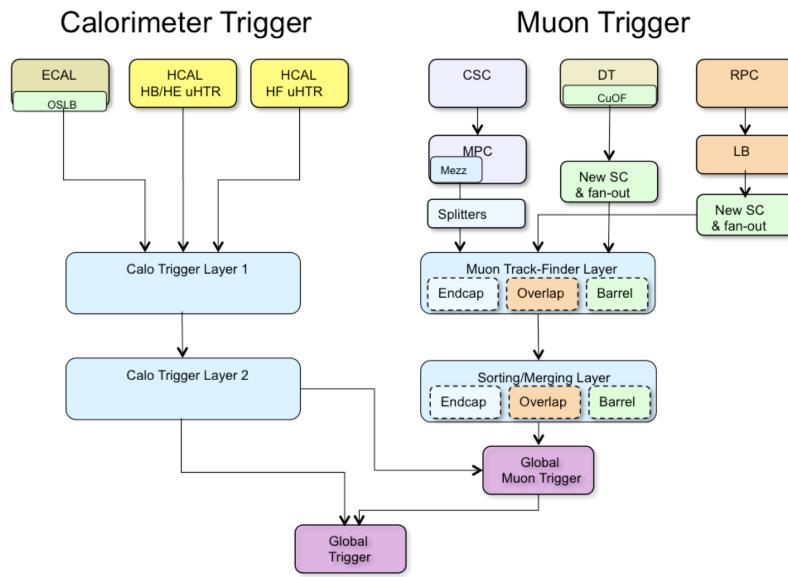


Figure 2.7: Dataflow for the Phase-1 Level-1 Trigger [28], which is implemented in custom hardware and is responsible for reducing the event rate from the LHC bunch crossing frequency of 400 MHz (bunch crossings every 25 ns) to a maximum rate of 100 kHz. In Phase-1, the Level-1 Trigger has access to information from the calorimeter and muon detectors.

The L1 calorimeter trigger begins with trigger tower energy sums formed by the individual calorimeter cell energies. In the original configuration, the ECAL energies were accompanied by a bit indicating the transverse extent of the electromagnetic energy deposits, and the HCAL energies were accompanied by a bit indicating the presence of minimum ionizing energy [29]. During Long Shutdowns 1 and 2 (LS1

919 and LS2), HF was upgraded to provide finer granularity information to the trigger,
920 and the HCAL barrel and endcap front-end electronics were upgraded to provide
921 high-precision timing information and depth segmentation information.

922 In the original design of the L1 calorimeter trigger, the trigger primitives are pro-
923 cessed by the Regional Calorimeter Trigger (RCT, upgraded to Calo Layer 1 after
924 LS2) which finds isolated and non-isolated electron/photon candidates [28]. At this
925 stage, electrons/photons candidates are treated together since they cannot be defini-
926 tively distinguished at this stage due to lack of tracking information in the L1 trigger.
927 The Global Calorimeter Trigger (GCT, upgraded to Calo Layer 2 after LS2) sorts
928 further the candidate electrons/photons, finds jets (classified as central, forward, and
929 tau) using the E_T sums and performs calibration of the clustered jet energies, and
930 calculates global quantities such as missing E_T . It sends the top four candidates of
931 each type to the Global Trigger [28].

932 During LS2 and before Run-2, the legacy calorimeter trigger was upgraded to be
933 more flexible, maintainable, and performant [30] [31] [32]. These upgrades included
934 the replacement of legacy VME-based electronics with the microTCA modern tele-
935 coms standard, and system-wide usage of the latest generation of FPGAs, Xilinx
936 Virtex 7. Parallel copper links were replaced in almost all cases with serial optical
937 links, allowing link speeds to increase from 1 Gb/s to 10 Gb/s [30]. A schematic of
938 the current calorimeter trigger is shown in Fig. 2.8. The calorimeter Layer-1 is imple-
939 mented in 18 Calorimeter Trigger Processor (CTP7) boards, with each card spanning
940 4 out of 72 towers in ϕ and all of η . Tower-level operations are performed in Layer-1,
941 such as the sum of ECAL and HCAL energies, energy calibration, and the compu-
942 tation of the ratio of HCAL to ECAL energies. The Layer-1 cards each transmit 48
943 output links at 10 Gb/s to the nine Layer-2 Master Processor cards (MP7) cards,
944 which host calorimeter algorithms that find particle candidates and compute global
945 energy sums. Each MP7 takes 72 input links and has access to the whole event at

trigger tower granularity, such that the algorithms are fully pipelined and start processing as soon as the minimum amount of data is received. The trigger candidates are sent to a demultiplexer (demux) board, also a MP7, which formats the data for the upgraded Global Trigger, also called the microGT (μ GT).

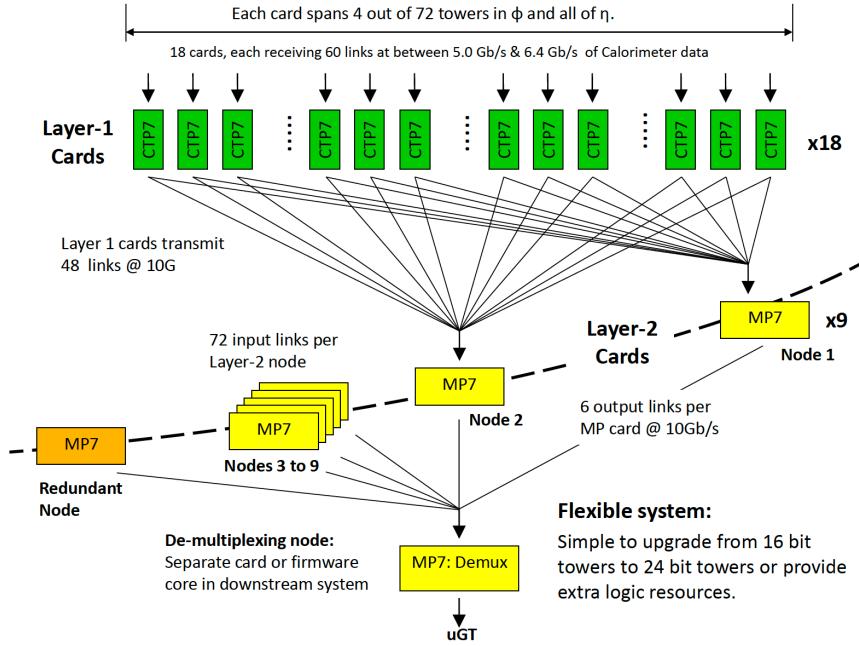


Figure 2.8: Schematic of the calorimeter trigger after Long Shutdown 2 [30]. The Layer-1 calorimeter trigger is implemented in CTP7 cards, which send time-multiplexed outputs to the Layer-2 MP7 cards. The Layer-2 cards handle the data in a round-robin style and the outputs are de-multiplexed, producing one output data stream to the Global Trigger.

Each of the L1 muon triggers has its own trigger logic [29]. The RPC strips are connected to a Pattern Comparator Trigger (PACT), which forms trigger segments that are used to build tracks and calculate p_T . The RPC logic also provides some hit data to the CSC trigger system to resolve ambiguities caused by two muons in the same CSC. The CSCs form local charged tracks (LCTs) from the cathode strips, which are combined with the anode wire information. LCTs are combined into full muon tracks and assigned p_T values.

The Global Muon Trigger (GMT) sorts the RPC, DT, and CSC muon tracks,

958 converts these tracks to the same η , ϕ , and p_T scale, and validates the muon sign [29].
959 It improves the trigger efficiency by merging muon candidates that were detected
960 in two complementary sub-systems (i.e. DT+RPC, or CSC+RPC). The GMT also
961 contains logic to correlate the found muon tracks with an $\eta-\phi$ grid of quiet calorimeter
962 towers to determine if the muons are isolated, as well as logic to remove duplicate
963 candidates originating in the overlap regions from both DT and CSC systems. The
964 final collection of muons are sorted based on their initial quality, correlation, and p_T ,
965 and the top four muons are sent to the Global Trigger [29].

966 The Global Trigger (GT) receives information from the GCT and GMT, and
967 makes the Level-1 Accept (L1A) decision to either discard or accept the bunch cross-
968 ing [29]. This is accomplished by sorting ranked trigger objects that are accompanied
969 by positional information in η and ϕ , permitting the trigger to applying criteria with
970 thresholds that can vary based on the location of the trigger objects, and/or to re-
971 quire trigger objects to be close to or opposite from each other. The GT L1A decision
972 arrives at the detector front end with a $3.8 \mu\text{s}$ latency after the interaction at a rate
973 which is required to be less than 100 kHz, and triggers a full readout of the detector
974 for further processing.

975 **2.5.6 The High-Level Trigger**

976 The HLT is implemented in software running on a large computer farm of fast com-
977 mercial processors [33] [34]. The algorithms in HLT have access to full data from
978 all CMS sub-detectors, including the tracker, with full granularity and resolution.
979 The HLT reconstruction software is similar to what is used offline for full CMS data
980 analysis. As a result, the HLT can calculate quantities with a resolution compara-
981 ble to the final detector resolution, compared to the L1 Trigger. The HLT performs
982 more computationally-intensive algorithms, such as combining tau-jet candidates in
983 the calorimeter with high- p_T stubs in the tracker, to form a hadronic tau trigger. The

984 maximum HLT input rate from the L1 Trigger is 100 kHz, and the HLT output rate
985 is approximately 100 Hz.

986 The HLT contains trigger paths, each corresponding to a dedicated trigger [35].
987 A path consists of several steps implemented as software modules. Each HLT trigger
988 path must be seeded by one or more L1 trigger bits: the first module always looks
989 for a L1 seed, consisting of L1 bit(s) and L1 object(s). Each module performs a well-
990 defined task such as unpacking (raw to digitized quantities), reconstruction of physics
991 objects (electrons, muons, jet, missing transverse energy, etc.), making intermediate
992 decisions that trigger more detailed reconstruction modules, and calculating the final
993 decision for the trigger path. If an intermediate filter decision is negative, the rest of
994 the path is not executed, and the trigger rejects the event.

995 **2.5.7 Particle reconstruction**

996 To build a description of the physics objects present in the particle collision, the
997 basic elements from the detector layers (tracks and clusters of energy) are correlated
998 to identify each particle in the final state. Measurements from different sub-detectors
999 are combined to reconstruct the particle properties. This approach is called particle-
1000 flow (PF) reconstruction [22]. Key to the success of the PF reconstruction is the
1001 fine spatial granularity of the detector layers. Coarse-grained detectors can cause
1002 the signals from different particles to merge, especially within jets. However, if the
1003 subdetectors are sufficiently segmented to separate individual particles, it becomes
1004 possible to produce a global event description that identifies all physics objects with
1005 high efficiencies and resolution.

1006 **2.5.8 Data storage and computational infrastructure**

1007 The LHC generates over 15 petabytes (15 million gigabytes) of data every year, neces-
1008 sitating a flexible computing system that can be accessed by researchers working at

1009 the four main LHC experiments: ALICE, ATLAS, CMS, and LHCb. The Worldwide
1010 LHC Computing Grid (WLCG) [36] is a global collaboration of computer centers that
1011 links thousands of computers and storage systems in over 170 centers across 41 coun-
1012 tries. These centers are arranged in “tiers”, and provide near real-time access to users
1013 processing, analyzing, and storing LHC data. One of the final stages of data analy-
1014 sis at LHC experiments is large-scale data processing taking place over distributing
1015 computing, for instance, with the use of Condor [37], a distributed, scalable, flexible
1016 batch processing system which accepts a computing job, allocates a resource to it,
1017 executes it, and returns the result back to a user transparently.

1018 **Chapter 3**

1019 **The Phase-2 Upgrade of CMS**

1020 This chapter gives an overview of the High-Luminosity LHC upgrade of the LHC in
1021 Section 3.1, and the upgrades for the Phase-2 CMS Level-1 (L1) Trigger in Section
1022 3.2. One of the major upgrades is the new availability of calorimeter crystal-level
1023 information to the L1 calorimeter trigger, compared to the current trigger which only
1024 has access to tower-level information (a tower being 5 by 5 in crystals). To capitalize
1025 on the increased spatial granularity of this information, an upgraded algorithm is
1026 presented which reconstructs and identifies electron and photon candidates in the the
1027 Layer-1 Calorimeter Trigger. A description of the algorithm and a validation of its
1028 performance in Phase-2 conditions is given in Section 3.3.

1029 **3.1 The High-Luminosity LHC**

1030 In order to sustain and extend the LHC’s physics discovery program and maintain
1031 operability for a decade or more, the LHC is undergoing a major upgrade to the High-
1032 Luminosity LHC (HL-LHC). In its final configuration, the HL-LHC will deliver a peak
1033 luminosity of $7.5 \times 10^{34} \text{ cm}^{-2} \text{ s}^{-1}$, potentially leading to total integrated luminosity
1034 of 4000 fb^{-1} after ten years of operations, scheduled to begin in 2027 [38]. This
1035 integrated luminosity is about ten times the predicted luminosity reach of the LHC

1036 in its initial configuration. To enable the CMS experiment to continue operations and
1037 data-taking and to maximize the discovery potential of the unprecedented amount
1038 of data, the CMS detector is undergoing Phase-2 upgrades in order to perform high-
1039 precision measurements and searches for physics beyond the Standard Model in the
1040 intense running conditions of the HL-LHC.

1041 **3.2 The Phase-2 Level-1 Trigger**

1042 To achieve the goals of the HL-LHC program and to ensure the collection of information-
1043 rich datasets in the HL-LHC, the Phase-2 upgrade of the CMS Level-1 Trigger [38]
1044 must be upgraded in conjunction with the CMS sub-detectors and their readouts, to
1045 maintain physics selectivity. The HL-LHC will produce an intense hadronic environ-
1046 ment corresponding to 200 simultaneous collisions per beam crossing, necessitating
1047 comprehensive upgrades of the trigger system outlined below.

1048 In order to cope with the increased pile-up and high occupancies of the HL-LHC,
1049 the latency of the L1 trigger system (time available to produce a L1 Accept signal) will
1050 be increased significantly from $3.8 \mu\text{s}$ to $12.5 \mu\text{s}$, with an increased maximum output
1051 bandwidth of 750 kHz [38]. With the increased latency, in addition to information
1052 from calorimeters and muon detectors (as in the Phase-1 system), information from
1053 the new tracker and high-granularity endcap calorimeter can also be included at L1
1054 for the first time. This is illustrated in the functional diagram of the architecture of
1055 the Phase-2 trigger system in Fig. 3.1.

1056 The key feature of the Phase-2 L1 Trigger is the introduction of a correlator layer,
1057 where algorithms produce higher-level trigger objects by combining information from
1058 sub-detectors, with a selectivity approaching that of offline reconstruction in the
1059 HLT [38]. Four independent data processing paths (grouped together in Fig. 3.1) are
1060 implemented: tracking, calorimetry, muon systems, and particle-flow techniques:

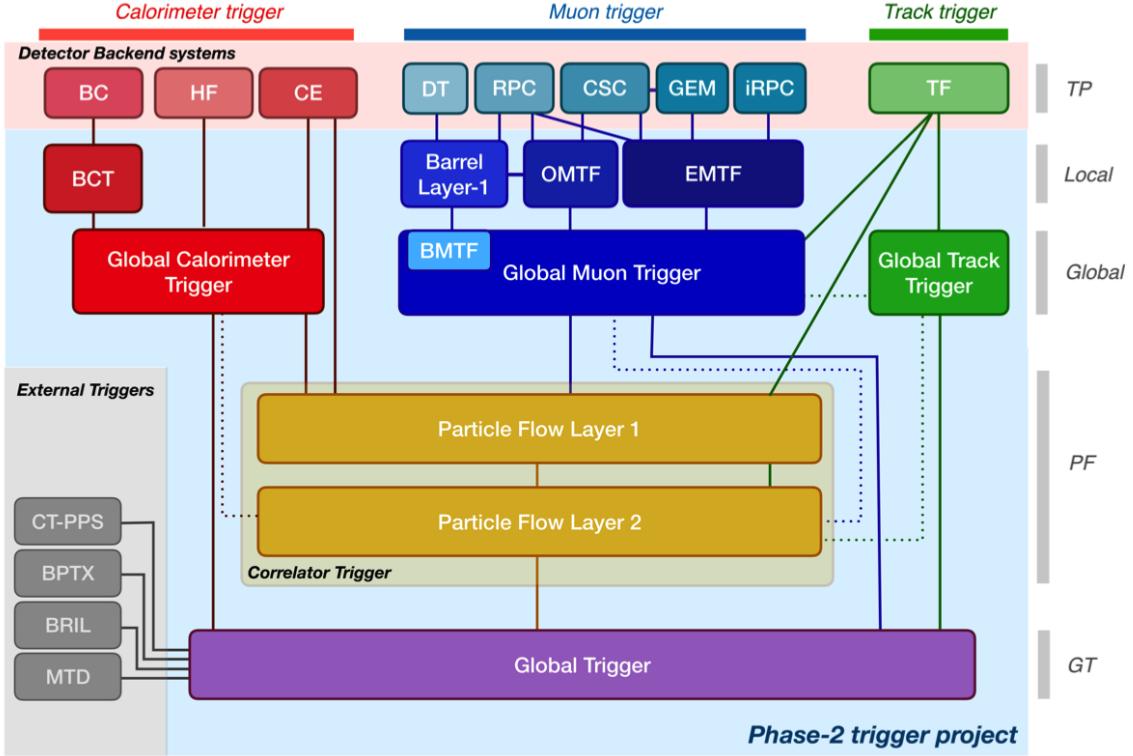


Figure 3.1: Functional diagram of the CMS L1 Phase-2 upgraded trigger design [38], showing the four trigger paths: calorimeter, muon, track, and Particle Flow. For the first time, tracking information will be available as early as the L1 Trigger.

- 1061 • **Calorimeter Trigger path:** (red, Fig. 3.1) A barrel calorimeter trigger (BCT)
 1062 and the HGCAL backend are used to process crystal-level information from the
 1063 calorimeters to produce high-resolution clusters and identification variables used
 1064 for later processing. Outputs from the BCT, HGCAL, and the HF are sent to
 1065 a global calorimeter trigger (GCT), where calorimeter-only objects such as e/γ
 1066 candidates, hadronically decaying tau lepton candidates, jets, and energy sums
 1067 are built.
- 1068 • **Track Trigger path:** (green, Fig. 3.1) Tracks from the Outer Tracker are
 1069 reconstructed in the track finder (TF) processors as part of the detector back-
 1070 end. A global track trigger (GTT) will reconstruct the primary vertices of the
 1071 event, along with tracker-only based objects, such as jets and missing transverse
 1072 momentum.

- **Muon Trigger path:** (*blue*, Fig. 3.1) Trigger primitives are processed by muon track finder algorithms, again separated into the barrel (barrel muon track finder, BMTF), overlap (overlap muon track finder, OMTF), and endcap (endcap muon track finder, EMTF). Standalone muons and stubs containing information such as position, bend angle, and timing, as well as L1 tracks, are sent to the global muon trigger (GMT).
- **Particle-Flow Trigger path:** (*yellow*, Fig. 3.1) The correlator trigger (CT) aims to approach the performance of offline Particle Flow, and is implemented in two layers. “Layer-1” produces the particle-flow candidates from matching calorimeter clusters and tracks. “Layer 2” builds and sorts final trigger objects and applies additional identification and isolation criteria.

The outputs from the above trigger paths are combined in the Global Trigger (GT) (*purple*, Fig. 3.1), which calculates the final trigger decision (Level-1 Accept), transmitting it to the Trigger Control and Distribution System (TCDS), which distributes it to the detector backend systems, initiating the readout to the DAQ. The GT also provides the interface to external triggers (*grey*, Fig. 3.1), such as triggers for the precision proton spectrometer (PPS), beam position and timing monitors (BPTX), and luminosity and beam monitoring (BRIL) detectors [38]. The design of the Phase-2 Level-1 Trigger allows for future inclusion of triggering information, for instance information about minimum ionizing particles (MIPs) from the MIP Timing Detector (MTD) [39].

3.3 Standalone Barrel Calorimeter electron/photon reconstruction

The reconstruction and identification of electrons and photons (e/γ) begin with the trigger primitives of the barrel ECAL and HCAL detectors and endcap HGCAL calorimeters, covering the pseudorapidity region $|\eta| < 3$. The barrel and endcap regions of the detector are intrinsically different enough to warrant different approaches to e/γ reconstruction. This work presents a firmware-based emulator for the standalone e/γ reconstruction in the barrel calorimeter (Fig. 3.2). “Standalone” refers to the fact that the tracker information is not used in this particular reconstruction chain. This firmware-based emulator is based on the parallelized, computational logic that will be deployed in the firmware of the Phase-2 Level-1 trigger. The emulator uses fixed-precision integers to represent all values, such as in the computation of cluster energies, and closely mimics the firmware logic which uses arrays and performs computations in flattened loops. It represents an improved, more realistic understanding of the trigger, compared to the previous emulator which used idealized logic such as vector operations, and floats to represent all values [38].

3.3.1 Electron/photon standalone barrel procedure

In Phase-2, the upgrade of both on-detector and off-detector electronics of the barrel calorimeters’ trigger primitive generator (TPG) will enable the streaming of single crystal data from the on-detector to the backend electronics. Currently in Phase-1, the ECAL and HCAL TPGs is restricted to providing lower-granularity information of trigger tower sums of 5×5 crystals to the Level-1 Trigger [38]. A schematic of the geometry of the ECAL barrel in the Phase-2 Regional Calorimeter Trigger (RCT) is shown in Fig. 3.3. The barrel is spanned by 36 RCT cards, each spanning 17×4 towers in $\eta \times \phi$. Each RCT card is subdivided into five “regions” as shown in Fig.

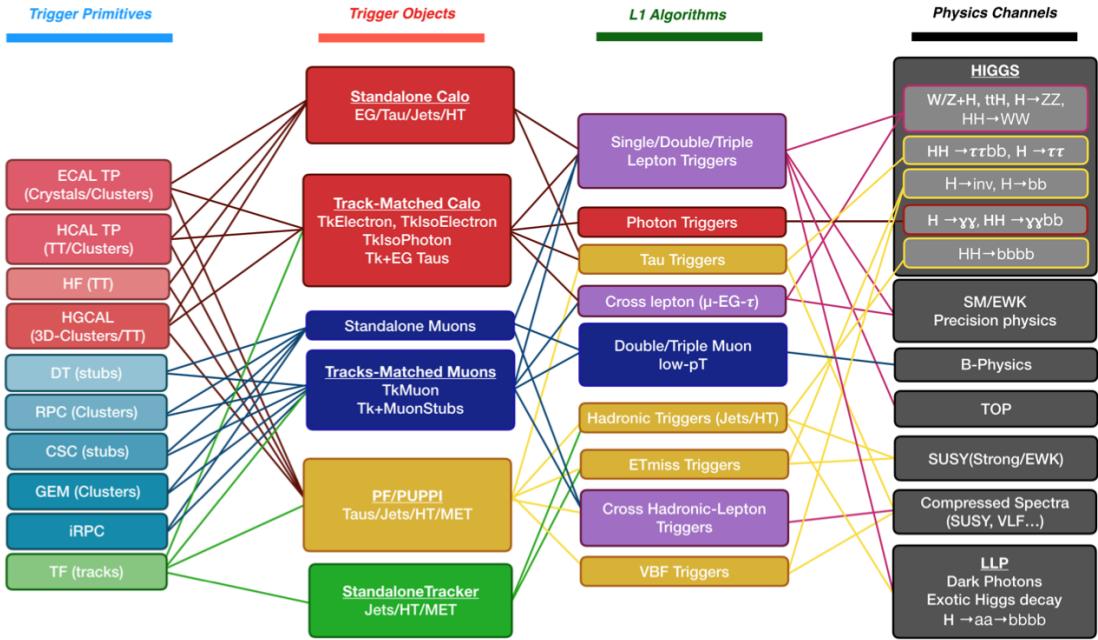


Figure 3.2: Summary of the links between the trigger primitives (*first column*), the trigger objects (*second column*), the Level-1 algorithms used in the menu (*3rd column*), and the physics channels (*4th column*), from [38], where a full description of the Phase-2 L1 algorithms can be found. This work focuses on developments for the Standalone Calorimeter electron and photon ("EG") reconstruction algorithm.

1119 3.4. After initial clustering and processing, the outputs of the RCT card are sent to
 1120 the Global Calorimeter (GCT) trigger, which is processed in three cards as shown in
 1121 Fig. 3.5. The reconstruction algorithm is detailed below.

1122 The standalone barrel algorithm for reconstructing and identifying electrons and
 1123 photons in the Phase-2 Level-1 Trigger takes as input the digitized response of each
 1124 crystal of the barrel ECAL, with a granularity 0.0175×0.0175 in $\eta \times \phi$, which is 25
 1125 times higher than the input to the Phase-1 trigger, which consisted of trigger towers
 1126 with a granularity of 0.0875×0.0875 . In HCAL the tower size of 0.0875×0.0875
 1127 is unchanged. The trigger algorithm is designed to closely reproduce the algorithm
 1128 used in the offline reconstruction, with limitations and simplifications due to trigger
 1129 latency.

1130 In the RCT, an initial requirement of $p_T > 0.5$ GeV is imposed on the input

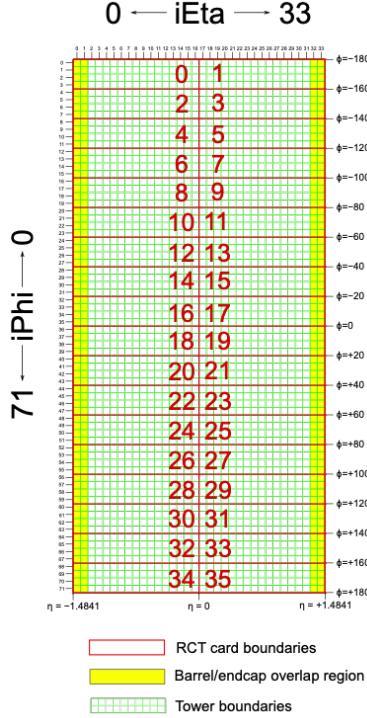


Figure 3.3: Schematic of the geometry of the Phase-2 ECAL barrel in the Regional Calorimeter Trigger (RCT), showing the division of the barrel region into 36 Regional Calorimeter Trigger (RCT) cards (*red*). Each card spans 17×4 towers in $\eta \times \phi$ (*green*), and each tower is 5×5 in single crystals in $\eta \times \phi$. Towers in the overlap region (*shaded yellow*) are read out to both the barrel and endcap.

trigger primitives (i.e. energies from the ECAL crystals and HCAL towers) to reject contribution from pile-up. In one of the regions inside a RCT card (Fig. 3.4), the crystal containing the highest energy deposit is identified as the seed crystal, as shown in Fig. 3.6. The energy in the crystals in a window of size 3×5 in $\eta \times \phi$ around the seed cluster is added into a cluster. The energy is considered “clustered”. The process is repeated with the remaining “unclustered” energy, until up to four clusters are produced in the region.

To improve e/γ identification and to reduce background contributions, identification and reconstruction algorithms are implemented at this stage:

- Shower shape: The energy deposit sums around the seed crystal is computed in windows of size 2×5 and 5×5 (Fig. 3.6, *dashed lines*), with true e/γ clusters

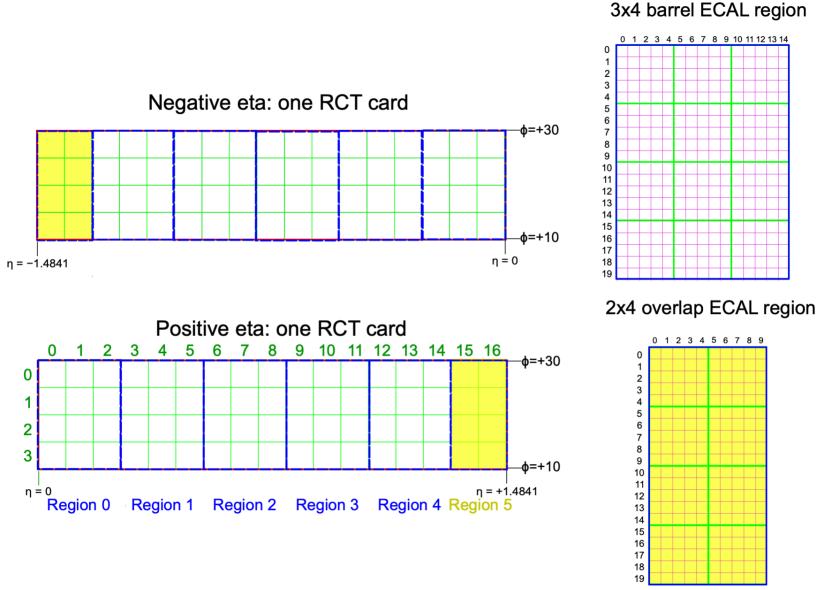


Figure 3.4: Schematic of two example RCT cards in the negative eta (*top left*) and positive eta (*bottom left*) regions of the ECAL barrel. Each RCT card is divided into six regions: five regions are of size 3×4 towers in $\eta \times \phi$ (*top right*), and a sixth smaller overlap region of size 2×4 towers (*bottom right*). Each tower is 5×5 ($\eta \times \phi$) in crystals.

1142 tending to produce showers that deposit most of their energy in a 2×5 region.

- 1143 • Bremsstrahlung recovery: e/γ tend to spread in the ϕ direction due to charged
 1144 particles being bent by the magnetic field of the CMS solenoid. If sufficient
 1145 energy comparable to the core 3×5 cluster is found in the adjacent 3×5
 1146 windows (Fig. 3.6, *shaded yellow*), the energy is added to the core cluster and
 1147 no longer considered unclustered energy.

1148 After parallel processing in the regions, the clusters in a RCT card are stitched
 1149 together if they are located directly along the borders of a region (Fig. 3.3). The
 1150 remaining unclustered ECAL energy is summed into ECAL towers.

1151 From each RCT card, the twelve highest-energy clusters, as well as any remaining
 1152 unclustered energy, are sent to the GCT. Since each GCT card has information from
 1153 sixteen RCT cards (Fig. 3.5), final stitching across the boundaries of the RCT cards
 1154 is performed. One more identification algorithm is performed at this stage:

GCT/“Layer 2”

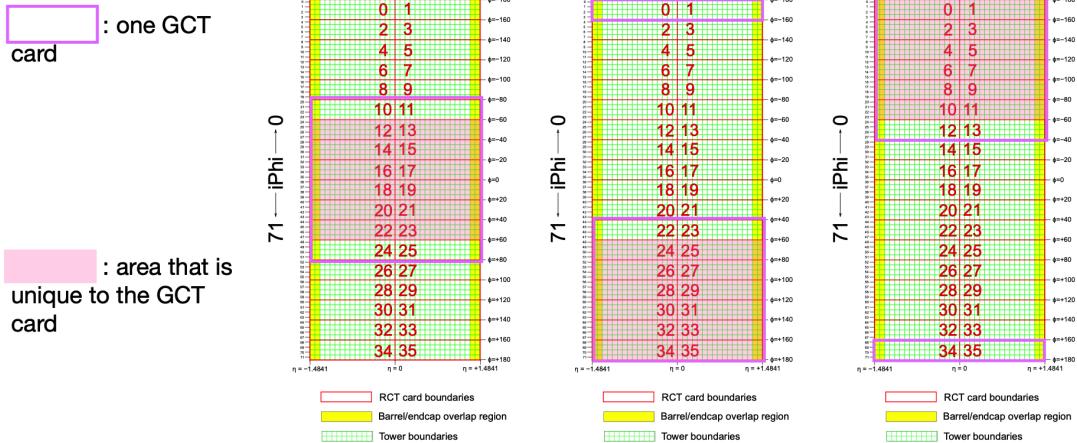


Figure 3.5: Schematic of the Phase-2 ECAL barrel in the Global Calorimeter Trigger (GCT), which will process the outputs of the Regional Calorimeter Trigger (RCT) in three GCT cards (*purple borders*). Each card in the GCT processes the equivalent of sixteen RCT cards, with the center twelve RCT cards being unique to that GCT card (*shaded pink*), and the remaining four RCT cards overlapping with one other GCT card.

- Isolation: One handle to reject backgrounds from e.g. pile-up, comes from the tendency for background to be spread more uniformly across a large area in the detector, whereas genuine e/γ are expected to produce showers concentrated in the 3×5 crystal window. The energy sum in a large window of 7×7 in towers is computed and used to reject background.

Flags that provide discrimination power between genuine e/γ and background, are computed using the relative isolation and shower shape quantities. The standalone working point (WP) is defined as the logical OR of the relative isolation and shower shape flags.

The information of the clusters in the event, including their energies, crystal-level position, the relative isolation flag, the shower shape flag, the standalone WP, and the ratio of the HCAL over ECAL energies, are sent in 64 bits to the Correlator Trigger and the Global Trigger. The towers in the event are computed as the sum of all unclustered energy in the ECAL with the corresponding HCAL energy at each

3x4 barrel ECAL region

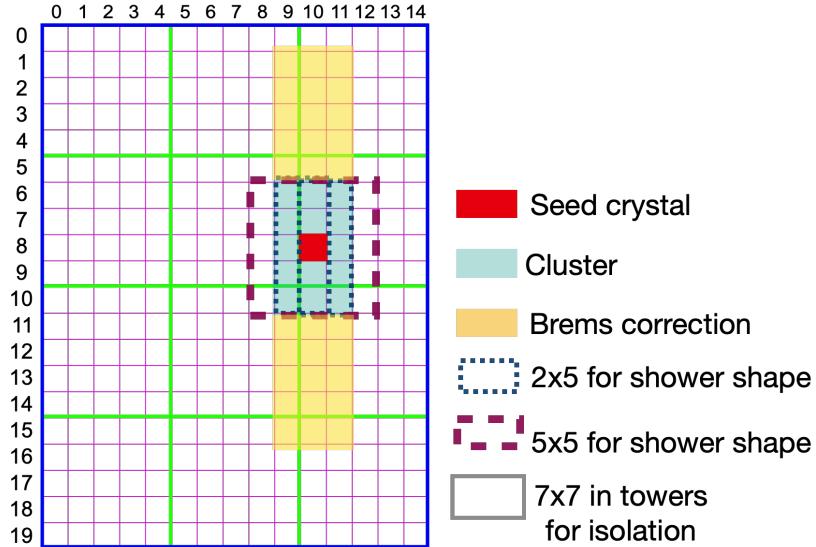


Figure 3.6: Illustration of an example electron/photon (e/γ) cluster in the Phase-2 Level-1 Trigger standalone barrel e/γ reconstruction, in a region of 15×20 crystals (3×4 towers) in $\eta \times \phi$. Each small pink square is one crystal, the highest-granularity ECAL trigger primitives available to the L1 Trigger in Phase-2. The core cluster consists of the energy sum in a 3×5 window of crystals (*shaded light blue*), centered around the seed crystal (*red*). The presence of energy lost to bremsstrahlung radiation is checked in the adjacent 3×5 windows in the ϕ direction (*shaded light yellow*). The ratio of the total energies in windows of size 2×5 and 5×5 in crystals (*dashed dark blue and dark red*) around the seed crystal, is computed and compared to the core cluster energy to obtain shower shape flags. Lastly, the isolation, defined as the sum of the energy in a large window of size 7×7 in towers (not shown in figure) is computed, and compared to the core cluster energy to obtain isolation flags.

¹¹⁶⁹ tower location, and their energies are sent to the Correlator Trigger.

