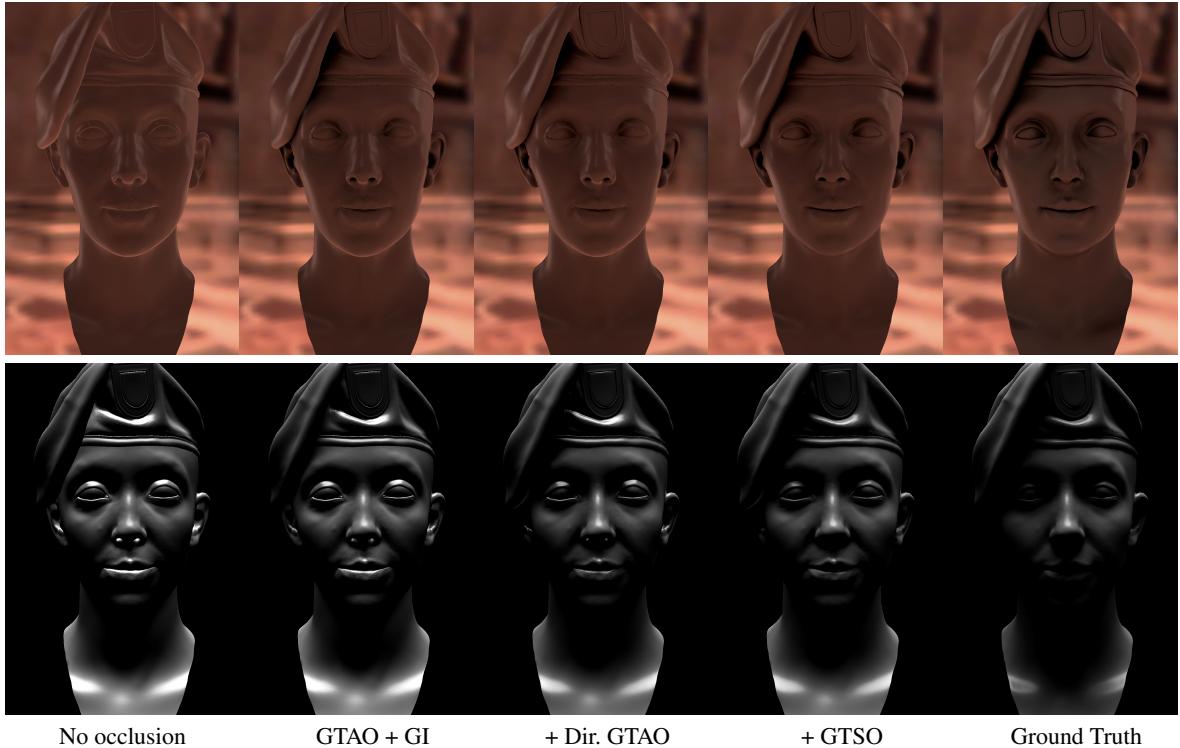


# Practical Real-Time Strategies for Accurate Indirect Occlusion

Technical Memo ATVI-TR-19-01

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**Figure 1:** Example renders with our practical occlusion techniques under illumination from the Grace light probe (top) and a high-frequency binary probe (bottom). From left to right: no occlusion, our GTAO with multiple-bounces, with spherical harmonics directional occlusion, with our GTSO modeling specular occlusion, and a ray-traced ground truth. Our techniques render high-quality occlusion matching the ground truth, with the baseline GTAO + GI rendering in just 0.5 ms on a PS4 at 1080p (for a standard halfres occlusion buffer).

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## Abstract

In this work we introduce a set of techniques for real-time ambient occlusion targeted to very tight budgets. We propose GTAO, a new formulation of screen-space ambient occlusion that allows the composited occluded illumination to match the ground truth reference in half a millisecond on current console hardware. This is done by using a radiometrically-correct formulation of the ambient occlusion equation, and an efficient implementation that distributes computation using spatio-temporal sampling. As opposed to previous methods, our technique incorporates the energy lost by missing interreflections by using an efficient, accurate physically-based parametric form, avoiding the use of ad-hoc approximations of indirect illumination. Then, we extend GTAO to account for directionally-resolved illumination, by fastly projecting coupled visibility and foreshortening factors into spherical harmonics, and thoroughly analyze with previous work. Finally, we introduce a novel model for specular occlusion formulation that accounts for the coupling between visibility and BRDF, closely matching the ground truth specular illumination from probe-based lighting, and propose GTSO, an efficient implementation of this concept based on tabulation. Our techniques are practical real-time, give results close to the ray-traced ground truth, and have been integrated in recent AAA console titles.

1 **1. Introduction**

2 Ambient occlusion (AO) is an approximation of global illumination,  
 3 that models the diffuse shadows produced by close, potentially small  
 4 occluders in a tight budget. It allows to preserve high-frequency details and contrast in low-frequency precomputed  
 5 indirect illumination via pre-baked illumination or light probes. Unfortunately, solving the ambient occlusion integral is still too expensive  
 6 to be practical in certain scenarios (e.g. 1080p or 4K rendering at 60 fps), so approximations have been developed in the past to  
 7 achieve the target performance budget.

8 We introduce a new set of screen-space occlusion techniques,  
 9 that target *practical* real-time performance while matching ray-traced ground truth solutions. We propose a novel technique for  
 10 ambient occlusion, that we call *ground truth-based ambient occlusion (GTAO)*, that decouples ambient occlusion from the near-range  
 11 indirect illumination. This allows us to solve efficiently the ambient  
 12 occlusion integral by avoiding piecewise integration (as required  
 13 when using obscurrence estimators), while recovering the lost multiple  
 14 scattered diffuse lighting by using an efficient physically-based  
 15 functional approximation. This allows to match not only ground  
 16 truth *occlusion*, but also *illumination* references. Then, we extend  
 17 our ambient occlusion model to directionally-resolved illumination  
 18 from distant probes, that uses our accurate ambient occlusion term  
 19 and our from-horizons bent normal calculations to derive an effi-  
 20 cient expansion in spherical harmonics, that can be used to effi-  
 21 ciently integrate ambient illumination. Finally, we generalize am-  
 22 bient occlusion for arbitrary specular materials and formulate it  
 23 by using a novel split-integral formulation that couples the BRDF  
 24 with the visibility. We propose an efficient implementation of this  
 25 formulation, that we call *ground truth-based specular occlusion (GTSO)*, to compute it in runtime by accessing a small precom-  
 26 puted table.

27 In particular, our contributions are:

- 28 • **GTAO:** An efficient ambient occlusion technique that matches  
 29 a radiometrically-correct ambient occlusion integral, and incor-  
 30 porates the lost energy due to close-range indirect illumination  
 31 using a simple closed-form analytical expression.
- 32 • **Directional GTA0:** an extension that accounts for directionally-  
 33 resolved distant illumination, which includes a ground truth  
 34 derivation of horizon-based bent normals.
- 35 • **Specular occlusion (SO):** A generalization of the standard am-  
 36 bient occlusion formulation for arbitrary specular BRDFs that  
 37 couples visibility and reflectance for efficiently computing spec-  
 38 ular reflection from distant probes. This formulation matches  
 39 ground truth references under the same assumptions as ambient  
 40 occlusion (uniform dome and a single bounce), and is one of the  
 41 principal results of our work.
- 42 • **GTSO:** An efficient implementation of this specular formulation  
 43 for microfacets-based BRDFs.

44 Figure 1 shows the effect of these techniques, and how their com-  
 45 bination match the Monte Carlo raytraced ground truth. We imple-  
 46 ment them efficiently, leveraging temporal reprojection and spatial  
 47 filtering to compute our baseline ambient occlusion in just 0.5 ms  
 48 per frame on a Sony Playstation 4, for a game running at 1080p (us-  
 49 ing a standard halfres occlusion buffer). Our results highlight that

50 for today hardware standards, performing ad-hoc occlusion calcu-  
 51 lations are no longer necessary for performance reasons anymore.

52 **2. Related Work**

53 Given the large amount of previous work on global illumination in  
 54 general, and ambient occlusion in particular, here we focus on the  
 55 most related works with ours. For a wider overview on the field we  
 56 refer to the surveys by Ritschel et al. [RDGK12], and Aalund and  
 57 Bærentzen [AB12].

58 **Screen-Space Ambient Occlusion** Ambient occlusion [ZIK98]  
 59 integrates the visibility from a point in the scene, to modulate the  
 60 ambient illumination term. It requires to perform expensive visibility  
 61 queries from the shaded point. In order to alleviate the overdarken-  
 62 ing resulting from ignoring interreflections the visibility is com-  
 63 monly modulated by an ad-hoc fall-off function; in these cases, it is  
 64 common to term AO as *ambient obscurrence*. In his seminal work,  
 65 Mittring [Mit07] proposed to move the visibility queries to screen-  
 66 space, assuming that only the geometry visible from the camera  
 67 acts as occluder. He approximated ambient occlusion by sampling  
 68 the depth map of the scene, and evaluated whether a point is oc-  
 69 cluded (behind) geometry in the depth map, effectively calculating  
 70 volumetric occlusion using point samples. Several works have im-  
 71 proved the sampling strategy [LS10, SKUT\*10, HSEE15], by in-  
 72 tegrating using line samples rather than points. While they obtain  
 73 high quality results, those methods simplify the integral function  
 74 resulting into radiometrically-incorrect ambient occlusion<sup>†</sup>. Bavoil  
 75 et al. [BSD08] proposed to perform line integrals based on the  
 76 horizon angles of the geometry around  $x$  using screen space ray  
 77 tracing. They termed their technique *horizon based AO* (HBAO).  
 78 This work is similar in spirit to volumetric line sampling ap-  
 79 proaches in that it realizes that any ray under the horizon will be  
 80 occluded if the horizon was already occluded. McGuire and col-  
 81 leagues [MOBH11, MML12] later simplified the ray tracing pro-  
 82 cess by assuming that  $x$  and any near-field position on the pos-  
 83 itive hemisphere are mutually visible. While HBAO and its im-  
 84 provements are efficient, they are not radiometrically correct, and  
 85 does not account for multiple scattering in the near field. Timo-  
 86 nen [Tim13a, ST15] improves HBAO by performing line sweeps  
 87 along all the image, finding the maximum horizon angle for a given  
 88 direction in constant time by amortizing samples over many pixels.  
 89 Closely related to our GTA0, the same author [Tim13b] proposed  
 90 a radiometrically-correct estimator for ambient obscurrence by line-  
 91 scanning and filtering the depth map, which is able to match the  
 92 raytraced *obscurrence* ground truth at small cost even for very large  
 93 gathering radii. While the technique yields impressive results, the  
 94 use of a obscurrence estimator prevents matching ground truth *il-  
 95 lumination* (rather matching obscurrence), and obligates the usage  
 96 of a piecewise inner integral for the occlusion computations, im-  
 97 plemented with a look up table. Our work efficiently computes  
 98 radiometrically-correct ambient occlusion based on visibility hori-  
 99 zons. It does not require ad-hoc fall-off functions to avoid over-  
 100 darkening, since indirect illumination is accounted by a physically-  
 101 based parametric formula. This allows to reduce AO computa-  
 102 tions