¹¹⁷⁰ 3.3.2 Electron/photon standalone barrel results

¹¹⁷¹ The performance of the current emulator of the standalone barrel e/γ algorithm in
¹¹⁷² Phase-2 conditions is quantified in efficiencies and rates. Efficiency is the fraction of
¹¹⁷³ true electrons that the algorithm can reconstruct and identify, and is evaluated in
¹¹⁷⁴ a Monte Carlo simulated sample containing electrons with transverse momentum p_T
¹¹⁷⁵ ranging from 1 to 100 GeV. The efficiencies of the current and previous emulators as

1176 a function of the electron generator-level p_T are shown in Fig. 3.7.

1177 The rates are the event rates that this reconstruction and identification algorithm
1178 would obtain if it were deployed in a trigger, assuming that proton-proton collisions
1179 are occurring at the 40 MHz event rate of the HL-LHC. The rate is reported as a
1180 function of the minimum energy threshold required by the trigger, and is estimated
1181 using a simulated sample of minimum bias events, i.e. generic proton-proton colli-
1182 sions without any specific physics selections. The rates for the current and previous
1183 emulator are shown in Fig. 3.8.

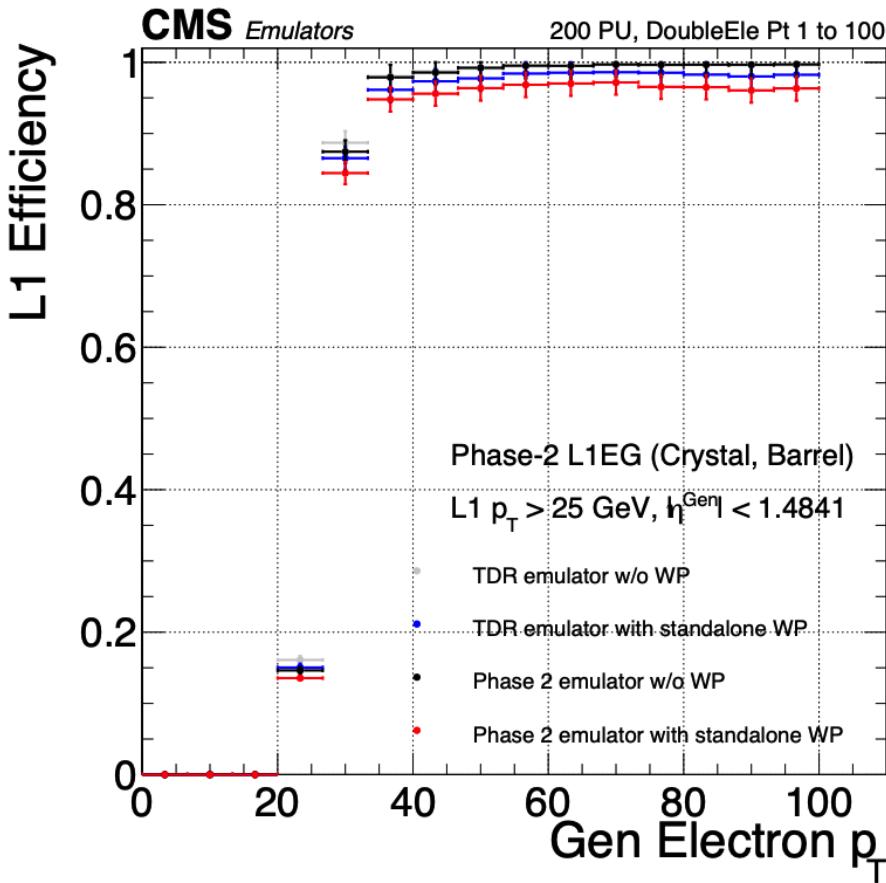


Figure 3.7: Efficiencies of the current and previous emulators of the standalone barrel e/γ algorithm for the Phase-2 Level-1 Trigger, evaluated in a simulated sample containing electrons, as a function of the electron's generator-level transverse momentum p_T . The standalone working point (WP) is defined as the logical OR of the isolation flag and shower shape flag. The efficiencies with and without requiring the standalone WP, are shown for the current emulator (labeled “Phase 2”, *black, red*) and the previous emulator (labeled “TDR”, *dark blue, grey*).

1184 The current emulator is incorporated into the full Phase-2 L1 menu, allowing an
1185 estimate of the rates produced by the standalone e/γ barrel trigger path and all
1186 other algorithms in the L1 Trigger. All rates are estimated with the assumption of
1187 an average pile-up of 200 and event rate of 40 MHz. The standalone working point
1188 single e/γ path with requirements on the e/γ candidate to have $|\eta| < 2.4$, offline p_T
1189 to be greater than 51 GeV, and online p_T to be greater than 41 GeV, is projected to
1190 have a rate of around 23 kHz. The standalone working point double e/γ path with
1191 requirements on the two e/γ candidates to have $|\eta| < 2.4$, offline p_T greater than 37
1192 and 24 GeV, and online p_T greater than 29 and 18 GeV, is projected to give a rate of
1193 around 6 kHz. For both paths, the objects efficiency plateau is 99%.

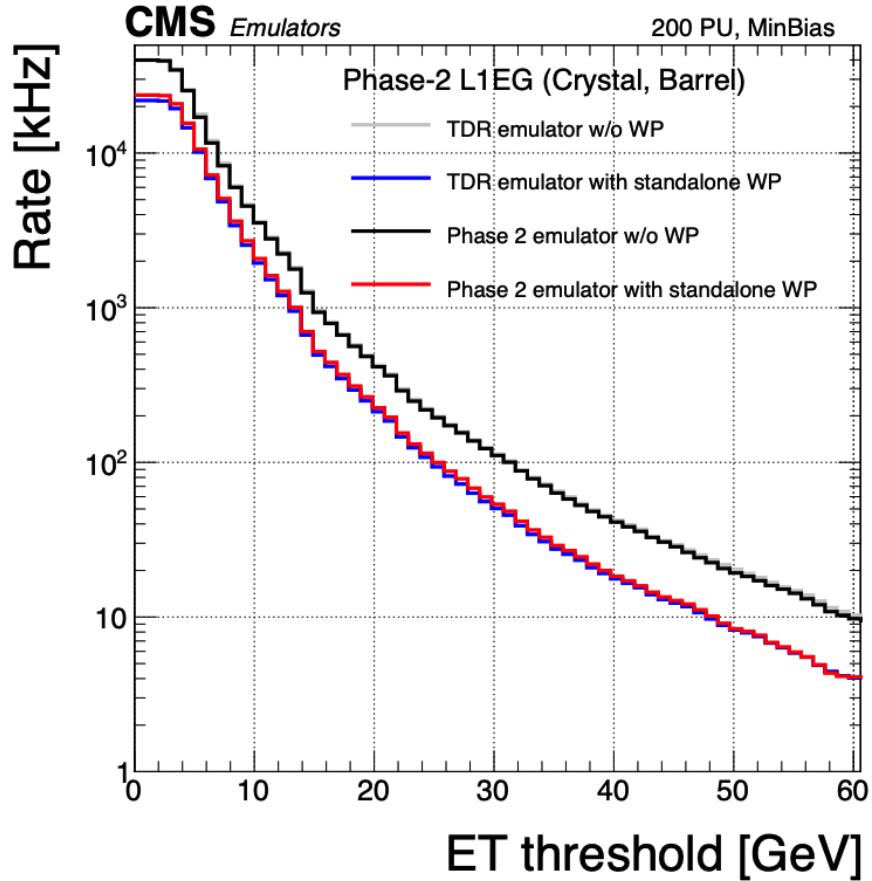


Figure 3.8: Rates in kHz of the current Phase-2 and previous (“TDR”) emulators of the standalone barrel e/γ algorithm for the Phase-2 Level-1 Trigger, evaluated on a minimum bias (MinBias) sample with 200 pile-up (PU), measured as a function of the minimum energy (E_T) required of the reconstructed e/γ object in each event. The standalone working point (standalone WP) is defined to be the logical OR of the isolation flag and the shower shape flag. The rates with and without requiring the standalone WP, are shown for the current emulator (labeled “Phase 2”, *black, red*) and the previous emulator (labeled “TDR”, *dark blue, grey*).

1194 **Chapter 4**

1195 **Datasets and Monte Carlo samples**

1196 The search for the exotic decay of the 125 GeV Higgs boson to two light neutral scalars
1197 decaying to a pair of bottom quarks and a pair of tau leptons ($h \rightarrow aa \rightarrow bb\tau\tau$) is
1198 based on proton-proton collision data at a center-of-mass energy 13 TeV collected
1199 in Run-2 of data-taking, spanning the data-taking years 2016, 2017, and 2018. The
1200 datasets used and the triggers used to collect the data are described in Section 4.1.
1201 Section 4.2 describes the Monte Carlo simulated samples that are used to model the
1202 $h \rightarrow aa \rightarrow bb\tau\tau$ signal and background Standard Model processes. Lastly, in order
1203 to obtain a better description of Standard Model backgrounds that contain two tau
1204 leptons, a data-Monte Carlo hybrid technique is used to generate embedded samples
1205 which model processes with genuine $\tau\tau$ in the final state, as detailed in Section 4.3.
1206 All samples are listed in Appendix A.

1207 **4.1 Datasets used**

1208 The $h \rightarrow aa \rightarrow bb\tau\tau$ analysis [40] is based on proton-proton collision data at a center-
1209 of-mass energy of 13 TeV collected in full Run-2 (2016-18) with the CMS detector.
1210 The data analyzed corresponds to a total integrated luminosity of 138 fb^{-1} (36.33 fb^{-1}
1211 for 2016, 41.53 fb^{-1} for 2017, and 59.74 fb^{-1} for 2018) [41] [42] [43]. The cumulative

1212 delivered and recorded luminosity versus time for 2015-2018 is shown in Fig. 4.1.

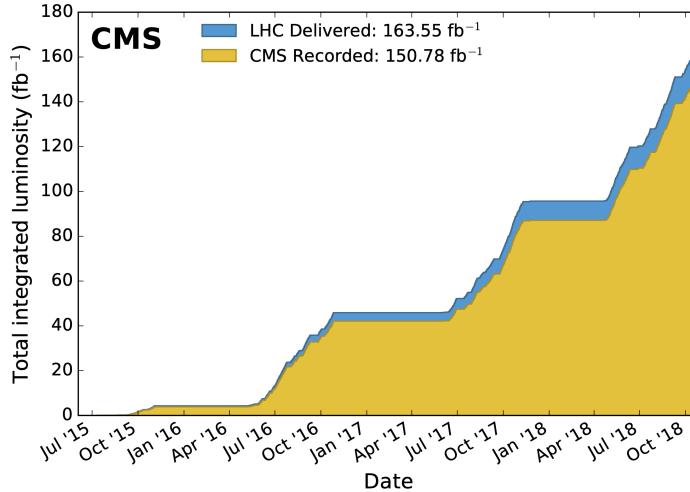


Figure 4.1: Cumulative delivered and recorded luminosity versus time for 2015-2018 at CMS, in proton-proton collision data only, at nominal center-of-mass energy [44].

1213 Data collected with the single muon trigger is used for the $\mu\tau_h$ channel. For the
1214 $e\tau_h$ channel, data collected with the single electron trigger is used; and for the $e\mu$
1215 channel, data collected with the electron + muon trigger is used. A more in-depth
1216 discussion of the triggers used follows in a later section. The datasets are listed in
1217 Appendix A in Tables A.1, A.2, and A.3.

1218 4.2 Monte Carlo samples

1219 Modeling and computing observables originating from arbitrary physics processes at
1220 the tree level and at next-to-leading order (NLO) is performed by Monte Carlo (MC)
1221 event generators, such as Powheg and MadGraph5_amCNLO [45] [46]. The informa-
1222 tion generated, e.g. the computation of the differential cross sections and kinematics
1223 of the final state particles, is saved in a compressed file and used to generate MC sam-
1224 ples that are used in physics analyses. The samples are digitized using GEANT4 [47],
1225 a platform used at the LHC and other facilities to comprehensively simulate the

passage of particles through matter. The digitized samples are passed through the same detector reconstruction as real data events collected in the detector. The MC background samples used in this analysis for 2016-2018 are listed in Appendix A in Tables A.7, A.8, and A.9.

The samples for modeling the signal ($h \rightarrow aa \rightarrow 2b2\tau$ and $h \rightarrow a_1a_2$) in the 2HDM+S and TRSM are generated at tree-level, for a range of masses of the light neutral scalar a . For $h \rightarrow aa$, the mass hypotheses for the a range from $m_a = (12\text{ GeV}, 62.5\text{ GeV})$. For $h \rightarrow a_1a_2$, the mass hypotheses for the two light scalars span combinations of m_{a1} , m_{a2} ranging from $(12\text{ GeV}, 62.5\text{ GeV})$ for the two scalars. The MC signal samples used in this analysis for 2016-2018 are listed in Appendix A in Tables A.10, A.11 and A.12.

4.3 Embedded samples

An important background for Higgs boson studies and searches for additional Higgs bosons is the decay of Z bosons into pairs of τ leptons ($Z \rightarrow \tau\tau$). An embedded technique was developed in the context of Standard Model Higgs to $\tau\tau$ measurements, to model $Z \rightarrow \tau\tau$ decays, and was expanded to also model all Standard Model processes that contain $\tau\tau$ [48]. The embedded technique has since been used successfully at CMS for the Standard Model $H \rightarrow \tau\tau$ measurement, as well as searches for minimal supersymmetric extensions to the Standard Model (MSSM) [49] [50].

The advantage of the embedded technique is that aspects of the event that are difficult to model and describe are directly taken from data, resulting in a better data description than can be achieved with only the $Z \rightarrow \tau\tau$ simulation [48]. The simulation must be tuned extensively to accurately model aspects of the data, such as time-dependent pile-up profiles, the production of additional jets, e.g. in multijet and vector boson fusion topologies, the number of reconstructed primary interaction

1251 vertices, and the missing transverse momentum p_T^{miss} . Since all events with genuine
 1252 $\tau\tau$ are estimated with samples made with the embedded technique (referred to as
 1253 embedded samples from here on), events in Monte Carlo simulation with genuine $\tau\tau$
 1254 are not used, in order to avoid double-counting.

1255 Fig. 4.2 shows a schematic of how embedded samples are produced. Data events
 1256 containing $Z \rightarrow \mu\mu$ decays are selected. In these events, all energy deposits of the
 1257 recorded muons are removed, and are replaced with simulated tau leptons with the
 1258 same kinematic properties as the removed muons. This results in a hybrid data format
 1259 containing information from both observed and simulated events, as illustrated in Fig.
 1260 4.2 [48]. The embedded samples used for the years 2016-2018 are listed in Appendix
 1261 A in Tables A.4, A.5, and A.6.

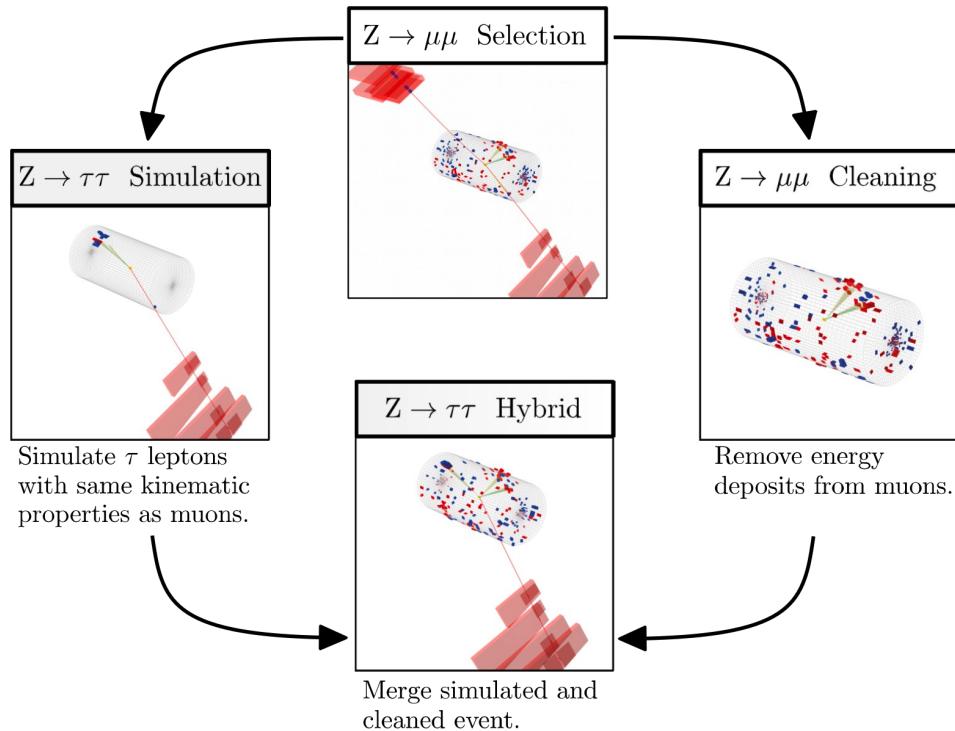


Figure 4.2: Schematic view of the four main steps of the embedding technique for τ leptons, as described in Section 4.3 [48]. A $Z \rightarrow \mu\mu$ event is selected in data ($Z \rightarrow \mu\mu$ selection), all of the energy deposits associated with the muons are removed ($Z \rightarrow \mu\mu$ cleaning), and two τ leptons and their decays are simulated in an empty detector ($Z \rightarrow \tau\tau$ simulation). Lastly, all energy deposits of the simulated τ decays are combined with the data event ($Z \rightarrow \tau\tau$ hybrid).

1262 In the selection step of the embedded technique, events are selected with one of a
 1263 set of $\mu\mu$ trigger paths, which require $p_T > 17(8)$ GeV for the leading (sub-leading)
 1264 muons, and a minimum requirement between 3.8 and 8.0 GeV on the invariant di-
 1265 muon mass $m_{\mu\mu}$ [48]. The offline reconstructed muons must match the objects at
 1266 trigger level and also have offline $p_T > 17(8)$ GeV. They must have $|\eta| < 2.4$ and be
 1267 located at a distance $|d_z| < 0.2$ cm to the primary vertex along the beam axis. To form
 1268 a Z boson candidate, each muon is required to originate from a global muon track.
 1269 The muon pairs must have opposite charges with an invariant mass of $m_{\mu\mu} > 20$ GeV.
 1270 If more than two di-muon pairs are found, the pair with the invariant mass closest to
 1271 the Z boson mass (91.19 GeV) is chosen.

1272 This selection is designed to be tight enough to ensure a high purity of genuine
 1273 $\mu\mu$ events, and also loose enough to minimize biases of the embedded event samples.
 1274 Isolation requirements are avoided, since they would introduce a bias towards less
 1275 hadronic activity in the vicinities of the embedded leptons that will appear more
 1276 isolated than expected in data. The selection results in an expected mixture of events
 1277 summarized in Table 4.1 from [48]. $Z \rightarrow \mu\mu$ is the dominant process modeled by the
 1278 embedded technique, with $t\bar{t}$, QCD, and diboson and single top processes becoming
 1279 more significant when considering events with b-tag jets.

Process	Fraction (%)		
	Inclusive	$m_{\mu\mu} > 70$ GeV	$N(\text{b-tag jets}) > 0$
$Z \rightarrow \mu\mu$	97.36	99.11	69.25
QCD	0.84	0.10	2.08
$t\bar{t}$	0.78	0.55	25.61
$Z \rightarrow \tau\tau$	0.71	0.05	0.57
Diboson, single t	0.17	0.17	2.35
W+jets	0.08	0.02	0.14

Table 4.1: Expected event composition after selecting two muons in the embedded technique [48], before additional cuts (i.e. inclusive, *column 2*), and after adding a requirement on the di-muon mass $m_{\mu\mu} > 70$ GeV (*column 3*), or a requirement on the number of b-tag jets in the event (*column 4*).

₁₂₈₀ **Chapter 5**

₁₂₈₁ **Object reconstruction and**
₁₂₈₂ **corrections applied**

₁₂₈₃ In the data processing workflow, data events and simulated events are analyzed to
₁₂₈₄ reconstruct physics objects of interest, and algorithms for distinguishing genuine par-
₁₂₈₅ ticle candidates from background, are employed. Section 5.1 describes the physical
₁₂₈₆ properties of the most important objects in the $h \rightarrow aa \rightarrow bb\tau\tau$ analysis: taus,
₁₂₈₇ muons, electrons, jets, and jets originating from b-quarks (b-flavor jets), as well as
₁₂₈₈ their reconstruction and identification in CMS. In this analysis, the full energy and
₁₂₈₉ momentum of the two tau leptons ($m_{\tau\tau}$) is estimated from the measured (i.e. visible)
₁₂₉₀ components of the tau leptons using the SVFit/FastMTT algorithm, which is de-
₁₂₉₁ scribed in Section 5.2. Corrections are applied to the simulated samples at the object
₁₂₉₂ level and the event level to account for known discrepancies between simulations and
₁₂₉₃ the data that the simulations are intended to model. These corrections are listed and
₁₂₉₄ detailed in Section 5.3.

₁₂₉₅ **5.1 Object reconstruction**

₁₂₉₆ **5.1.1 Taus**

₁₂₉₇ The tau (τ) is the heaviest known lepton. With a rest mass of 1776.86 MeV, it can
₁₂₉₈ decay to not only electrons and muons, but also hadrons. In two thirds of the cases, τ
₁₂₉₉ leptons decay hadronically, typically into one or three charged mesons (predominantly
₁₃₀₀ π^+ , π^-), often accompanied by neutral pions (that decay $\pi^0 \rightarrow \gamma\gamma$), and a ν_τ . These
₁₃₀₁ hadronic decays are denoted τ_h . In the remainder of the decays, the tau decays to
₁₃₀₂ the lighter leptons (electron or muon), termed leptonic decays. The mean lifetime of
₁₃₀₃ the τ is $\tau = 290 \times 10^{-15}$ seconds, corresponding to $c\tau = 87.03 \mu\text{m}$, which is short
₁₃₀₄ enough that taus decay in the CMS detector before reaching the detector elements,
₁₃₀₅ but also long enough that some decay length variables can help with hadronic tau
₁₃₀₆ identification. The tau's largest decay branching ratios (proportional to probability
₁₃₀₇ of decay) are listed below [26]:

- ₁₃₀₈ • 17.8% decay to $e^- \bar{\nu}_e \nu_\tau$
- ₁₃₀₉ • 17.4% decay to $\mu^- \bar{\nu}_\mu \nu_\tau$
- ₁₃₁₀ • 25.5% decay to $\pi^- \pi^0 \nu_\tau$ (ρ^- resonance at 770 MeV)
- ₁₃₁₁ • 10.8% decay to $\pi^- \nu_\tau$
- ₁₃₁₂ • 9.3% decay to $\pi^- \pi^0 \pi^0 \nu_\tau$ (a_1^- resonance at 1200 MeV)
- ₁₃₁₃ • 9.0% decay to $\pi^- \pi^- \pi^+ \nu_\tau$ (a_1^- resonance at 1200 MeV)

₁₃₁₄ In all cases, at least one neutrino is produced. The neutrinos escape undetected
₁₃₁₅ from the CMS detector, resulting in missing transverse energy. Charged hadrons leave
₁₃₁₆ tracks in the tracking detector before being absorbed in the hadronic calorimeter; in
₁₃₁₇ CMS tau reconstruction terminology, they are often called “prongs”, i.e. the dominant

₁₃₁₈ τ_h decay modes are termed “1 prong” (π^\pm), “1 prong + π^0 (s)”, and “3-prong”. Neutral
₁₃₁₉ pions decay to two photons which lose their energy in the electromagnetic calorimeter.
₁₃₂₀ Taus that decay to electrons and muons, are typically triggered on and reconstructed
₁₃₂₁ as electrons and muons respectively.

₁₃₂₂ **Hadron plus strips (HPS) reconstruction of τ_h**

₁₃₂₃ At CMS, hadronically decaying tau leptons are reconstructed with the hadron plus
₁₃₂₄ strips (HPS) algorithm [51] [52]. The HPS algorithm capitalizes on photon conversions
₁₃₂₅ in the CMS tracker material, which originate from the neutral pion (π^0) decaying
₁₃₂₆ to two photons. The bending of electron/positron tracks due to the CMS solenoid
₁₃₂₇ magnetic field leads to a spread of the neutral pions’ calorimeter signatures in the ϕ
₁₃₂₈ direction. This motivates the reconstruction of photons in “strips”: objects that are
₁₃₂₉ built out of PF photons and electrons. The strip reconstruction starts with centering
₁₃₃₀ a strip on the most energetic electromagnetic particle in a PF jet. Among other
₁₃₃₁ electromagnetic particles located in a window of size $\Delta\eta = 0.05$ and $\Delta\phi = 0.20$
₁₃₃₂ around the strip center, the most energetic one is associated with the strip and its
₁₃₃₃ momentum is added to the strip momentum. This is repeated iteratively until no
₁₃₃₄ further particles can be associated. Lastly, strips satisfying a requirement of $p_T^{\text{strip}} > 1$
₁₃₃₅ GeV are combined with charged hadrons to reconstruct individual τ_h decay modes,
₁₃₃₆ where h stands for both π and K :

- ₁₃₃₇ • *Single hadron:* $h^- \nu_\tau$ and $h^- \pi^0 \nu_\tau$ decay modes, in which the neutral pions have
₁₃₃₈ too little energy to be reconstructed as strips.
- ₁₃₃₉ • *One hadron + one strip:* $h^- \pi^0 \nu_\tau$ decay modes, where the photons from the π^0
₁₃₄₀ decay are close together in the calorimeter.
- ₁₃₄₁ • *One hadron + two strips:* $h^- \pi^0 \nu_\tau$ decay modes, where the photons from the π^0
₁₃₄₂ decay are well separated.

- 1343 • *Three hadrons:* $h^- h^+ h^- \nu_\tau$ decay modes. The three charged hadrons are re-
1344 quired to originate from the same secondary vertex.

1345 The $h^- \pi^0 \pi^0 \nu_\tau$ and $h^- h^+ h^- \pi^0 \nu_\tau$ decay modes do not have their own treatment are
1346 reconstructed with the above topologies.

1347 In the HPS algorithm, the direction of the reconstructed tau momentum \vec{p}^{τ_h}
1348 is required to fall within a distance of $\Delta R = 0.1$ from the original PF jet. All
1349 charged hadrons and strips are required to be contained within a cone of size $\Delta R =$
1350 $(2.8 \text{ GeV})/p_T^{\tau_h}$, from the τ_h as reconstructed by the HPS.

1351 All charged hadrons are assumed to be pions, and they are required to be consis-
1352 tent with the masses of the intermediate meson resonances (if applicable), with the
1353 following allowed windows for candidates: 50-200 MeV for π^0 , 0.3-1.3 GeV for ρ , and
1354 0.8-1.5 GeV for a_1 . If the τ_h decay is compatible with more than one hypothesis, the
1355 one giving the highest $p_T^{\tau_h}$ is chosen. Lastly, an isolation requirement is applied: aside
1356 from the τ_h decay products, no charged hadrons or photons can be present within
1357 an isolation cone of size $\Delta R = 0.5$ around the direction of the τ_h . The outputs of
1358 the HPS algorithm are the reconstructed decay mode and the visible four-momentum
1359 (i.e. the four-momenta of all decay products excluding the neutrinos).

1360 **DeepTau for identifying τ_h**

1361 The identification of τ_h candidates in CMS has historically been divided into separate
1362 discriminators against jets, electrons, and muons. Discriminators versus jets and
1363 electrons use information from derived quantities, such as the p_T sum of particles
1364 near the τ_h axis. Building on the previous multivariate analysis (MVA) classifier [53]
1365 based on a boosted decision tree (BDT), DeepTau is a more recent classifier based on a
1366 deep neural network (DNN) that simultaneously discriminates against jets, electrons,
1367 and muons. The DNN uses a combination of high-level inputs, similar to previous
1368 algorithms, and also uses convolutional layers in η - ϕ space to process information

1369 from all reconstructed particles near the τ_h axis. Convolutional layers are based on
1370 the principle that an image can be processed independently of its position.

1371 The final DeepTau discriminators against jets, muons, and electrons are given by

$$D_\alpha(y) = \frac{y_\tau}{y_\tau + y_\alpha} \quad (5.1)$$

1372 where y_τ (y_α) are estimates of the probabilities for the τ_h candidate to come from
1373 a genuine τ_h (jet, μ , e). Working points for each discriminator with different τ_h
1374 identification efficiencies are defined for D_e , D_μ , and D_{jet} , for usage in physics analyses
1375 and derivation of data-to-simulation corrections [54].

1376 5.1.2 Muons

1377 Muons are the next lightest lepton after taus, with a mass of 105.66 MeV and a
1378 mean lifetime of $\tau = 2.20 \times 10^{-6}$ seconds, or $c\tau = 658.64$ m. At CMS, muons are
1379 identified with requirements on the quality of the track reconstruction and on the
1380 number of measurements in the tracker and the muon systems [55]. In the standard
1381 CMS reconstruction, tracks are first reconstructed independently in the inner tracker
1382 (tracker track) and in the muon system (standalone-muon track). Next, these tracks
1383 are processed in two different methods.

1384 The first is Global Muon reconstruction (outside-in) [55], which fits combined hits
1385 from the tracker track and standalone-muon track, using the Kalman-filter technique.
1386 At large transverse momenta, $p_T \gtrsim 200$ GeV, the global-muon fit can improve the
1387 momentum resolution compared to the tracker-only fit.

1388 The second is Tracker Muon reconstruction (inside-out) [55], which starts with
1389 tracker tracks with $p_T > 0.5$ GeV and total momentum $p_T > 2.5$ GeV. These tracks
1390 are extrapolated outwards to the muon system and matched to detector segments
1391 there, taking into account the magnetic field, expected energy losses, and multiple

1392 Coulomb scattering in the detector material. Tracker Muon reconstruction is more
 1393 efficient than the Global Muon reconstruction at low momenta, $p \lesssim 5$ GeV, because
 1394 it only requires a single muon segment in the muon system, whereas Global Muon
 1395 reconstruction typically requires segments in at least two muon stations.

1396 To further suppress fake muons from decay in flight, isolation cuts are used. A
 1397 relative isolation variable is defined to quantify the energy flow of particles near the
 1398 muon trajectory. A relative isolation is defined similarly for muons and electrons:

$$I^\ell \equiv \frac{\sum_{\text{charged}} p_T + \max\left(0, \sum_{\text{neutral}} p_T - \frac{1}{2} \sum_{\text{charged, PU}} p_T\right)}{p_T^\ell} \quad (5.2)$$

1399 where $\sum_{\text{charged}} p_T$ is the scalar sum of the p_T of the charged particles originating from
 1400 the primary vertex and located in a cone of size $\Delta R = \sqrt{(\Delta\eta)^2 + (\Delta\phi)^2} = 0.4(0.3)$
 1401 centered on the direction of the muon (electron). The sum $\sum_{\text{neutral}} p_T$ is the equivalent
 1402 for neutral particles. The sum $\sum_{\text{charged, PU}} p_T$ is the scalar sum of the p_T of the
 1403 charged hadrons in the cone originating from pile-up vertices. The factor 1/2 comes
 1404 from simulation estimations, which find that the ratio of neutral to charged hadron
 1405 production in the hadronization process of inelastic pp collisions is 1/2. Thus the
 1406 subtracted term is intended to subtract contribution from pile-up, from the neutral
 1407 particle contribution to the isolation sum. Finally, this is divided by the lepton
 1408 transverse momentum, p_T^ℓ .

1409 **5.1.3 Electrons**

1410 Electrons are the lightest lepton with a mass of 0.511 MeV. At CMS, electrons are
 1411 reconstructed by associating a track reconstructed in the silicon tracking detector
 1412 with a cluster of energy in the ECAL. Performance is maximized via a combination
 1413 of a stand-alone approach and the complementary global particle-flow approach [56].

1414 In the stand-alone approach, the electron energy, which is typically spread over

several crystals of the ECAL, is clustered with the “hybrid” algorithm in the barrel and the “multi- 5×5 ” in the endcaps [56]. The hybrid algorithm collects energy in a small window in η and an extended window in ϕ . It identifies a seed crystal, and adds arrays of 5×1 crystals in $\eta \times \phi$ in a range of $N = 17$ crystals in both directions of ϕ , if their energies exceed a minimum threshold, thus forming a supercluster (SC). In the endcap, crystals are not arranged in an $\eta \times \phi$ geometry; instead clusters are build around seed crystals in clusters of 5×5 crystals that can partly overlap. Nearby clusters are grouped into a supercluster, and energy is recovered from associated deposits in the preshower.

In the PF reconstruction [56], PF clusters are reconstructed by aggregating around a seed all contiguous crystals with energies two standard deviations above the electronic noise observed at the beginning of a data-taking run. The energy of a given crystal can be shared among two or more clusters.

The electron track reconstruction is performed in two ways [56]: the ECAL-based seeding, which begins with the SC energy and positioning, and the tracker-based seeding (part of the PF reconstruction algorithm), which uses tracks reconstructed from the general algorithm for charged particles, extrapolated towards the ECAL and matched to an SC. Kalman filter (KF) tracks with a small number of hits or that are not well-fitted, are re-fitted with a dedicated Gaussian sum Filter (GSF).

A global identification variable [56] is defined using a multivariate analysis (MVA) technique that combines information on track observables (kinematics, quality of the KF track and GSF track), the electron PF cluster observables (shape and pattern), and the association between the two (geometric and kinematic observables). For electrons seeded only through the tracker-based approach, a weak selection is applied on this MVA variable. For electrons seeded through both approaches, a logical OR is taken.

Electron isolation, i.e. the presence of energy deposits near the electron trajectory,

1442 is a separate key handle in rejecting significant background. Compared to isolated
 1443 electrons, electrons from misidentified jets or genuine electrons within a jet resulting
 1444 from semileptonic decays of b or c quarks tend to have significant energy deposits
 1445 near the primary trajectory [56]. Offline analyses benefit from the PF technique
 1446 for defining isolation, which sums the PF candidates reconstructed located within a
 1447 specified isolation cone around the electron candidate, as in Eqn. 5.2.

1448 5.1.4 Jets

1449 The vast majority of processes of interest at the LHC contains quarks or gluons in
 1450 the final state, but these particles cannot be observed directly. In a process called
 1451 hadronization, they fragment into spatially-grouped collections of particles called jets,
 1452 which can be detected in the tracking and calorimeter systems. Hadronization and
 1453 the subsequent decays of unstable hadrons can produce hundreds of nearby particles
 1454 in the CMS detector. Jets are reconstructed by the PF algorithm (PF jets), or from
 1455 the sum of the ECAL and HCAL energies deposited in the calorimeter towers (Calo
 1456 jets). In PF jets, typically used in offline analyses, jets are built using the anti- k_T
 1457 (AK) clustering algorithm [57]. The anti- k_T algorithm iterates over particle pairs and
 1458 finds the two that are closest in a distance measure d , and determines whether to
 1459 combine them:

$$d_{ij} = \min(p_{T,i}^{-2}, p_{T,j}^{-2}) \frac{\Delta_{ij}^2}{R^2}, \text{ combine when } d_{ij} < p_{T,i}^{-2}; \text{ stop when } d_{ij} > p_{T,i}^{-2} \quad (5.3)$$

1460 where $\Delta_{ij}^2 = (\eta_i - \eta_j)^2 + (\phi_i - \phi_j)^2$ and $p_{T,i}$, η_i , ϕ_i are the transverse momentum, rapid-
 1461 ity, and azimuthal angle of particle i . The power -2 means that higher-momentum
 1462 particles are clustered first, leading to jets that tend to be centered on the hardest
 1463 (highest p_T) particle.

1464 There are several methods to remove contributions of pile-up collisions from jet

1465 clustering [58]:

- 1466 • Charged hadron subtraction (CHS), which removes all charged hadron candi-
1467 dates associated with a track that is not associated with the primary vertex.
- 1468 • PileUp Per Particle Identification (PUPPI), which weighs input particles based
1469 on their likelihood of arising from pile-up. QCD particles tend to have a collinear
1470 structure, compared to soft diffuse radiation coming from pile-up. The local
1471 shape for charged pile-up, used as a proxy for all pile-up particles, is used on an
1472 event-by-event basis to calculate a weight for each particle. PUPPI is deployed
1473 in Run-2 and is more performant than CHS in high pile-up scenarios.

1474 5.1.5 B-flavored jets

1475 Jets that arise from bottom-quark hadronization (b-flavor jets) have overwhelming
1476 background from processes involving jets from gluons (g) and light-flavor quarks (u, d,
1477 s), and from c-quark fragmentation. The ability to identify b-flavor jets, or b-tagging,
1478 exploits the b hadrons' relatively large masses, long lifetimes, and daughter particles
1479 with hard momentum spectra [57].

1480 The impact parameter (IP) of a track is the 3-dimensional distance between the
1481 track and the primary vertex (PV) at the point of closest approach. The IP is positive
1482 if the track originates from the decay of particles travelling along the jet axis. The
1483 resolution of the IP depends on the p_T and η of the track, motivating the use of the
1484 impact parameter significance S_{IP} (ratio of the IP to its estimated uncertainty) as an
1485 observable [57].

1486 Because of the large but finite lifetimes of the b hadrons, b hadrons tend to
1487 travel a short distance before decaying at a secondary vertex (SV), which can be
1488 measured and reconstructed separately from the primary vertex due to the excellent
1489 position resolution of the pixel detector [57]. Previous b-tagging algorithms (e.g.

1490 CSV, cMVAv2, and DeepCSV) have capitalized on variables such as the presence of
1491 a SV, the flight distance and direction (computed from the vector between the PV
1492 and the SV), and kinematics of the system of associated secondary tracks (e.g. track
1493 multiplicity, mass, and energy).

1494 The DeepJet (formerly known as DeepFlavour) algorithm [59] is a deep-neural-
1495 network multi-classification algorithm, which uses 16 properties of up to 25 charged
1496 and 6 properties of 25 neutral particle-flow jet constituents, as well as 17 properties
1497 from up to 4 secondary vertices associate with the jet. Compared to the previous clas-
1498 sifying algorithm DeepCSV, DeepJet has been demonstrated to have higher efficiency
1499 with lower misidentification probability in Phase-1 data [60].

1500 5.2 Reconstruction of the di-tau mass

1501 The final signal extraction is done to the total di-tau ($\tau\tau$) mass, which is estimated
1502 from the visible $\tau\tau$ mass using the FastMTT algorithm [61]. FastMTT is based on the
1503 SVFit algorithm, originally developed for the Standard Model $H \rightarrow \tau\tau$ analysis [62].
1504 Both the SVFit algorithms, and the FastMTT algorithm, are described below, to give
1505 a complete picture of how the algorithms attempt to reconstruct the true invariant
1506 mass of a Higgs or Z boson decay.

1507 To specify a hadronic τ decay, six parameters are needed [62]: the polar and
1508 azimuthal angles of the visible decay product system in the τ rest frame, the three
1509 boost parameters from the τ rest frame to the laboratory frame, and the invariant
1510 mass m_{vis} of the visible decay products. For a leptonic τ decay, two neutrinos are
1511 produced, and a seventh parameter, the invariant mass of the two-neutrino system, is
1512 necessary. The unknown parameters are constrained by four observables that are the
1513 components of the four-momentum of the system formed by the visible decay products
1514 of the τ lepton, measured in the laboratory frame. The remaining unconstrained

1515 parameters for hadronic and leptonic τ decays are thus:

- 1516 • The fraction of the τ energy in the laboratory frame carried by the visible decay
1517 products,
- 1518 • ϕ , the azimuthal angle of the τ direction in the laboratory frame,
- 1519 • $m_{\nu\nu}$, the invariant mass of the two-neutrino system in leptonic τ decays (for
1520 hadronic τ decays, $m_{\nu\nu}$ is set to 0).

1521 E_x^{miss} and E_y^{miss} , the x and y components of the missing transverse energy E_T^{miss}
1522 provide two further constraints.

1523 5.2.1 Original SVFit ‘‘standalone’’: maximum likelihood

1524 In one of the original versions of SVFit, called ‘‘standalone’’ SVFit [62], a maximum
1525 likelihood fit method is used to reconstruct the mass $m_{\tau\tau}$ by combining the measured
1526 observables E_x^{miss} and E_y^{miss} with a likelihood model that includes terms for the τ
1527 decay kinematics and the E_T^{miss} resolution [62]. The likelihood function $f(\vec{z}, \vec{y}, \vec{a}_1 \vec{a}_2)$
1528 of the parameters $\vec{z} = (E_x^{\text{miss}}, E_y^{\text{miss}})$ in an event is constructed, where the remaining
1529 parameters are the kinematics of the two τ decays, denoted $\vec{a}_1 = (x_1, \phi_1, m_{\nu\nu,1})$ and
1530 $\vec{a}_2 = (x_2, \phi_2, m_{\nu\nu,2})$, and the four-momenta of the visible decay products with the
1531 measured values $\vec{y} = (p_1^{\text{vis}}, p_2^{\text{vis}})$.

1532 The likelihood f is the product of three likelihood functions. The first two likelihood
1533 functions model the decay parameters \vec{a}_1 and \vec{a}_2 of the two τ leptons. For leptonic
1534 decays, the likelihood function is modeled using matrix elements for τ decays,
1535 and integrated over the allowed phase space $0 \leq x \leq 1$ and $0 \leq m_{\nu\nu} \leq m_\tau \sqrt{1-x}$. For
1536 hadronic τ decays, a model based on the two-body phase space is used and integrated
1537 over $m_{\text{vis}}^2/m_{\tau\tau}^2 \leq x \leq 1$. The third likelihood function quantifies the compatibility of
1538 a τ decay hypothesis with the reconstructed \vec{E}_T^{miss} in an event, assuming the neutrinos
1539 are the only source of missing transverse energy. The expected \vec{E}_T^{miss} resolution

1540 is represented by a covariant matrix, estimated on an event-by-event basis using a
1541 significance algorithm [63].

1542 5.2.2 “Classic SVFit” with matrix element

1543 Classic SVFit is an improved algorithm of the original “standalone” SVFit using the
1544 formalism of the matrix element (ME) method [61]. In the ME method, an estimate
1545 for the unknown model parameter Θ (here, the mass $m_{\tau\tau}$) is obtained by maximizing
1546 the probability density \mathcal{P} . The key ingredients of the probability density are the
1547 squared modulus of the matrix element $|\mathcal{M}(\mathbf{p}, \Theta)|^2$ and the transfer function $W(\mathbf{y}|\mathbf{p})$
1548 (probability density to observe the measured observables \mathbf{y} given the phase space
1549 point \mathbf{p}). The best estimate $m_{\tau\tau}$ is obtained by computing the probability density \mathcal{P}
1550 for a range of mass hypotheses and finding the value of $m_{\tau\tau}$ that maximizes \mathcal{P} .

1551 Distributions illustrating the performance of the classic matrix element SVFit
1552 algorithm are shown in Fig. 5.1 from [61], showing the di-tau mass after and before
1553 application of SVFit to recover energy lost to neutrinos. The SVFit algorithm is
1554 found to improve the sensitivity of the Standard Model $H \rightarrow \tau\tau$ analysis performed
1555 by CMS by about 30%, compared to performing the same analysis using only the
1556 visible mass m_{vis} .

1557 5.2.3 FastMTT: optimized SVFit

1558 FastMTT [64] is a further simplification to the matrix element method of Classic
1559 SVFit which has comparable performance but is about 100 times faster. FastMTT
1560 drops the matrix element component of the computation without significant impact
1561 on the final mass resolution, and simplifies the computation of the transfer functions.
1562 The opening angle of the τ decay products with respect to the initial τ momenta ap-
1563 proaches 0 for τ with high $\gamma = E_\tau/m_\tau$, with typical τ decays from the Z boson decays
1564 already satisfying this condition. In this collinear approximation, the dimensionality

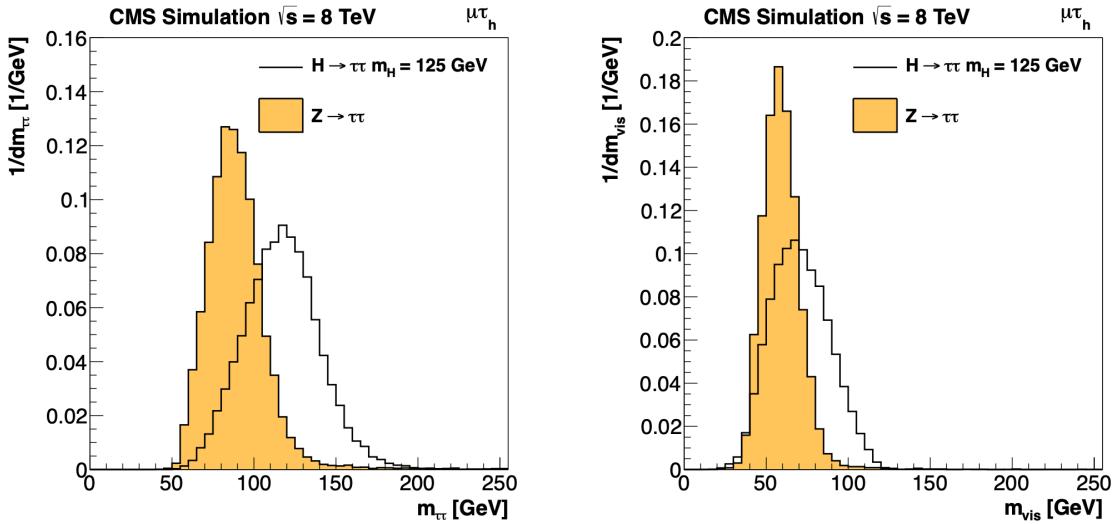


Figure 5.1: Distributions from [61], of $m_{\tau\tau}$ after reconstruction with the original SVFit algorithm (*left*), and before SVFit with only the visible tau decay products (*right*), for $H \rightarrow \tau\tau$ signal events of mass $m_H = 125$ GeV (*black line*) and the $Z/\gamma^* \rightarrow \tau\tau$ background (*orange, solid*), in the decay channel $\tau\tau \rightarrow \mu\tau_h$.

1565 of the transfer function can be reduced in the computation of FastMTT, while still
1566 yielding similar results to Classic SVFit [64].

1567 5.3 Corrections applied to simulation

1568 Corrections are applied to simulated samples to account for known effects in the event
1569 modeling and reconstruction and data-taking, and are intended to bring simulations
1570 in closer agreement with data. Corrections fall into two broad categories: *energy*
1571 *scale corrections* applied to physics objects, and *event-level corrections*. Energy scale
1572 corrections are multiplicative factors applied to the energy and transverse momentum
1573 p_T of simulated objects (e.g. leptons or jets), and bring the average reconstructed en-
1574 ergies of simulated particles into better agreement with those of objects reconstructed
1575 from data. Event-level corrections are applied as a per-event multiplicative weight,
1576 and account for effects such as mis-modeling in simulations of the underlying physics
1577 process, or changing detector operating conditions during data-taking. Event-level

1578 corrections change the shapes of the distributions of all the physical observables.

1579 Uncertainties in scale factors and corrections are also sources of systematic errors
1580 in the analysis, detailed in Chapter 8. Systematic uncertainties in the tau, muon, and
1581 electron energy scales can shift the p_T of the leptons up or down, which can change
1582 whether events pass or fail the offline p_T thresholds for the trigger paths described in
1583 the previous section, i.e. change the number of events in the signal region.

1584 5.3.1 Tau energy scale

1585 An energy scale is applied to the transverse momentum p_T and mass of the hadronic
1586 tau τ_h in the $\mu\tau_h$ and $e\tau_h$ channels, to correct for a deviation of the average recon-
1587 structed τ_h energy from the generator-level energy of the visible τ_h decay products.
1588 These correction factors are derived centrally [53], by fitting to events in $e\tau_h$ and $\mu\tau_h$
1589 final states in Z/γ^* events separately for the h^\pm , $h^\pm\pi^0$, and $h^\pm h^\mp h^\pm$ decays. The
1590 values used are shown in Table 5.1.

1591 When applying the energy scale to the τ_h , the 4-momentum of the missing trans-
1592 verse energy (MET) is adjusted such that the total 4-momenta of the τ_h and the MET
1593 remains unchanged [65].

Tau energy scale factor				
Decay mode	2018	2017	2016 pre-VFP	2016 post-VFP
0	0.991 ± 0.008	0.986 ± 0.009	0.987 ± 0.01	0.993 ± 0.009
1	1.004 ± 0.006	0.999 ± 0.006	0.998 ± 0.006	0.991 ± 0.007
10	0.998 ± 0.007	0.999 ± 0.007	0.984 ± 0.008	1.001 ± 0.007
11	1.004 ± 0.009	0.996 ± 0.01	0.999 ± 0.011	0.997 ± 0.016

Table 5.1: Energy scales applied to genuine hadronic tau decays τ_h by data-taking year/era and decay mode, along with systematic errors.

1594 **5.3.2 Muon energy scale**

1595 An energy scale is applied to the p_T and mass of genuine muons from τ decays in the
1596 $e\mu$ and $\mu\tau_h$ channels [66]. The applied values are the same for MC and embedded
1597 samples and are shown in Table 5.2. Following the SM $H \rightarrow \tau\tau$ analysis, Rochester
1598 corrections are not applied, and instead prescriptions from [67] are followed.

Muon energy scale factor	
Eta range	Value for all years
$ \eta \in [0.0, 1.2)$	1.0 ± 0.004
$ \eta \in [1.2, 2.1)$	1.0 ± 0.009
$ \eta \in [2.1, 2.4)$	1.0 ± 0.027

Table 5.2: Energy scales and systematic errors applied to genuine muons. The values are the same for MC and embedded for all years [68] [67].

1599 **5.3.3 Electron energy scale**

1600 Corrections to the electron energy scale are applied to genuine e from τ decays, and
1601 are binned in two dimensions by electron p_T and η for barrel vs. endcap [69]. The
1602 scale factors are binned in p_T and η for MC samples: e.g. values for 2018 are shown
1603 in Fig. 5.2 from [70]. For embedded samples the electron energy scale is taken as
1604 only binned in η (Table 5.3).

Electron energy scale factor for embedded samples			
Eta range	2018	2017	2016
$ \eta \in [0.0, 1.479)$	0.973 ± 0.005	0.986 ± 0.009	0.9976 ± 0.0050
$ \eta \in [1.479, 2.4)$	0.980 ± 0.0125	0.887 ± 0.0125	0.993 ± 0.0125

Table 5.3: Energy scales and systematic errors applied to electrons in embedded samples, binned in the electron η , by data-taking year [71] [72] [73].

1605 **5.3.4 τ_h identification efficiency**

1606 The τ_h identification efficiency can differ in data and MC [65]. Recommended correc-
1607 tions are provided by the Tau POG, and we use the medium DeepTau vs. jet working

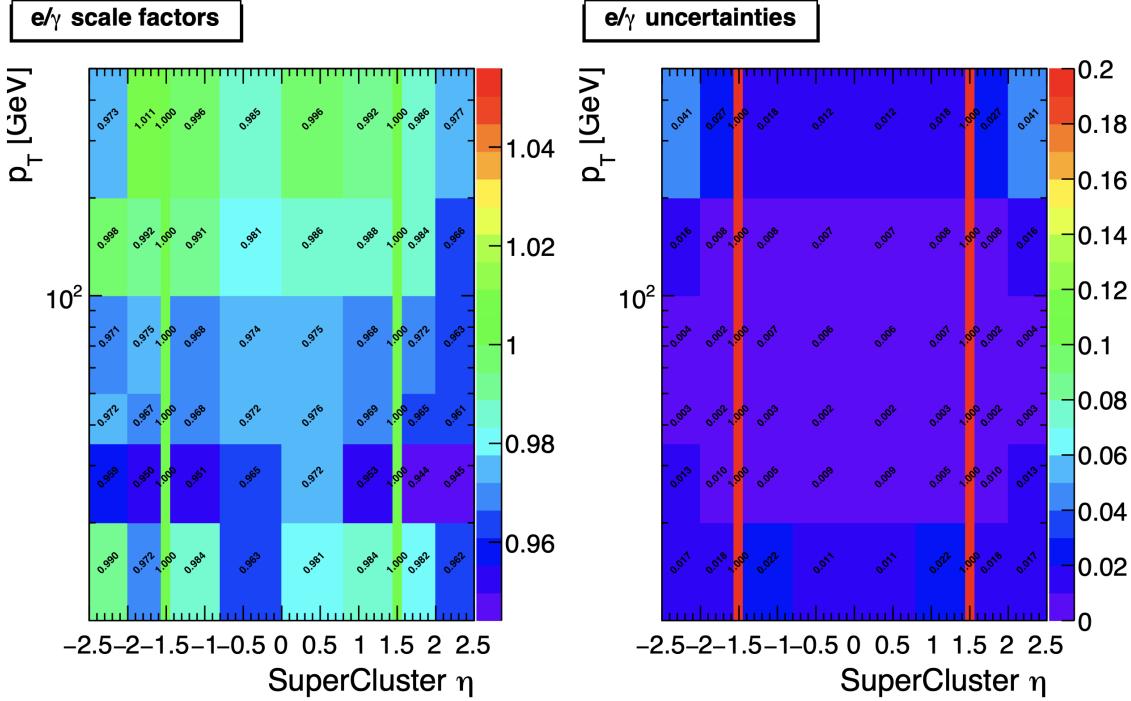


Figure 5.2: Electron/photon energy scale factors (*left*) and corresponding uncertainties (*right*) binned in the electron η and p_T , for the data-taking year 2018 [70].

1608 point values. The identification efficiency is measured in $Z \rightarrow \tau\tau$ events in the $\mu\tau_h$
 1609 final state, and is binned in p_T due to clear p_T dependence of the DeepTau ID.

Tau ID efficiency for DeepTau Medium vs. jet WP in 2018						
p_T (GeV)	< 20	(20, 25]	(25, 30]	(30, 35]	(35, 40]	(40, 500]
Central value	0	0.945	0.946	0.916	0.921	1.005
Up value	0	1.001	0.981	0.946	0.950	1.035
Down value	0	0.888	0.981	0.883	0.893	0.953

Table 5.4: Tau ID efficiency for the DeepTau vs. jet medium working point, with central, up, and down values for 2018, binned in the tau p_T [65].

1610 5.3.5 Trigger efficiencies

1611 Scale factors are applied to correct for differences in trigger efficiencies between MC
 1612 and embedded vs. data, with values taken from tools provided by the Standard Model
 1613 $H \rightarrow \tau\tau$ working group which uses the same trigger paths [68]. In the following

1614 sections we review relevant trigger efficiencies in data, which form the basis of the
1615 trigger efficiency corrections applied to MC and embedded.

1616 **5.3.6 Tau trigger efficiencies**

1617 The efficiencies in data of the single- τ_h leg in $\mu\tau_h$, $e\tau_h$, and di- τ_h triggers is computed
1618 centrally per using a Tag and Probe (TnP) method [74] which is outlined here. In
1619 this method, $Z \rightarrow \tau\tau \rightarrow \mu\tau_h$ are selected in data and a Drell-Yan simulated sample
1620 ($Z \rightarrow \ell\ell, \ell = e, \mu, \tau_h$) with high purity. Cuts are applied to reject events not in this
1621 final state, e.g. suppressing $Z \rightarrow \mu\mu$ by vetoing events with a single loose ID muon.
1622 An isolated muon candidate (the tag) with online $p_T > 27$ GeV and $|\eta| < 2.1$ is
1623 identified and matched to an offline μ . An offline τ_h candidate (the probe) is selected,
1624 which is separated from the tag μ , and has $p_T > 20$ GeV and $|\eta| < 2.1$. The probe
1625 τ_h must pass anti-muon and anti-electron discriminators to avoid fakes from muons
1626 and electrons, and must pass the medium MVA tau isolation to suppress fakes from
1627 QCD jets. The trigger efficiency in the TnP method is calculated as

$$\text{Efficiency} = \frac{\text{Number of events passing the TnP selection with fires the HLT path}}{\text{Number of events passing the TnP selection}} \quad (5.4)$$

1628 The efficiencies for the hadronic tau legs in the relevant channels of this analyses
1629 ($\mu\tau_h$ and $e\tau_h$) as a function of the offline tau p_T and η , are shown for data taken in
1630 2016, 2017, and 2018 in Figures 5.3a and 5.3b [74] [75]. In both figures, the different
1631 HLT thresholds and differences in the L1 seed result in higher efficiencies in 2016 and
1632 differences in shapes of the 2016 efficiencies compared to 2017 and 2018. The low
1633 pile-up in 2016 also leads to higher efficiencies in that year.

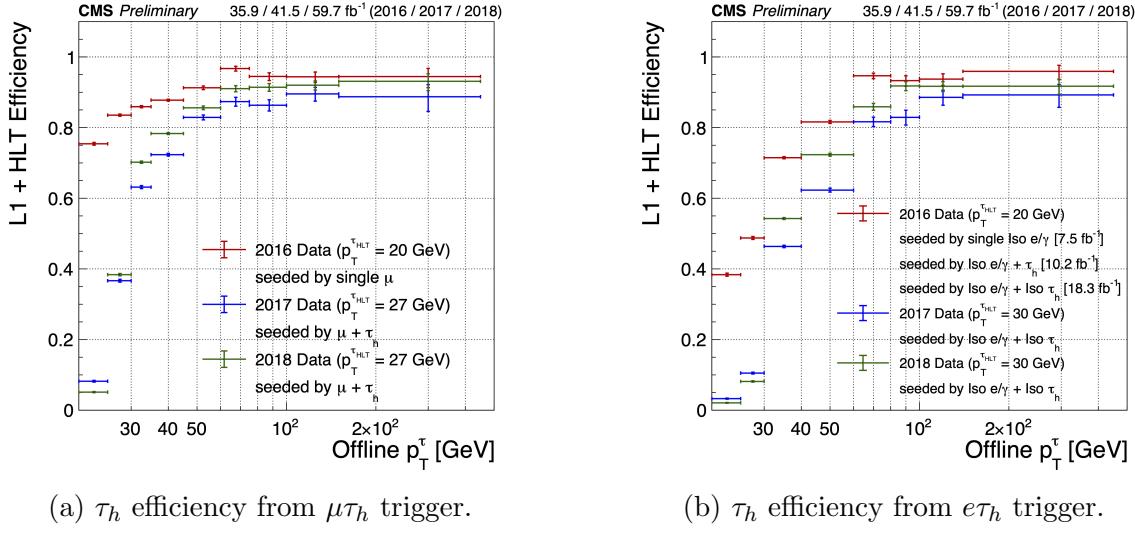


Figure 5.3: Hadronic tau leg efficiency of the cross-triggers for $\mu\tau_h$ (left) and $e\tau_h$ (right) triggers as a function of offline tau p_T for the years 2016 (red), 2017 (blue) and 2018 (green), from [75]. HLT p_T thresholds and L1 seeds are indicated in the legends.

5.3.7 Single muon trigger efficiencies

The efficiencies for the single isolated muon trigger with $p_T > 24$ GeV used in this analysis, is shown for the data-taking year 2018 in Fig. 5.4a as a function of the muon p_T and as a function of the muon $|\eta|$ in Fig. 5.4b from [76]. The data is split with respect to a HLT muon reconstruction update that was deployed on 15/05/2018. A small asymmetry in efficiencies between negative and positive η in Fig. 5.4b is due to disabled muon chambers (CSCs). The efficiencies shown are estimated using a Tag and Probe method using $Z \rightarrow \mu\mu$ events, with the tag being an offline muon with $p_T > 29$ GeV and $|\eta| < 2.4$ passing a tight ID criteria, and the probe is an online (L1) trigger object with $\Delta R < 0.3$ and passing tight ID and Particle Flow based isolation requirements with $p_T > 26$ GeV.

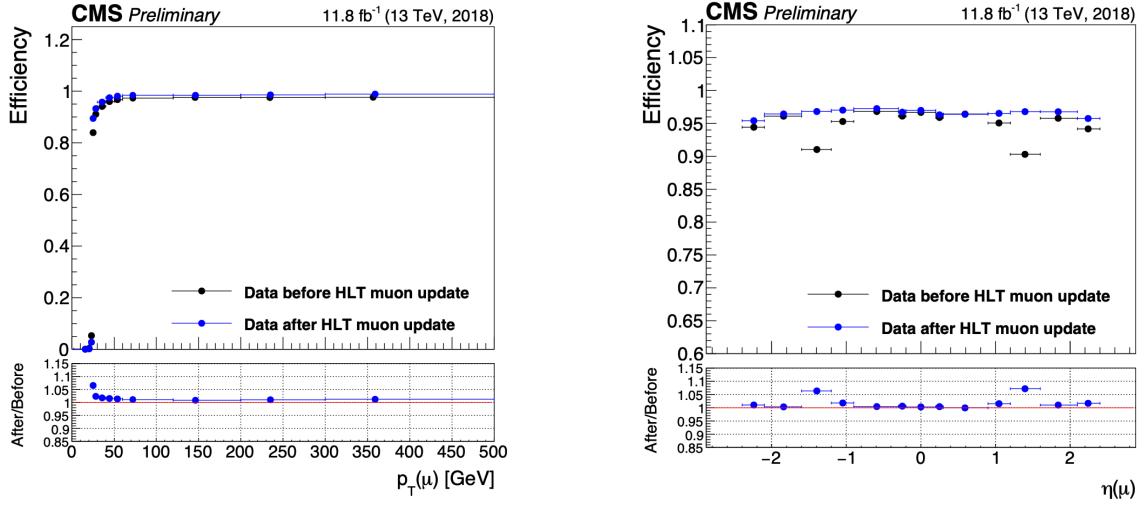
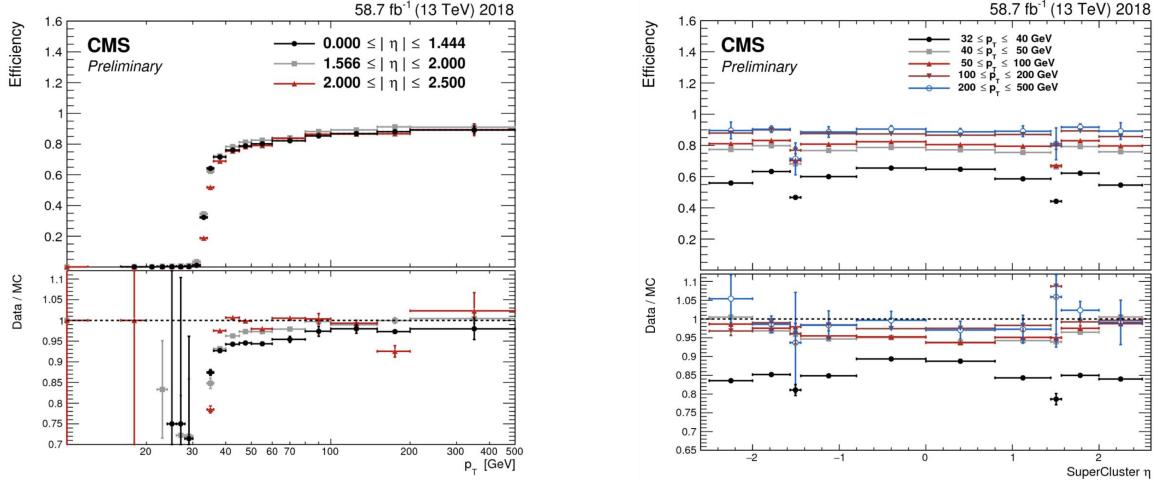
(a) Muon efficiency vs p_T for SingleMuon.(b) Muon efficiency vs $|\eta|$ for SingleMuon.

Figure 5.4: Trigger efficiencies in data (*top panels*) and ratio of efficiencies after/before a HLT muon reconstruction update (*bottom panels*) for the muon in the isolated single muon trigger with threshold $p_T > 24$ GeV in the data-taking year 2018, as functions of the muon p_T (*left*) and muon $|\eta|$ (*right*). Only statistical errors are shown [76].

1645 5.3.8 Single electron trigger efficiencies

1646 The efficiencies in data, and the ratio between data and MC, of the single electron
 1647 HLT trigger with p_T threshold 32 GeV used in this analysis are shown for 2018,
 1648 as a function of the electron p_T in Fig. 5.5a and of the electron $|\eta|$ in Fig. 5.5b,
 1649 from [77]. In the Tag and Probe method used for the 2018 dataset, the tag is an
 1650 offline reconstructed electron with $|\eta| \leq 2.1$ and not in the barrel and endcap overlap
 1651 region, with $p_T > 35$ GeV with tight isolation and shower shape requirements, firing
 1652 the tag trigger. The probe is an offline reconstructed electron with $|\eta| \leq 2.5$ with
 1653 $E_T^{\text{ECAL}} > 5$ GeV with no extra identification criteria [77].

1654 The disagreement between data and MC, particularly at low transverse momentum,
 1655 is in part due to detector effects that are difficult to simulate, such as crys-
 1656 tal transparency losses in the ECAL and the evolution of dead regions in the pixel
 1657 tracker [77].



(a) Electron efficiency vs p_T for single electron.

(b) Electron efficiency vs $|\eta|$ for single electron.

Figure 5.5: Trigger efficiencies in data, and the data/MC ratio for the electron in the single electron trigger with threshold $p_T > 32$ GeV in the data-taking year 2018, as functions of the electron p_T (*left*) and electron $|\eta|$ (*right*) [77]. In the plot vs. p_T , the region $1.442 \leq |\eta| \leq 1.566$ is not included as it corresponds to the transition between barrel and endcap parts of the ECAL.

1658 5.3.9 $e\mu$ cross-trigger efficiencies

1659 The efficiencies of the electron and muons for the cross-trigger with leading muon
 1660 used in the $e\mu$ channel are shown for data in 2016, 2017, and 2018 in Figures 5.6a and
 1661 5.6b [78]. These efficiencies were measured centrally using a Tag and Probe in events
 1662 with Z to dileptons with the same flavor and opposite charge, where the tags are an
 1663 isolated muon or electron, and the probe (offline) candidate is required to satisfy the
 1664 same lepton selection as that of the tag candidate, be matched within $\Delta R < 0.1$ with
 1665 a corresponding online trigger object, and also to pass the cross-trigger. The trigger
 1666 efficiency is then:

$$\text{Efficiency} = \frac{\text{Events passing lepton pair selections and probe passing trigger}}{\text{Events passing lepton pair selections}} \quad (5.5)$$

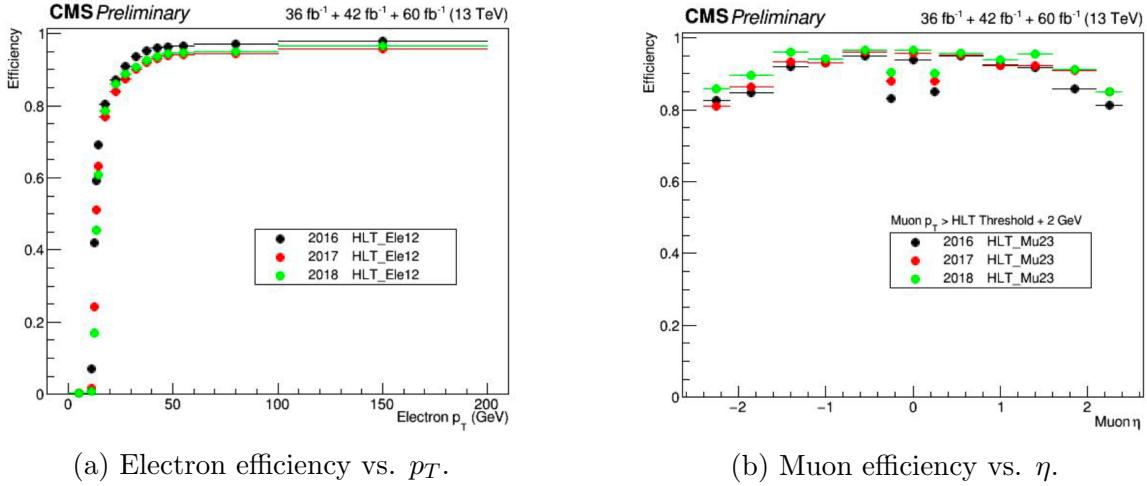


Figure 5.6: Efficiencies of the electron leg vs. p_T (*left*) and the muon log vs. η (*right*), for the HLT path with online thresholds of 12 GeV for the electron and 23 GeV for the muon, for the data-taking years 2016 (*black*), 2017 (*red*), and 2018 (*green*) [78].

1667 5.3.10 Electrons and muons faking τ_h : energy scales

1668 Energy scales for electrons misidentified as hadronic tau decays (e faking τ_h) are
 1669 provided by the Tau POG, and were measured in the $e\tau_h$ channel with the visible
 1670 invariant mass of the electron and hadronic tau system [68]. This energy scale is
 1671 applied for τ_h with $p_T > 20$ GeV regardless of which DeepTau vs. electron working
 1672 point was used. Values for 2018 are shown in Table 5.5.

Electrons faking τ_h energy scale factor in 2018	
Reconstructed decay mode of the fake τ_h	Central value and (up, down) shifts
0	1.01362 (+0.00474, -0.00904)
1	1.01945 (+0.01598, -0.01226)
10	0.96903 (+0.0125, -0.03404)
11	0.985 (+0.04309, -0.05499)

Table 5.5: Energy scales and up/down systematic uncertainties applied to electrons misidentified as hadronic taus for 2018, binned in decay mode of the fake τ_h [68].

1673 No nominal energy scale is applied for muons mis-reconstructed as τ_h , and the
 1674 uncertainty is treated as $\pm 1\%$ and uncorrelated in the reconstructed decay mode [68].

1675 **5.3.11 Electrons and muons faking τ_h : misidentification effi-**
 1676 **ciencies**

1677 Corrections on identification efficiencies are applied to genuine electrons and muons
 1678 misidentified as τ to account for differences in data and MC.

1679 The specific values depend on the vs. electron and vs. muon discriminator working
 1680 points used. For misidentified $\mu \rightarrow \tau_h$, the scale factors are split into different $|\eta|$
 1681 regions, determined by the CMS muon and tracker detector geometries, as shown in
 1682 Table 5.6 for 2018 [65].

Tau ID efficiency for DeepTau vs. muon WPs in 2018		
$ \eta $	Tight working point	VLoose working point
(0.0, 0.2)	0.767 ± 0.127	0.954 ± 0.069
(0.2, 0.6)	1.255 ± 0.258	1.009 ± 0.098
(0.6, 1.0)	0.902 ± 0.203	1.029 ± 0.075
(1.0, 1.45)	0.833 ± 0.415	0.928 ± 0.145
(1.45, 2.0)	4.436 ± 0.814	5.000 ± 0.377
(2.0, 2.53)	1.000 ± 0.000	1.000 ± 0.000

Table 5.6: Tau mis-identification efficiency for the DeepTau Tight and Very Loose (VLoose) working points vs. muons in 2018, binned in the muon $|\eta|$ [65].

1683 For misidentified $e \rightarrow \tau_h$, the scale factors are split into barrel and endcap regions,
 1684 dictated by the ECAL detector geometry, as shown in Table 5.7 for 2018.

Tau ID efficiency for DeepTau vs. electron WPs in 2018		
$ \eta $	Tight working point	VLoose working point
(0.0, 0.73)	1.47 ± 0.27	0.95 ± 0.07
(0.73, 1.509)	1.509 ± 0.0	1.00 ± 0.0
(1.509, 1.929)	1.929 ± 0.2	0.86 ± 0.1
(1.929, 2.683)	2.683 ± 0.9	2.68 ± 0.0

Table 5.7: Tau mis-identification efficiency for the DeepTau Tight and Very Loose (VLoose) working points vs. electrons in 2018, binned in the electron $|\eta|$ [65].