<sup>†</sup> By "radiometrically-correct" we mean that foreshortening is taken into account in the ambient occlusion integral.

108 to its bare bones by solving the inner integral analytically. In addition,  
 109 we generalize GTAO to directional and non-Lambertian oc-  
 110 clusion.

111 **Directional Occlusion** While AO has received significant attention,  
 112 only a few works have focused on introducing the directional dependence  
 113 of ambient illumination encoded in e.g. probes. Rather than relying on a pre-filtered probe, Ritschel et al. [REG\*09]  
 114 approximate directional diffuse lighting by evaluating the render  
 115 equation on the fly. Despite of using an approximate visibility test,  
 116 it is too slow for practical real-time environments. Landis [Lan02]  
 117 proposed to use bent normals to fetch from the ambient probe in the  
 118 most visible direction, in order to increase the directional fidelity  
 119 of ambient occlusion. Since then, bent normals have observed a  
 120 widespread usage both for off-line and real-time rendering. Klehm  
 121 et al. [KRES11] extended SSAO [Mit07] to handle bent normals,  
 122 averaging the directions to visible samples. They also propose a  
 123 variant for HBAO using a similar rationale, although averaging  
 124 horizon directions does not match 3d ray-traced bent normals. Oat  
 125 and Sander [OS07] precalculated visibility by means of ambient oc-  
 126 clusion and bent normals, then using this information during real-  
 127 time rendering by calculating the spherical cap intersection with the  
 128 light source aperture, effectively applying visibility to diffuse light-  
 129 ing. Ramamoorthi and Hanrahan [RH01] proposed to encode light  
 130 probes into spherical harmonics, allowing to efficiently convolve  
 131 light and the foreshortening factor in real-time. Green [Gre03] con-  
 132 volved with the visibility as well by means of the triple SH product.  
 133 We build on these ideas, and propose an efficient projection into SH  
 134 of the coupled visibility and foreshortening factor in run-time that  
 135 avoids the more expensive and less accurate triple SH product.  
 136

137 **Specular Occlusion** While ambient occlusion as been largely  
 138 studied over the last two decades, specular occlusion has not re-  
 139 ceived similar attention, despite of possibly being as important,  
 140 specially with the adoption of physically-based shading models.  
 141 Gotanda [Got12] derived empirically specular occlusion from am-  
 142 bient occlusion. He noted that ambient occlusion does not con-  
 143 sider the BRDF lobe shape, resulting in a mismatching occlu-  
 144 sion scale. Lagarde [Ld14] adopted a similar empirical approach,  
 145 adapted to GGX-based microfacets using the roughness of the sur-  
 146 face to shape the resulting specular occlusion. Jimenez and von  
 147 der Pahlen. [Jv13] highlighted the importance of specular occlusion  
 148 for rendering photorealistic characters, even when using ad-hoc ap-  
 149 proaches. In contrast, we formally derive a specular occlusion term  
 150 analogous to ambient occlusion, that couples visibility and specu-  
 151 lar BRDF. In addition, we propose an efficient model for rendering  
 152 with this specular occlusion term.

### 153 3. Background & Overview

The reflected radiance  $L_r(x, \omega_0)$  from a point  $x$  with normal  $\mathbf{n}$  to-  
 154 wards a direction  $\omega_0$  can be modeled as

$$155 L_r(x, \omega_0) = \int_{\mathcal{H}^2} L(x, \omega_i) f_r(x, \omega_i, \omega_0) \langle \mathbf{n}, \omega_i \rangle^+ d\omega_i, \quad (1)$$

where  $\mathcal{H}^2$  is the hemisphere centered in  $x$  and having  $\mathbf{n}$  as its axis,  $L(x, \omega_i)$  is the incoming radiance at  $x$  from direction  $\omega_i$ ,  $f_r(x, \omega_i, \omega_0)$  is the BRDF at  $x$ , and  $\langle \mathbf{n}, \omega_i \rangle^+$  models foreshortening. Ambient occlusion [ZIK98] approximates Equation (1), by in-

troducing a set of assumptions: *i*) all surfaces around  $x$  are purely absorbing (i.e. do not bounce light), *ii*) all light comes from an infinite uniformly white environment light (or generalizing, of any uniform color), which might be occluded by the geometry around  $x$ ; and *iii*) the surface at  $x$  is a Lambertian surface. This transforms Equation (1) into

$$156 L_r(x, \omega_0) \approx L_i \frac{\rho(x)}{\pi} \int_{\mathcal{H}^2} V(x, \omega_i) \langle \mathbf{n}, \omega_i \rangle^+ d\omega_i \\ 157 = L_i \frac{\rho(x)}{\pi} \mathcal{A}(x), \quad (2)$$

158 where  $\mathcal{A}(x)$  is the *ambient occlusion* term at point  $x$ ,  $\frac{\rho(x)}{\pi}$  is the dif-  
 159 fuse BRDF with albedo  $\rho(x)$ , and  $V(x, \omega_i)$  is the visibility term at  $x$   
 160 in direction  $\omega_i$ . Previous works [ZIK98, Mit07, BSD08] have mod-  
 161 eled this visibility term  $V(x, \omega_i)$  as an attenuation function with  
 162 respect to the distance to the occluder, referring to  $\mathcal{A}(x)$  as *obscu-  
 163 rance*. This attenuation function was used as an ad-hoc solution to  
 164 avoid the typical over-darkening in AO produced by ignoring near-  
 165 field interreflections.

Ambient occlusion is only exact for uniform illumination. However, it is often used in practice for *any* illumination stored in a light probe. In these cases, the illumination is approximated as

$$166 L_r(x, \omega_0) \approx \frac{\rho(x)}{\pi} \mathcal{A}(x) \int_{\mathcal{H}^2} L(x, \omega_i) f_r(x, \omega_i, \omega_0) \langle \mathbf{n}, \omega_i \rangle^+ d\omega_i \\ 167 = \frac{\rho(x)}{\pi} \mathcal{A}(x) \mathcal{L}(x, \omega_h), \quad (3)$$

168 where  $\mathcal{L}(x, \omega_h)$  is the light probe pre-convolved with the BRDF,  
 169 and  $\omega_h$  is the query direction at the probe. Several works base on  
 170 the bent normals  $\omega_h = \mathbf{b}$  [Lan02] to fetch the probe, which is later  
 171 attenuated by the ambient occlusion term  $\mathcal{A}(x)$ . While this incor-  
 172 porates some degree of directionality in the incoming radiance, the  
 173 visibility  $V$  and lighting  $L$  terms remain decoupled in Equation (3).

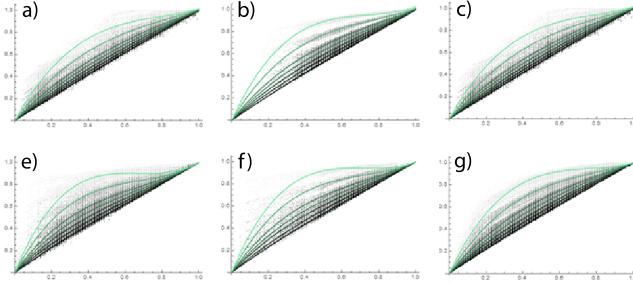
An alternatively common approach is to encode the light probe and visibility as a  $n^{\text{th}}$  order spherical harmonics (SH) expansion [RH01]. This allows to compute  $L_r(x, \omega_0)$  efficiently, as a SH double product

$$174 L_r(x, \omega_0) \approx \frac{\rho(x)}{\pi} \sum_{j=0}^n \hat{L}_j \hat{V}_j, \quad (4)$$

175 where  $\hat{L}_j$  and  $\hat{V}_j$  are the  $j^{\text{th}}$  SH coefficient for  $L$  and  $V$  respectively.  
 176 Unfortunately, introducing the foreshortening or adding a BRDF  
 177 requires an expensive triple SH product, which can limit the appli-  
 178 cability of this approach.

179 **Objectives** In this work we have two main goals: On one hand,  
 180 we propose a technique that matches the radiometrically-correct  
 181 ambient occlusion definition, while being efficient enough to be  
 182 used in demanding real-time applications. On the other hand, we  
 183 want to extend the amount of global illumination effects that can  
 184 be efficiently approximated, to not only match ground truth *occlu-  
 185 sion* but rather ground truth *illumination*, for an extended set of  
 186 material BRDFs and input lighting configurations including non-  
 187 uniform dome illuminations.

The first goal imposes severe limitations in terms of input data, number of passes, and number of instructions. Bounded by these



**Figure 2:** Mapping between the ambient occlusion (x-axis) and the global illumination (y-axis) for the scenes in Figure 17 and different albedos. A cubic polynomial (drawn in green) fits the data very well, suggesting a functional relationship between AO and GI. We develop such relationship in Section 4.2.

183 limitations, we describe in Section 4 a technique that works in  
 184 screen space, taking as inputs only the depth buffer and surface  
 185 normals (which can be derived from it by differentiation or can be  
 186 supplied separately), and that can coexist and enhance other sources  
 187 of global illumination (specifically baked irradiance).

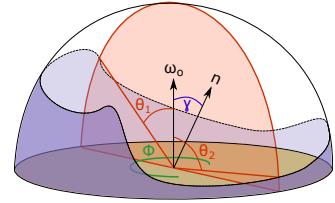
188 In order to achieve the second goal, we will relax all of the as-  
 189 sumptions done for traditional ambient occlusion. In particular: a)  
 190 we include diffuse interreflections of near-field occluders, for ad-  
 191 dressing assumption *i* (Section 4.2); b) we introduce a fast and ac-  
 192 curate Lambertian directional occlusion approach, for the uniform  
 193 dome assumption *ii* (Section 5); and c) we propose a formal direc-  
 194 tional and specular occlusion formulation, for relaxing the purely  
 195 Lambertian surface assumption *iii* (Section 6).

#### 4. GTAO: Ground Truth-based Ambient Occlusion

196 To develop an efficient model that accounts for near-field indi-  
 197 rect illumination, we make the key observation that there is a  
 198 functional relationship between the total ambient occlusion, the  
 199 surface’s albedo, and the indirect illumination reflected from  $x$ ,  
 200 as shown in Figure 2. This allows us to build a GI-aware am-  
 201 bient occlusion technique in two parts: First we compute the  
 202 radiometrically-correct ambient occlusion at  $x$  assuming binary vis-  
 203 ibility (Section 4.1), and then we reintroduce the lost indirect illu-  
 204 mination based on the computed ambient occlusion (Section 4.2).  
 205 This has two main benefits: 1) We model interreflections based on  
 206 a physically-plausible approximation, instead of an heuristic ob-  
 207 scurcence term; and 2) eliminating the empirical obscurcence term  
 208 allows us to reduce complexity by removing the piecewise inner  
 209 integration, and consequently, for computations to be performed  
 210 once per direction, rather than once per sample.

##### 4.1. Computing ambient occlusion

212 Our formulation of ambient occlusion follows the horizon-based  
 213 approach of Bavoil et al. [BSD08], which under the height field as-  
 214 sumption computes Equation (2) as an integral along an azimuthal



**Figure 3:** Diagram of our reference frame when computing horizon-based ambient occlusion. Horizons angles  $\theta_1$  and  $\theta_2$  are drawn in red, slice angle  $\phi$  is drawn in green, view direction  $\omega_o$  and normal  $\mathbf{n}$  in black, and  $\gamma$  the angle between  $\omega_o$  and  $\mathbf{n}$  in blue.

angle  $\phi$  as

$$\mathcal{A}(x) = \frac{1}{\pi} \int_0^\pi \int_{-\pi/2}^{\pi/2} V(\phi, \theta) \cos(\theta - \gamma)^+ |\sin(\theta)| d\theta d\phi, \quad (5)$$

where  $\theta$  is the polar angle along the view vector  $\omega_o$ ,  $\gamma$  is the angle between the normal  $\mathbf{n}$  and the view vector  $\omega_o$  [TW10],  $\cos(\theta)^+ = \max(\cos(\theta), 0)$ , and  $V(\phi, \theta)$  is the visibility attenuation function. Note that unlike [BSD08], this integral is written here on its radiometrically-correct form, and hence accounting for the foreshortening factor. The coordinate system has also been changed by defining  $(\phi, \theta)$  with respect to the view vector  $\omega_o$  instead of the tangent vector, which requires introducing *abs* values to account for the *sin* term of the differential solid angle. Assuming a binary visibility function  $V(\phi, \theta)$  that returns 1 when  $\theta$  is above the horizon angles  $\theta_1(\phi)$  and  $\theta_2(\phi)$ , and 0 below them (see Figure 3 for the reference system), and consequently not having per-sample attenuation, Equation (5) can be transformed as

$$\mathcal{A}(x) = \frac{1}{\pi} \int_0^\pi \underbrace{\int_{\theta_1(\phi)}^{\theta_2(\phi)} \cos(\theta - \gamma)^+ |\sin(\theta)| d\theta}_{\hat{a}} d\phi. \quad (6)$$

Given the horizon angles  $\theta_1$  and  $\theta_2$  we can solve analytically the inner integral  $\hat{a}$  in Equation (6) as

$$\begin{aligned} \hat{a}(\theta_1, \theta_2, \gamma) &= \frac{1}{4} (-\cos(2\theta_1 - \gamma) + \cos(\gamma) + 2\theta_1 \sin(\gamma)) \\ &\quad + \frac{1}{4} (-\cos(2\theta_2 - \gamma) + \cos(\gamma) + 2\theta_2 \sin(\gamma)). \end{aligned} \quad (7)$$

It is important to note that this formulation requires the normal  $\mathbf{n}$  to lay in the plane  $P$  defined by the horizon vectors, which does not hold in general. However, it can be shown that the following identity holds [Tim13b]:

$$\int_{-\pi/2}^{\pi/2} \langle \mathbf{n}, \omega_i \rangle^+ |\sin(\theta)| d\theta = \|\bar{\mathbf{n}}\| \int_{-\pi/2}^{\pi/2} \langle \frac{\bar{\mathbf{n}}}{\|\bar{\mathbf{n}}\|}, \omega_i \rangle^+ |\sin(\theta)| d\theta, \quad (8)$$

where  $\bar{\mathbf{n}}$  is the *projected* normal in  $P$ . Combining with Equation (6) we obtain

$$\mathcal{A}(x) = \frac{1}{\pi} \int_0^\pi \|\bar{\mathbf{n}}\| \hat{a}(\theta_1(\phi), \theta_2(\phi), \gamma') d\phi, \quad (9)$$

where  $\gamma' = \arccos(\langle \frac{\bar{\mathbf{n}}}{\|\bar{\mathbf{n}}\|}, \omega_o \rangle)$ .

This analytic integral can be efficiently executed only once per

direction. Additionally, after optimization only two  $\cos$  and one  $\sin$  are needed per sample, plus three additional  $\arccos$  functions per direction for setting up the integration domain, which can be efficiently approximated [Dro14].

**Computing maximum horizon angles** Core to the solution of Equation (9) is to find the maximum horizon angles  $\theta_1(\phi)$  and  $\theta_2(\phi)$  for a direction  $\hat{\mathbf{t}}(\phi)$  in the image plane, where  $\phi$  is the uniformly distributed azimuthal angle. We compute  $\theta_1(\phi)$  by ray-tracing in screen-space from the projected pixel  $\hat{x}$  of point  $x$  using  $\hat{s}(r) = \hat{x} + \hat{\mathbf{t}}(\phi) \cdot r$ , with  $r \in [0, 1]$  the parametrization of the ray. For each camera space point  $s(r)$  we compute  $\omega_s(r) = \frac{s(r) - x}{\|s(r) - x\|}$ . The maximum horizon angle with respect to the view vector  $\omega_0$  is then

$$\theta_1(\phi) = \arccos \left( \max_{r \in [0, 1]} (\langle \omega_s(r), \omega_0 \rangle^+) \right). \quad (10)$$

We compute a fixed number of discrete samples per direction. Angle  $\theta_2(\phi)$  is computed analogously with  $\hat{s}(r) = \hat{x} - \hat{\mathbf{t}}(\phi) \cdot r$ . The maximum screen-space ray tracing distance  $r$  is scaled depending on the distance from the camera; this is necessary to make  $\mathcal{A}(x)$  view-independent. We clamp the maximum  $r$  to avoid large gathering areas in objects close to the camera, which would trash the GPU cache. Algorithm 1 details GTAO computations.

## 4.2. Adding indirect illumination

Equation (9) matches the ground truth if we assume that the neighborhood of  $x$  only occludes light, and therefore no interreflections are present. This results into an energy lost, visible as an overdarkening at e.g. corners. In other words, Equation (9) computes Equation (2), but not the physically-accurate Equation (1), for which no analytical closed-form solution exists. However, as shown by Nayar et al. [NIK91], if we assume that the neighborhood  $S(x)$  of  $x$  has constant albedo  $\rho(x)$  and diffuse reflectance, we can express Equation (1) as a Neumann series as

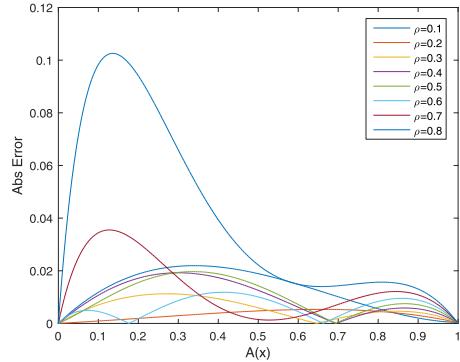
$$L_r(x, \omega_0) = L_i \frac{\rho(x)}{\pi} \mathcal{A}(x) + \sum_{m=1}^{\infty} \rho^m \int_{S(x)} K_m(x \leftarrow x') L_i(x') dx', \quad (11)$$

where  $K_m(x \leftarrow x')$  is the transfer function between  $x'$  and  $x$ . This equation relates the light incoming at the points in  $S(x)$ , the geometric relationship  $K_m(x \leftarrow x')$  between  $x$  and the points  $x' \in S(x)$ , and the ground truth total light reflected at  $x$ . Introducing the assumption of uniform illumination  $L_i$  at  $S(x)$ , and following Stewart and Langer [SL96], we can find a closed-form solution for Equation (11) as

$$\begin{aligned} L_r(x, \omega_0) &= L_i \frac{\rho}{\pi} \frac{\pi^{-1} \int_{\mathcal{H}^2} V(x, \omega_i) \langle \mathbf{n}, \omega_i \rangle^+ d\omega_i}{1 - \rho (1 - \pi^{-1} \int_{\mathcal{H}^2} V(x, \omega_i) \langle \mathbf{n}, \omega_i \rangle^+ d\omega_i)} \\ &= L_i \frac{\rho}{\pi} \frac{\mathcal{A}(x)}{1 - \rho (1 - \mathcal{A}(x))}. \end{aligned} \quad (12)$$

Equation (12) accurately computes the indirect illumination as a function of the ambient occlusion at  $x$ , with just a few algebraic operations. This allows to eliminate the need of ad-hoc obscuration operators when computing  $\mathcal{A}(x)$ .