1685 5.3.12 Electron ID and tracking efficiency

1686 Scale factors are applied to MC to correct for differences between MC and data in
 1687 the performance of electron identification (ID) and tracking.

1688 Electron and photon identification, as discussed earlier, use variables with good
 1689 signal vs. background discrimination power such as lateral shower shape and ratio
 1690 of energy deposited in the HCAL to energy deposited in the ECAL at the position
 1691 of the electron. The cut-based electron identification efficiencies in data and ratio of
 1692 efficiencies in data to MC are shown in Fig. 5.7a for the multivariate analysis (MVA)
 1693 identification working point.

1694 The tracking efficiencies in data and the data/MC ratio are shown in Fig. 5.7b
 1695 for the Gaussian-sum filter (GSF) tracking [79].

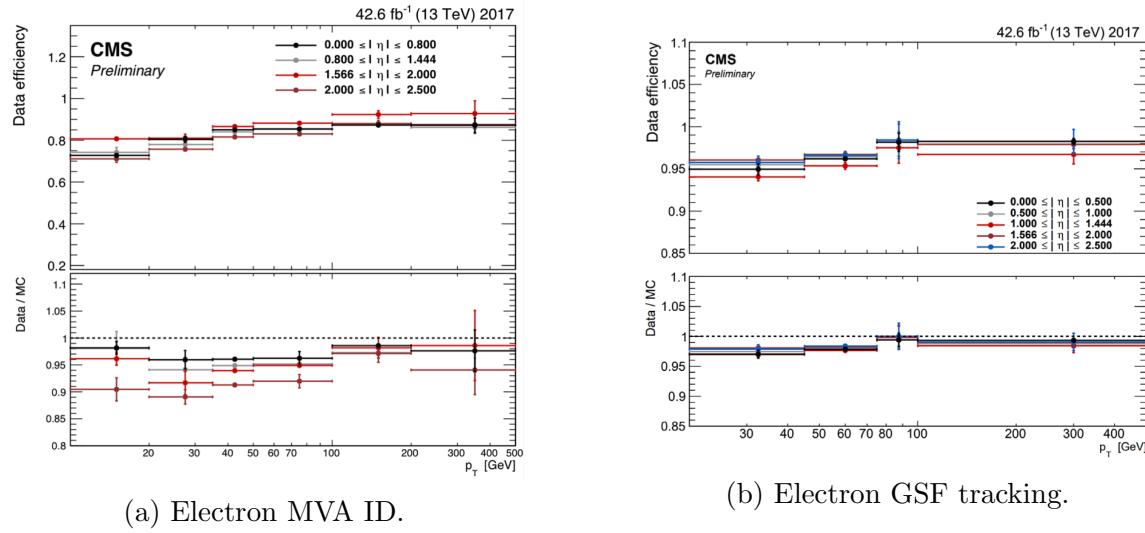


Figure 5.7: Efficiencies in data (*top panels*) and the ratio of efficiencies in data/MC (*bottom panels*), for the electron multivariate analysis (MVA) identification (*left*) and for the Gaussian-sum filter (GSF) tracking (*right*) [79]. Error bars represent statistical and systematic uncertainties.

1696 5.3.13 Muon ID, isolation, and tracking efficiencies

1697 Scale factors are applied to MC to correct for differences between MC and data in
 1698 the performance of muon identification, isolation, and tracking, as detailed below.

1699 The efficiencies for muon identification measured in 2015 data and MC simulation
 1700 are shown in Figures 5.8a and 5.8b for the loose ID and tight ID respectively [80]. The
 1701 loose ID is chosen such that efficiency exceeds 99% over the full η range, and the data
 1702 and simulation agree to within 1%. The tight ID is chosen such that efficiency varies
 1703 between 95% and 99% as a function of η , and the data and simulation agree to within
 1704 1-3%. The muon identification working point used in this analysis is the medium ID,
 1705 which has an efficiency of 98% for all η and an agreement within 1-2% [80].

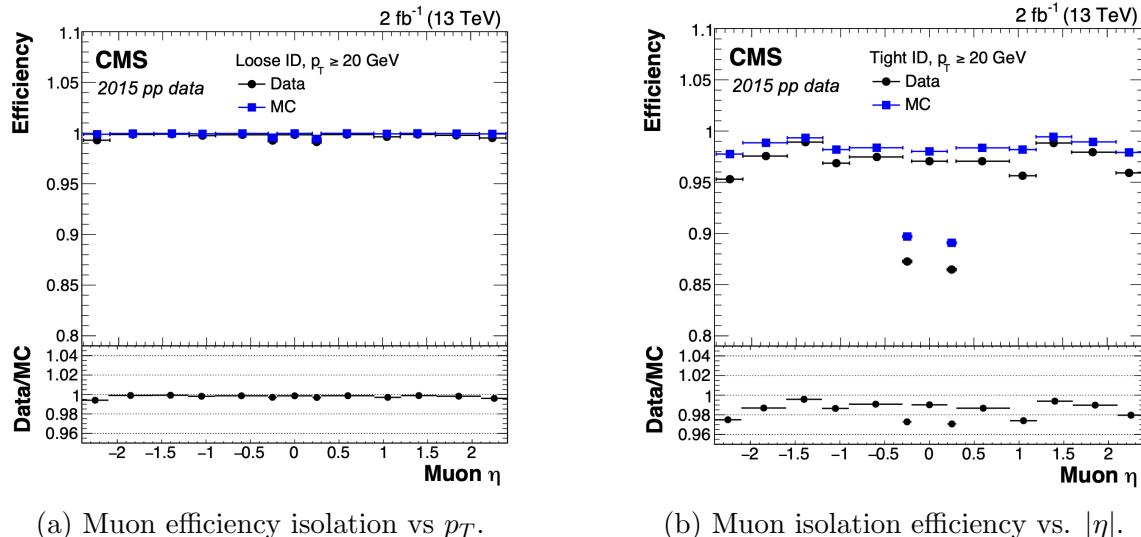
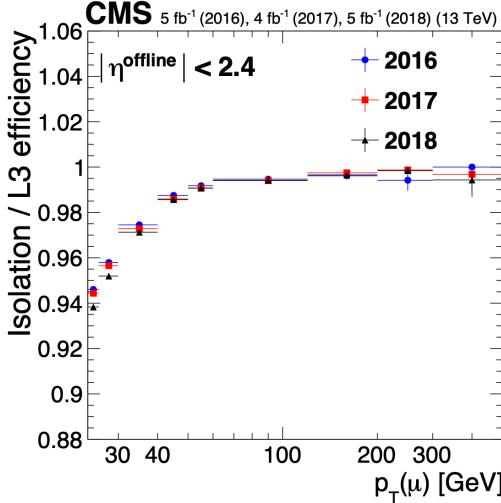


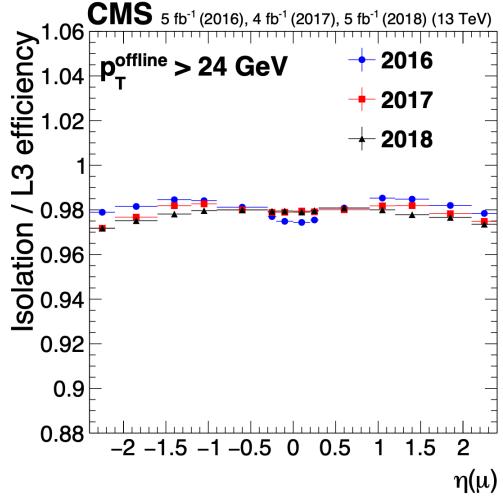
Figure 5.8: Muon identification efficiencies in 2015 data and MC as a function of the muon p_T for the loose ID (*left*) and tight ID (*right*) working points [80].

1706 The efficiencies in data for the muon isolation, as measured in Level-3 muons
 1707 (muons in one of the final stages of reconstruction in the HLT), as a function of the
 1708 muon p_T and $|\eta|$ are shown in Figures 5.9a and 5.9b [80]. The HLT muon reconstruc-
 1709 tion consists of two steps: Level-2 (L2), where the muon is reconstructed in the muon
 1710 subdetectors only, and Level-3 (L3) which is a global fit of tracker and muon hits (i.e.
 1711 the global muon reconstruction as described in Section 5.1.2) [81].

1712 The muon tracking efficiencies as a function of $|\eta|$ for standalone muons (i.e. tracks
 1713 from only the muon system, i.e. DT, CSC, and RPC, as discussed in Section 5.1.2),
 1714 is shown for data and simulated Drell-Yan samples in Fig. 5.10 [82].



(a) Muon efficiency isolation vs p_T .



(b) Muon isolation efficiency vs. $|\eta|$.

Figure 5.9: Muon isolation efficiencies in Run-2 data with respect to Level-3 muons (one of the final stages of HLT muon reconstruction) as a function of the muon p_T (*left*) and $|\eta|$ (*right*) [80].

1715 5.3.14 Recoil corrections

1716 In proton-proton collisions, W and Z bosons are predominantly produced through
1717 quark-antiquark annihilation. Higher-order processes can induce radiated quarks or
1718 gluons that recoil against the boson, imparting a non-zero transverse momentum to
1719 the boson [83]. Recoil corrections accounting for this effect are applied to samples
1720 with W+jets, Z+jets, and Higgs bosons [68]. The corrections are performed on the
1721 vectorial difference between the measured missing transverse momentum and the total
1722 transverse momentum of neutrinos originating from the decay of the W, Z, or Higgs
1723 boson. This vector is projected onto the axes parallel and orthogonal to the boson
1724 p_T . This vector, and the resulting correction to use, is measured in $Z \rightarrow \mu\mu$ events,
1725 since these events have leptonic recoil that do not contain neutrinos, allowing the
1726 4-vector of the Z boson to be measured precisely. The corrections are binned in
1727 generator-level p_T of the parent boson and also the number of jets in the event.

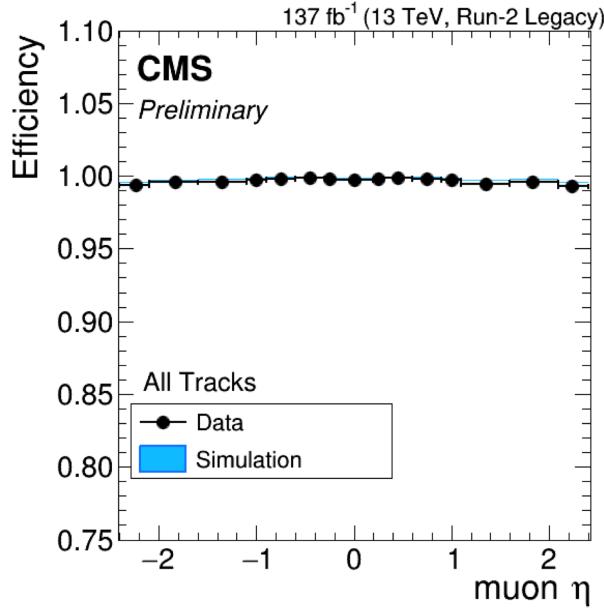


Figure 5.10: Muon tracking efficiencies as a function of $|\eta|$ for standalone muons in Run-2 data (*black*) and Drell-Yan MC simulation (*blue*) [82]. All Tracks refers to tracks which exploit the presence of muon candidates in the muon system to seed the track reconstruction in the inner tracker, in contrast to tracks that use tracker-only hits for seeding. Uncertainties shown are statistical.

5.3.15 Drell-Yan corrections

The Z boson transverse momentum distribution disagrees between leading-order (LO) simulations and data in a $Z \rightarrow \mu\mu$ control region with at least one b-tag jet [84]. Per-event weights derived by the 2016 data-only version of this analysis [84] are applied to $Z \rightarrow \tau\tau/\ell\ell$ events, as a function of the generator-level Z boson p_T to provide better matching of MC to data.

5.3.16 Pile-up reweighting

Reweighting is performed to rescale MC events to account for differences between MC and data, in the distribution of the pile-up (number of additional proton-proton interactions per bunch crossing). A tool for calculating the pile-up reweighting for the MC samples used is provided centrally by the Luminosity POG [85].

₁₇₃₉ **5.3.17 Pre-firing corrections**

₁₇₄₀ In 2016 and 2017 data-taking, a gradual timing shift of ECAL was not properly
₁₇₄₁ propagated to L1 trigger primitives (TPs), resulting in a large fraction of high η
₁₇₄₂ TPs being incorrectly associated with the previous bunch crossing. L1 trigger rules
₁₇₄₃ prevent two consecutive bunch crossings from firing, causing events to be rejected if
₁₇₄₄ significant ECAL energy was deposited in $2.0 < |\eta| < 3.0$. To account for this issue,
₁₇₄₅ MC simulations for 2016 and 2017 are corrected using an event-dependent weight.
₁₇₄₆ Embedded samples are not corrected [50].

₁₇₄₇ **5.3.18 Top p_T spectrum reweighing**

₁₇₄₈ In Run-1 and Run-2 it was observed that the p_T spectra of top quarks in $t\bar{t}$ data
₁₇₄₉ was significantly softer than those predicted by MC simulations [86]. Possible sources
₁₇₅₀ of this discrepancy are higher order QCD and/or electroweak corrections, and non-
₁₇₅₁ resonant production of $t\bar{t}$ -like final states. To account for this, corrections derived
₁₇₅₂ from Run-2 data by the Top Physics Analysis Group (PAG) are applied to the p_T
₁₇₅₃ of the top and anti-top quarks in MC simulations, computed as a function of their
₁₇₅₄ generator-level p_T [86].

₁₇₅₅ **5.3.19 B-tagging efficiency**

₁₇₅₆ In order to predict correct b-tagging discriminant distributions and event yields in
₁₇₅₇ data, the weight of selected MC events is reweighed according to recommendations by
₁₇₅₈ the BTV POG [87]. The reweighing depends on the jet p_T , η , and the b-tagging dis-
₁₇₅₉ criminant. In this method, there is no migration of events from one b-tag multiplicity
₁₇₆₀ bin to another.

1761 5.3.20 Jet energy resolution and jet energy smearing

1762 Calibration of jet energies, i.e. ensuring that the energy and momentum of the recon-
1763 structed jet matches that of the quark/gluon-initiated jet, is a challenging task due
1764 to time-dependent changes in the detector response and calibration and high pile-
1765 up [88] [89]. Jet calibration is done via jet energy corrections (JECs) applied to the
1766 p_T of jets in MC samples, accounting successively for the effects of pile-up, uniformity
1767 of the detector response, and residual data-simulation jet energy scale differences [90].
1768 Typical jet energy resolutions reported at $\sqrt{s} = 8$ TeV in the central rapidities are
1769 15-20% at 30 GeV and about 10% at 100 GeV [88]. Jet energy corrections are also
1770 propagated to the missing transverse energy.

1771 Measurements show that the jet energy resolution (JER) in data is worse than
1772 in simulation, and so the jets in MC need to be smeared to describe the data. JER
1773 corrections are applied after JEC on MC simulations, and adjust the width of the p_T
1774 distribution based on pile-up, jet size, and jet flavor [91]. Tools for applying JEC and
1775 JER are provided centrally by the JER Corrections group.

1776 Chapter 6

1777 Event selection

1778 This chapter describes how events in data and simulated samples are selected in the
1779 search for $h \rightarrow aa \rightarrow bb\tau\tau$. The event selection is motivated by optimization checks
1780 aimed at maximizing the final expected limit, and is also based on recommendations
1781 from CMS Physics Objects Groups. As described in the previous chapter, the tau
1782 lepton can decay to electrons (e), muons (μ), or hadronic states (τ_h). As a result,
1783 several different final states of the $\tau\tau$ system are possible, and are here referred to
1784 as “channels” since they are mutually exclusive. The three $\tau\tau$ final states studied in
1785 this analysis are muon and hadronic tau ($\mu\tau_h$), electron and hadronic tau ($e\tau_h$), and
1786 electron and muon ($e\mu$). The procedure for dividing events into these three channels
1787 begins with checking the High-Level Trigger paths passed by the events as detailed
1788 in Section 6.1. Events are further accepted or rejected based on criteria applied to
1789 the leptons in the event. These event selections are described for the $\mu\tau_h$ channel in
1790 Section 6.2, the $e\tau_h$ channel in Section 6.3, and the $e\mu$ channel in Section 6.4.

1791 6.1 General procedure for all channels

1792 For the search for $h \rightarrow aa \rightarrow bb\tau\tau$, three final states of the $\tau\tau$ system are considered:
1793 $\mu\tau_h$, $e\tau_h$, and $e\mu$. The $\tau_h\tau_h$ final state is not considered because signal events in the

1794 $\tau_h\tau_h$ channel would typically produce hadronic taus with momenta below data-taking
1795 trigger thresholds. In all three final states, events are required to have at least one
1796 b-tag jet passing the medium working point of the DeepFlavour tagger, with $p_T > 20$
1797 GeV, and $|\eta| < 2.4$. A second b-tag jet is not required because such a requirement
1798 would reduce signal acceptance by 80% compared to only requiring one b-tag jet.

1799 Events in MC samples are sorted into one of the three $\tau\tau$ channels if they pass the
1800 following trigger requirements and requirements on the offline reconstructed objects
1801 in the event, first checking the HLT paths for the $\mu\tau_h$ channel, then $e\tau_h$, and finally $e\mu$.
1802 The two leading leptons (e.g. muon and hadronic tau for the $\mu\tau_h$ channel) that were
1803 determined to have originated from the $\tau\tau$ decay, are called the $\tau\tau$ “legs”. For events
1804 in data and embedded samples, the HLT paths requirements for the corresponding
1805 channel are checked.

1806 After sorting events by HLT paths and identifying the leading tau legs in the offline
1807 reconstructed objects, the p_T of the offline objects is checked against the online trigger
1808 thresholds. Trigger matching is also performed, which checks the correspondence
1809 between each offline reconstructed object used in the analysis (e.g. a muon), and a
1810 trigger object in the HLT (e.g. a HLT muon). An offline object is considered to be
1811 matched, if it corresponds to a trigger object of the same object type, with $\Delta R < 0.5$.
1812 This matched trigger object is also required to pass the filter(s) of the HLT trigger.
1813 The trigger thresholds used for the $bb\mu\mu$ final state and the $bb\tau\tau$ final state (the focus
1814 of this work) are summarized in Tables 6.1.

1815 After checking the HLT paths and trigger objects in each channel, events are
1816 subject to further selection to ensure that they contain leptons and b-tag jet(s) of in-
1817 terest. These requirements are summarized in Table 6.2, and detailed in the following
1818 sections.

Year	Single/dilepton trigger p_T	$bb\mu\mu$	$bb\tau\tau$					
			$e\mu$		$e\tau_h$		$\mu\tau_h$	
		μ	e	μ	e	τ_h	μ	τ_h
2016	Single lepton	24	–	–	25	–	22	–
	p_T -leading lepton	17	23	23	–	–	–	20
	p_T -subleading lepton	8	12	8	–	–	19	–
2017	Single lepton	24	–	–	27, 32	–	24, 27	–
	p_T -leading lepton	17	23	23	–	30	–	27
	p_T -subleading lepton	8	12	8	24	–	20	–
2018	Single lepton	24	–	–	32, 35	–	24, 27	–
	p_T -leading lepton	17	23	23	–	30	–	27
	p_T subleading lepton	8	12	8	24	–	20	–

Table 6.1: Trigger thresholds used for the leptons in the $bb\mu\mu$ analysis and the $bb\tau\tau$ analysis (the focus of this work). The thresholds for the three $bb\tau\tau$ channels ($e\mu$, $e\tau_h$, and $\mu\tau_h$) are listed separately, with some channels and years taking the logical OR of two triggers with different thresholds.

6.2 Event selection in the $\mu\tau_h$ channel

In all three years, a single muon trigger is used if the muon has sufficiently high p_T , otherwise a dilepton $\mu\tau_h$ cross-trigger is used (Tables 6.3, 6.4, and 6.5). For data taken in 2017-2018 (2016), the logical OR of the single muon triggers with online p_T thresholds 24 and 27 (23) GeV is used, with the corresponding offline muon required to have with p_T 1 GeV above the online threshold. For data taken in 2017-2018 (2016), a dilepton $\mu + \tau_h$ cross-trigger with p_T thresholds of 20 (19) and 27 (20) GeV for the muon and tau respectively, is used. The τ_h is required to have $|\eta| < 2.3$ if the single trigger is fired, $|\eta| < 2.1$.

The muon and τ_h are required to have opposite charge and be separated by $\Delta R > 0.4$. The muon is required to have $|\eta| < 2.4$, and the τ_h is required to have $|\eta| < 2.3$ unless a cross-trigger is required, in which case we require $|\eta| < 2.1$ as discussed above.

The muon is required to pass the medium identification (ID) working point [92], which is defined by the Muon POG as a loose muon (i.e. a Particle Flow muon that is either a global or a tracker muon - see Section 5.1.2) with additional requirements

All years (2016, 2017, 2018) and eras				
Kinematic variable	$bb\mu\mu$		$bb\tau\tau$	
	μ	$e\mu$	$e\tau_h$	$\mu\tau_h$
ΔR between leptons	>0.4	>0.3	>0.4	>0.4
$ \eta $ of electron	-	<2.4	<2.1	-
$ \eta $ of muon	<2.4	<2.4	-	<2.1
$ \eta $ of hadronic tau	-	-	<2.3/< 2.1	<2.3/< 2.1
Relative isolation of electron	-	<0.10	-	<0.15
Relative isolation of muon	<0.25	<0.15	-	<0.15
Leading b-tag jet p_T	>15 GeV		>20 GeV	
Leading b-tag jet $ \eta $	<2.4		<2.4	
Leading b-tag jet WP	Tight		Medium	
Sub-leading b-tag jet p_T	>15 GeV		-	
Sub-leading b-tag jet $ \eta $	<2.4		-	
Sub-leading b-tag jet WP	Loose		-	
ΔR between jet(s) and leptons	>0.4		>0.5	

Table 6.2: Summary of requirements applied to the leptons in the $bb\mu\mu$ analysis and the $bb\tau\tau$ analysis (the focus of this work). $\Delta R = \sqrt{(\Delta\eta)^2 + (\Delta\phi)^2}$ is a measure of spatial separation. Relative isolation is defined in Eqn. 5.2 and Section 5.1.2. The b-tag jets are required to pass the listed DeepFlavour working points (WP), which are described in Section 5.1.5. In the $bb\tau\tau$ analysis, the required $|\eta|$ of the hadronic taus are listed for the single and cross triggers respectively. The $bb\mu\mu$ analysis requires two b-tag jets in all events, while the $bb\tau\tau$ analysis only requires one.

1835 on track quality and muon quality. This identification criteria is designed to be
1836 highly efficiently for prompt muons and for muons from heavy quark decays. In
1837 addition to the ID, for prompt muons it is recommended to apply cuts on the impact
1838 parameter [92]: we apply $|\Delta(z)| < 0.2$ and $|\Delta(xy)| < 0.045$. A cut is applied on the
1839 muon relative isolation (defined in Section 5.1.2), to be less than 0.15 in a cone size of
1840 $\Delta R = 0.4$, which corresponds to the Tight Particle Flow isolation requirement [92].

1841 The τ_h is required to pass a cut on its impact parameter of $|\Delta(z)| < 0.2$. The τ_h
1842 is also required to pass the VLoose (Very Loose) DeepTau working point vs. elec-
1843 tron, the Tight DeepTau working point vs. muons, and the VVVLoose and Medium
1844 DeepTau working point vs. jets. Events with taus reconstructed in two of the decay
1845 modes (labeled 5 and 6) are rejected, since these decay modes are meant to recover
1846 3-prong taus, but are only recommended for use in analyses where the benefits in final
1847 significance outweigh the resulting increase in background [65]. Decays reconstructed
1848 with 2 prongs are not considered as they are only recommended for taus with a very
1849 high transverse momentum, where the prongs may overlap.

1850 For the estimation of the background from jets faking τ_h , which is described in Sec-
1851 tion 7.7, anti-isolated events are selected, by requiring events to pass all the selections
1852 described above, except failing the Medium DeepTau working point vs. jets.

1853 6.3 Event selection in the $e\tau_h$ channel

1854 The HLT trigger paths for the $e\tau_h$ channel are summarized in Tables 6.3, 6.4, and
1855 6.5. Similarly to the $\mu\tau_h$ channel, a single electron trigger is used if the electron has
1856 sufficiently high p_T in 2018 and 2017. For data taken in 2018 (2017), the OR of the
1857 single electron triggers with online p_T thresholds at 32 and 35 (27 and 32) GeV are
1858 used, with the corresponding offline electrons required to have p_T greater than 33
1859 (28) GeV. A $e + \tau_h$ cross-trigger is used for electrons with lower offline p_T between

1860 25 and 33 GeV (25 and 28 GeV). For the 2016 dataset, there is no cross trigger but
1861 only a single electron trigger with online p_T threshold at 25 GeV, which is used if the
1862 offline electron has p_T greater than 26 GeV.

1863 The electron and τ_h are required to have opposite charge and be separated by
1864 $\Delta R > 0.4$. The electron is required to be within $|\eta| < 2.3$ when no cross trigger is
1865 used, and $|\eta| < 2.1$ when the cross trigger is fired. The τ_h is required to have $|\eta| < 2.3$
1866 if no cross trigger is fired, and have $|\eta| < 2.1$ if the cross trigger is fired.

1867 The electron is required to have a relative isolation (same definition as in Section
1868 5.1.2) of less than 0.1 in a cone size of $\Delta R = 0.3$, which is the standard recommended
1869 cone size giving minimal pile-up dependence and reduced probability of other objects
1870 overlapping with the cone. The isolation quantity used includes an “effective area”
1871 (EA) correction to remove the effect of pile-up in the barrel and endcap parts of the
1872 detector [93].

1873 The electron is also required to pass cuts on its impact parameter of $|\Delta(z)| < 0.2$
1874 and $|\Delta(xy)| < 0.045$. It is also required to pass the non-isolated MVA working point
1875 corresponding to 90% efficiency. The electron’s number of missing hits, which are
1876 gaps in its trajectory through the inner tracker [93], must be less than or equal to
1877 1. The electron must pass a conversion veto, which rejects electrons coming from
1878 photon conversions in the tracker, which should instead be reconstructed as part of
1879 the photon [93].

1880 The impact parameter cut for the τ_h is $|\Delta(z)| < 0.2$. In contrast to the $\mu\tau_h$ event
1881 selection, the vs. electron and vs. muon DeepTau working points are flipped, to
1882 reject muons faking the τ_h leg. The τ_h is required to pass the Tight DeepTau working
1883 point vs. electrons, the VLoose DeepTau working point vs. muons, and the Medium
1884 DeepTau working point vs. jets.

1885 As in the $\mu\tau_h$ channel, for the estimation of the background from jets faking τ_h ,
1886 which is described in Section 7.7, anti-isolated events are selected, by requiring events

1887 to pass all the selections described above, except failing the Medium DeepTau working
1888 point vs. jets.

1889 6.4 Event selection in the $e\mu$ channel

1890 The HLT trigger paths for the $e\mu$ channel are summarized in Tables 6.3, 6.4, and
1891 6.5. Events are selected with the logical OR of two $e + \mu$ cross triggers, where either
1892 the electron or muon can have larger p_T : (1) leading electron, where the electron has
1893 online $p_T > 23$ GeV and muon has online $p_T > 8$ GeV, or (2) leading muon, where
1894 electron has online $p_T > 12$ GeV and muon has online $p_T > 23$ GeV.

1895 The leading and sub-leading leptons are required to have an offline p_T greater
1896 than 1 GeV above the online threshold (i.e. $p_T > 24$ GeV). If the sub-leading lepton
1897 is the electron, the offline p_T threshold is 1 GeV above the online threshold ($p_T > 13$
1898 GeV), but if it is a muon, the offline p_T threshold is required to be at least 5 GeV
1899 greater than the online threshold (i.e. $p_T > 13$ GeV). This is because of poor data
1900 and simulation agreement for low- p_T muons with p_T between 9 GeV and 13 GeV, and
1901 the higher probability of mis-identifying jets as muons at lower p_T . With no effect on
1902 the expected limits, the offline p_T threshold for muons is raised to 13 GeV instead of
1903 9 GeV, even though it may lead to loss in signal acceptance. Both the electron and
1904 muon are required to be have $|\eta| < 2.4$.

1905 The electron and muon are required to have opposite charge and be separated
1906 by $\Delta R > 0.3$ (note the decreased separation requirement compared to the other
1907 two channels). The electron is required to pass the non-isolated MVA identification
1908 working point corresponding to 90% efficiency, and to have a relative isolation less
1909 than 0.1 for a cone size of $\Delta R = 0.3$ with the EA pile-up subtraction correction.
1910 The electron must have one or fewer missing hits and pass the conversion veto (both
1911 described previously in Section 6.3).

1912 The muon is required to pass the medium identification working point (described
 1913 earlier in 6.2), and to have a relative isolation less than 0.15 for a cone size of $\Delta R =$
 1914 0.4. The muon impact parameter is required to have $|\Delta(z)| > 0.2$ and $|\Delta(xy)| < 0.045$.
 1915 For the QCD multijet background estimation described in Section 7.8, the same-
 1916 sign region is selected by requiring all the above selections, except the legs are required
 1917 to have the same electric charge rather than opposite.

2016 $\mu\tau_h$ trigger paths	
Notes	HLT Path
	HLT_IsoMu22_v
	HLT_IsoMu22_eta2p1_v
	HLT_IsoTkMu22_v
	HLT_IsoTkMu22_eta2p1_v
	HLT_IsoMu19_eta2p1_LooseIsoPFTau20_v
	HLT_IsoMu19_eta2p1_LooseIsoPFTau20_SingleL1_v
2016 $e\tau_h$ trigger paths	
Notes	HLT Path
	HLT_Ele25_eta2p1_WPTight_Gsf_v
2016 $e\mu$ trigger paths	
Notes	HLT Path
runs B-F and MC	HLT_Mu23_TrkIsoVVL_Ele12_CaloIdL_TrackIdL_IsoVL_v
runs B-F and MC	HLT_Mu8_TrkIsoVVL_Ele23_CaloIdL_TrackIdL_IsoVL_v
runs G-H	HLT_Mu23_TrkIsoVVL_Ele12_CaloIdL_TrackIdL_IsoVL_DZ_v
runs G-H	HLT_Mu8_TrkIsoVVL_Ele23_CaloIdL_TrackIdL_IsoVL_DZ_v

Table 6.3: High-Level Trigger (HLT) paths used to select data and simulation events in 2016 for the three $\tau\tau$ channels.

1918 6.5 Extra lepton vetoes in all channels

1919 Events containing a third lepton (electron or muon) that is neither of the leading $\tau\tau$
 1920 legs are rejected, and events with di-muons and di-electrons are vetoed, with criteria
 1921 taken from the Standard Model $H \rightarrow \tau\tau$ working group [68]. These vetoes on extra
 1922 leptons also ensure orthogonality of events to analyses such as the $bb\mu\mu$ final state,
 1923 whose results are combined with this $bb\tau\tau$ final state as described in Section 10.2.

2017 $\mu\tau_h$ trigger paths	
Notes	HLT Path
	HLT_IsoMu24_v
	HLT_IsoMu27_v
	HLT_IsoMu20_eta2p1_LooseChargedIso_PFTau27_eta2p1_CrossL1_v
2017 $e\tau_h$ trigger paths	
Notes	HLT Path
	HLT_Ele32_WPTight_Gsf_v
	HLT_Ele35_WPTight_Gsf_v
	HLT_Ele24_eta2p1_WPTight_Gsf_Loose_ChargedIsoPFTau30_eta2p1_CrossL1_v
2017 $e\mu$ trigger paths	
Notes	HLT Path
	HLT_Mu23_TrkIsoVVL_Ele12_CaloIdL_TrackIdL_IsoVL_DZ_v
	HLT_Mu8_TrkIsoVVL_Ele23_CaloIdL_TrackIdL_IsoVL_DZ_v

Table 6.4: High-Level Trigger (HLT) paths used to select data and simulation events in 2017 for the three $\tau\tau$ channels.

1924 The event is vetoed if a third electron is found with the following properties:
 1925 $p_T > 10$ GeV, $|\eta| < 2.5$, impact parameter $|\Delta(z)| < 0.2$ and $|\Delta(xy)| < 0.045$, passing
 1926 non-isolation MVA identification with 90% efficiency, conversion veto, ≤ 1 missing
 1927 hits, and relative isolation < 0.3 with cone size $\Delta R = 0.3$. The event is also vetoed if
 1928 a third muon is found with the following properties: $p_T > 10$ GeV, $|\eta| < 2.4$, impact
 1929 parameter $|\Delta(z)| < 0.2$ and $|\Delta(xy)| < 0.045$, medium ID, and isolation < 0.3 with
 1930 cone size $\Delta R = 0.4$.

1931 A di-muon veto is applied, which rejects events containing a pair of muons with
 1932 opposite charge and separation of $\Delta R > 0.15$, that both pass the following selections:
 1933 $p_T > 15$ GeV, $|\eta| < 2.4$, flag for global muons, flag for tracker muon, flag for Particle
 1934 Flow muon, $|\Delta(z)| < 0.2$, $|\Delta(xy)| < 0.045$, and isolation < 0.3 with cone size $\Delta R =$
 1935 0.4. A similar di-electron veto is applied to reject events containing a pair of electrons
 1936 with opposite charge and separation of $\Delta R > 0.15$, that both pass the following
 1937 selections: $p_T > 15$ GeV, $|\eta| < 2.5$, a dedicated electron ID (cut-based) for vetoing
 1938 third leptons, $|\Delta(z)| < 0.2$, $|\Delta(xy)| < 0.045$, with pile-up corrected relative isolation
 1939 < 0.3 with cone size $\Delta R = 0.3$.

2018 $\mu\tau_h$ trigger paths	
Notes	HLT Path
	HLT_IsoMu24_v
	HLT_IsoMu27_v
only data run < 317509	HLT_IsoMu20_eta2p1_ (contd.)
	LooseChargedIsoPFTauHPS27_eta2p1_CrossL1_v
MC and data run \geq 317509	HLT_IsoMu20_eta2p1_ (contd.)
	LooseChargedIsoPFTauHPS27_eta2p1_TightID_CrossL1_v
2018 $e\tau_h$ trigger paths	
Notes	HLT Path
	HLT_Ele32_WPTight_Gsf_v
	HLT_Ele35_WPTight_Gsf_v
only data run < 317509	HLT_Ele24_eta2p1_WPTight_Gsf_ (contd.)
	LooseChargedIsoPFTauHPS30_eta2p1_CrossL1_v
MC and data run \geq 317509	HLT_Ele24_eta2p1_WPTight_Gsf_ (contd.)
	LooseChargedIsoPFTauHPS30_eta2p1_TightID_CrossL1_v
2018 $e\mu$ trigger paths	
Notes	HLT Path
	HLT_Mu23_TrkIsoVVL_Ele12_CaloIdL_TrackIdL_IsoVL_DZ_v
	HLT_Mu8_TrkIsoVVL_Ele23_CaloIdL_TrackIdL_IsoVL_DZ_v

Table 6.5: High-Level Trigger (HLT) paths used to select data and simulation events in 2018 for the three $\tau\tau$ channels. In 2018 a HLT trigger path using the hadron plus strips (HPS) tau reconstruction algorithm became available.