**Accuracy analysis** We analyze the accuracy of Equation (12) by



**Figure 4:** Error comparison between the analytic approximation for indirect illumination in Equation (12) and the polynomial data-driven fitting derived from a set of simulations, with respect to the amount of ambient occlusion  $\mathcal{A}(x)$ , for different albedos  $\rho$ . Even for high albedo values in highly occluded areas (low  $\mathcal{A}(x)$ ), where indirect illumination dominates, the introduced error is below 10%.



**Figure 5:** Comparison between the samples computed on a single pixel (left), adding the spatial occlusion gathering using a bilateral reconstruction filter (middle), and adding the temporal reprojection using an exponential accumulation buffer (right). In each image we use 1, 16 and 96 effective sample directions per pixel respectively.

comparing it to Monte Carlo-based measurements of global illumination with respect to ambient occlusion (Figure 2). To analyze the average reflected radiance on that scenario and to reduce the effect of simulation variance, we fit a polynomial relating the albedo  $\rho(x)$ , the ambient occlusion term  $\mathcal{A}(x)$ , and the total reflected radiance illuminated by a uniform dome. Details can be found in Appendix B. As we can observe in Figure 4, Equation (12) provides a very accurate approximation of global illumination based on the surface's albedo and ambient occlusion for a reasonable range of surface albedos.

## 4.3. Implementation details

For a game running at 60 frames per second, around half a millisecond is a reasonable screen-space ambient occlusion budget, which makes optimization mandatory. Similarly, working in screen space imposes some limitations.

**Spatio-temporal sampling approach** We compute our ambient occlusion on half-resolution, which is later bilaterally upscaled to full resolution. Moreover, in order to compute as many samples as possible without harming the performance, we distribute the occlusion integral over both space and time: We sample the horizon



**Figure 6:** Effect of using our thickness heuristic (right) in comparison to not using it (left). In screen-space methods, thin occluders such as leaves or branches cast an unrealistic amount of occlusion, which is not temporally consistent. Our simple heuristic allows for significantly diminishing the effect of such thin occluders.



**Figure 7:** Comparison between the ground truth-based ambient occlusion computed with Monte Carlo ray-tracing (left) and our method (right) without multiple scattering (Equation (9)). Our method closely matches the ground truth, while being significantly faster to compute.

in only one direction per pixel (including both sides of a direction, with 12 steps in total) but use the information gathered on a neighborhood of  $4 \times 4$  using a bilateral filter for reconstruction, using uniform convolution weights. To generate per-pixel directions we use a tileable spatial uniform noise of  $4 \times 4$ . In addition, we make aggressive use of temporal coherency by alternating between 6 different rotations and reprojecting the results, using an exponential accumulation buffer. All this gives a total of  $4 \times 4 \times 6 = 96$  effective sampled directions per pixel. Figure 5 shows the effect of the spatial and temporal gathering on the final reconstruction. We opted for a regular sampling approach rather than using line sweeps [Tim13a] because it fitted better our tight budget and target quality (single direction of 12 steps per pixel). Line-sweep ambient occlusion achieves very high quality results, but unfortunately requires a high scan direction count to avoid banding, given the impossibility of randomizing directions per-pixel. Silvennoinen et al. [ST15] reported a cost of 1.6ms for a  $1280 \times 720$  image on the Xbox One, representing a different tradeoff than our approach.

**Bounding the sampling area** As opposed to ambient obscurance techniques, in our formulation we do not use an attenuation function. However, we only want to calculate local ambient occlusion, as larger-range low-frequency occlusion can be computed using baked irradiance or occlusion. We compute near-field occlusion using our formulation, and combine it with baked far-field occlusion by calculating their minimum. In order to minimize artifacts we employ a conservative attenuation strategy. We linearly interpolate the current sample horizon angle cosine ( $\cos(\theta)$ ) towards  $-1$  when we exceed the near-field occlusion area, meaning that any sample outside of the near-field will be progressively attenuated, and with all the other samples remaining unmodified.

#### 4.4. Results

We implement our GTAO in an DirectX stand-alone application. In all cases, we compare against a Monte Carlo ray-traced reference. Figure 7 compares our GTAO without global illumination (ambient occlusion only) against the ground truth: Our technique matches the ray-traced results, while being practical for games at 1080p and 60 fps. Similarly, we compare our approximation to near-field global illumination against a path traced ground truth. Figure 8 shows a scene rendered with ambient occlusion only,

in only one direction per pixel (including both sides of a direction, with 12 steps in total) but use the information gathered on a neighborhood of  $4 \times 4$  using a bilateral filter for reconstruction, using uniform convolution weights. To generate per-pixel directions we use a tileable spatial uniform noise of  $4 \times 4$ . In addition, we make aggressive use of temporal coherency by alternating between 6 different rotations and reprojecting the results, using an exponential accumulation buffer. All this gives a total of  $4 \times 4 \times 6 = 96$  effective sampled directions per pixel. Figure 5 shows the effect of the spatial and temporal gathering on the final reconstruction. We opted for a regular sampling approach rather than using line sweeps [Tim13a] because it fitted better our tight budget and target quality (single direction of 12 steps per pixel). Line-sweep ambient occlusion achieves very high quality results, but unfortunately requires a high scan direction count to avoid banding, given the impossibility of randomizing directions per-pixel. Silvennoinen et al. [ST15] reported a cost of 1.6ms for a  $1280 \times 720$  image on the Xbox One, representing a different tradeoff than our approach.

**Height-field assumption considerations** Screen-space techniques assume that the depth map is a height-field, which generally does not hold. As a result, thin features at depth discontinuities cast too much occlusion. While this could be solved with e.g. depth peeling, it is impractical in our case. Instead, we introduce a conservative heuristic derived from the assumption that the thickness of an object is similar to its size in screen space. We introduce this heuristic by modifying the horizon search (Equation (10)): For each iteration  $i$ , the cosine of the maximum horizon angle  $\bar{\theta}^i(\phi) = \cos(\theta^i(\phi))$  is updated using the sample at distance  $r_i$  as

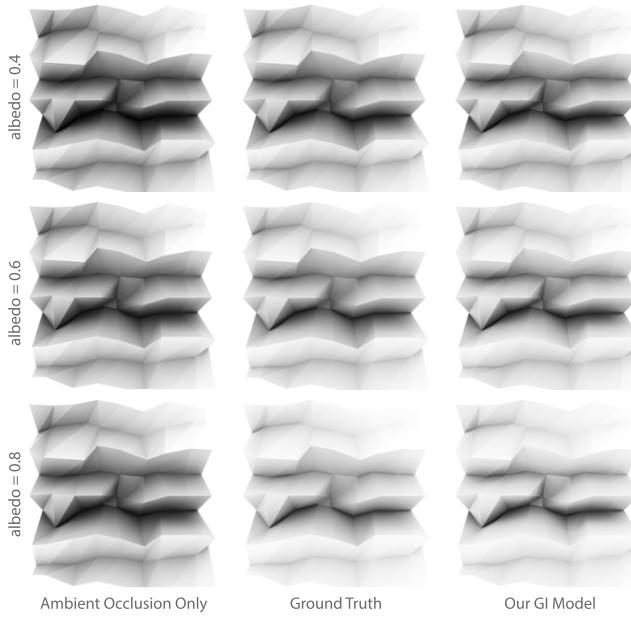
$$\bar{\theta}^i(\phi) = \begin{cases} \langle \omega_s(r_i), \omega_0 \rangle^+ & \text{if } \langle \omega_s(r_i), \omega_0 \rangle^+ \geq \bar{\theta}^{i-1}(\phi) \\ \bar{\theta}^{i-1}(\phi) - \beta & \text{if } \langle \omega_s(r_i), \omega_0 \rangle^+ < \bar{\theta}^{i-1}(\phi) \end{cases} \quad (13)$$

where  $\beta$  is a correction constant, and  $\bar{\theta}^0 = -1$ . Note that the superscript notation is used here to distinguish the iteration process from horizon angles  $\theta_1$  and  $\theta_2$  notation, with the overline being used in this equation to indicate cosine of the horizon angle. A single sample that is behind the horizon will not significantly decrease the computed horizon, but many of them (in e.g. a thin feature) will considerably attenuate it. This allows to progressively attenuate the occlusion on convex features by reducing the horizon angle, while leaving concavities unmodified in e.g. simple corners in indoor settings. For the correction to not affect small convex objects, such as facial features in a human, it is only applied when the sample distance to current maximum horizon is sufficient, and when it is not too far away from the sampling hemisphere base. Figure 6 shows the effect of this heuristic.

Screen-space techniques assume that the depth map is a height-field, which generally does not hold. As a result, thin features at depth discontinuities cast too much occlusion. While this could be solved with e.g. depth peeling, it is impractical in our case. Instead, we introduce a conservative heuristic derived from the assumption that the thickness of an object is similar to its size in screen space. We introduce this heuristic by modifying the horizon search (Equation (10)): For each iteration  $i$ , the cosine of the maximum horizon angle  $\bar{\theta}^i(\phi) = \cos(\theta^i(\phi))$  is updated using the sample at distance  $r_i$  as



**Figure 8:** Adding near-field global illumination to ambient occlusion: From left to right, HBAO [BSD08], our GTAO with ambient occlusion only, GTAO with our global illumination approximation for gray albedo, our GTAO with colored indirect illumination, and path traced Monte Carlo ground truth in a surface with colored albedo. Our approximation model for diffuse interreflections based on ambient occlusion matches very closely the ground truth, and it is able to recover the energy lost by assuming one-bounce illumination only.



**Figure 9:** Effect of albedo in our ambient occlusion-based global illumination approximation, for the groove scene. From left to right: GTAO only, Monte Carlo ground truth, and our approximation based on GTAO, for albedos 0.4, 0.6 and 0.8.

305 and then including global illumination with the analytical model  
306 proposed in Section 4.2, both for gray and colored albedos. Fig-  
307 ure 9 shows the same comparison in an abstract groove-like shape,  
308 with increasing values of gray albedo. In both cases, for a uni-  
309 form distant illumination our technique delivers similar results to  
310 the ground truth, while rendering it in a tight practical real-time  
311 budget.

### 312 5. Directional GTAO

So far we have assumed a uniform infinite light source (i.e. a colored probe). Unfortunately, this approach is too simplifying in practical conditions, specially with the widespread use of light probes for ambient illumination. To account for that, we need to recover the directional component of the light in Equation (1), while still being able to retain real-time performance. Let us approximate Equation (1) for diffuse reflectance and distant illumination as

$$L_r(x, \omega_0) \approx \frac{\rho(x)}{\pi} \int_{\mathcal{H}^2} V(x, \omega_i) L(\omega_i) \langle \mathbf{n}, \omega_i \rangle^+ d\omega_i, \quad (14)$$

with  $L(\omega_i)$  the light incoming from an infinitely far lighting environment (light probe), and  $V(x, \omega_i)$  its visibility. To solve Equation (14), we project the terms of the integral as their spherical harmonics expansion [RH01] as

$$L_r(x, \omega_0) \approx \frac{\rho(x)}{\pi} \int_{\mathcal{H}^2} \left( \sum_j \hat{L}_j y_j(\omega_i) \right) \left( \sum_j \hat{V}'_j y_j(\omega_i) \right) d\omega_i \quad (15)$$

$$= \sum_j \hat{L}_j \hat{V}'_j, \quad (16)$$

313 where  $\hat{L}_j$  and  $\hat{V}'_j$  are the  $j$ -th term of the SH expansion of  $L(\omega_i)$   
314 and  $V'(x, \omega_i)$  respectively, with  $V'(x, \omega_i) = V(x, \omega_i) \langle \mathbf{n}, \omega_i \rangle^+$ , and  
315  $y_j$  is the  $j$ -th spherical harmonics basis function. Assuming that the  
316 visibility  $V(x, \omega_i)$  can be approximated by a cone centered at the

bent normal  $\mathbf{b}$  [Lan02] with aperture angle  $\alpha_v$  defined as a function of the AO term  $\mathcal{A}$ , we can project both the visibility and the dot product in *zonal harmonics* [Slo08], which can be computed efficiently in runtime, and from their expansion compute  $\hat{V}'$ . This allows to compute Equation (14) as a simple dot product between the expansion of  $L(\omega_i)$  and  $V'(x, \omega_i)$ .

**Zonal Harmonics** Zonal harmonics [Slo08] are the projection on spherical harmonics for functions that have rotational symmetry around an axis. They only contain non-zero information for the central coefficients of the expansion (i.e. for  $m = 0$  in  $y_l^m$ ). The key advantage is that they can be efficiently rotated to a new direction  $\omega_i$  at runtime following:

$$f_l^m = \sqrt{\frac{4\pi}{2l+1}} z_l y_l^m(\omega_i), \quad (17)$$

where  $f_l^m$  is the rotated spherical harmonics coefficient,  $z_l$  is the zonal harmonic coefficient of level  $l$ , and  $y_l^m(\omega_i)$  is the spherical harmonics basis for direction  $\omega_i$ .

**Computing the bent normal  $\mathbf{b}$**  We compute  $\mathbf{b}$  using a radiometric formulation weighted by the cosine as

$$\mathbf{b} = \int_{\mathcal{H}^2} \omega_i V(x, \omega_i) \langle \mathbf{n}, \omega_i \rangle^+ d\omega_i. \quad (18)$$

We compute Equation (18) using a similar approach to Equation (6), following the horizon-based approximation as

$$\mathbf{b} = \int_0^\pi \int_{\theta_1(\phi)}^{\theta_2(\phi)} \omega_i(\theta, \phi) \cos(\theta - \gamma)^+ |\sin(\theta)| d\theta d\phi, \quad (19)$$

with  $\omega_i(\theta, \phi)$  the direction defined by the polar coordinates  $(\theta, \phi)$ . We integrate the vertical slices defined by the rotation angle  $\phi$ , while the inner integral over  $\theta$  can be solved analytically for each component of  $\mathbf{b}$  as

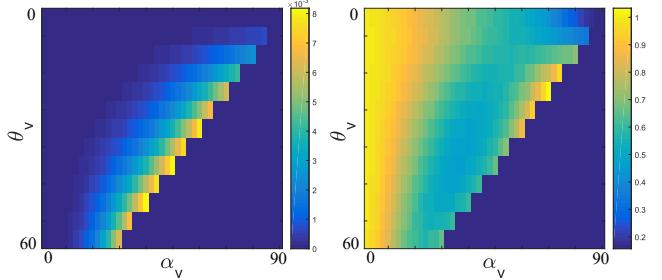
$$\begin{aligned} \mathbf{b}_x &= \cos(\phi) \int_{\theta_1(\phi)}^{\theta_2(\phi)} \sin(\theta) \cos(\theta - \gamma)^+ |\sin(\theta)| d\theta = \cos(\phi) \hat{v}_{xy}(\phi), \\ \mathbf{b}_y &= \sin(\phi) \int_{\theta_1(\phi)}^{\theta_2(\phi)} \sin(\theta) \cos(\theta - \gamma)^+ |\sin(\theta)| d\theta = \sin(\phi) \hat{v}_{xy}(\phi), \\ \mathbf{b}_z &= \int_{\theta_1(\phi)}^{\theta_2(\phi)} \cos(\theta) \cos(\theta - \gamma)^+ |\sin(\theta)| d\theta \\ &= \frac{1}{12} (-\cos(3\theta_1(\phi) - \gamma) - \cos(3\theta_2(\phi) - \gamma) + 8\cos(\gamma) \\ &\quad - 3(\cos(\theta_1(\phi) + \gamma) + \cos(\theta_2(\phi) + \gamma))), \end{aligned} \quad (20)$$

where  $\hat{v}_{xy}(\phi)$  can be analytically solved as

$$\begin{aligned} \hat{v}_{xy}(\phi) &= \frac{1}{12} (6\sin(\theta_1(\phi) - \gamma) - \sin(3\theta_1(\phi) - \gamma) \\ &\quad + 6\sin(\theta_2(\phi) - \gamma) - \sin(3\theta_2(\phi) - \gamma) + 16\sin(\gamma) \\ &\quad - 3(\sin(\theta_1(\phi) + \gamma) + \sin(\theta_2(\phi) + \gamma))). \end{aligned} \quad (21)$$

Note that  $\mathbf{b}$  needs to be normalized after it is computed. We calculate Equation (20) at the same time as the ambient occlusion term  $\mathcal{A}$ , following the approach and implementation described in Section 4, and detailed in Algorithm 2.

**Computing  $V'(x, \omega_i)$**  In order to compute the spherical harmonics expansion of  $V'(x, \omega_i)$  efficiently we leverage the speed of zonal harmonics. We approximate visibility as a visibility cone; thus,



**Figure 10:** Left: Error introduced by computing the cosine term in  $V'(x, \omega_i)$  with respect to the bent normal in zonal harmonics, for a 3 levels SH expansion, for the visibility cone aperture  $\alpha_v$  and the angle between the bent normal and normal of  $\theta_v$ . Right: Ratio of the error introduced by our technique, with respect to the error introduced by the triple product approximation, for a 3 levels SH expansion, where 1 is equal performance, and below 1 means that our technique has less error. Our technique introduces less error than the triple product for most cases, while being more efficient.

both  $V(x, \omega_i)$  and  $\langle \mathbf{n}, \omega_i \rangle^+$  are radially symmetric with respect a particular axis. In particular, the visibility cone approximating  $V(x, \omega_i)$  is symmetric with respect the bent normal [Lan02], and the dot product is symmetric with respect the normal at  $x$ .