¹⁹⁴⁰ Chapter 7

¹⁹⁴¹ Background estimation

¹⁹⁴² This section describes methods used to estimate sources of background from Standard
¹⁹⁴³ Model processes in the search for $h \rightarrow aa \rightarrow bb\tau\tau$. Similar background estimation
¹⁹⁴⁴ methods are being used for the $h \rightarrow a_1a_2$ analysis. The background contributions
¹⁹⁴⁵ directly taken from MC are described in Sections 7.1 to 7.6. Section 7.7 describes
¹⁹⁴⁶ the data-driven method for estimating backgrounds from jets faking hadronic tau
¹⁹⁴⁷ decays ($\text{jet} \rightarrow \tau_h$), which is used in the $\mu\tau_h$ and $e\tau_h$ channels. Section 7.8 describes
¹⁹⁴⁸ the data-driven method for estimating background from quantum chromodynamic
¹⁹⁴⁹ (QCD) processes in the $e\mu$ channel.

¹⁹⁵⁰ 7.1 Z+jets

¹⁹⁵¹ A major source of background for $\tau\tau$ analyses is the Drell-Yan (DY) process (Z+jets).
¹⁹⁵² The Z boson decays to $\tau\tau/\mu\mu/ee$ with equal probability of 3.4% each, with the domi-
¹⁹⁵³ nant decay modes being to hadrons (around 70%) and neutrinos (invisible) (20%) [26].
¹⁹⁵⁴ The Drell-Yan contribution with genuine taus, $Z \rightarrow \tau\tau$, is estimated using embed-
¹⁹⁵⁵ ded samples, described in Section 4.3. To avoid double-counting between embedded
¹⁹⁵⁶ and MC samples, in all MC samples, events with legs that originated from genuine τ
¹⁹⁵⁷ are discarded.

1958 The other decays of the Z , $Z \rightarrow ee$ and $Z \rightarrow \mu\mu$, are estimated from MC simulation,
1959 and are hereafter referred to as simply the Drell-Yan background. These MC samples
1960 are generated to leading order (LO) with different numbers of jets (jet multiplicity) in
1961 the matrix element: $Z+1$ jet, $Z+2$ jets, $Z+3$ jets, $Z+4$ jets, and inclusive $Z+jets$. The
1962 cross-sections of the samples with ≥ 1 jets are normalized to next-to-NLO (NNLO)
1963 in QCD.

1964 For the inclusive Drell-Yan sample, two samples are used with different thresholds
1965 for the di-lepton invariant mass ($m_{\ell\ell}$) at the generator level: one with $m_{\ell\ell} > 50$ GeV
1966 and the other with $10 < m_{\ell\ell} < 50$.

1967 7.2 W+jets

1968 The dominant W boson decay modes are to hadrons (67.4%), $e + \nu_e$ (10.7%), $\mu + \nu_\mu$
1969 (10.6%), and $\tau + \nu_\tau$ (11.4%) [26]. The $W+jets$ background is estimated from MC
1970 simulation. Similarly to the $Z+jets$, the $W+jets$ samples are generated with different
1971 jet multiplicities in the matrix element. LO samples are used for greater statistics
1972 and are normalized to NNLO cross sections.

1973 7.3 $t\bar{t} + jets$

1974 In hadron collisions, top quarks are produced singly with the weak interaction, or in
1975 pairs via the strong interaction, with interference between these leading-order pro-
1976 cesses possible in higher orders of the perturbation theory. The top quark is the
1977 heaviest fermion in the Standard Model and has a short lifetime ($\sim 10^{-25}$ s), decay-
1978 ing without hadronization into a bottom quark and a W boson [26], with the decay
1979 modes of the W boson as listed in the previous section. With two top quarks, the
1980 final states of the two resulting W bosons can be described as fully leptonic, semilep-
1981 tonic, and fully hadronic. These three final states are modeled separately with MC

1982 simulation in 2018 and 2017, while for 2016 the sample used is inclusive.

1983 7.4 Single top

1984 There are three main production modes of the single top in pp collisions [94]: the
1985 exchange of a virtual W boson (t channel), the production and decay of a virtual W
1986 boson (s channel), and the associated production of a top quark and W boson (tW ,
1987 or W-associated) channel. As the s channel process is rare and only 3% of the total
1988 production, the dominant production mode of the t -channel and the tW production
1989 are considered and modeled with MC.

1990 7.5 Diboson

1991 In pp collisions, the production of dibosons (pairs of electroweak gauge bosons, i.e.
1992 WW, WZ, and ZZ) is dominated by quark-antiquark annihilation, with a small con-
1993 tribution from gluon-gluon interaction [95]. MC is used to model the pair production
1994 and decays of VV to $2\ell 2\nu$, WZ to $2q 2\ell$ and $3\ell\nu$, and ZZ to 4ℓ and $2q 2\ell$ (q being
1995 quarks and ℓ being leptons).

1996 7.6 Standard Model Higgs

1997 MC is used to simulate backgrounds from major production modes of the Standard
1998 Model 125 GeV Higgs boson: gluon-gluon fusion (ggH), vector boson fusion (VBF),
1999 associated production with a W or Z (WH, ZH), and associated production with a
2000 top pair (ttH) (see Fig. 7.1 for leading-order diagrams). For these production modes,
2001 samples with the Higgs decaying to $\tau\tau$ or to WW are used. Samples made with
2002 higher-order diagrams for WH and ZH that include the production of a jet, with the
2003 Higgs decaying to WW, are also used.

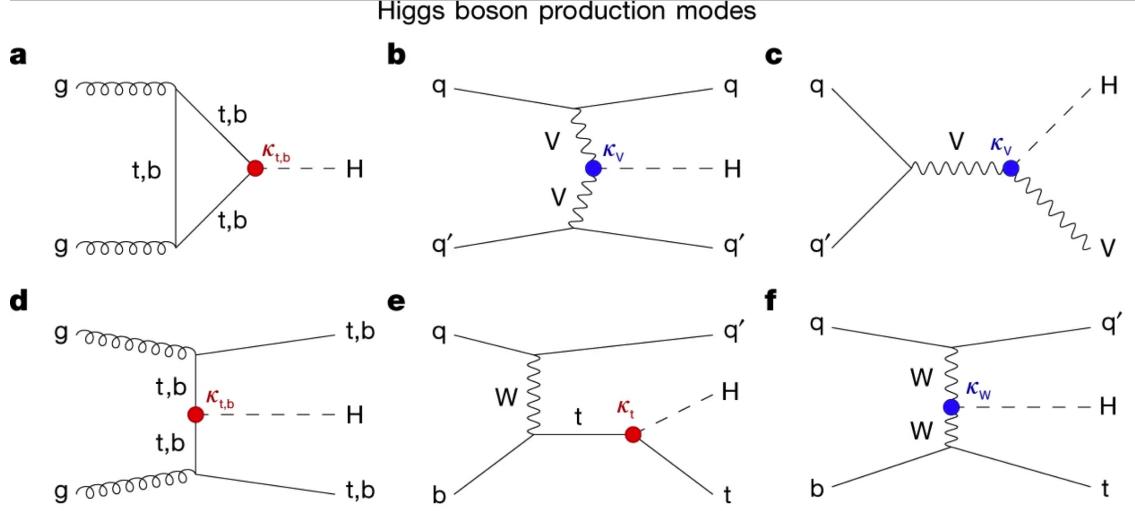


Figure 7.1: Leading-order Feynman diagrams of Higgs production from [96], in ggH (a) and vector boson fusion (VBF; b), associated production with a W or Z (V) boson (VH; c), associated production with a top or bottom quark pair (ttH or bbH); d, and associated production with a single top quark (tH; e, f).

2004 7.7 Jet faking τ_h

Events with a jet mis-reconstructed as the hadronic tau leg τ_h are a major source of background in the $\mu\tau_h$ and $e\tau_h$ channels. The main processes contributing to jet $\rightarrow \tau_h$ events are QCD multijet, W+jets, and $t\bar{t}$ production. These events are estimated using a data-driven method adapted from past analyses [50] [84]. This background includes contributions from W+jets, QCD multijets, and $t\bar{t}$ +jets. To estimate this background, a sideband region is constructed, where events are required to pass all baseline $\mu\tau_h/e\tau_h$ selection criteria, but fail the τ_h isolation criteria. The events in this sideband region are reweighed with a factor $f/(1 - f)$, where f is the probability for a jet to be misidentified as a τ_h . The jet $\rightarrow \tau_h$ background is the anti-isolated, reweighed MC and embedded events subtracted from the anti-isolated, reweighted data events.

The fake factor is measured in $Z \rightarrow \mu\mu + \text{jets}$ events in data in the $\mu\mu\tau_h$ final state, as any reconstructed τ_h in these events must originate from a jet. The two muons are required to be isolated (< 0.15), have opposite electric charge, and have

2019 an invariant mass between 76 and 106 GeV (close to the Z mass). These events are
2020 selected with a double muon trigger, with the leading muon having offline $p_T > 20$
2021 GeV and the subleading muon $p_T > 10$ GeV. Simulated diboson (ZZ and WZ) events
2022 are subtracted to avoid contamination from events with real τ_h . The denominator of
2023 the fake rate corresponds to fake taus passing the VVVLoose working point of the
2024 discriminator vs. jets, while the numerator corresponds to those passing the Medium
2025 working point, i.e. $f = N_{\text{jet passing tight}} / N_{\text{jet passing loose}}$.

2026 f is measured as a function of the τ_h transverse momentum and is 8% - 10% in
2027 each of the data-taking years. f is derived separately for the $\mu\tau_h$ and $e\tau_h$ channels
2028 because the channels use different anti-lepton identification working points.

2029 7.8 QCD multijet background

2030 In the $e\mu$ channel, events with jets faking electrons or muons originating from QCD
2031 multijet, is estimated from data events with the same baseline selection as in the
2032 signal region, except with same-signed (SS) charged $e + \mu$, ensuring orthogonality
2033 with the signal region which requires opposite-sign (OS) $e\mu$ pairs. All same-sign MC
2034 events (both events with real and fake $e + \mu$) are subtracted from same-sign data
2035 events to remove contamination from other backgrounds. i.e. $\text{QCD}_{\text{SS}} = \text{Data}_{\text{SS}} -$
2036 MC_{SS} .

2037 Three scale factors are applied to the QCD_{SS} events to compute the QCD multijet
2038 background [84] [40]:

- 2039 • *OS-to-SS scale factor*: This scales the SS QCD to the OS region, and is mea-
2040 sured from an orthogonal region with an isolated electron and an anti-isolated
2041 muon. Only the muon is chosen to be anti-isolated because this scale factor was
2042 observed to depend more strongly on electron isolation than that of the muon.
2043 This scale factor is treated as a function of the ΔR separation of the trajectories

2044 of the electron and muon, and is measured separately for events with 0 jets, 1,
2045 jet, and greater than 1 jet.

- 2046 • *2D closure correction for the lepton p_T :* This factor accounts for subleading
2047 dependencies of the first scale factor on the p_T of the two leptons. A 2D weight
2048 is derived in a similar fashion, as a ratio of QCD_{OS} events to QCD_{SS} events,
2049 but parameterized by both electron and muon p_T , where the SS events have the
2050 previous scale factor applied.
- 2051 • *Isolation correction for the muon:* The third and final factor is an isolation
2052 correction, which is a bias correction to account for the fact that the fake
2053 factor was determined for less-isolated muons. This factor is obtained as the
2054 ratio of the OS-to-SS scale factors measured in two other control regions: (1)
2055 events where the electron is anti-isolated ($0.15 < \text{iso} < 0.5$) and the muon is
2056 isolated, and (2) events where both leptons are anti-isolated.

2057

Chapter 8

2058

Systematic uncertainties

2059 Uncertainties in the measurement of a physical observable can be statistical or sys-
2060 tematic in nature. Statistical uncertainties originate from limitations on the number
2061 of events and experiments that can be performed. Systematic uncertainties arise
2062 from the dependence of the physical observable on quantities whose exact values are
2063 unknown and which can only be modeled imperfectly.

2064 The handling of systematic uncertainties is separated into normalization uncer-
2065 tainties (those that affect the total yield of a variables' distribution) and shape un-
2066 certainties (those that shift the distribution of events). Normalization uncertainties
2067 are expressed as multiplicative factors, while shape uncertainties are represented as
2068 up and down shifts of a variable's distribution.

2069 Up/down shifts of shape uncertainties can change the number of background
2070 events in a distribution. For instance, hadronic taus receive corrections from the
2071 nominal tau energy scale, with the nominal, up, and down energy scales provided
2072 centrally by CMS. For the $\mu\tau_h$ channel, an event could have a τ_h with p_T just below
2073 the offline threshold of 20 GeV (for instance, 19.5 GeV), so in the nominal distribution
2074 of $m_{\tau\tau}$ (or any other variable for this channel), the event is excluded. However, when
2075 we build our distributions with the tau energy scale “up” shift, the energy of this τ_h

2076 may be scaled up to, say, 20.5 GeV, and now the event passes the offline p_T threshold
2077 for the single muon trigger, leading to the event’s inclusion in the distributions made
2078 with the tau energy scale “up” shift.

2079 In evaluating the up and down shifts of a specific source of uncertainty, all other
2080 corrections and scale factors are held at their nominal values, and the full chain of
2081 object and event selection and event categorization is performed to obtain the observ-
2082 able distributions. Any “downstream” variables that depend on the shifted variable,
2083 e.g. the invariant di-tau mass $m_{\tau\tau}$, must be computed for the nominal case, and then
2084 re-computed separately for each up and down shift of the tau legs’ energy scale. The
2085 objective of this process is to quantify the effect of a single source of uncertainty on
2086 the resulting observable distributions. Each scale factor and correction described in
2087 Section 5.3 has an associated uncertainty. The binning of the uncertainties follows
2088 that of the nominal scale factor value.

2089 Sections 8.1 to 8.5 describe uncertainties associated with physics objects, and
2090 Sections 8.6 and 8.7 describe uncertainties associated with sample-level effects. The
2091 pulls and impacts for the top sixty most important systematics are shown in Section
2092 8.8.

2093 8.1 Uncertainties in the lepton energy scales

2094 The uncertainties in the tau energy scales [65] are binned by the tau decay mode and
2095 are taken as shape uncertainties treated as uncorrelated across the tau decay modes
2096 and years. Same as with the application of the nominal scale factor, when applying
2097 the up or down shifts, the missing transverse energy (p_T^{miss}) of the event is adjusted
2098 so that the 4-vector sum of the tau p_T^{miss} is unchanged.

2099 The uncertainties in the muon energy scale [66] are 0.4% for $|\eta| < 1.2$, 0.9% for
2100 $1.2 < |\eta| < 2.1$, and 2.7% for $2.1 < |\eta| < 2.4$, and are treated as shape uncertainties,

2101 fully uncorrelated between embedded and MC samples.

2102 The uncertainties in the electron energy scale [69] in MC are binned in the electron
2103 $|\eta|$ and p_T , and are shown in Fig. 5.2. The uncertainties range from 0.5% to 2.2% in
2104 the barrel, and 0.3% to 4.1% in the endcap, across the p_T range. The uncertainties
2105 for the embedded sample are binned only in $|\eta|$ and are on the order of 0.5% and
2106 1.25% for the barrel and endcap [73].

2107 There are also uncertainties in the energy scales for electrons and muons misiden-
2108 tified as τ_h . The uncertainty for muons misidentified as τ_h is 1% [65]. For electrons
2109 misidentified as τ_h , the uncertainty is binned in barrel/endcap η and by 1-prong and
2110 1-prong + π_0 decays. The probability for e/μ faking a 3-prong decay mode is much
2111 lower.

2112 8.2 Uncertainties from other lepton corrections

2113 Uncertainties associated with the τ_h identification efficiencies are treated as shapes,
2114 uncorrelated across the seven p_T bins and years. The shape uncertainties in the
2115 embedded samples are taken as 50% correlated with those of the MC samples.

2116 The uncertainties on electron and muon identification efficiencies are taken as
2117 normalization uncertainties of 2% each, with a 50% correlation between embedded
2118 and MC samples.

2119 In the $e\tau_h$ channel, there is an additional uncertainty for the vs. jet discrimination
2120 efficiency [65], because the analysis uses a looser anti-lepton working point (VLoose
2121 WP) than the working points used in the measurement of the efficiency (namely,
2122 VLoose WP vs e, and Tight WP vs mu). For nominal $\tau_h p_T < 100$ GeV, an additional
2123 uncertainty of 3% (5%) is used in MC (embedded), and for high p_T an uncertainty of
2124 15% is used for both.

2125 The uncertainties in trigger efficiencies are taken as shapes [65]. In the $e\tau_h$ and $\mu\tau_h$

2126 channels, there are uncertainties for the single and cross lepton triggers, and in the
2127 $e\mu$ channel there is one uncertainty each for the two $e + \mu$ triggers, and one combined
2128 uncertainty since their trigger phase spaces are not mutually exclusive.

2129

8.3 Uncertainties from jet energy scale and reso- 2130 lution

2131 The jet energy scale uncertainties are taken as shape uncertainties: there are eleven
2132 in total, with seven correlated across years (labeled “Year” below) and the remainder
2133 uncorrelated across years. They affect the b-tag jet p_T and mass, and hence the
2134 missing transverse energy p_T^{miss} . The shifts are propagated through the b-tagging
2135 scale factor calculation and b-tag jet counting.

2136 The uncertainties in the jet energy correction and resolution [88] [97] are as follows:

- 2137
 - *Absolute, AbsoluteYear*: flat absolute scale uncertainties.
- 2138
 - *BBEC1, BBEC1Year*: for sub-detector regions, with barrel “BB” in $|\eta| < 1.3$
2139 and endcap region 1 “EC1”: $1.3 < |\eta| < 2.5$.
- 2140
 - *EC2, EC2 year*: for sub-detector regions, with endcap region 2 “EC2” in $2.5 <$
2141 $|\eta| < 3.0$.
- 2142
 - *HF, HF year*: for sub-detector regions, with hadron forward “HF” in $|\eta| > 3$.
- 2143
 - *FlavorQCD*: for uncertainty in jet flavor (uds/c/b-quark and gluon) estimates
2144 based on comparing Pythia and Herwig (different MC generator) predictions.
- 2145
 - *RelativeBal*: account for difference between log-linear fits of the two methods
2146 used to study the jet energy response: MPF (missing transverse momentum
2147 projection fraction) and p_T balance.

- 2148 • *RelativeSample*: account for η -dependent uncertainty due to a difference be-
 2149 tween relative residuals, observed with dijet and Z+jets in Run D of 2018 data.
- 2150 • *JetResolution*: uncertainty in the jet energy resolution.

2151 8.4 Uncertainties from b-tagging scale factors

2152 The b-tagging scale factor has its own set of associated uncertainties (not to be
 2153 confused with shifts in the b-tagging scale factor due to the propagation of the jet
 2154 energy scale uncertainties described in the previous section 8.3). They are:

- 2155 • *hf*: contamination from heavy flavor (b+c) jets in the light flavor region.
- 2156 • *hfstats1, hfstats2*: linear and quadratic statistical fluctuations from b-flavor jets.
- 2157 • *lf*: contamination from light flavor (udsg+c jets) in the heavy flavor region.
- 2158 • *lfstats1, lfstats2*: linear and quadratic statistical fluctuations from udsg jets.
- 2159 • *cferr, cferr2*: uncertainty for charm jets.

2160 The variations for “lf, hf, hfstats1/2, lfstats1/2” are applied to both b and udsg jets.
 2161 For c-flavor jets, only “cferr1/2” is applied.

2162 8.5 Uncertainties from MET

2163 Samples where recoil corrections were applied (Z+jets, W+jets, and Standard Model
 2164 Higgs, as described in Section 5.3) have uncertainties from the response and resolution
 2165 of the hadronic recoil against the leptonic system. These are each binned in jet
 2166 multiplicity.

2167 8.6 Uncertainties associated with samples used

2168 Normalization uncertainties related to the samples used are:

- 2169 • *Cross-section uncertainties*: $\sigma(t\bar{t})$: 4.2%, $\sigma(\text{diboson})$: 5%, $\sigma(\text{single top})$: 5%,
2170 $\sigma(\text{ggH})$: 3.2%, $\sigma(\text{qqH})$: 2.1%, $\sigma(\text{WH})$: 1.9%, $\sigma(\text{ZH})$: 1.3%, $\sigma(\text{ttH})$: 3.6%
- 2171 • *Uncertainties in QCD renormalization scale*: QCD scale(qqH): +0.43%-0.33%,
2172 QCD scale(WH): +0.5%-0.7%, QCD scale(ttH): +5.8%-9.2%
- 2173 • *Branching ratio uncertainties*: $\text{BR}(\text{H} \rightarrow \tau\tau)$: 1.8%, and $\text{BR}(\text{H} \rightarrow \text{WW})$: 1.5%.
- 2174 • *Normalization uncertainties*: 2% for Drell-Yan, 4\$ for embedded, 20% pre-fit
2175 for the QCD multijet background in the $e\mu$ channel, 20% pre-fit for the jet
2176 faking background.

2177 The $t\bar{t}$ process has additional acceptance uncertainties from QCD scale variation
2178 and parton shower uncertainties [98]. Parton shower uncertainties originate from
2179 the modeling of perturbative and non-perturbative QCD effects handled in parton
2180 shower MC generators. The scale variations are determined from the envelope of the
2181 6 provided shapes due to variations in the factorization scale, renormalization scale,
2182 and their combined variation [98].

2183 The uncertainty in the Z p_T reweighting in Drell-Yan samples is taken as a shape
2184 uncertainty and the up and down values are 0.9 and 1.1 times the nominal reweighting.
2185 This 10% uncertainty is sufficient to cover uncertainties in the weights derived from
2186 the discrepancies between LO simulations and data in the di-muon mass in $Z \rightarrow \mu\mu$
2187 events.

2188 The weight applied to anti-isolated events in the $\mu\tau_h$ and $e\tau_h$ channels to estimate
2189 the background from jets faking τ_h , has shape uncertainties covering uncertainties in
2190 the derivation of the weight. There are six shape uncertainties corresponding to the
2191 binning of the fake rate in the τ_h transverse momentum. For the weight applied to

2192 scale up anti-isolated events in cross-trigger regions, 20% of the nominal weight is
2193 taken as a shape uncertainty.

2194 8.7 Other uncertainties

2195 A 3.6% yield uncertainty in the signal is used to cover uncertainties in the parton
2196 distribution functions (PDFs), knowledge of the α_s (fine structure constant), and
2197 QCD scale. The size of these uncertainties was estimated by a different analysis
2198 searching for two light scalars decaying to four muons, which compared the PDFs
2199 from different model libraries using recommendations from the PDF4LHC Working
2200 Group [99] [100].

2201 Uncertainties in the luminosity measurements can originate from uncertainties
2202 in the luminosity calibration in the van de Meer scan procedure and from detector
2203 operations [43]. Some effects are fully uncorrelated (e.g. if the systematic error is
2204 limited by the statistical uncertainty in the calibration scans taken independently in
2205 each year), and some are correlated, for example in the 2017 and 2018 measurements
2206 which used a method with the same systematic bias. The luminosity normalization
2207 uncertainties are applied all MC samples, divided into those uncorrelated across years
2208 (0.26% for 2016, 0.60% for 2017, and 0.65% for 2018), one correlated between 2017
2209 and 2018 (0.27%), and one correlated between all three years (1.30%) [41] [42] [43] [85].

2210 8.8 Pulls and impacts

2211 The top impacts and pulls computed for the combination of all channels and years is
2212 shown in Fig. 8.1. The top impacts are related to uncertainty in the signal sample and
2213 cross-section of the $t\bar{t}$ cross-section, and also the yields of the jet faking τ_h background,
2214 which is a major background in all channels and expected to be constrained due to
2215 the yield uncertainty which is taken to be 20% pre-fit.

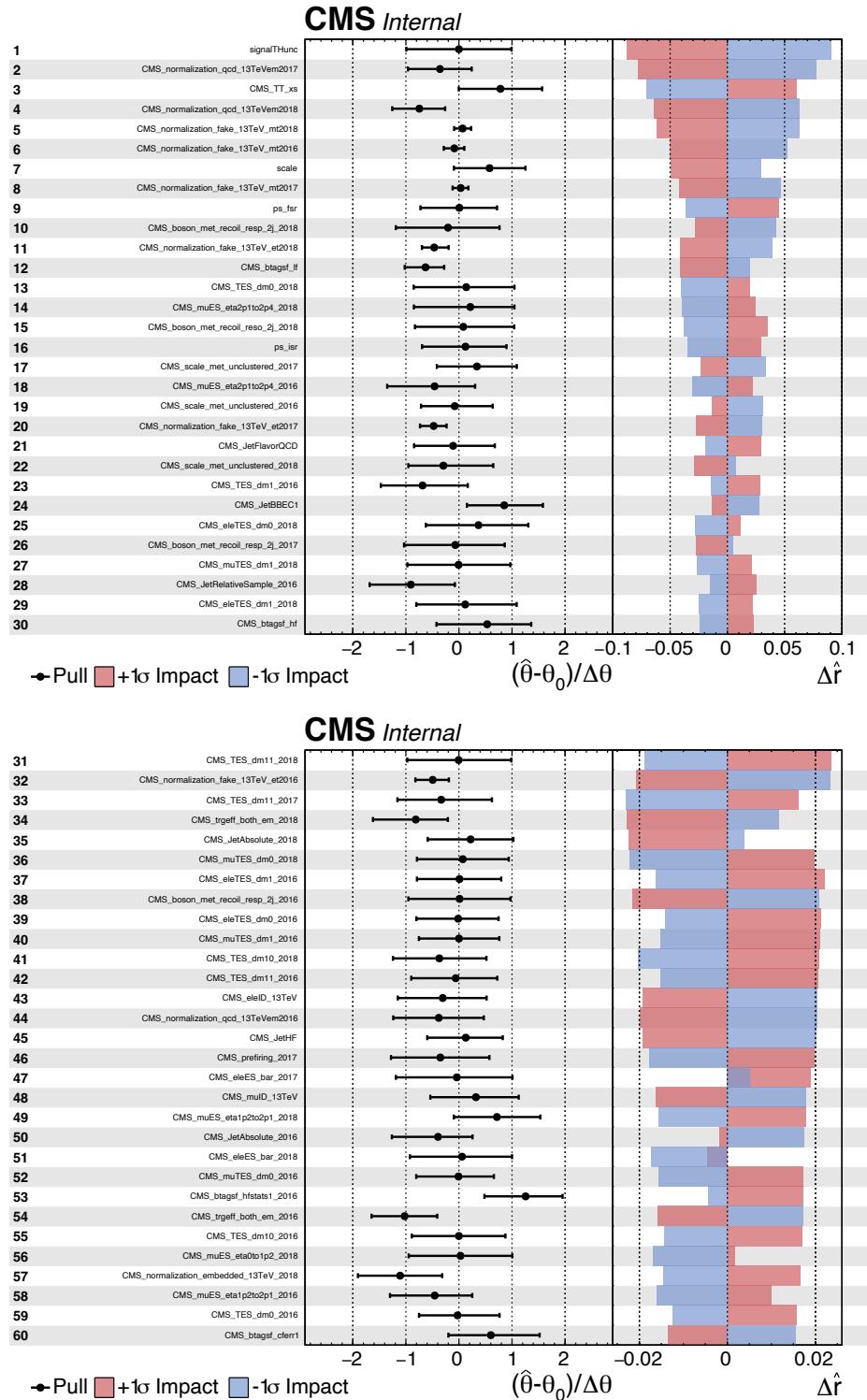


Figure 8.1: Top sixty pulls and impacts for the combination of all channels and years [101].

2216 **Chapter 9**

2217 **Event categorization and signal
2218 extraction**

2219 Measured events are divided into categories, based on cuts on values of observables
2220 in the event, or some derived quantity based on the observables in the event. The
2221 objective of event categorization is to divide events into signal regions, where the
2222 signal is enhanced and the background is suppressed, and control regions, which are
2223 signal-poor and used to check that the background estimation methods employed in
2224 the analysis in fact accurately models the data. In this analysis, events in each tau-tau
2225 channel are selected to contain one or more b-tag jets reconstructed in the event as
2226 described in Section 9.1. Events are further divided into signal and control regions
2227 using a deep learning-based approach described in Section 9.2. The signal is extracted
2228 from the di-tau mass distribution in the signal region using the statistical procedure
2229 described in Section 9.3.

2230 **9.1 B-tag jet multiplicity**

2231 The increased statistics of the full Run-2 dataset enables the separation of events into
2232 events with exactly 1 b-tag jet and events with greater than 1 b-tag jet. Further event

2233 categorization is performed with deep neural networks (DNNs) described below. The
2234 DNNs are used only for separating events into signal and control regions in the 1
2235 b-tag and 2 b-tag jets scenarios. The final results are extracted from the statistical
2236 fitting to the mass of the $\tau\tau$, $m_{\tau\tau}$.

2237 9.2 DNN-based event categorization

2238 Neural networks for event categorization are trained for each of the $\mu\tau_h$, $e\tau_h$, and $e\mu$
2239 channels, for 1 and 2 b-tag jets, giving $3 \times 2 = 6$ networks in total. In the training,
2240 the signal is taken to be all of the possible pseudoscalar mass m_a hypotheses together.
2241 The backgrounds for each DNN are taken to be a representative combination of the
2242 three major backgrounds: $Z \rightarrow \tau\tau$, $t\bar{t} + \text{jets}$, and fake backgrounds. The proportions of
2243 each background for each channel and b-tag jet multiplicity are taken from the yields
2244 in the $m_{\tau\tau}$ distribution. For instance, in the $\mu\tau_h$ 1 b-tag jet category, the composition
2245 of the background for training is 17.4% from $Z \rightarrow \tau\tau$, 42.4% from $t\bar{t} + \text{jets}$, and 40.2%
2246 fakes.

2247 The input variables capture the key differences between the signal and the back-
2248 ground:

- 2249 • Transverse momentum p_T of the electron and muon in the $e\tau_h$ and $\mu\tau_h$ channels,
2250 where the signal tends to have a softer p_T spectrum (lower energy) than the
2251 background.
- 2252 • p_T of the b-tag jet(s). The signal sample b-tag jet(s) tend to have softer p_T .
- 2253 • Invariant masses of the various objects ($\tau\tau$ legs and the b-tag jet(s)), which
2254 tend to be smaller for the signal samples.
- 2255 • The angular separation ΔR between pairs of the objects, where signal samples
2256 peak at smaller ΔR values.

- 2257 • The transverse mass between the missing transverse energy p_T^{miss} and each of
 2258 the four objects [84], defined as

$$m_T(\ell, p_T^{\text{miss}}) \equiv \sqrt{2p_T^\ell \cdot p_T^{\text{miss}}[1 - \cos(\Delta\phi)]} \quad (9.1)$$

2259 where p_T^ℓ is the transverse momentum of the object ℓ , and $\Delta\phi$ is the difference
 2260 in azimuthal angle between the object and the p_T^{miss} . Events from $t\bar{t}$ +jets and
 2261 jets faking τ_h backgrounds have larger p_T^{miss} resulting in larger transverse mass
 2262 values compared to the signal, which tends to have smaller p_T^{miss} that is also
 2263 more aligned with the lepton legs.

- 2264 • The variable D_ζ [84], defined as

$$D_\zeta \equiv p_\zeta - 0.85p_\zeta^{\text{vis}} \quad (9.2)$$

2265 where the ζ axis is the bisector of the transverse directions of the visible τ decay
 2266 products. p_ζ is the component of the p_T^{miss} along the ζ axis, and p_ζ^{vis} is the sum
 2267 of the components of the lepton p_T along the same axis. This variable captures
 2268 the fact that in signal the p_T^{miss} is small and approximately aligned with the $\tau\tau$.
 2269 In contrast, the $Z \rightarrow \tau\tau$ background tends towards large D_ζ values because the
 2270 p_T^{miss} is collinear to the $\tau\tau$, and the $t\bar{t}$ +jets events tend to have small D_ζ due to
 2271 a large p_T^{miss} not aligned with the $\tau\tau$.

- 2272 • For events with 2 b-tag jets, one additional variable is defined to capture the
 2273 difference in the invariant mass of the bb and the $\tau\tau$:

$$\Delta m_{a_1} \equiv (m_{bb} - m_{\tau\tau})/m_{\tau\tau} \quad (9.3)$$

2274 This variable peaks at zero for the $h \rightarrow aa \rightarrow 2b2\tau$ signal.

After training, events in data, MC, and embedded are evaluated with the six DNNs and assigned a raw score between 0 and 1 (background-like or signal-like). In order to flatten the distribution of the score and define score thresholds for categorizing events, the raw output scores are transformed with the function $\tilde{p}(n) = \text{arctanh}(p \times \tanh(n))/n$ where n is a positive integer. The thresholds of the DNN score used for signal/control region definition are determined using scans that optimize the signal sensitivity and are shown in Tables 9.1 and 9.2.

1bNN $\tilde{p}(n = 1.5)$				
	SR1	SR2	SR3	CR
$\mu\tau_h$ 2018	> 0.98	$\in [0.95, 0.98]$	$\in [0.90, 0.95]$	< 0.90
$\mu\tau_h$ 2017	> 0.97	$\in [0.94, 0.97]$	$\in [0.90, 0.94]$	< 0.90
$\mu\tau_h$ 2016	> 0.97	$\in [0.94, 0.97]$	$\in [0.89, 0.94]$	< 0.89
1bNN $\tilde{p}(n = 1.5)$				
	SR1	SR2	SR3	CR
$e\tau_h$ 2018	> 0.97	$\in [0.945, 0.97]$	$\in [0.90, 0.945]$	< 0.90
$e\tau_h$ 2017	> 0.985	$\in [0.965, 0.985]$	$\in [0.93, 0.965]$	< 0.93
$e\tau_h$ 2016	> 0.985	$\in [0.965, 0.985]$	$\in [0.93, 0.965]$	< 0.93
1bNN $\tilde{p}(n = 2.5)$				
	SR1	SR2	SR3	CR
$e\mu$ 2018	> 0.99	$\in [0.95, 0.99]$	$\in [0.85, 0.95]$	< 0.85
$e\mu$ 2017	> 0.985	$\in [0.95, 0.985]$	$\in [0.85, 0.95]$	< 0.85
$e\mu$ 2016	> 0.99	$\in [0.95, 0.99]$	$\in [0.85, 0.95]$	< 0.85

Table 9.1: Event categorization based on DNN scores for events with exactly 1 b-tag jet (1bNN), for the three $\tau\tau$ channels and three eras.

9.3 Methodology for signal extraction

After events are divided into categories, the data is compared to the expected backgrounds in the signal region categories. Here, we describe the fundamental concepts behind hypothesis testing in high-energy physics, as well as how exclusion limits can be set on parameters whose true values we cannot measure, culminating in the modified frequentist method CL_S which is used to perform signal extraction in this

	2bNN $\tilde{p}(n = 1.5)$		
	SR1	SR2	CR
$\mu\tau_h$ 2018	> 0.99	$\in [0.96, 0.99]$	< 0.96
$\mu\tau_h$ 2017	> 0.98	$\in [0.94, 0.98]$	< 0.94
$\mu\tau_h$ 2016	> 0.97	$\in [0.93, 0.97]$	< 0.93
	2bNN $\tilde{p}(n = 1.5)$		
	SR1	SR2	CR
$e\tau_h$ 2018	> 0.96	NA	< 0.96
$e\tau_h$ 2017	> 0.985	NA	< 0.985
$e\tau_h$ 2016	> 0.96	NA	< 0.96
	2bNN $\tilde{p}(n = 2.5)$		
	SR1	SR2	CR
$e\mu$ 2018	> 0.98	$\in [0.94, 0.98]$	< 0.94
$e\mu$ 2017	> 0.97	$\in [0.93, 0.97]$	< 0.93
$e\mu$ 2016	> 0.98	$\in [0.94, 0.98]$	< 0.94

Table 9.2: Event categorization based on DNN scores for events with 2 b-tag jets (2bNN), for the three $\tau\tau$ channels and three eras.

2288 analysis.