We compute the visibility cone in run-time, by computing the bent normal, and an ambient occlusion term. These two terms allow us to compute the visibility cone centered at the bent normal, with an aperture angle derived from the visibility as (see Appendix A for details)

$$\alpha_v(x) = \arccos \left( \sqrt{1 - \mathcal{A}(x)} \right). \quad (22)$$

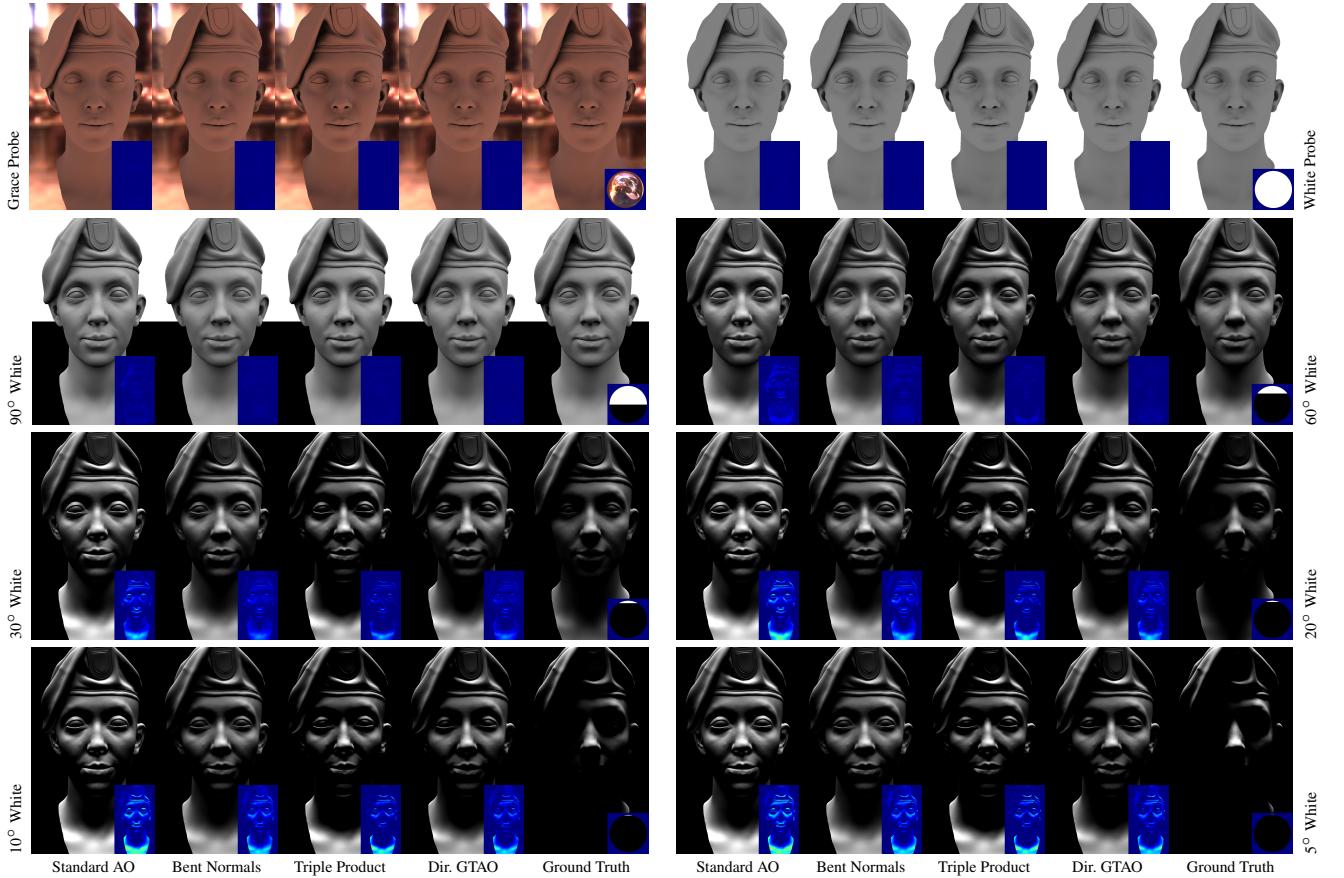
Then, by assuming for efficiency that the dot product is computed with respect to the bent normal instead of the geometric normal, we can compute the zonal harmonics expansion of  $V'(x, \omega_i)$  as

$$\begin{aligned} z_0 &= \frac{\sqrt{\pi}}{2} \sin(\alpha_v(x))^2, \\ z_1 &= \frac{\sqrt{3\pi}}{3} (1 - \cos(\alpha_v(x))^3), \\ z_2 &= \frac{\sqrt{5\pi}}{16} \sin(\alpha_v(x))^2 (2 + 6\cos(\alpha_v(x))^2), \end{aligned} \quad (23)$$

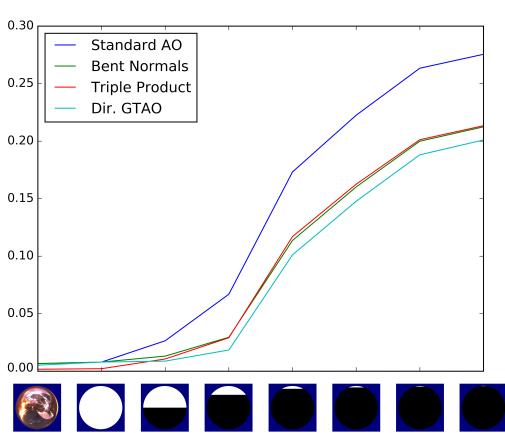
where only the first three coefficients are shown. This formulation introduces error, since we use the bent normal to compute the cosine term, instead of the normal at  $x$ . However, we observed that the divergence between them is not very large (see Figure 10, left). Moreover, computing  $V'(x, \omega_i)$  instead of multiplying the SH projections of  $V(x, \omega_i)$  and the cosine term, gives significant more accurate approximation for a practical low-order SH expansion: The cosine term smooths the step function  $V(x, \omega_i)$ , that is better approximated using SH (see Figure 10, right).

## 5.1. Results

We compare our method for directional occlusion against the standard non-directional AO approximation, the bent normal approx-



**Figure 11:** Error analysis for directional ambient occlusion under illumination with varying frequency, for standard AO (Section 4), standard AO using the bent normal to fetch the environment map [Lan02], the triple product approximation [Sny06], and our work (Dir. GTAO). The insets show the per-pixel error (MSRE for each method is shown in Figure 12).



**Figure 12:** MRSE for the results in Figure 11

imation [Lan02], and the triple product approximation [Sny06]. The former two are computed by fetching a pre-filtered environment map, while the triple product and our directional GTAO use a three-levels SH expansion (SH9) of the probe. Figure 11 shows the results of such comparison for different probes, with increasing frequency. The standard AO approximation quickly fails to capture the directional behaviour of light, while the un-expensive bent normals approximation performs similarly as the triple product; our directional GTAO performs the best in all scenarios both qualitatively and quantitatively (MRSE, see Figure 12).

In terms of cost, we evaluate the performance on a GCN platform, by measuring the final ISA instructions for a dedicated pixel shader running only the occlusion code. The bent normal is given as a pixel shader input, and could come in practice from either our GTAO screen-space approach, or baked offline. Table 1 shows the results. Not surprisingly, the non-directional AO and the bent normal approximation on a prefiltered probe are the cheapest options. Modeling the probe using a SH expansion almost doubles their cost; the cost of our technique is comparable with these simpler techniques, while introducing significantly less error. Note that the actual cost of those techniques is not directly proportional to the

	#ISA	VMMEM	Cycles	VGPR
Standard AO (Probe)	39	1	357	8
Bent Normal (Probe)	39	1	357	8
Standard AO (SH9)	74	0	642	12
Bent Normal (SH9)	74	0	642	12
Triple Product (SH9)	452	0	2053	44
<b>Dir. GTAO (SH9)</b>	<b>90</b>	<b>0</b>	<b>722</b>	<b>12</b>

**Table 1:** Performance comparison for different techniques approximating directional AO on a GCN platform on a dedicated pixel shader. We measure the number of ISA instructions, use of VMMEM, total cycles, and register pressure (VGPR). Our directional GTAO is comparable to simpler ones working with SH.

390 is the full microfacet BRDF at the pixel under uniform light, that  
 391 can be stored in a pre-computed lookup table (typically referred  
 392 to as *environment lut*). The first integral, on the other hand, is  
 393 the convolution of the distant environment light  $L(x, \omega_i)$  with a  
 394 circularly symmetric kernel that approximates the NDF of the  
 395 microfacets. This first integral can be precomputed by convolving  
 396 the distant illumination (e.g. a cubemap) with lobes from different  
 397 surfaces roughness, making it very efficient for rendering glossy  
 398 materials. However, most approximations ignore occlusion or  
 399 approximate it with heuristics.

In order to account for occlusion in specular lighting, we opt for an approach similar to the split-integral approximation in Equation (25). We separate the visibility term from the first integral as a constant, to modulate the amount of illumination reaching  $x$ . This allows us to transform Equation (25) into a product of three integrals, or our *triple-split-integral* approximation:

$$L_r(x, \omega_0) \approx \mathcal{S}(x, \omega_0) \cdot \mathcal{L}(x) \cdot \mathcal{F}(x, \omega_0), \quad (26)$$

where  $\mathcal{S}$  is our *specular occlusion* term modeling visibility. It is computed as

$$\mathcal{S}(x, \omega_0) = \frac{1}{C_V} \int_{\mathcal{H}^2} V(x, \omega_i) f_r(x, \omega_i, \omega_0) \langle \mathbf{n}, \omega_i \rangle^+ d\omega_i, \quad (27)$$

with the normalization term  $C_V = \int_{\mathcal{H}^2} f_r(x, \omega_i, \omega_0) \langle \mathbf{n}, \omega_i \rangle^+ d\omega_i$  ensuring that the specular occlusion  $\mathcal{S}$  ranges into  $[0, 1]$ . Our definition of specular occlusion is weighted by the BRDF, and thus is directionally dependent. This weight was carefully chosen for Equation (26) to match the ground truth for uniform illumination. The normalization factor  $C_V$  is the same as the latter integral  $\mathcal{F}$ , and thus it cancels out if substituting  $\mathcal{S}$  into Equation (26) reducing  $L_r(x, \omega_0)$  to

$$L_r(x, \omega_0) \approx \mathcal{L}(x) \cdot \int_{\mathcal{H}^2} V(x, \omega_i) f_r(x, \omega_i, \omega_0) \langle \mathbf{n}, \omega_i \rangle^+ d\omega_i. \quad (28)$$

In this form, we can observe that for a uniform distant illumination ( $\mathcal{L}(x) = 1$ ) it matches exactly the ground truth expressed by Equation (24).