2289 9.3.1 Model building and parameter estimation

In the frequentist interpretation of probability, an experiment measuring an observable can be repeated, resulting in different values of the observable, e.g. the invariant mass of a candidate Higgs boson in a search for the Higgs [102]. The ensemble of values of the observable x gives rise to the probability density function (PDF) $f(x)$, which has the important property that it is normalized to unity:

$$\int f(x) dx = 1.$$

A parametric family of PDFs

$$f(x|\alpha),$$

2290 read “ f of x given α ”, is referred to as a probability model or model. The parameters α
 2291 typically represent parameters of the theory or an unknown property of the detector’s
 2292 response. The parameters are not frequentist in nature, unlike x . Out of all the
 2293 parameters, typically only a few are of interest, and are called the parameters of
 2294 interest (POI), labeled μ here. The remaining are referred to as nuisance parameters
 2295 (NP) [102] and are labeled θ .

2296 $f(x)$ is the probability density for the observable in one event and we wish to
 2297 describe the probability density for a dataset with many events, $\mathcal{D} = \{x_1, \dots, x_n\}$,
 2298 called the total probability model \mathbf{f} . For instance, if we also have a prediction for
 2299 the total number of events expected, called ν , we also account for the overall Poisson
 2300 probability for observing n events given ν expected:

$$\mathbf{f}(\mathcal{D}|\nu, \alpha) = \text{Poisson}(n|\nu) \prod_{e=1}^n f(x_e|\alpha) \quad (9.4)$$

The likelihood function $L(\alpha)$ is numerically equivalent to $f(x|\alpha)$ for fixed x , or
 $\mathbf{f}(\mathcal{D}|\alpha)$ with \mathcal{D} fixed [102]. The likelihood function is not a probability density for α
 and is not normalized to unity:

$$\int L(\alpha) d(\alpha) \neq 1.$$

2301 i.e. the likelihood function is the value of f as a function of α given a fixed value of
 2302 x .

2303 To estimate the parameter α we use an estimator, which is a function of the
 2304 data. Take for example the measurement of data distributed according to a Gaussian
 2305 probability density $f(x|\mu, \sigma) = \text{Gauss}(x|\mu, \sigma)$. One possible estimator of the mean μ ,
 2306 is the mean of the measured data points $\bar{x} = \sum_{i=1}^n x_i/n$ [102].

2307 A commonly used estimator in physics is the maximum likelihood estimator
 2308 (MLE), defined as the value α which maximizes the likelihood function $L(\alpha)$. This

2309 value, labeled $\hat{\alpha}$, also maximizes $\ln L(\alpha)$ and minimizes $-\ln L(\alpha)$. By convention the
2310 $-\ln L(\alpha)$ is minimized, in a process called “fitting”, and the maximum likelihood
2311 estimate is called the “best fit value”.

2312 9.3.2 Hypothesis testing

2313 In this section we next introduce concepts related to hypothesis testing such as the
2314 test statistic constructed from the ratio of likelihood functions.

2315 The objective of a likelihood analysis is to distinguish different models repre-
2316 senting the various hypotheses, and determine the one that best explains the ex-
2317 perimental outcome. In a search for new physics, a signal is additive on top of the
2318 background. The background-only hypothesis is the null hypothesis, and the signal-
2319 plus-background hypothesis is the alternative.

2320 As a simple example, take the p -value test, for an experiment where we count
2321 events in the signal region, n_{SR} , and expect ν_B background events and ν_S events from
2322 the signal [102]. Then

2323 1. The null hypothesis (H_0), i.e. the background-only hypothesis in this experi-
2324 ment, with the probability modeled by $\text{Poisson}(n_{SR}|\nu_B)$.

2325 2. The alternate hypothesis (H_1), i.e. signal-plus-background hypothesis, with the
2326 probability modeled by $\text{Poisson}(n_{SR}|(\nu_B + \nu_S))$.

2327 The compatibility of the observed data n_{SR}^0 and the null hypothesis, is quantified as
2328 the probability that the background-only hypothesis would produce at least as many
2329 events as was observed. This probability is the p -value:

$$p = \sum_{n=n_{SR}^0}^{\infty} \text{Poisson}(n|\nu_B). \quad (9.5)$$

2330 If the p -value is very small, we might reject the null hypothesis. The p -value is not the

2331 probability of the null hypothesis given the data; rather, it expresses the probability
 2332 that data with a certain property was obtained, assuming the null hypothesis [102].

2333 The *p*-value is an example of a test statistic T , which maps the data to a single
 2334 real number. The Neyman-Pearson lemma states that out of the infinite possibilities
 2335 of choices of test statistic, the uniformly most powerful test statistic is the likelihood
 2336 ratio T_{NP} [102]:

$$T_{NP}(\mathcal{D}) = \frac{L(\mathcal{D}|H_1)}{L(\mathcal{D}|H_0)} \quad (9.6)$$

To reiterate, the test statistic T is a real-valued function of the data, implying that a particular probability model $\mathbf{f}(\mathcal{D}|\boldsymbol{\alpha})$ implies a distribution of the test statistic, $f(T|\boldsymbol{\alpha})$, which depends on the value of $\boldsymbol{\alpha}$. With this distribution in hand, the *p*-value can be evaluated in the following equivalent formulations:

$$p(\boldsymbol{\alpha}) = \int_{T_0}^{\infty} f(T|\boldsymbol{\alpha}) dT \quad (9.7)$$

$$= \int \mathbf{f}(\mathcal{D}|\boldsymbol{\alpha}) \theta(T(\mathcal{D}) - T_0) d\mathcal{D} \quad (9.8)$$

$$= P(T \geq T_0|\boldsymbol{\alpha}) \quad (9.9)$$

2337 where T_0 is the value of T based on the observed data, and $\theta()$ is the Heaviside
 2338 function. The size of the test is conventionally chosen to be 10%, 5%, or 1%. As
 2339 the *p*-value depends on $\boldsymbol{\alpha}$ (both the POI and NP), the null hypothesis should not be
 2340 rejected if the *p*-value is larger than the size of the test for any value of the nuisance
 2341 parameters.

2342 9.3.3 Confidence intervals

2343 In an example of the measurement of the Standard Model Higgs boson, $\boldsymbol{\alpha}_{POI} =$
 2344 $(\sigma/\sigma_{SM}, M_H)$, with σ/σ_{SM} is the ratio of the production cross-section for Higgs with

respect to its value in the SM, and M_H is the unknown mass of the Higgs, values
 of these parameters outside specific bounds are said to be “excluded at the 95%
 confidence level”. These allowed regions are called confidence levels or confidence
 regions, and the parameter values outside of them are considered excluded [102]. A
 95% confidence interval does not mean that there is a 95% chance that the true value
 of the parameter is inside the interval. Rather, a 95% confidence interval covers the
 true value 95% of the time (even though we do not know the true value).

To construct a confidence interval for a parameter α , the Neyman Construction
 is used to invert a series of hypothesis tests; i.e. for each possible value of α , the null
 hypothesis is treated as α , and we perform a hypothesis test based on a test statistic.
 To construct a 95% confidence interval, we construct a series of hypothesis tests with
 size of 5%. The confidence interval $I(\mathcal{D})$ is constructed by taking the set of parameter
 values $\boldsymbol{\alpha}$ where the null hypothesis is accepted:

$$I(\mathcal{D}) = \{\boldsymbol{\alpha} | P(T(\mathcal{D}) > k_\alpha | \boldsymbol{\alpha}) < \alpha\}, \quad (9.10)$$

where $T(\mathcal{D})$ is the test statistic, and the last α (not bolded) and the subscript k_α
 refer to the size of the test. A schematic of the Neyman construction is shown in Fig.
 9.1. In a more generalized case, the x -axis is the test statistic T .

9.3.4 Profile likelihood ratio

In this section we describe a frequentist statistical procedure based on the profile
 likelihood ratio test statistic, which is implemented using asymptotic distributions.

With a multi-parameter likelihood function $L(\boldsymbol{\alpha})$, the the maximum likelihood of
 one specific parameter α_p with other parameters $\boldsymbol{\alpha}_o$ fixed, is called the conditional
 maximum likelihood estimate and is denoted $\hat{\alpha}_p(\boldsymbol{\alpha}_0)$. The process of choosing specific
 values of the nuisance parameters for a given value of μ , $\mathcal{D}_{\text{simulated}}$, and value of global

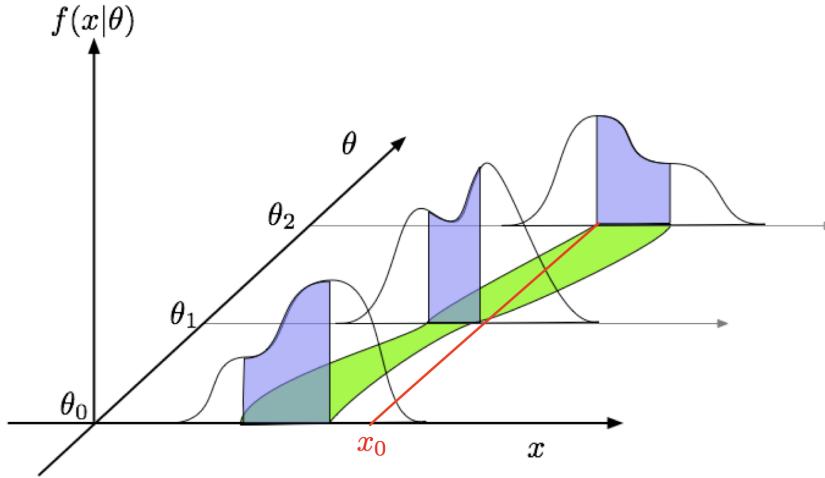


Figure 9.1: Schematic of the Neyman construction for confidence intervals [102]. For each value of θ , we find a region in x where $\int f(x|\theta)dx$ satisfies the size of the test (blue). These regions form a confidence belt (green). The intersection of the observation x_0 (red) with the confidence belt defines the confidence interval $[\theta_1, \theta_2]$ [102].

2368 observables \mathcal{G} is called profiling. From the full list of parameters $\boldsymbol{\alpha}$, we denote the
2369 parameter of interest μ , and the nuisance parameters $\boldsymbol{\theta}$.

2370 We construct the profile likelihood ratio,

$$\lambda(\mu) = \frac{L(\mu, \hat{\boldsymbol{\theta}}(\mu))}{L(\mu, \hat{\boldsymbol{\theta}})} \quad (9.11)$$

2371 which depends explicitly on the parameter of interest μ , implicitly on the data \mathcal{D}_{sim}
2372 and global observables \mathcal{G} , and is independent of the nuisance parameters $\boldsymbol{\theta}$, which
2373 have been eliminated in profiling [102].

2374 The main conceptual reason for constructing the test statistic from the profile
2375 likelihood ratio is that asymptotically (i.e. for measurements with many events) the
2376 distribution of the profile likelihood ratio $\lambda(\mu = \mu_{\text{true}})$ is independent of the values of
2377 the nuisance parameters [102].

2378 The following p -value is used to quantify the consistency with the hypothesis of a
2379 signal strength of μ :

$$p_\mu = \int_{\tilde{q}_{\mu,\text{obs}}}^{\infty} f(\tilde{q}_\mu | \mu, \hat{\boldsymbol{\theta}}(\mu, \text{obs})) d\tilde{q}_\mu \quad (9.12)$$

2380 **9.3.5 Modified frequentist method: CL_S**

2381 In the modified frequentist method called CL_S , to test a hypothesis with signal, we
2382 define p'_μ as a ratio of p -values [102]:

$$p'_\mu = \frac{p_\mu}{1 - p_b} \quad (9.13)$$

2383 where p_b is the p -value derived under the background-only hypothesis:

$$p_b = 1 - p_0 \equiv 1 - \int_{\tilde{q}_{\mu,\text{obs}}}^{\infty} f(\tilde{q}_\mu | 0, \hat{\theta}(\mu = 0, \text{obs})) d\tilde{q}_\mu. \quad (9.14)$$

2384 The CL_S upper limit on μ , denoted μ_{up} , is obtained by solving for $p'_{\mu_{up}} = 5\%$.
2385 If testing the compatibility of the data with the background-only hypothesis, we
2386 consider the p_b value defined above and conventionally convert it into the quantile
2387 or “sigma” of a unit Gaussian. z standard deviations (e.g. $z = 5$ in “ 5σ ”) means
2388 that the probability of falling above these standard deviations, equals p_b (e.g. 3σ
2389 corresponds to $p_b = 2.7 \times 10^{-3}$ or 95.43%, and 5σ corresponds to $p_b = 5.7 \times 10^{-7}$ or
2390 99.999943%).

2391 **Chapter 10**

2392 **Results**

2393 In this chapter, Section 10.1 presents the results from the $h \rightarrow aa \rightarrow bb\tau\tau$ analysis
2394 performed on 137 fb^{-1} of data from the full CMS Run-2 dataset in the years 2016 to
2395 2018, with interpretations provided for different 2HDM+S scenarios. This analysis
2396 was combined with a different search in the $h \rightarrow aa \rightarrow bb\mu\mu$ final state, which was
2397 also performed on the full Run-2 dataset. The combination procedure and results
2398 from the combined analyses ($h \rightarrow aa \rightarrow bb\ell\ell$, with $\ell = \mu, \tau$) are detailed in 10.2.
2399 The combined analysis places some of the most stringent limits to date at CMS for
2400 2HDM+S scenarios in the light scalar mass range $m_a = 12 \text{ GeV}$ to 60 GeV .

2401 **10.1 Results from $bb\tau\tau$**

2402 In each of the three $\tau\tau$ channels studied ($\mu\tau_h$, $e\tau_h$, and $e\mu$), events are divided based
2403 on whether they contain exactly 1 or 2 b-tag jets, and further divided into signal
2404 and control regions (SRs and CRs) using the DNN categorization score as described
2405 in Section 9.2. The control regions demonstrate good agreement between observed
2406 events in data, and the sum of the contributions from expected backgrounds that
2407 are modeled in simulated and embedded samples. The signal regions are defined to
2408 be sensitive to the $h \rightarrow aa \rightarrow bb\tau\tau$ signal. The postfit final observed and expected

2409 distributions of the di-tau invariant mass $m_{\tau\tau}$ reconstructed with SVFit (described
2410 in Section 5.2) are shown in Fig. 10.1 for the $\mu\tau_h$ channel, Fig. 10.2 for the $e\tau_h$
2411 channel, and Fig. 10.3 for the $e\mu$ channel. In all figures, the hypothesized yield for
2412 the $h \rightarrow aa \rightarrow bb\tau\tau$ signal is shown for the pseudoscalar mass $m_a = 35$ GeV and
2413 assuming a branching fraction $B(H \rightarrow aa \rightarrow bb\tau\tau) = 10\%$.

2414 The 95% CL expected and observed exclusion limits on the signal strength of the
2415 branching fraction $B(h \rightarrow aa \rightarrow bb\tau\tau)$ as a function of the pseudoscalar mass m_a
2416 ranging from 12 GeV to 60 GeV, are shown for the three $\tau\tau$ channels and all three
2417 channels combined in Fig. 10.4. The limits are shown as percentages and normalized
2418 to the production cross-section of the Standard Model Higgs boson. No excess of
2419 events above the Standard Model expectations is observed. In the limits for the three
2420 $\tau\tau$ channels combined, expected (observed) limits range from 1.4 to 5.6% (1.7 to
2421 7.6%) for pseudoscalar masses between 12 and 60 GeV.

2422 The $e\mu$ channel is the only channel that has signal sensitivity to the $m_a = 12$
2423 GeV pseudoscalar mass hypothesis, because the minimum required spatial separation
2424 $\Delta R = \sqrt{(\Delta\eta)^2 + (\Delta\phi)^2}$ between the two τ legs is smaller than the other two channels
2425 ($\Delta R < 0.3$ for $e\mu$, compared to $\Delta R < 0.4$ for the other two channels). This decreased
2426 ΔR requirement results in better signal acceptance for low mass signals for the $e\mu$
2427 channel. The $\mu\tau_h$ and $e\tau_h$ channels are most sensitive to the intermediate mass points
2428 studied, since the analysis targets a resolved signature: at low mass points, the tau
2429 legs are boosted, and at high mass points, the $m_{\tau\tau}$ distributions in signal have larger
2430 overlap with background distributions. In the combination of the three $\tau\tau$ channels,
2431 the limit for $m_a = 12$ GeV comes only from the $e\mu$ channel, and the best sensitivity
2432 is attained at intermediate mass points around $m_a = 20$ GeV to 45 GeV.

2433 To set limits on the branching fraction of the 125 GeV Higgs to the two pseu-
2434 doscalars, $B(h \rightarrow aa)$, we interpret the results in four types of 2HDM+S, which were
2435 introduced in Section 1.4. In 2HDM+S, the theorized branching fraction of the pseu-

2436 doscalars depends on the 2HDM+S model type, the pseudoscalar mass m_a , and the
2437 ratio of the two Higgs doublets' vacuum expectation values $\tan \beta$. In Type I models,
2438 the branching fraction is independent of $\tan \beta$, while in Types II, III, and IV, it is
2439 a function of m_a and $\tan \beta$. Limits for the $bb\tau\tau$ final state as a function of m_a for
2440 2HDM+S Type I (valid for all $\tan \beta$ values), Type II with $\tan \beta = 2.0$, Type III with
2441 $\tan \beta = 2.0$, and Type IV with $\tan \beta = 0.6$ are overlaid and shown in Fig. 10.5a.

2442 10.2 Combination with $bb\mu\mu$ final state

2443 Results from this analysis for the $h \rightarrow aa \rightarrow bb\tau\tau$ final state are combined with the
2444 analysis for the $h \rightarrow aa \rightarrow bb\mu\mu$ final state [103]. While the predicted branching ratio
2445 for $aa \rightarrow bb\mu\mu$ is comparatively small, the $bb\mu\mu$ final state has competitive results
2446 due to the excellent di-muon resolution measured by CMS. The $bb\mu\mu$ analysis uses
2447 an unbinned fit to the data using the di-muon mass $m_{\mu\mu}$ distribution. Details can be
2448 found in [103].

2449 Combining the results is possible since the $bb\tau\tau$ analysis explicitly rejects events
2450 with extra leptons, so there is no overlap between the events studied in the $bb\tau\tau$
2451 analysis and the $bb\mu\mu$ analysis. In the statistical combination, several systematic
2452 uncertainties are treated as correlated: the integrated luminosity normalization, the
2453 b-tagging scale factor, the scale factors related to muon reconstruction, identifica-
2454 tion, and trigger efficiencies, the inefficiency in the ECAL trigger readout, and the
2455 theoretical uncertainties related to signal modeling.

2456 Since the results in both final states are statistically limited, the combination ben-
2457 efits from the additional data. For $m_a = 35$ GeV, all systematic uncertainties amount
2458 to around 6% of the total uncertainty, with the dominant systematic uncertainties
2459 coming from jet energy systematics in the $bb\mu\mu$ final state, theoretical uncertainties
2460 in the signal, and uncertainties in the QCD multijet backgrounds in the $e\mu$ channel

2461 of the $bb\tau\tau$ final state.

2462 The mass distributions of the di-muon and di-tau objects ($m_{\mu\mu}$ and $m_{\tau\tau}$) are
2463 compared to the data in a combined maximum likelihood fit to derive upper limits
2464 on $B(h \rightarrow aa)$. The observed limits at 95% CL on $B(h \rightarrow aa)$ for different 2HDM+S
2465 scenarios, are shown for the search for $h \rightarrow aa \rightarrow bb\mu\mu$ in Fig. 10.5b, and the
2466 combined analyses $h \rightarrow aa \rightarrow bb\ell\ell$ in Fig. 10.6.

2467 Exclusion limits in a two-dimensional plane as a function of $\tan\beta$ and m_a are
2468 set for 2HDM+S Types II, III, and IV in Fig. 10.7. The most stringent constraints
2469 are observed for 2HDM+S type III because of large branching fractions predicted in
2470 theory, with predicted branching fractions between 0.47 and 0.42 for $\tan\beta = 2.0$ and
2471 values of m_a between 15 and 60 GeV, compared to the observed 95% CL upper limits
2472 which are between 0.08 and 0.03. For 2HDM+S type IV, the predicted branching
2473 fractions from theory are between 0.26 and 0.20 for $\tan\beta = 0.6$ for values of m_a
2474 between 15 and 60 GeV, and the 95% CL observed upper limits are between 0.12 and
2475 0.05.

2476 The combined results from $h \rightarrow aa \rightarrow bb\ell\ell$ are compared with CMS results in
2477 other final states as a function of the pseudoscalar mass m_a : for 2HDM+S type I in
2478 Fig. 10.8, type II with $\tan\beta = 2.0$ in Fig. 10.9, and type III with $\tan\beta = 2.0$ in Fig.
2479 10.10. In other scenarios, e.g. type III with $\tan\beta = 5.0$, more stringent limits are set
2480 by analyses in other final states, $\mu\mu\tau\tau$ in this case. Other summary plots for other
2481 model types and $\tan\beta$ values can be found at [104].

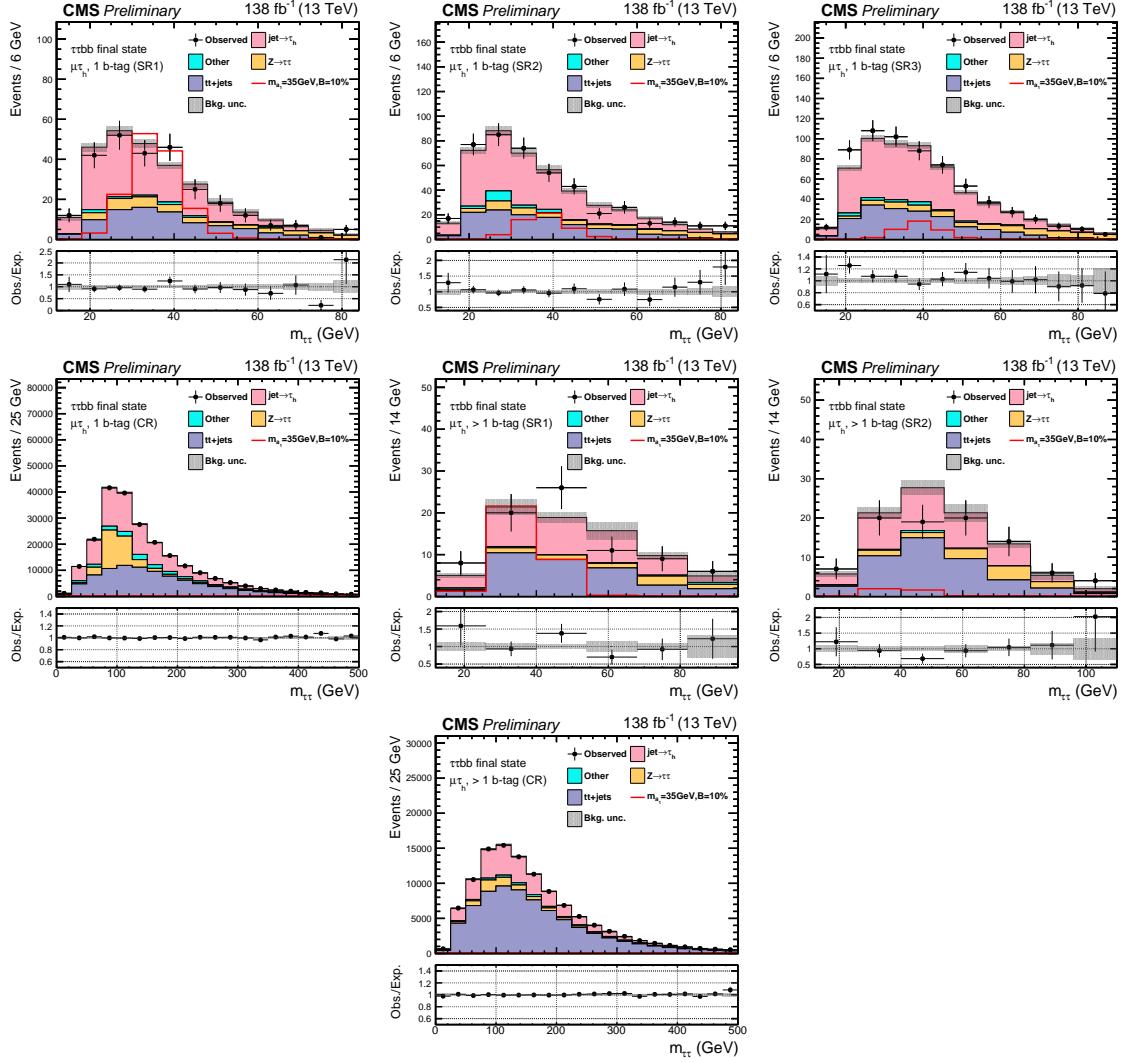


Figure 10.1: Postfit final $m_{\tau\tau}$ observed and expected distributions, and the observed/expected ratios, in the $\mu\tau_h$ channel [101]. Events are divided into the 1 b-tag jet signal regions (SR1, SR2, SR3) (*top row*), 1 b-tag jet control region (*middle row*), 2 b-tag jet signal regions (SR1, SR2) (*middle row*), and lastly the 2 b-tag jet control region (CR) (*bottom*). Statistical and systematic sources of uncertainties in the expected events are added in quadrature and labeled “Bkg. unc” (*shaded gray*). The dominant backgrounds in all categories are jets faking the τ_h leg (*pink*), $Z \rightarrow \tau\tau$ (*orange*), and $t\bar{t}+j$ ets (*purple*). For illustrative purposes, the beyond-Standard Model signal yield from $h \rightarrow aabb\tau\tau$ is shown for the pseudoscalar mass hypothesis $m_a = 35$ GeV, assuming a branching fraction $B(h \rightarrow aa \rightarrow bb\tau\tau) = 10\%$ (*red line*).

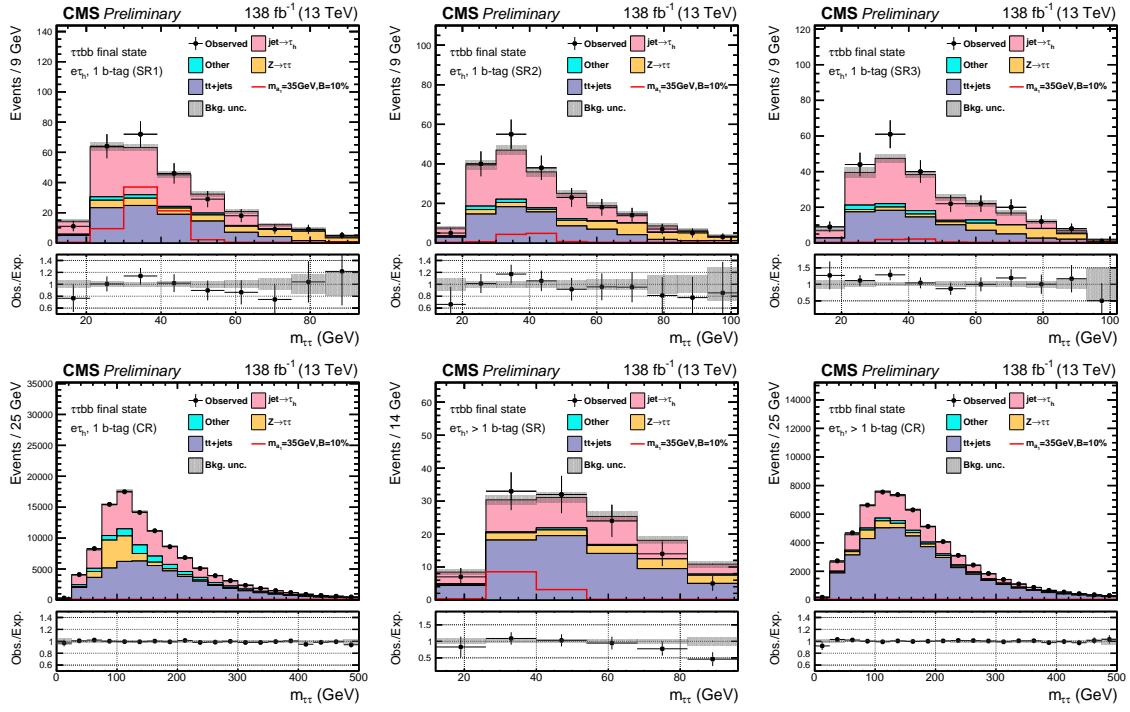


Figure 10.2: Postfit final observed and expected $m_{\tau\tau}$ distributions, and the observed/expected ratios, in the $e\tau_h$ channel [101]. Events are divided into the 1 b-tag jet signal regions (SR1, SR2, SR3) (*top row*), the 1 b-tag jet control region (CR) (*bottom row*), and 2 b-tag jet signal region (SR) and control region (CR) (*bottom row*). Statistical and systematic sources of uncertainties in the expected events are added in quadrature and labeled “Bkg. unc” (*shaded gray*). In this channel, the dominant backgrounds are jets faking the τ_h leg (*pink*), $Z \rightarrow \tau\tau$ (*orange*), and $t\bar{t}+{\rm jets}$ (*purple*). For illustrative purposes, the beyond-Standard Model signal yield from $h \rightarrow aabb\tau\tau$ is shown for the pseudoscalar mass hypothesis $m_a = 35$ GeV, assuming a branching fraction $B(h \rightarrow aa \rightarrow bb\tau\tau) = 10\%$ (*red line*).

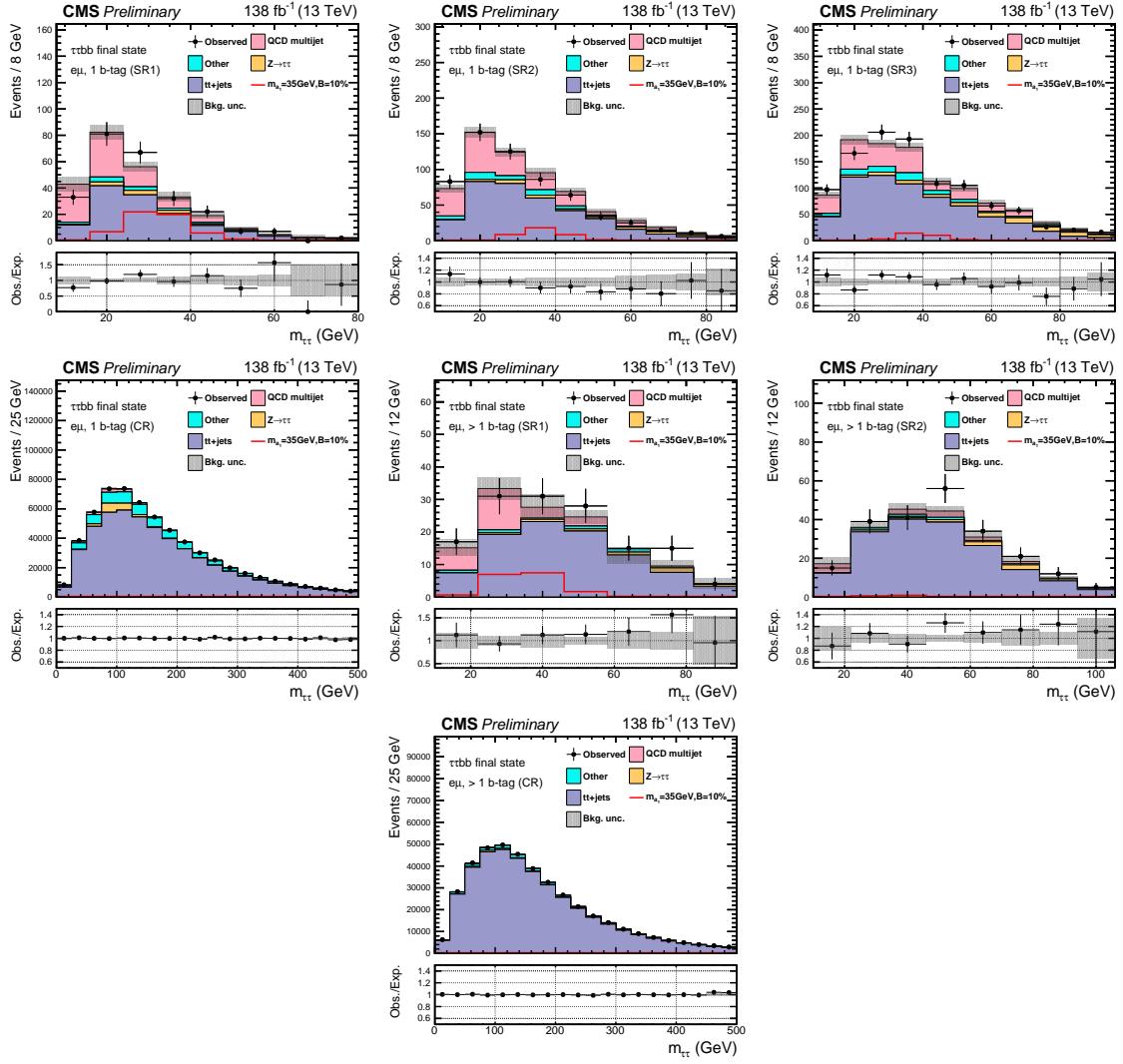


Figure 10.3: Postfit final observed and expected $m_{\tau\tau}$ distributions, and the observed/expected ratios, in the $e\mu$ channel [101]. Events are divided into the 1 b-tag jet signal regions (SR1, SR2, and SR3) (*top row*), 1 b-tag jet control region (CR) (*middle row*), 2 b-tag jet signal regions (SR1 and SR2) (*middle row*), and 2 b-tag jet control region (CR) (*bottom row*). Statistical and systematic sources of uncertainties in the expected events are added in quadrature and labeled “Bkg. unc” (*shaded gray*). The $t\bar{t}+j$ process (*purple*) is a major background, and in the signal regions the QCD multijet (*pink*) is also a major background. TFor illustrative purposes, the beyond-Standard Model signal yield from $h \rightarrow aabb\tau\tau$ is shown for the pseudoscalar mass hypothesis $m_a = 35$ GeV, assuming a branching fraction $B(h \rightarrow aa \rightarrow bb\tau\tau) = 10\%$ (*red line*).