Figure 13 shows the differences between our approximation in Equation (26) and the raytraced ground truth: For a constant probe, our formulation of specular occlusion models Equation (1) exactly, while for environment probes, it results into a faithful approximation of the rendering equation, specially for specular materials. In the following section, we describe a technique for solving Equation (27) practically for highly demanding applications.

## 7. GTSO: Ground Truth-based Specular Occlusion

Our key idea to compute specular occlusions  $\mathcal{S}(x, \omega_0)$  efficiently is to model an approximation for both the visibility and the BRDF lobes, and then compute the intersection between these two as the specular occlusion. With that in mind, the problem reduces to the question on how representing both the visibility and the BRDF compactly, and on how to compute the intersection between both.

For the **visibility**, we follow the same procedure as in Section 5, and build a visibility cone centered in the bent normal  $\mathbf{b}$  and with amplitude angle derived from the ambient occlusion term  $\mathcal{A}(x)$  using Equation (22). Similarly, we can model the **specular** lobe as a

370 cycles they consume on isolation and will depend on which shader  
 371 (and where) they are located, as often cost can be hidden by other  
 372 operations. It is interesting to observe that the bent normal approx-  
 373 imation results in low error; this suggests that when using a pre-  
 374 filtered environment map it could be the technique of choice. How-  
 375 ever, in cases when using SH9 to encode the light probe, our tech-  
 376 nique introduces minimal overhead over simpler techniques, while  
 377 reducing significantly the error.

378

### 6. Specular Occlusion

379 Here we generalize classic Lambertian-based ambient occlusion,  
 380 by proposing its specular counterpart. We develop an illumination  
 381 model where the near-field occlusion modulates the distant lighting  
 382 while supporting arbitrary BRDF models (e.g. microfacets). More-  
 383 over, for the specific cases of uniform dome illumination, our model  
 384 delivers ground truth results.

385

Let us assume that all light incides from an infinitely far lighting environment (light probe) to express Equation (1) as

$$L_r(x, \omega_0) = \int_{\mathcal{H}^2} V(x, \omega_i) L(x, \omega_i) f_r(x, \omega_i, \omega_0) \langle \mathbf{n}, \omega_i \rangle^+ d\omega_i. \quad (24)$$

Computing this integral by numerical integration is too expensive for real-time applications. Equation (24) is the generalization of Equation (14) to arbitrary BRDFs, and could be computed following a similar procedure as in Section 5; however, in order to support all-frequency BRDFs a large number of coefficients in the SH expansion would be required, reducing significantly the performance. The current state-of-the-art assumes uniform perfect visibility ( $\forall \omega_i | V(x, \omega_i) = 1$ ) and uses a split-integral approximation [Laz13, Kar13] as

$$\begin{aligned} L_r(x, \omega_0) &\approx \mathcal{L}(x) \cdot \mathcal{F}(x, \omega_0), \\ \mathcal{L}(x) &= \frac{1}{C_L} \int_{\mathcal{H}^2} \overbrace{V(x, \omega_i)}^{=1} L(x, \omega_i) D(x, \omega_h) \langle \mathbf{n}, \omega_i \rangle^+ d\omega_i, \\ \mathcal{F}(x, \omega_0) &= \int_{\mathcal{H}^2} f_r(x, \omega_i, \omega_0) \langle \mathbf{n}, \omega_i \rangle^+ d\omega_i, \end{aligned} \quad (25)$$

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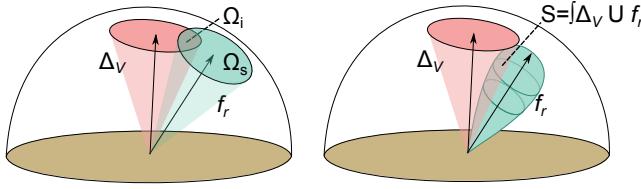
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**Figure 13:** Comparison between ground truth specular illumination and our specular occlusion model under two different illumination setups, for increasing roughness of the GGX microfacet BRDF. From top to bottom: Our specular occlusion (Equation (26)) under environment lighting, ground truth result under the same environment light, our specular occlusion with constant illumination, and ground the rendering equation under the same white probe. For constant illumination, our specular occlusion model exactly models the rendering equation.



**Figure 14:** Geometry of our specular occlusion, assuming that both the visibility and the specular are modeled as cones (left) and with accurate specular lobe (right).

express  $\mathcal{S}$  as a four dimensional function:

$$\mathcal{S}(\alpha_v, \beta, r, \theta_o) \approx \frac{1}{C_V} \int_{\mathcal{H}^2} \Delta_V(\alpha_v, \beta) f_r(\omega_i, \theta_o, r) \langle \mathbf{n}, \omega_i \rangle^+ d\omega_i. \quad (30)$$

This function can be compactly baked as a four-dimensional table if assuming a reflectance at normal incidence of 0.04. Given that the function is relatively smooth, we can encode it to a four-dimensional 32<sup>4</sup> BC4 8-bit look up table, which can be efficiently accessed in runtime. While not explored in this work, this lookup table for  $\mathcal{S}(x, \omega_o)$  could be merged with the lookup table often used for  $\mathcal{F}(x, \omega_o)$ .

## 7.1. Results

Figure 15 compares the results of our GTSO implementation using a 4D look-up table for computing the specular occlusion, compared against the ground truth and the empirically-based technique described in Lagarde [Ld14]. For any roughness parameter of the microfacet BRDF, the introduced error is minimal.

## 8. Conclusions

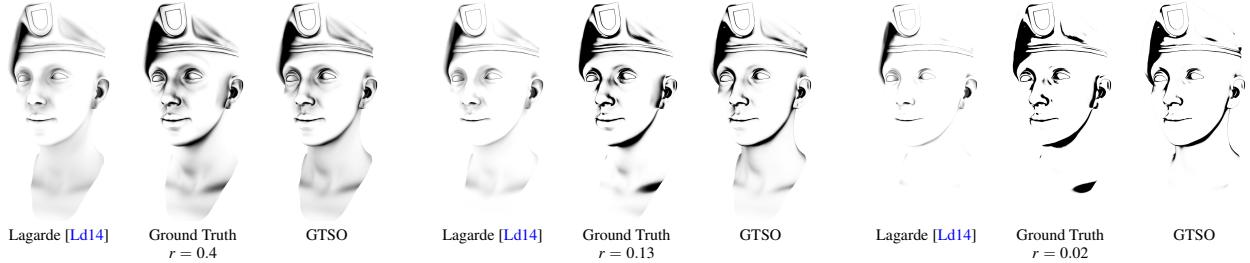
In this work we have presented several contribution to real-time ambient occlusion. In the first place, we have presented GTAO: an efficient formulation of ambient occlusion that matches the Monte Carlo ground truth within a very tight budget. We implement our technique efficiently, by aggressively making use of both spatial and temporal coherence to effectively integrate almost 100 samples per pixel while computing only one each frame. GTAO goes together with a simple but effective technique that simulates near-field diffuse inter-reflections based on the ambient occlusion at the shading point. The technique bases on the observation that these inter-reflections can be modeled as a function of the local albedo and the ambient occlusion. Then, we have generalized our GTAO to

cone centered on the reflection direction  $\omega_r$ , and compute  $\mathcal{S}(x, \omega_o)$  as the intersection of the visibility and BRDF cones (see Figure 14, left, and Appendix C for more details on this approach).

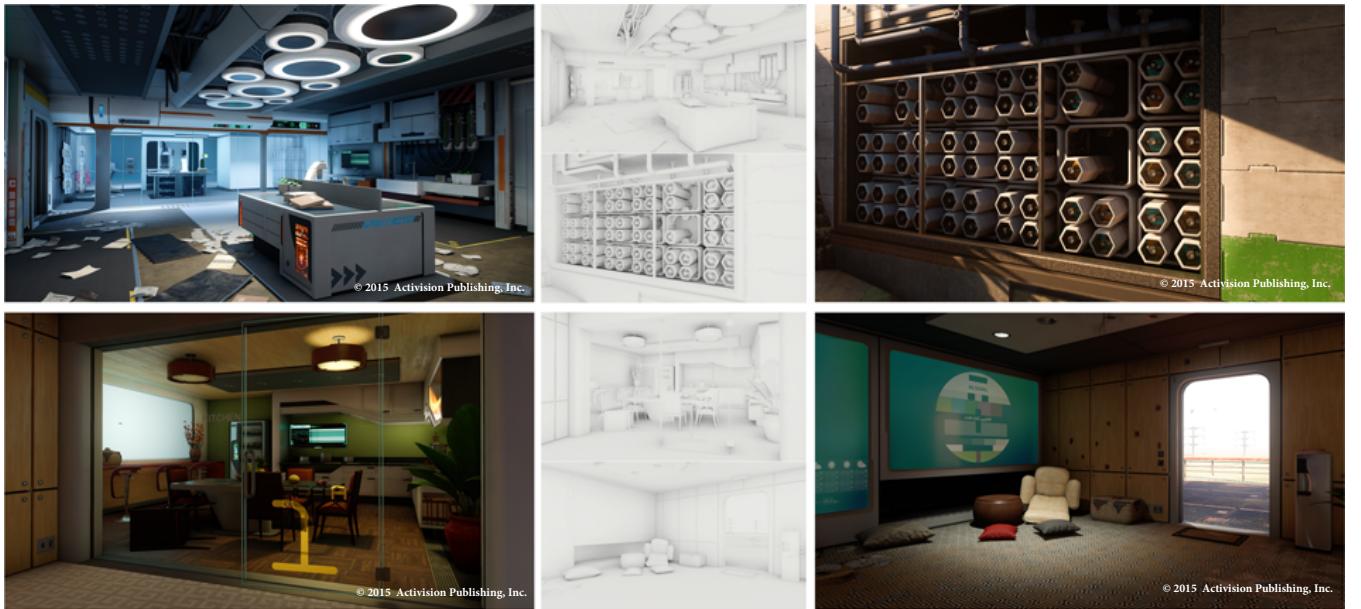
Unfortunately, in a tight-bounded real-time application, these computations are still expensive. Furthermore, we have found specular lobes to be poorly represented by cones. To improve on both qualities, we opt for a more accurate approximation by precomputing the specular occlusion  $\mathcal{S}$  as the product of the visibility cone  $\Delta_V$  and the actual BRDF  $F$  (Figure 14, right):