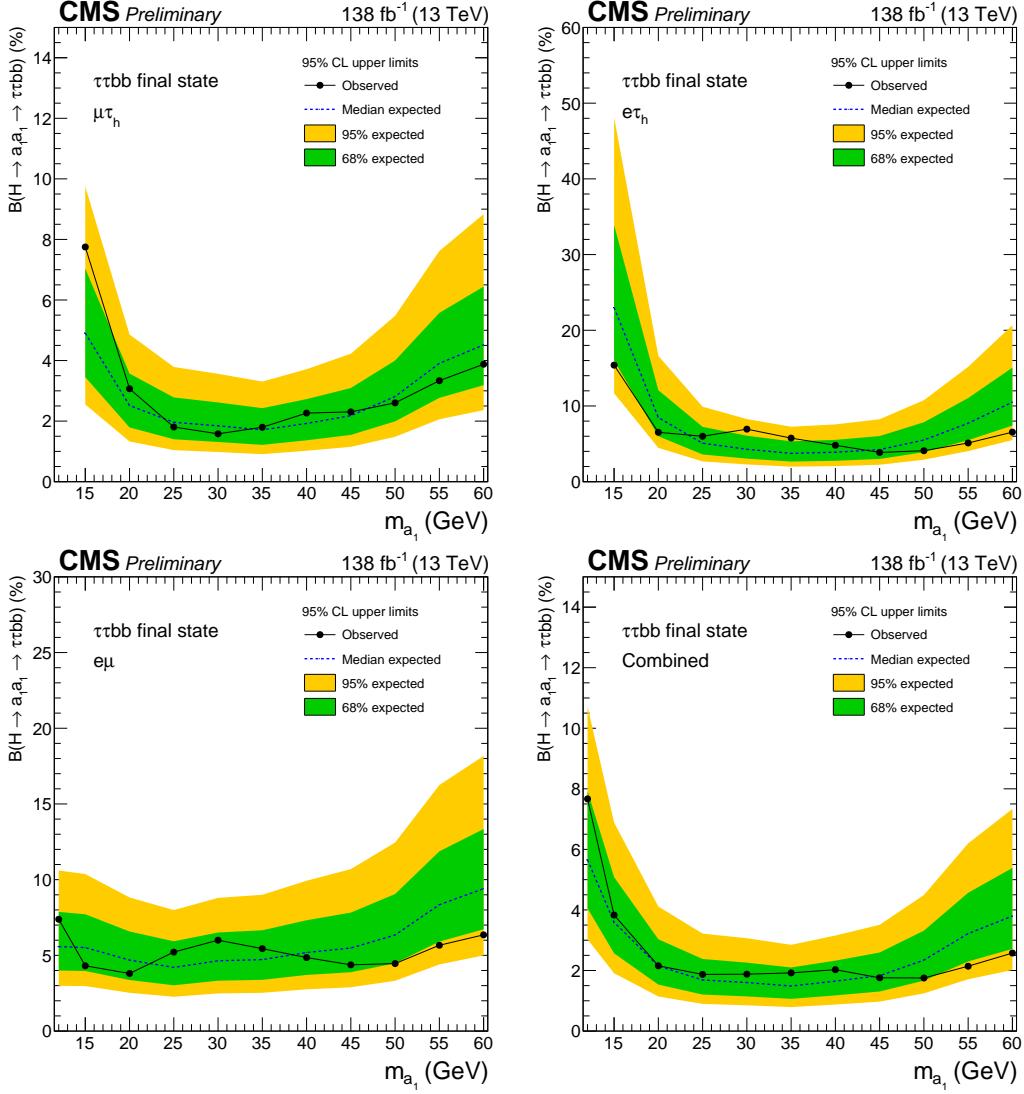
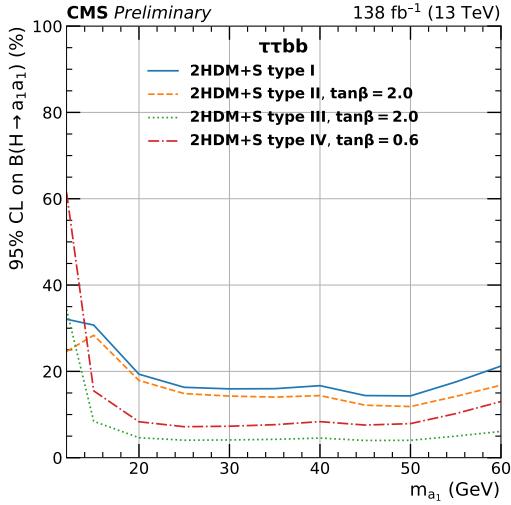
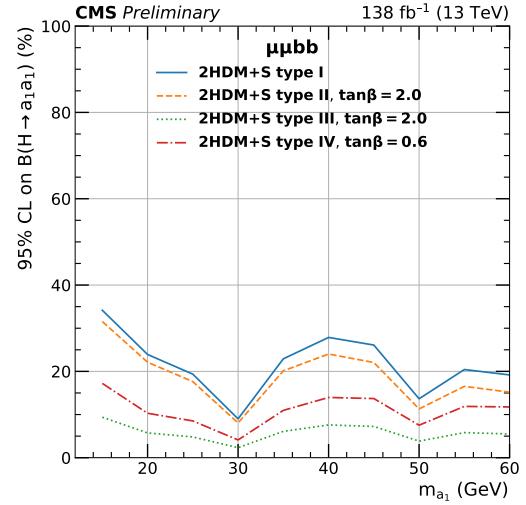


Figure 10.4: Observed 95% CL exclusion limits (*black, solid lines*) and expected 95% CL and 68% CL limits (*shaded yellow and green*) on the branching fraction $B(h \rightarrow aa \rightarrow bb\tau\tau)$ in percentages, assuming the Standard Model production for the 125 GeV Higgs (h). Limits are shown for the $\mu\tau_h$ channel (*top left*), the $e\tau_h$ channel (*top right*), and the $e\mu$ channel (*bottom left*), and lastly the combination of all three channels (*bottom right*) [101]. The dataset corresponds to 138 fb^{-1} of data collected in the years 2016-2018 at a center-of-mass energy 13 TeV. Only the $e\mu$ channel has sensitivity to the mass hypothesis $m_a = 12 \text{ GeV}$. The best sensitivity is attained at intermediate mass points.



(a) $bb\tau\tau$ final state.



(b) $bb\mu\mu$ final state.

Figure 10.5: Observed 95% CL upper limits on $B(h \rightarrow aa)$ in %, for the $bb\tau\tau$ final state (*left*) and $bb\mu\mu$ final state (*right*) using the full Run 2 integrated luminosity of 138 fb^{-1} in 2HDM+S type I (blue), type II with $\tan\beta = 2.0$ (orange dashed), type III with $\tan\beta = 2.0$ (dotted green), and type IV with $\tan\beta = 0.6$ (red dashed) [101]. Linear interpolation is used between points in the graphs. The $\tan\beta$ values chosen here correspond to the most stringent limits in each model.

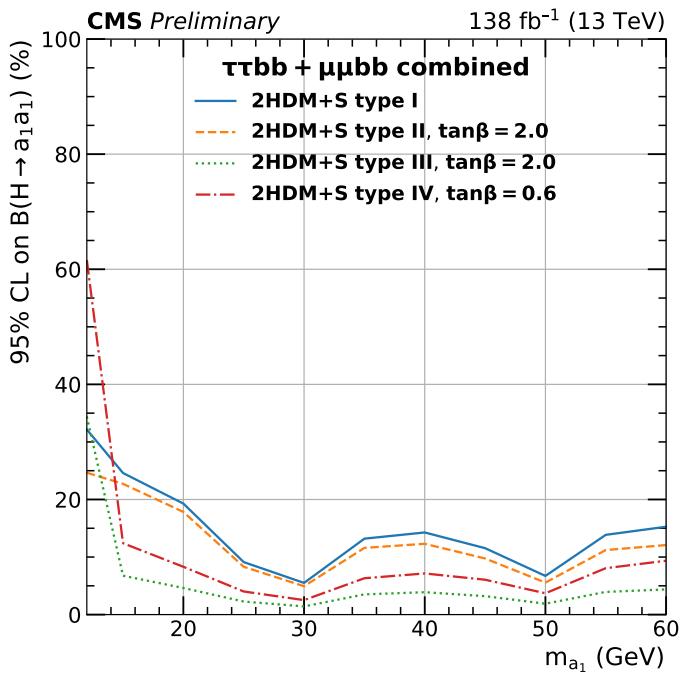


Figure 10.6: Observed 95% CL upper limits on the branching fraction of the 125 GeV Higgs boson to two pseudoscalars, $B(h \rightarrow aa)$, in percentages, as a function of the pseudoscalar mass m_a , in 2HDM+S type I (blue), type II with $\tan\beta = 2.0$ (orange dashed), type III with $\tan\beta = 2.0$ (dotted green), and type IV with $\tan\beta = 0.6$ (red dashed), for the combination of $bb\mu\mu$ and $bb\tau\tau$ channels using the full Run 2 integrated luminosity of 138 fb^{-1} [101].

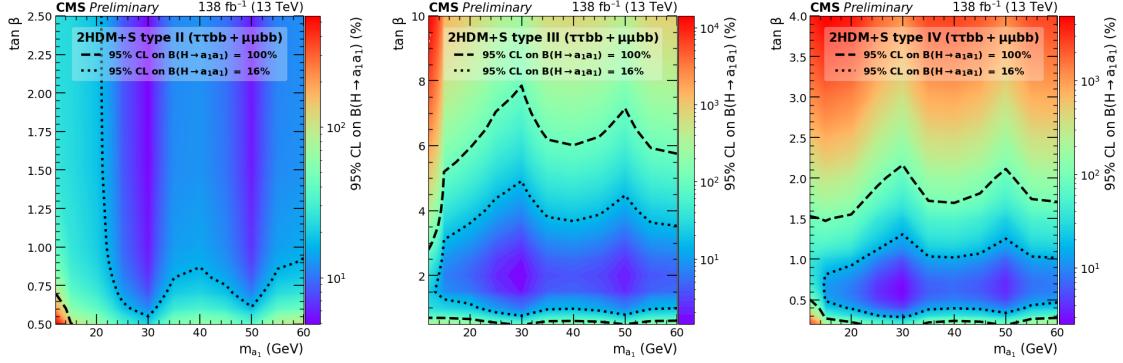


Figure 10.7: Observed 95% CL upper limits on $\mathcal{B}(h \rightarrow aa)$ in %, for the combination of $bb\mu\mu$ and $bb\tau\tau$ channels using the full Run 2 integrated luminosity of 138 fb^{-1} for Type II (*left*), Type III (*middle*), and Type IV (*right*) 2HDM+S in the $\tan \beta$ vs. m_a phase space. The contours (*dashed black*) correspond to branching fractions of 100% and 16%, where 16% is the combined upper limit on Higgs boson to undetected particle decays from previous Run-2 results. All points inside the contour are allowed within that upper limit. Linear extrapolation has been used between different points on the figures [101].

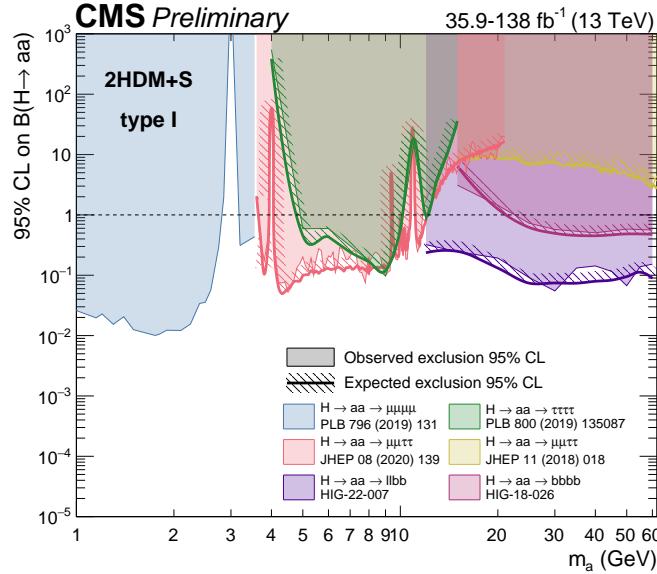


Figure 10.8: Summary plot of current 95% limits on the branching ratio of the 125 GeV Higgs boson to two pseudoscalars, normalized to the Standard Model Higgs production cross-section, $\frac{\sigma(h)}{\sigma_{\text{SM}}} \times B(H \rightarrow aa)$ in the 2HDM+S type I scenario performed with data collected at 13 TeV [104]. Results from different final states studied at CMS are overlaid on this figure: $\mu\mu\mu\mu$ (blue), $\tau\tau\tau\tau$ (green), boosted $2\mu 2\tau$ (red), resolved $2\mu 2\tau$ (yellow), $bbbb$ (magenta), and the combined result for $\ell\ell bb$ ($\ell = \mu, \tau$) (purple).

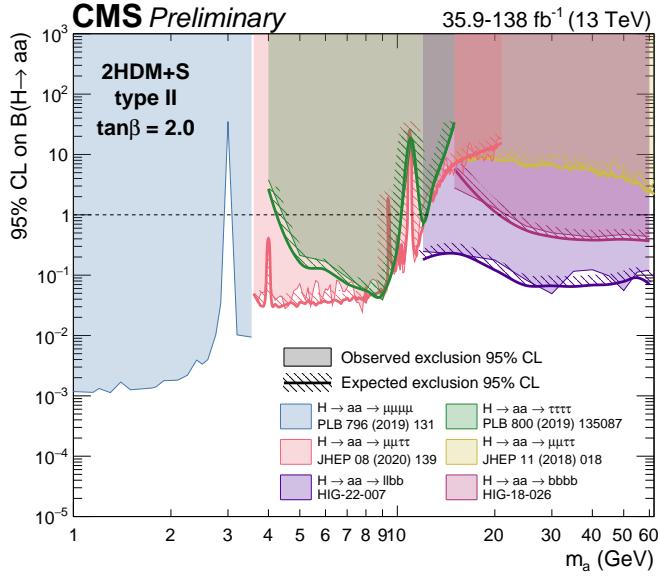


Figure 10.9: Summary plot of current observed and expected 95% CL limits on the branching ratio of the 125 GeV Higgs boson to two pseudoscalars, normalized to the Standard Model Higgs production cross-section, $\frac{\sigma(h)}{\sigma_{SM}} \times B(h \rightarrow aa)$, in the 2HDM+S type II scenario with $\tan \beta = 2.0$, obtained at CMS with data collected at 13 TeV [104]. Results from different final states studied at CMS are overlaid on this figure: $\mu\mu\mu\mu$ (blue), $\tau\tau\tau\tau$ (green), boosted $2\mu 2\tau$ (red), resolved $2\mu 2\tau$ (yellow), $bbbb$ (magenta), and the combined result for $\ell\ell bb$ ($\ell = \mu, \tau$) (purple).

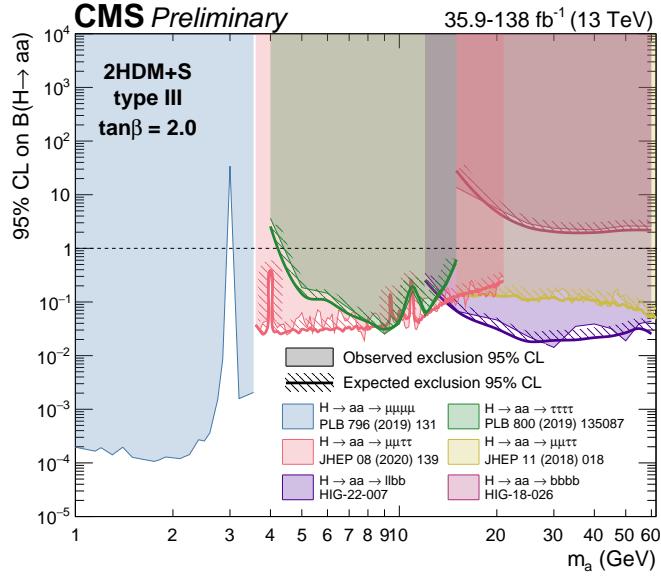


Figure 10.10: Summary plot of current observed and expected 95% CL limits on the branching ratio of the 125 GeV Higgs boson to two pseudoscalars, normalized to the Standard Model Higgs production cross section, $\frac{\sigma(h)}{\sigma_{SM}} \times B(h \rightarrow aa)$ in the 2HDM+S type-III scenario with $\tan \beta = 2.0$, obtained at CMS with data collected at 13 TeV [104]. Results from different final states studied at CMS are overlaid on this figure: $\mu\mu\mu\mu$ (blue), $\tau\tau\tau\tau$ (green), boosted $2\mu 2\tau$ (red), resolved $2\mu 2\tau$ (yellow), $bbbb$ (magenta), and the combined result for $\ell\ell bb$ ($\ell = \mu, \tau$) (purple).

2482 **Chapter 11**

2483 **Asymmetric exotic Higgs decays**

2484 This chapter presents progress towards a search for exotic Higgs decays to two light
2485 scalars with unequal mass ($h \rightarrow a_1 a_2$) final states with bottom quarks and τ leptons,
2486 with plans to interpret the results in the context of Two Real Singlet Models (TRSMs),
2487 described in Section 1.5. Compared to the symmetric decay scenario $h \rightarrow aa$ which
2488 has been studied in multiple final states at CMS with stringent limits set on the
2489 various 2HDM+S scenarios, this asymmetric decay scenario has not been directly
2490 searched for at the CMS experiment. Section 11.1 lists the mass hypotheses of the
2491 new particles a_1 and a_2 that will be studied. Section 11.2 describes the studies on
2492 which channels the analysis will be carried out in. Section 11.3 shows the control
2493 plots produced using the analysis framework that will be used for this analysis.

2494 **11.1 Signal masses**

2495 As discussed in Section 1.5, $h \rightarrow a_1 a_2$ can result in a “cascade” decay if one of the
2496 scalars, a_2 is sufficiently heavy ($m_{a_2} > 2m_{a_1}$). The “non-cascade” case is where the
2497 light scalars decay directly to Standard Model particles.

2498 The mass hypotheses (mass points) (m_{a_1}, m_{a_2}) studied here are:

- 2499 • *Cascade mass points:* (15, 30), (15, 40), (15, 50), (15, 60), (15, 70), (15, 80),
2500 (15, 90), (15, 100), (15, 110), (20, 40), (20, 50), (20, 60), (20, 70), (20, 80), (20,
2501 90), (20, 100), (30, 60), (30, 70), (30, 80), and (30, 90) GeV

- 2502 • *Non-cascade mass points:* (15, 20), (15, 30), (20, 30), (20, 40), (30, 40), (30,
2503 50), (30, 60), (40, 50), (40, 60), (40, 70), (40, 80), (50, 60), and (50, 70) GeV

2504 Samples were produced using the MadGraph5_aMCatNLO event generator, for each
2505 signal mass point in the gluon-gluon fusion (ggF) and vector boson fusion (VBF)
2506 production modes of the 125 GeV Higgs boson. In the sample generation, the decays
2507 of a to Standard Model particles were specified to be decays to bottom quarks or τ
2508 leptons.

2509 11.2 Cascade scenario signal studies

2510 Studies of the signal phenomenology in the cascade scenario were performed to de-
2511 termine the viability of the $4b2\tau$ and/or $2b4\tau$ channels.

2512 Cross sections and branching fractions of the $4b2\tau$ and $2b4\tau$ final states were
2513 compared using cross-section predictions provided by the authors of [7]. For an
2514 example mass point $m_{a_2} = 80$ GeV, $m_{a_1} = 30$ GeV, the branching fractions to
2515 $4b2\tau$ is ten times larger than $2b4\tau$: $B(h \rightarrow a_1 a_2 \rightarrow 3a_1 \rightarrow 4b2\tau) = 0.00857$, vs.
2516 $B(h \rightarrow a_1 a_2 \rightarrow 3a_1 \rightarrow 2b4\tau) = 0.00068$. The $4b2\tau$ final state is chosen for this
2517 analysis.

2518 In general the four b-flavor jets have low p_T at generator level, as illustrated for
2519 example mass points (100, 15) GeV and (40, 20) GeV in Fig. 11.1. The p_T distribution
2520 of the sub-leading jet peaks at an energy below 20 GeV, with the third and fourth
2521 jets tending to have even softer energies.

2522 An event category with three or more b-tag jets was determined to be infeasible
2523 due to low statistics in this category, due to the difficulties in reconstructing the third

2524 and fourth b-flavor jets which have very low transverse momenta p_T . Event categories
 2525 with exactly 1 b-tag jet and ≥ 2 b-tag jets will be used.

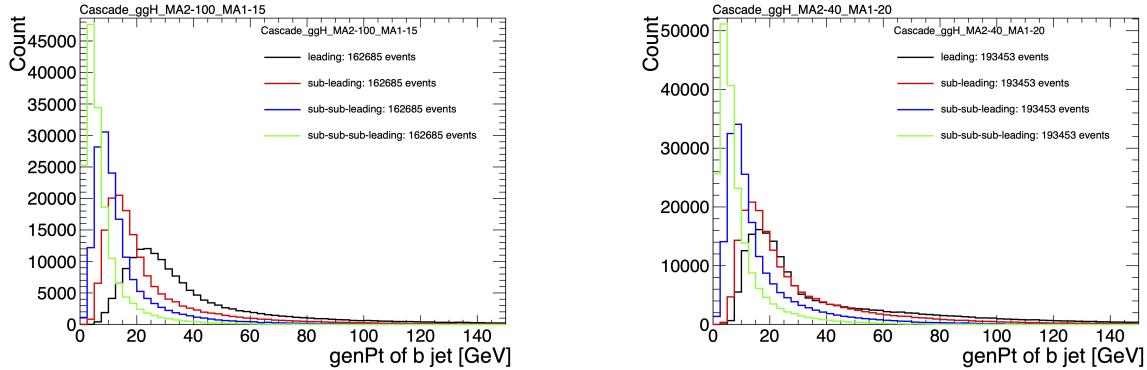


Figure 11.1: Generator-level b-flavor jet transverse momenta p_T , for $h \rightarrow a_1 a_2$ cascade scenario in the $4b2\tau$ final state, for mass hypotheses $(m_{a_1}, m_{a_2}) = (100, 15)$ GeV (*left*) and $(40, 20)$ GeV (*right*). In each plot the generator-level p_T of the leading (*black*), sub-leading (*red*), third (*blue*), and fourth (*light green*) are overlaid.

2526 In the $4b2\tau$ final state, the possibility of the leading and sub-leading b-tag jets
 2527 being sufficiently close in ΔR to require boosted jet reconstruction techniques was
 2528 explored. In the $4b2\tau$ case, the two b-flavor-jets in the generated event that were
 2529 spatially closest in ΔR were considered as one object. This two b-flavor jet object was
 2530 spatially matched in ΔR to the jets reconstructed with the standard AK4 algorithm
 2531 which uses a cone size of $\Delta R = 0.4$. The quality of the p_T resolution (computed as
 2532 $(p_{T,\text{reconstructed}} - p_{T,\text{gen}})/p_{T,\text{gen}}$) and closeness in distance ΔR of the reconstructed jet
 2533 to the nearest generator-level jets, was seen to depend on the absolute and relative
 2534 masses of the light scalars. The best (worst) performance occurred in samples with
 2535 large (small) mass differences between the heavier scalar a_2 and the lighter scalar a_1 ,
 2536 as illustrated for the mass hypotheses (m_{a_1}, m_{a_2}) (100, 15) GeV and (40, 20) GeV in
 2537 Fig. 11.2.

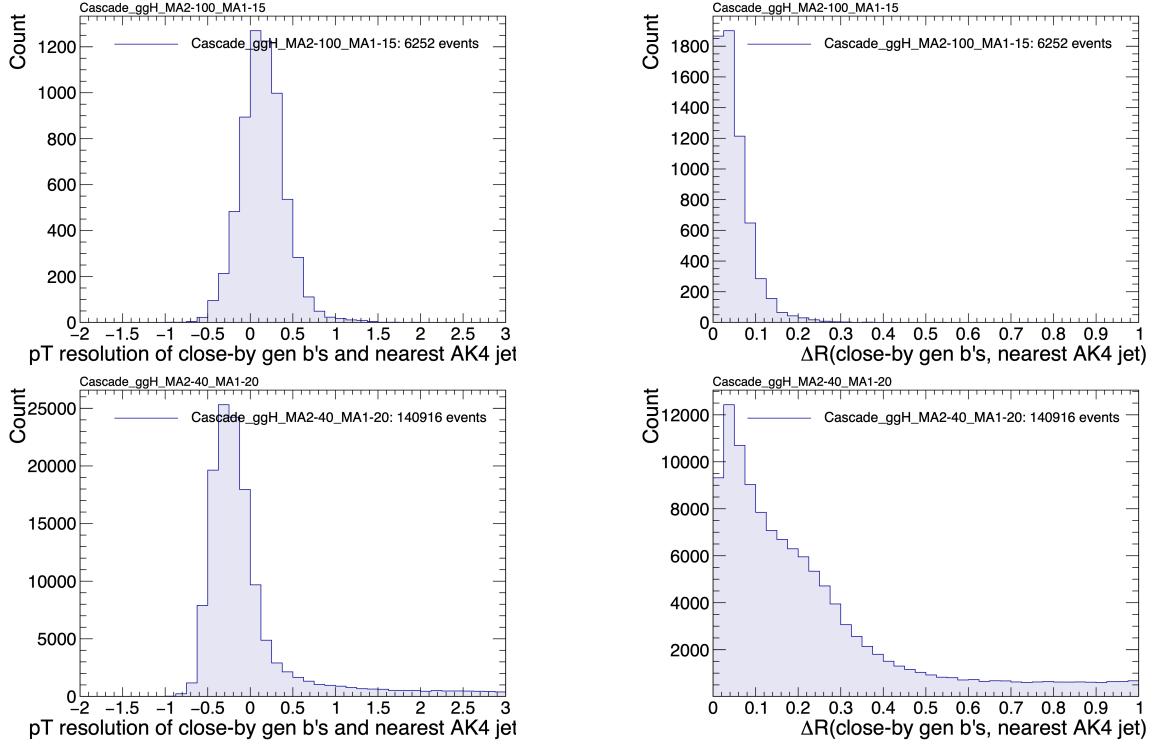


Figure 11.2: Distributions (arbitrary units) of transverse momentum p_T resolution and ΔR between the two closest generator-level b jets, treated as one object, and the nearest reconstructed AK4 jet, for two different $h \rightarrow a_1 a_2$ mass hypotheses (m_{a_1}, m_{a_2}) = (100, 15) GeV (top left, top right) and (40, 20) GeV (bottom left, bottom right) in the ggH production of the 125 GeV h . In the (40, 20) GeV mass point, the longer p_T resolution tail (bottom left) indicates that the reconstructed jet underestimates the generator b -flavor jets' energy, and the significant fraction of events with larger ΔR values (bottom right) indicate worse matching.

11.3 Current control plots for $\mu\tau_h$ channel

The $\tau\tau$ states for the $h \rightarrow a_1 a_2$ to $4b2\tau$ analysis will be similar to those studied in $h \rightarrow aa \rightarrow bb\tau\tau$. For the $\mu\tau_h$ channel, histograms of the key kinematic variables are made for data and the sum of the expected backgrounds, which are estimated from Monte Carlo samples, embedded samples, and the data-driven method for estimating jets faking τ_h as described in Chapter 7. Nominal values of the scale factors and event reweighting are applied, as described in Chapter ???. The errors shown in the figures only include statistical errors and only several of the full set of systematic errors (only those associated with the lepton energy scales and τ_h identification efficiency,

2547 described in Sections 5.3.1, 5.3.2, and 5.3.4).

2548 The p_T , η , and ϕ of the leading muon and hadronic tau τ_h , and the di-tau visible
2549 mass m_{vis} and momentum $p_{T,\text{vis}}$, are shown in Fig. 11.3. The p_T , η , and ϕ of the the
2550 leading and sub-leading b-tag jets, and the missing transverse energy magnitude and
2551 azimuthal direction, are shown in Fig. 11.4.

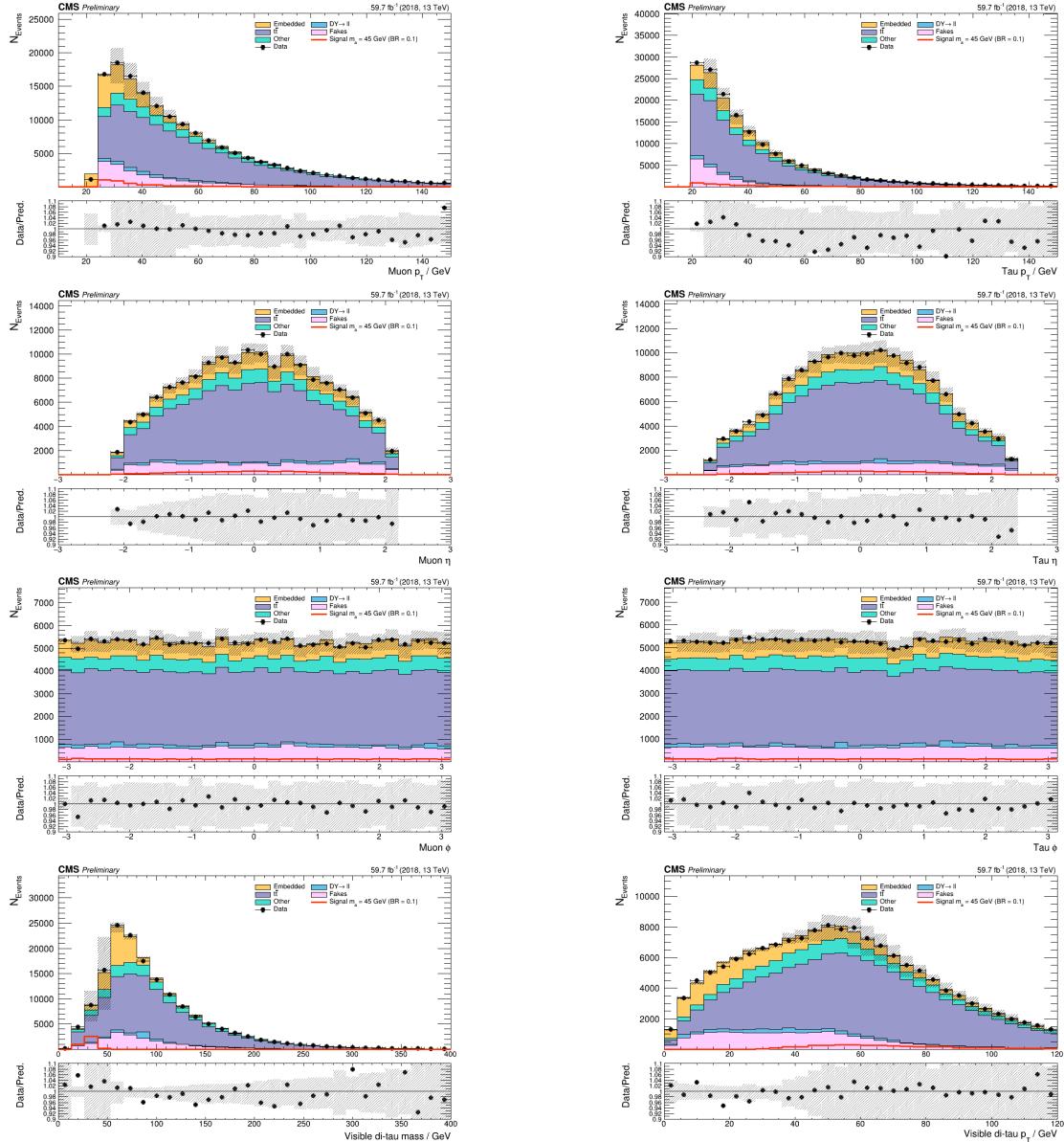


Figure 11.3: Kinematic properties of the leading muon and τ_h in the $\mu\tau_h$ channel: p_T (top row), η (second row), and ϕ (third row). The visible 4-momenta of the muon and τ_h are summed, giving the visible di-tau mass m_{vis} and transverse momentum $p_{T,\text{vis}}$. The errors shown in the figures only include statistical errors and only several of the full set of systematic errors (only those associated with the lepton energy scales and τ_h identification efficiency).

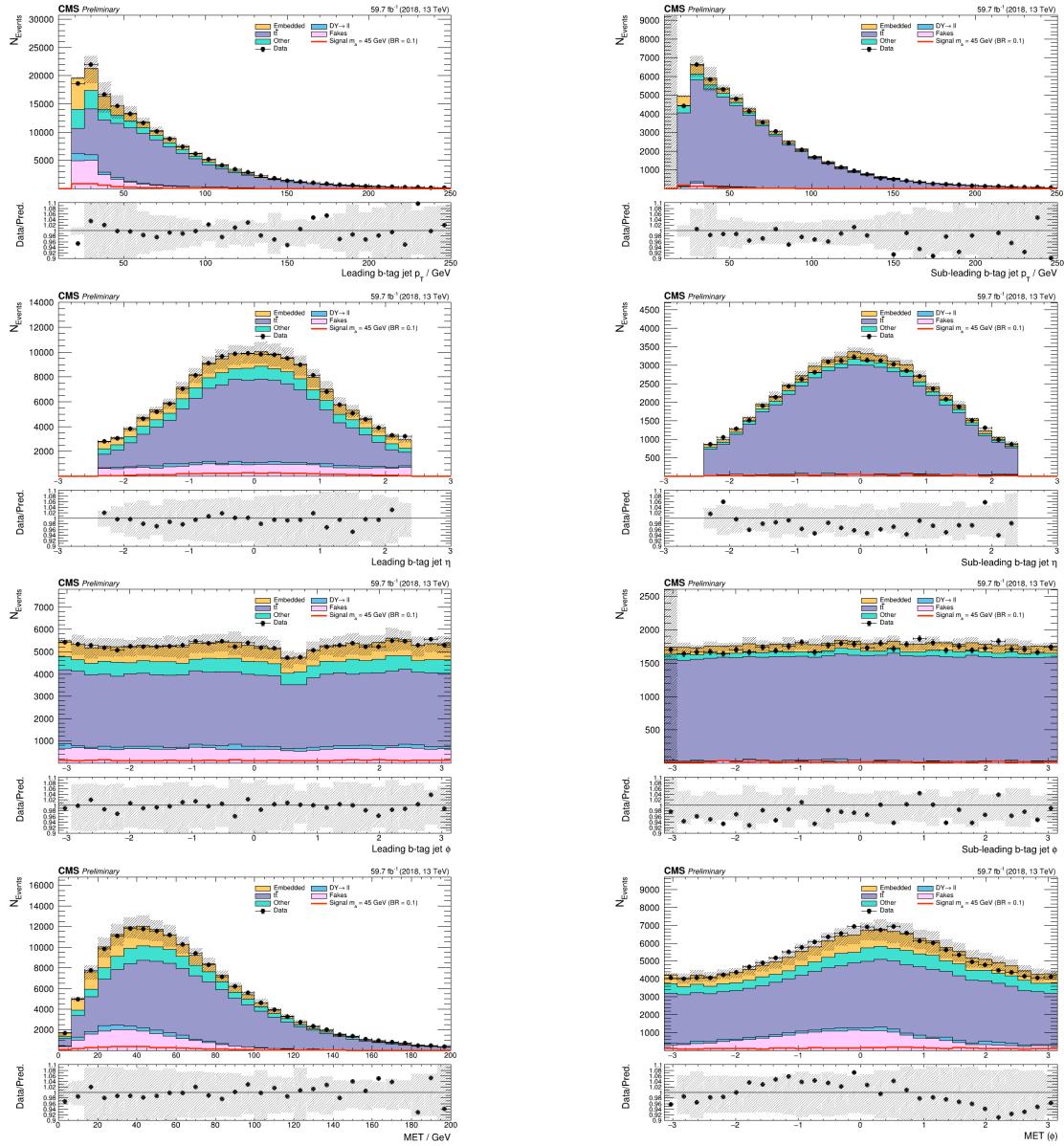


Figure 11.4: Kinematic properties of the leading and sub-leading b-tag jets in the $\mu\tau_h$ final state: jet p_T (*top row*), η (*second row*), ϕ (*third row*), as well as the missing transverse energy magnitude and azimuthal direction (*bottom row*). The errors shown in the figures only include statistical errors and only several of the full set of systematic errors (only those associated with the lepton energy scales and τ_h identification efficiency).

2552 **Chapter 12**

2553 **Conclusion and outlook**

2554 With the discovery of a Higgs boson with mass 125 GeV at the LHC in 2012, the LHC
2555 and CMS physics program has evolved to include the precise characterization of the
2556 125 GeV Higgs boson and searching for evidence of additional Higgs particles in an
2557 extended Higgs sector. This thesis presents a direct search at CMS for exotic decays
2558 of the Higgs boson with mass 125 GeV in data collected in the years 2016-2018 in
2559 proton-proton collisions at center-of-mass energy 13 TeV, to two light neutral scalar
2560 particles that decay to two bottom quarks and two tau leptons ($h \rightarrow aa \rightarrow bb\tau\tau$). The
2561 results are combined with another search that was performed in the $h \rightarrow aa \rightarrow bb\mu\mu$
2562 final state, giving the most stringent limits to date for theories with Two Higgs
2563 Doublet Models extended with a singlet scalar (2HDM+S), for pseudoscalar masses
2564 m_a ranging from 15 GeV to 60 GeV, in a number of 2HDM+S scenarios such as type
2565 II and III with $\tan\beta = 2.0$.

2566 As the rich physics program of CMS has set stringent limits on the exotic decay
2567 $h \rightarrow aa$, we turn our attention to direct searches for decays to light neutral scalars
2568 with potentially unequal mass, $h \rightarrow a_1a_2$, which has not been performed at CMS
2569 to date. Preliminary studies on $h \rightarrow a_1a_2$ signals in the Two Real Singlet Model
2570 (TRSM) are shown, and work is ongoing to develop the analysis for $h \rightarrow a_1a_2$ in final

2571 states with bottom quarks and tau leptons.

2572 To ensure the continued performance of the CMS detector and to enhance its
2573 data-taking capabilities in the intense pile-up conditions of the Phase-2 upgrade of
2574 the High-Luminosity LHC, upgrades of the Level-1 Trigger are paramount for filtering
2575 the increased data rate of the HL-LHC. This thesis presents work on the standalone
2576 barrel calorimeter algorithm for reconstructing and identifying electron and photon
2577 candidates, using high granularity crystal-level information from the ECAL subdetec-
2578 tor. For Phase-2, the increase in the granularity of information sent from the electro-
2579 magnetic calorimeter to the Level-1 trigger, from energy sums over towers (which are
2580 5×5 in crystals) to crystal-level information, allows for the implementation of a more
2581 sophisticated clustering algorithm that can exploit the fact that genuine electrons
2582 and photons tend to leave energies concentrated a 3×5 window in crystals, and use
2583 shape and isolation information to distinguish genuine electrons and photons from
2584 noise. Electrons and photons are key to characterizing Standard Model processes and
2585 performing searches for new physics, and this represents one of the many upgrades of
2586 the CMS detector in preparation for Phase-2. With the ongoing Run-3 data collecting
2587 period, and wealth of ongoing and scheduled upgrades, there remains an abundance
2588 of directions for detector development and physics at CMS heading into Phase-2 of
2589 the LHC.

2590

Appendix A

2591

Samples used

2592 The datasets used in the MiniAOD-based framework for the $h \rightarrow aa \rightarrow bb\tau\tau$ analysis
2593 are listed in this appendix. The NanoAOD-based framework uses the NanoAOD ver-
2594 sions of these datasets. The data used for the years 2016-2018 are listed in Tables A.1,
2595 A.2, and A.3 respectively. The embedded samples used for the years 2016-2018 are
2596 listed in Tables A.4, A.5, and A.6 respectively. The Monte Carlo simulated samples
2597 used to estimate backgrounds for the years 2016-2018 are listed in Tables A.7, A.8,
2598 and A.9 respectively.

2599 The $h \rightarrow aa \rightarrow bb\tau\tau$ signal samples are generated for 11 pseudoscalar masses
2600 between 12 GeV and 60 GeV for gluon fusion (ggF) and vector boson fusion (VBF)
2601 Higgs production. The 2016-2018 signal samples are listed in Tables A.10, A.11 and
2602 A.12 respectively. A filter is applied at the generator level for each $\tau\tau$ final state:

- 2603 • ee final state: $p_T(e_1) > 22$ GeV, $p_T e_2 > 10$ GeV, $|\eta(e_1)| < 2.6$, and $|\eta(e_2)| < 2.6$.
- 2604 • $e\tau_h$ final state: $p_T(e) > 22$ GeV, $p_T(\tau_h) > 16$ GeV, $|\eta(e)| < 2.6$, and $|\eta(\tau_h)| < 2.7$.
- 2605 • $e\mu$ final state: $p_T(e) > 11$ GeV, $p_T(\mu) > 7$ GeV, $|\eta(e)| < 2.6$, and $|\eta(\mu)| < 2.5$.
- 2606 • $\tau_h\tau_h$ final state: $p_T(\tau_{h1}) > 28$ GeV, $p_T(\tau_{h2}) > 28$ GeV, $|\eta(\tau_{h1})| < 2.5$, and
2607 $|\eta(\tau_{h2})| < 2.5$.

- 2608 • $\mu\tau_h$ final state: $p_T(\mu) > 19 \text{ GeV}$, $p_T(\tau_h) > 16 \text{ GeV}$, $|\eta(\mu)| < 2.5$, and $|\eta(\tau_h)| <$
 2609 2.7.
- 2610 • $\mu\mu$ final state: $p_T(\mu_1) > 17 \text{ GeV}$, $p_T(\mu_2) > 8 \text{ GeV}$, $|\eta(\mu_1)| < 2.5$, and $|\eta(\mu_2)| <$
 2611 2.5.

2612 The tables also show for each sample the filter efficiencies, which is the percentage
 2613 of events that pass the above filters, and the number of events that were generated
 2614 after applying the filters.

Channel	Data sets (2016)	Run range
$e\mu$	/MuonEG/Run2016B-17Jul2018_ver1-v1/MINIAOD	272760-273017
	/MuonEG/Run2016B-17Jul2018_ver2-v1/MINIAOD	273150-275376
	/MuonEG/Run2016C-17Jul2018-v1/MINIAOD	275656-276283
	/MuonEG/Run2016D-17Jul2018-v1/MINIAOD	276315-276811
	/MuonEG/Run2016E-17Jul2018-v2/MINIAOD	276831-277420
	/MuonEG/Run2016F-17Jul2018-v1/MINIAOD	277932-278808
	/MuonEG/Run2016G-17Jul2018-v1/MINIAOD	278820-280385
	/MuonEG/Run2016H-17Jul2018-v1/MINIAOD	281613-284044
$e\tau_h$	/SingleElectron/Run2016B-17Jul2018_ver1-v1/MINIAOD	272760-273017
	/SingleElectron/Run2016B-17Jul2018_ver2-v1/MINIAOD	273150-275376
	/SingleElectron/Run2016C-17Jul2018-v1/MINIAOD	275656-276283
	/SingleElectron/Run2016D-17Jul2018-v1/MINIAOD	276315-276811
	/SingleElectron/Run2016E-17Jul2018-v1/MINIAOD	276831-277420
	/SingleElectron/Run2016F-17Jul2018-v1/MINIAOD	277932-278808
	/SingleElectron/Run2016G-17Jul2018-v1/MINIAOD	278820-280385
	/SingleElectron/Run2016H-17Jul2018-v1/MINIAOD	281613-284044
$\mu\tau_h$	/SingleMuon/Run2016B-17Jul2018_ver1-v1/MINIAOD	272760-273017
	/SingleMuon/Run2016B-17Jul2018_ver2-v1/MINIAOD	273150-275376
	/SingleMuon/Run2016C-17Jul2018-v1/MINIAOD	275656-276283
	/SingleMuon/Run2016D-17Jul2018-v1/MINIAOD	276315-276811
	/SingleMuon/Run2016E-17Jul2018-v1/MINIAOD	276831-277420
	/SingleMuon/Run2016F-17Jul2018-v1/MINIAOD	277932-278808
	/SingleMuon/Run2016G-17Jul2018-v1/MINIAOD	278820-280385
	/SingleMuon/Run2016H-17Jul2018-v1/MINIAOD	281613-284044

Table A.1: Data sets used in the $h \rightarrow aa \rightarrow bb\tau\tau$ analysis for the 2016 era.

Channel	Data sets (2017)	Run range
$e\mu$	/MuonEG/Run2017B-31Mar2018-v1/MINIAOD	297047-299329
	/MuonEG/Run2017C-31Mar2018-v1/MINIAOD	299368-302029
	/MuonEG/Run2017D-31Mar2018-v1/MINIAOD	302031-302663
	/MuonEG/Run2017E-31Mar2018-v1/MINIAOD	303824-304797
	/MuonEG/Run2017F-31Mar2018-v1/MINIAOD	305040-306460
$e\tau_h$	/SingleElectron/Run2017B-31Mar2018-v1/MINIAOD	297047-299329
	/SingleElectron/Run2017C-31Mar2018-v1/MINIAOD	299368-302029
	/SingleElectron/Run2017D-31Mar2018-v1/MINIAOD	302031-302663
	/SingleElectron/Run2017E-31Mar2018-v1/MINIAOD	303824-304797
	/SingleElectron/Run2017F-31Mar2018-v1/MINIAOD	305040-306460
$\mu\tau_h$	/SingleMuon/Run2017B-31Mar2018-v1/MINIAOD	297047-299329
	/SingleMuon/Run2017C-31Mar2018-v1/MINIAOD	299368-302029
	/SingleMuon/Run2017D-31Mar2018-v1/MINIAOD	302031-302663
	/SingleMuon/Run2017E-31Mar2018-v1/MINIAOD	303824-304797
	/SingleMuon/Run2017F-31Mar2018-v1/MINIAOD	305040-306460

Table A.2: Data sets used in the $h \rightarrow aa \rightarrow bb\tau\tau$ analysis for the 2017 era.

Channel	Data sets (2018)	Run range
$e\mu$	/MuonEG/Run2018A-17Sep2018-v1/MINIAOD	315257-316995
	/MuonEG/Run2018B-17Sep2018-v1/MINIAOD	317080-319310
	/MuonEG/Run2018C-17Sep2018-v1/MINIAOD	319337-320065
	/MuonEG/Run2018D-PromptReco-v2/MINIAOD	320500-325175
$e\tau_h$	/EGamma/Run2018A-17Sep2018-v2/MINIAOD	315257-316995
	/EGamma/Run2018B-17Sep2018-v1/MINIAOD	317080-319310
	/EGamma/Run2018C-17Sep2018-v1/MINIAOD	319337-320065
	/EGamma/Run2018D-PromptReco-v2/MINIAOD	320497-325175
$\mu\tau_h$	/SingleMuon/Run2018A-17Sep2018-v2/MINIAOD	315257-316995
	/SingleMuon/Run2018B-17Sep2018-v1/MINIAOD	317080-319310
	/SingleMuon/Run2018C-17Sep2018-v1/MINIAOD	319337-320065
	/SingleMuon/Run2018D-PromptReco-v2/MINIAOD	320500-325175

Table A.3: Data sets used in the $h \rightarrow aa \rightarrow bb\tau\tau$ analysis for the 2018 eras.

Channel	Embedded samples (2016)
$e\mu$	/EmbeddingRun2016B/ElMuFinalState-inputDoubleMu_94X_Legacy_miniAOD-v5
	/EmbeddingRun2016C/ElMuFinalState-inputDoubleMu_94X_Legacy_miniAOD-v5
	/EmbeddingRun2016D/ElMuFinalState-inputDoubleMu_94X_Legacy_miniAOD-v5
	/EmbeddingRun2016E/ElMuFinalState-inputDoubleMu_94X_Legacy_miniAOD-v5
	/EmbeddingRun2016F/ElMuFinalState-inputDoubleMu_94X_Legacy_miniAOD-v5
	/EmbeddingRun2016G/ElMuFinalState-inputDoubleMu_94X_Legacy_miniAOD-v5
	/EmbeddingRun2016H/ElMuFinalState-inputDoubleMu_94X_Legacy_miniAOD-v5
$e\tau_h$	/EmbeddingRun2016B/ElTauFinalState-inputDoubleMu_94X_Legacy_miniAOD-v5
	/EmbeddingRun2016C/ElTauFinalState-inputDoubleMu_94X_Legacy_miniAOD-v5
	/EmbeddingRun2016D/ElTauFinalState-inputDoubleMu_94X_Legacy_miniAOD-v5
	/EmbeddingRun2016E/ElTauFinalState-inputDoubleMu_94X_Legacy_miniAOD-v5
	/EmbeddingRun2016F/ElTauFinalState-inputDoubleMu_94X_Legacy_miniAOD-v5
	/EmbeddingRun2016G/ElTauFinalState-inputDoubleMu_94X_Legacy_miniAOD-v5
	/EmbeddingRun2016H/ElTauFinalState-inputDoubleMu_94X_Legacy_miniAOD-v5
$\mu\tau_h$	/EmbeddingRun2016B/MuTauFinalState-inputDoubleMu_94X_Legacy_miniAOD-v5
	/EmbeddingRun2016C/MuTauFinalState-inputDoubleMu_94X_Legacy_miniAOD-v5
	/EmbeddingRun2016D/MuTauFinalState-inputDoubleMu_94X_Legacy_miniAOD-v5
	/EmbeddingRun2016E/MuTauFinalState-inputDoubleMu_94X_Legacy_miniAOD-v5
	/EmbeddingRun2016F/MuTauFinalState-inputDoubleMu_94X_Legacy_miniAOD-v5
	/EmbeddingRun2016G/MuTauFinalState-inputDoubleMu_94X_Legacy_miniAOD-v5
	/EmbeddingRun2016H/MuTauFinalState-inputDoubleMu_94X_Legacy_miniAOD-v5

Table A.4: Embedded samples used in the analysis for the 2016 era.

Channel	Embedded samples (2017)
$e\mu$	/EmbeddingRun2017B/ElMuFinalState-inputDoubleMu_94X_miniAOD-v2 /EmbeddingRun2017C/ElMuFinalState-inputDoubleMu_94X_miniAOD-v2 /EmbeddingRun2017D/ElMuFinalState-inputDoubleMu_94X_miniAOD-v2 /EmbeddingRun2017E/ElMuFinalState-inputDoubleMu_94X_miniAOD-v2 /EmbeddingRun2017F/ElMuFinalState-inputDoubleMu_94X_miniAOD-v2
$e\tau_h$	/EmbeddingRun2017B/ElTauFinalState-inputDoubleMu_94X_miniAOD-v2 /EmbeddingRun2017C/ElTauFinalState-inputDoubleMu_94X_miniAOD-v2 /EmbeddingRun2017D/ElTauFinalState-inputDoubleMu_94X_miniAOD-v2 /EmbeddingRun2017E/ElTauFinalState-inputDoubleMu_94X_miniAOD-v2 /EmbeddingRun2017F/ElTauFinalState-inputDoubleMu_94X_miniAOD-v2
$\mu\tau_h$	/EmbeddingRun2017B/MuTauFinalState-inputDoubleMu_94X_miniAOD-v2 /EmbeddingRun2017C/MuTauFinalState-inputDoubleMu_94X_miniAOD-v2 /EmbeddingRun2017D/MuTauFinalState-inputDoubleMu_94X_miniAOD-v2 /EmbeddingRun2017E/MuTauFinalState-inputDoubleMu_94X_miniAOD-v2 /EmbeddingRun2017F/MuTauFinalState-inputDoubleMu_94X_miniAOD-v2

Table A.5: Embedded samples used in the analysis for the 2017 era.

Channel	Embedded samples (2018)
$e\mu$	/EmbeddingRun2018A/ElMuFinalState-inputDoubleMu_102X_miniAOD-v1 /EmbeddingRun2018B/ElMuFinalState-inputDoubleMu_102X_miniAOD-v1 /EmbeddingRun2018C/ElMuFinalState-inputDoubleMu_102X_miniAOD-v1 /EmbeddingRun2018D/ElMuFinalState-inputDoubleMu_102X_miniAOD-v1
$e\tau_h$	/EmbeddingRun2018A/ElTauFinalState-inputDoubleMu_102X_miniAOD-v1 /EmbeddingRun2018B/ElTauFinalState-inputDoubleMu_102X_miniAOD-v1 /EmbeddingRun2018C/ElTauFinalState-inputDoubleMu_102X_miniAOD-v1 /EmbeddingRun2018D/ElTauFinalState-inputDoubleMu_102X_miniAOD-v1
$\mu\tau_h$	/EmbeddingRun2018A/MuTauFinalState-inputDoubleMu_102X_miniAOD-v1 /EmbeddingRun2018B/MuTauFinalState-inputDoubleMu_102X_miniAOD-v1 /EmbeddingRun2018C/MuTauFinalState-inputDoubleMu_102X_miniAOD-v1 /EmbeddingRun2018D/MuTauFinalState-inputDoubleMu_102X_miniAOD-v1

Table A.6: Embedded samples used in the analysis for the 2018 era.

Process	Simulated background samples (2016)	Cross section (pb)
DY	/DY1JetsToLL_M-50_TuneCUETP8M1	1012.5 (LO)
	/DY2JetsToLL_M-50_TuneCUETP8M1	332.8 (LO)
	/DY3JetsToLL_M-50_TuneCUETP8M1	101.8 (LO)
	/DY4JetsToLL_M-50_TuneCUETP8M1	54.8 (LO)
	/DYJetsToLL_M-50_TuneCUETP8M1	4963.0 (LO)
	/DY1JetsToLL_M-10to50_TuneCUETP8M1	730.3 (LO)
	/DY2JetsToLL_M-10to50_TuneCUETP8M1	387.4 (LO)
	/DY3JetsToLL_M-10to50_TuneCUETP8M1	95.0 (LO)
	/DY4JetsToLL_M-10to50_TuneCUETP8M1	36.7 (LO)
	/DYJetsToLL_M-10to50_TuneCUETP8M1	16290.0 (LO)
Top	/TTTo2L2Nu_TuneCP5_PSweights	88.29
	/TTToHadronic_TuneCP5_PSweights	377.96
	/TTToSemiLeptonic_TuneCP5_PSweights	365.35
	/ST_t-channel_antitop_4f_inclusiveDecays [†]	26.23
	/ST_t-channel_top_4f_inclusiveDecays [†]	44.07
	/ST_tW_antitop_5f_inclusiveDecays_TuneCUETP8M1	35.6
	/ST_tW_top_5f_inclusiveDecays_TuneCUETP8M1	35.6
VV	/VVTTo2L2Nu_13TeV_amcatnloFXFX_madspin_pythia8	13.84
	/WZTo2L2Q_13TeV_amcatnloFXFX_madspin_pythia8	5.52
	/WZTo3LNu_TuneCUETP8M1_13TeV-amcatnloFXFX-pythia8	4.43
	/ZZTo2L2Q_13TeV_amcatnloFXFX_madspin_pythia8	3.38
	/ZZTo4L_13TeV-amcatnloFXFX-pythia8	1.212
W	/W1JetsToLNu_TuneCUETP8M1	8104.0 (LO)
	/W2JetsToLNu_TuneCUETP8M1	2793.0 (LO)
	/W3JetsToLNu_TuneCUETP8M1	992.5 (LO)
	/W4JetsToLNu_TuneCUETP8M1	544.3 (LO)
	/WJetsToLNu_TuneCUETP8M1	52940.0 (LO)
H	/GluGluHToTauTau_M125	48.58*0.0627
	/GluGluHToWWTo2L2Nu_M125	48.58*0.2137*0.3258*0.3258
	/GluGluZH_HToWW_M125	0.1227*0.2137
	/HWminusJ_HToWW_M125	0.5328*0.2137
	/HWplusJ_HToWW_M125	0.840*0.2137
	/HZJ_HToWW_M125	0.7612*0.2137
	/VBFHToTauTau_M125	3.782*0.0627
	/VBFHToWWTo2L2Nu_M125	3.782*0.2137*0.3258*0.3258
	/WminusHToTauTau_M125	0.5328*0.0627
	/WplusHToTauTau_M125	0.840*0.0627
	/ZHToTauTau_M125	0.7612*0.0627
	/ggZH_HToTauTau_ZToLL_M125	0.1227*0.0627*3*0.033658
	/ggZH_HToTauTau_ZToNuNu_M125	0.1227*0.0627*0.2000
	/ggZH_HToTauTau_ZToQQ_M125	0.1227*0.0627*0.6991
	/ttHToNonbb_M125_TuneCUETP8M2_ttHtranche3	0.5071*(1-0.5824)
	/ttHTobb_M125_TuneCP5	0.5071*0.5824

Table A.7: Background MC samples used in the analysis for the 2016 era. Samples marked with a [†] are generated with the powhegV2-madspin-pythia8 tag.

Process	Simulated background samples (2017)	Cross section (pb)
DY	DY1JetsToLL_M-50_TuneCP5	877.8 (LO)
	DY2JetsToLL_M-50_TuneCP5	304.4 (LO)
	DY3JetsToLL_M-50_TuneCP5	111.5 (LO)
	DY4JetsToLL_M-50_TuneCP5	44.0 (LO)
	DYJetsToLL_M-50_TuneCP5	5343.0 (LO)
	DYJetsToLL_M-10to50_TuneCP5	15810.0 (LO)
Top	TTTo2L2Nu_TuneCP5	88.29
	TTToHadronic_TuneCP5	377.96
	TTToSemileptonic_TuneCP5	365.35
	ST_t-channel_antitop_4f_inclusiveDecays_TuneCP5 [†]	80.94
	ST_t-channel_top_4f_inclusiveDecays_TuneCP5 [†]	136.02
	ST_tW_antitop_5f_inclusiveDecays_TuneCP5	35.85
	ST_tW_top_5f_inclusiveDecays_TuneCP5	35.85
VV	VVTo2L2Nu_13TeV_amcatnloFXFX_madspin_pythia8	13.84
	WZTo2L2Q_13TeV_amcatnloFXFX_madspin_pythia8	5.52
	WZTo3LNu_TuneCP5_13TeV-amcatnloFXFX-pythia8	4.43
	ZZTo2L2Q_13TeV_amcatnloFXFX_madspin_pythia8	3.38
	ZZTo4L_TuneCP5_13TeV-amcatnloFXFX-pythia8	1.212
W	W1JetsToLNu_TuneCP5	8104.0 (LO)
	W2JetsToLNu_TuneCP5	2793.0 (LO)
	W3JetsToLNu_TuneCP5	992.5 (LO)
	W4JetsToLNu_TuneCP5	544.3 (LO)
	WJetsToLNu_TuneCP5	52940.0 (LO)
H	GluGluHToTauTau_M125	48.58*0.0627
	GluGluHToWWTo2L2Nu_M125 ^{††}	48.58*0.2137*0.3258*0.3258
	GluGluZH_HToWW_M125	0.1227*0.2137
	HWminusJ_HToWW_M125	0.5328*0.2137
	HWplusJ_HToWW_M125	0.840*0.2137
	HZJ_HToWW_M125 ^{††}	0.7612*0.2137
	VBFHToTauTau_M125	3.782*0.0627
	VBFHToWWTo2L2Nu_M125 ^{††}	3.782*0.2137*0.3258*0.3258
	WminusHToTauTau_M125	0.5328*0.0627
	WplusHToTauTau_M125	0.840*0.0627
	ZHToTauTau_M125	0.7612*0.0627
	ggZH_HToTauTau_ZToLL_M125	0.1227*0.0627*3*0.033658
	ggZH_HToTauTau_ZToNuNu_M125	0.1227*0.0627*0.2000
	ggZH_HToTauTau_ZToQQ_M125	0.1227*0.0627*0.6991
	ttHToNonbb_M125_TuneCP5	0.5071*(1-0.5824)
	ttHTobb_M125_TuneCP5	0.5071*0.5824

Table A.8: Background MC samples used in the analysis for the 2017 era. All samples use powheg, except the DYJets and WJets samples, which use madgraphMLM. Samples marked with a [†], ^{††}, or ^{†††} had the tags powhegV2-madspin-pythia8, powheg2_JHUGenV714_pythia8, or powheg_jhugen724_pythia8 respectively.

Process	Simulated background samples (2018)	Cross section (pb)
DY	DY1JetsToLL_M-50_TuneCP5	877.8 (LO)
	DY2JetsToLL_M-50_TuneCP5	304.4 (LO)
	DY3JetsToLL_M-50_TuneCP5	111.5 (LO)
	DY4JetsToLL_M-50_TuneCP5	44.0 (LO)
	DYJetsToLL_M-50_TuneCP5	5343.0 (LO)
	DYJetsToLL_M-10to50_TuneCP5	15810.0 (LO)
Top	TTTo2L2Nu_TuneCP5	88.29
	TTToHadronic_TuneCP5	377.96
	TTToSemiLeptonic_TuneCP5	365.35
	ST_t-channel_antitop_4f_InclusiveDecays_TuneCP5 [†]	80.94
	ST_t-channel_top_5f_TuneCP5 [†]	136.02
	ST_tW_antitop_5f_inclusiveDecays_TuneCP5	35.85
VV	ST_tW_top_5f_inclusiveDecays	35.85
	VVTo2L2Nu_13TeV_amcatnloFXFX_madspin	13.84
	WZTo2L2Q_13TeV_amcatnloFXFX_madspin	5.52
	WZTo3LNu_TuneCP5_13TeV-amcatnloFXFX-pythia8	4.43
	ZZTo2L2Q_13TeV_amcatnloFXFX_madspin	3.38
W	ZZTo4L_TuneCP5_13TeV-amcatnloFXFX-pythia8	1.212
	W1JetsToLNu_TuneCP5	8104.0 (LO)
	W2JetsToLNu_TuneCP5	2793.0 (LO)
	W3JetsToLNu_TuneCP5	992.5 (LO)
	W4JetsToLNu_TuneCP5	544.3 (LO)
H	WJetsToLNu_TuneCP5	52940.0 (LO)
	GluGluHToTauTau_M125	48.58*0.0627
	GluGluHToWWTo2L2Nu_M125 ^{††}	48.58*0.2137*0.3258*0.3258
	GluGluZH_HToWW_M125	0.1227*0.2137
	HWminusJ_HToWW_M125 ^{†††}	0.5328*0.2137
	HWplusJ_HToWW_M125 ^{†††}	0.840*0.2137
	HZJ_HToWW_M125 ^{††}	0.7612*0.2137
	VBFHToTauTau_M125	3.782*0.0627
	VBFHToWWTo2L2Nu_M125 ^{†††}	3.782*0.2137*0.3258*0.3258
	WminusHToTauTau_M125	0.5328*0.0627
	WplusHToTauTau_M125	0.840*0.0627
	ZHToTauTau_M125	0.7612*0.0627
	ggZH_HToTauTau_ZToLL_M125	0.1227*0.0627*3*0.033658
	ggZH_HToTauTau_ZToNuNu_M125	0.1227*0.0627*0.2000
	ggZH_HToTauTau_ZToQQ_M125	0.1227*0.0627*0.6991
	ttHTobb_M125_TuneCP5	0.5071*(1-0.5824)
	ttHTobb_M125_TuneCP5	0.5071*0.5824

Table A.9: Background Monte Carlo samples used in the analysis for the 2018 era. All samples listed are generated for 13 TeV collisions and use pythia8. All samples use powheg, except the DYJets and WJets samples, which use madgraphMLM. Samples marked with a [†], ^{††}, or ^{†††}, had the tags powhegV2-madspin-pythia8, powheg2_JHUGenV714_pythia8, and powheg_jhugen724_pythia8 respectively.

Signal samples (2016)	# events	Filter eff.
/SUSYGluGluToHToAA_AToBB_AToTauTau_M-12_FilterTauTauTrigger	0.4M	3.81%
/SUSYGluGluToHToAA_AToBB_AToTauTau_M-15_FilterTauTauTrigger	0.4M	3.54%
/SUSYGluGluToHToAA_AToBB_AToTauTau_M-20_FilterTauTauTrigger	1M	3.37
/SUSYGluGluToHToAA_AToBB_AToTauTau_M-25_FilterTauTauTrigger	0.2M	3.56%
/SUSYGluGluToHToAA_AToBB_AToTauTau_M-30_FilterTauTauTrigger	0.2M	3.16%
/SUSYGluGluToHToAA_AToBB_AToTauTau_M-35_FilterTauTauTrigger	0.2M	3.30%
/SUSYGluGluToHToAA_AToBB_AToTauTau_M-40_FilterTauTauTrigger	1M	3.30%
/SUSYGluGluToHToAA_AToBB_AToTauTau_M-45_FilterTauTauTrigger	0.2M	3.23%
/SUSYGluGluToHToAA_AToBB_AToTauTau_M-50_FilterTauTauTrigger	0.2M	3.42%
/SUSYGluGluToHToAA_AToBB_AToTauTau_M-55_FilterTauTauTrigger	0.2M	3.65%
/SUSYGluGluToHToAA_AToBB_AToTauTau_M-60_FilterTauTauTrigger	1M	3.73
/SUSYVBFHToAA_AToBB_AToTauTau_M-12_FilterTauTauTrigger	0.2M	7.94%
/SUSYVBFHToAA_AToBB_AToTauTau_M-15_FilterTauTauTrigger	0.2M	7.38%
/SUSYVBFHToAA_AToBB_AToTauTau_M-20_FilterTauTauTrigger	0.2M	7.27%
/SUSYVBFHToAA_AToBB_AToTauTau_M-25_FilterTauTauTrigger	0.2M	7.21%
/SUSYVBFHToAA_AToBB_AToTauTau_M-30_FilterTauTauTrigger	0.2M	6.87%
/SUSYVBFHToAA_AToBB_AToTauTau_M-35_FilterTauTauTrigger	0.2M	6.80%
/SUSYVBFHToAA_AToBB_AToTauTau_M-40_FilterTauTauTrigger	0.2M	6.78%
/SUSYVBFHToAA_AToBB_AToTauTau_M-45_FilterTauTauTrigger	0.2M	6.56%
/SUSYVBFHToAA_AToBB_AToTauTau_M-50_FilterTauTauTrigger	0.2M	6.40%
/SUSYVBFHToAA_AToBB_AToTauTau_M-55_FilterTauTauTrigger	0.2M	6.54%
/SUSYVBFHToAA_AToBB_AToTauTau_M-60_FilterTauTauTrigger	0.2M	6.55%

Table A.10: Signal samples used in the analysis for the 2016 era. All belong to the RunIISummer16MiniAODv3 campaign and are produced with the tags TuneCUETP8M1_13TeV_madgraph_pythia8 PUMoriond17_94X_mcRun2_asymptotic_v3. The second column is the number of events after the generator-level filter is applied, and the third column is the filter efficiency (percentage of all events that pass the generator-level filter).

Signal samples (2017)	# events	Filter eff.
/SUSYGluGluToHToAA_AToBB_AToTauTau_M-12_FilterTauTauTrigger	0.4M	3.78%
/SUSYGluGluToHToAA_AToBB_AToTauTau_M-15_FilterTauTauTrigger	0.4M	3.55%
/SUSYGluGluToHToAA_AToBB_AToTauTau_M-20_FilterTauTauTrigger	1M	3.40%
/SUSYGluGluToHToAA_AToBB_AToTauTau_M-25_FilterTauTauTrigger	0.2M	3.32%
/SUSYGluGluToHToAA_AToBB_AToTauTau_M-30_FilterTauTauTrigger	0.2M	3.36%
/SUSYGluGluToHToAA_AToBB_AToTauTau_M-35_FilterTauTauTrigger	0.2M	3.27%
/SUSYGluGluToHToAA_AToBB_AToTauTau_M-40_FilterTauTauTrigger	1M	3.03%
/SUSYGluGluToHToAA_AToBB_AToTauTau_M-45_FilterTauTauTrigger	0.2M	3.03%
/SUSYGluGluToHToAA_AToBB_AToTauTau_M-50_FilterTauTauTrigger	0.2M	3.31%
/SUSYGluGluToHToAA_AToBB_AToTauTau_M-55_FilterTauTauTrigger	0.2M	3.56%
/SUSYGluGluToHToAA_AToBB_AToTauTau_M-60_FilterTauTauTrigger	1M	3.95%
/SUSYVBFHToAA_AToBB_AToTauTau_M-12_FilterTauTauTrigger	0.2M	7.73%
/SUSYVBFHToAA_AToBB_AToTauTau_M-15_FilterTauTauTrigger	0.2M	7.35%
/SUSYVBFHToAA_AToBB_AToTauTau_M-20_FilterTauTauTrigger	0.2M	7.33%
/SUSYVBFHToAA_AToBB_AToTauTau_M-25_FilterTauTauTrigger	0.2M	7.23%
/SUSYVBFHToAA_AToBB_AToTauTau_M-30_FilterTauTauTrigger	0.2M	6.84%
/SUSYVBFHToAA_AToBB_AToTauTau_M-35_FilterTauTauTrigger	0.2M	6.97%
/SUSYVBFHToAA_AToBB_AToTauTau_M-40_FilterTauTauTrigger	0.2M	6.17%
/SUSYVBFHToAA_AToBB_AToTauTau_M-45_FilterTauTauTrigger	0.2M	6.67%
/SUSYVBFHToAA_AToBB_AToTauTau_M-50_FilterTauTauTrigger	0.2M	6.61%
/SUSYVBFHToAA_AToBB_AToTauTau_M-55_FilterTauTauTrigger	0.2M	6.51%
/SUSYVBFHToAA_AToBB_AToTauTau_M-60_FilterTauTauTrigger	0.2M	6.71%

Table A.11: Signal samples used in the analysis for the 2017 era. All belong to the RunIIFall17MiniAODv2 campaign and are produced with the tags TuneCP5_13TeV_madgraph_pythia8 and PU2017_12Apr2018_94X_mc2017_realistic_v14. The second column is the number of events after the generator-level filter is applied, and the third column is the filter efficiency (percentage of all events that pass the generator-level filter).

Signal samples (2018)	# events	Filter eff.
/SUSYGluGluToHToAA_AToBB_AToTauTau_M-12_FilterTauTauTrigger	0.4M	3.78%
/SUSYGluGluToHToAA_AToBB_AToTauTau_M-15_FilterTauTauTrigger	0.4M	3.49%
/SUSYGluGluToHToAA_AToBB_AToTauTau_M-20_FilterTauTauTrigger	1M	3.36%
/SUSYGluGluToHToAA_AToBB_AToTauTau_M-25_FilterTauTauTrigger	0.2M	3.46%
/SUSYGluGluToHToAA_AToBB_AToTauTau_M-30_FilterTauTauTrigger	0.2M	3.18%
/SUSYGluGluToHToAA_AToBB_AToTauTau_M-35_FilterTauTauTrigger	0.2M	3.28%
/SUSYGluGluToHToAA_AToBB_AToTauTau_M-40_FilterTauTauTrigger	1M	3.10%
/SUSYGluGluToHToAA_AToBB_AToTauTau_M-45_FilterTauTauTrigger	0.2M	3.21%
/SUSYGluGluToHToAA_AToBB_AToTauTau_M-50_FilterTauTauTrigger	0.2M	3.14%
/SUSYGluGluToHToAA_AToBB_AToTauTau_M-55_FilterTauTauTrigger	0.2M	3.56%
/SUSYGluGluToHToAA_AToBB_AToTauTau_M-60_FilterTauTauTrigger	1M	3.38%
/SUSYVBFHToAA_AToBB_AToTauTau_M-12_FilterTauTauTrigger	0.2M	7.78%
/SUSYVBFHToAA_AToBB_AToTauTau_M-15_FilterTauTauTrigger	0.2M	7.52%
/SUSYVBFHToAA_AToBB_AToTauTau_M-20_FilterTauTauTrigger	0.2M	6.87%
/SUSYVBFHToAA_AToBB_AToTauTau_M-25_FilterTauTauTrigger	0.2M	7.21%
/SUSYVBFHToAA_AToBB_AToTauTau_M-30_FilterTauTauTrigger	0.2M	6.51%
/SUSYVBFHToAA_AToBB_AToTauTau_M-35_FilterTauTauTrigger	0.2M	6.95%
/SUSYVBFHToAA_AToBB_AToTauTau_M-40_FilterTauTauTrigger	0.2M	6.81%
/SUSYVBFHToAA_AToBB_AToTauTau_M-45_FilterTauTauTrigger	0.2M	6.62%
/SUSYVBFHToAA_AToBB_AToTauTau_M-50_FilterTauTauTrigger	0.2M	6.56%
/SUSYVBFHToAA_AToBB_AToTauTau_M-55_FilterTauTauTrigger	0.2M	6.64%
/SUSYVBFHToAA_AToBB_AToTauTau_M-60_FilterTauTauTrigger	0.2M	6.75%

Table A.12: Signal samples used in the analysis for the 2018 era. All belong to the RunIIAutumn18MiniAOD campaign and are produced with `_TuneCP5_13TeV_madgraph_pythia8`, with the tag `102X_upgrade2018_realistic_v15`. The second column is the number of events after the generator-level filter is applied, and the third column is the filter efficiency (percentage of all events that pass the generator-level filter).

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