$$\mathcal{S}(x, \omega_o) \approx \frac{1}{C_V} \int_{\mathcal{H}^2} \Delta_V(\alpha_v(x), \beta(\mathbf{b}(x), \omega_r)) f_r(x, \omega_i, \omega_o) \langle \mathbf{n}, \omega_i \rangle^+ d\omega_i, \quad (29)$$

with  $\beta = \arccos(\langle \mathbf{b}, \omega_r \rangle)$  the angle between the bent normal and the reflection vector  $\omega_r$ , and  $\Delta_V(\alpha, \beta)$  is a binary function returning 1 if  $\beta \leq \alpha$  and 0 elsewhere. Assuming a isotropic microfacet-based BRDF with a GGX NDF [WMLT07] parametrized by the roughness  $r$ , we model the reflected direction  $\omega_r$  as a single angle  $\theta_o = \arccos(\langle \mathbf{n}, \omega_r \rangle)$  with respect to the normal  $\mathbf{n}$ . With these assumptions, and omitting the spatial dependence for clarity, we can



**Figure 15:** Comparison of Lagarde specular occlusion [Ld14], the Monte Carlo ground truth and our GTSO using a four-dimensional look-up table, for a GGX microfacet BRDF with roughness  $r = 0.4, 0.13$  and  $0.02$ . The ground truth case shows  $\mathcal{S}(x, \omega_0)$  for traced visibility  $V(x, \omega_i)$ , whereas GTSO shows it for our cone-based, lookup table approximation.



**Figure 16:** Screenshots of our GTao being used in-game for accurate and efficient ambient occlusion, in scenes with high-quality physically-based shading and high geometric complexity. Our GTao computes the ambient occlusion layer (in the insets) in just  $0.5$  ms for PS4.

450 a directional spherical harmonics-based generalization, that lever-  
 451 ages zonal harmonics and efficient on-line computation of SH-  
 452 based Lambertian occlusion. Finally, we have introduced an ap-  
 453 proximation of specular occlusion with our *Ground-Truth Specu-*  
 454 *lar Occlusion*, which generalizes the ambient occlusion operator to  
 455 deal with specular surfaces, and introduced an efficient technique  
 456 based on a precomputed look-up table to efficiently compute the  
 457 specular reflection from uniform and non-uniform probe-based il-  
 458 lumination.

459 As shown in Figure 1 combining all our techniques results into a  
 460 complete solution for efficient probe-based illumination, allowing  
 461 to match the raytraced ground truth. The near-field indirect illumi-  
 462 nation, directional GTao, and GTSO base on the results obtained  
 463 using our efficient implementation of ambient occlusion (GTao),  
 464 resulting in very optimized techniques targeting very tight time  
 465 budgets, like videogames, even for current console platforms (Fig-  
 466 ure 16).

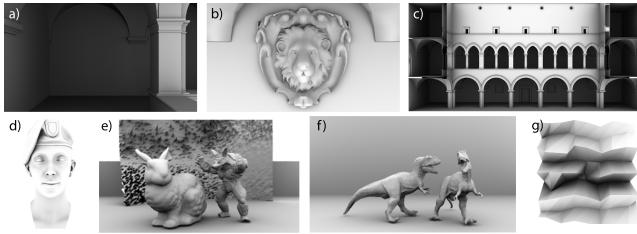
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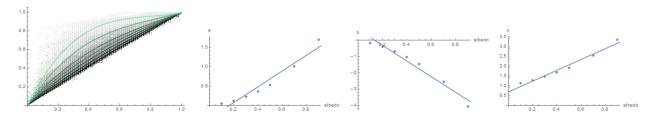
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**Figure 17:** Input scenes used for computing the mapping between the ambient occlusion and the near-field global illumination, rendered using only ambient occlusion.



**Figure 18:** Cubic fit for our mapping between the ambient occlusion and the three-bounce global illumination for different albedos (left). We observed that a linear fit between the coefficients of the polynomial wrt the albedo gives a good continuous fit, as shown in the three rightmost figures. The combination of these fits give form to our model (Equation (33)).

## 572 Appendix A: Aperture Calculation

The visibility equation for a cone can be calculated as follows:

$$\mathcal{A}(x) = \frac{1}{\pi} \int_0^{2\pi} \left( \int_0^{\alpha_v(x)} \cos(\theta) \sin(\theta) d\theta \right) d\phi = 1 - \cos(\alpha_v(x))^2 \quad (31)$$

where solving for  $\alpha_v(x)$  yields the equation to convert from occlusion to aperture angles:

$$\alpha_v(x) = \arccos(\sqrt{1 - \mathcal{A}(x)}) \quad (32)$$

## 573 Appendix B: Polynomial Fitting of Global Illumination

Based on the observation that there is a relationship between ambient occlusion and global illumination exists (Figure 2), and assuming that the albedo  $\rho(s)$  at all points  $s$  around  $x$  is  $\rho(s) = \rho(x)$ , we want to design a mapping between the albedo and ambient occlusion at  $x$  and the reflected global illumination at  $x$ . To build this function  $\mathcal{G}(\mathcal{A}(x), \rho(x))$  we compute seven simulations with different albedos ( $\rho = [0.1, 0.2, 0.3, 0.4, 0.5, 0.7, 0.9]$ ) in a set of scenes showing a variety of different types of occlusion conditions (see Figure 17). We compute both the ambient occlusion and multi-bounce indirect illumination (in our case, up to three bounces). By taking the combination of all points, we fit this mapping using a cubic polynomial for each albedo (Figure 18 (left)), generating a set of polynomial coefficient for each scene albedo. We then observed that said coefficients were well approximated by a linear fit as a function of the input albedo (Figure 18). This last observation allows us to build a bidimensional mapping between the albedo  $\rho$  and ambient occlusion  $\mathcal{A}$ :

$$\begin{aligned} \mathcal{G}(\mathcal{A}, \rho) &= a(\rho) \mathcal{A}^3 - b(\rho) \mathcal{A}^2 + c(\rho) \mathcal{A}, \\ a(\rho) &= 2.0404 \rho - 0.3324, \\ b(\rho) &= 4.7951 \rho - 0.6417, \\ c(\rho) &= 2.7552 \rho + 0.6903. \end{aligned} \quad (33)$$

574

## 575 Appendix C: Analytical Cone-to-Cone Specular Occlusion

Specular occlusion can be computed as the ratio between the intersection of the visibility and specular cones  $\Omega_i$ , and the specular cone  $\Omega_s$  (see Figure 14, left):

$$\mathcal{S}(x, \omega_0) = \frac{\Omega_i(x, \omega_0)}{\Omega_s(x, \omega_0)}, \quad (34)$$

We then need to compute the visibility and specular cones, defined by a direction and an aperture, and their intersection solid angle  $\Omega_i$ . To leverage previous work on mappings from specular lobes to cones, Phong is used instead of GGX on the experiments described in this section.

The visibility cone is explained in Section 6 (Equation (22)). In the case of the specular cone, its direction is defined by the reflection vector  $\omega_r$ . Its aperture  $\alpha_s$ , on the other hand, is defined by the roughness  $r$  (or specular power  $p$  in the case of a Phong BRDF). Since there are no exact solution for this, we opt of an approach similar to the one by Uludag [Ulu14], which uses the Phong importance sampling routine by Walter et al. [WMLT07] to relate the aperture with the Phong power  $p$ :

$$\alpha_s = \arccos\left(u^{\frac{1}{p+2}}\right), \quad (35)$$

where  $u$  is a constant. As opposed to Uludag, we do not obtain  $u$  by fitting the cone to lobes ( $u = 0.244$ ), but minimize differences between resulting GTSO and Monte Carlo ground truth references, getting  $u = 0.01$ .

Once we have both cones, we can compute their intersection solid angle  $\Omega_i$ . This intersection has analytical solution [OS07, Maz12], as a function of the cone apertures and the angle between their respective directions, the bent normal  $\mathbf{b}$  and the reflection direction  $\omega_r$ .

---

**Algorithm 1** Computes the ambient occlusion term  $\mathcal{A}(x)$ .

---

```

1: cPosV ← VIEWSPACEPOSFROMDEPTHBUFFER(cTexCoord)                                ▷ We will abbreviate center with c
2: viewV ← NORMALIZE(-cPosV)
3: visibility ← 0
4: for slice ∈ [0, sliceCount) do
5:   φ ← ( $\pi$ /sliceCount) * slice
6:   ω ← {cos φ, sin φ}
7:
8:   directionV ← {ω[0], ω[1], 0}
9:   orthoDirectionV ← directionV - DOT(directionV, viewV) * viewV
10:  axisV ← CROSS(directionV, viewV)
11:  projNormalV ← normalV - axisV * DOT(normalV, axisV)
12:
13:  sgnN ← SIGN(DOT(orthoDirectionV, projNormalV))
14:  cosN ← SATURATE(DOT(projNormalV, viewV)/LEN(projNormalV))
15:  n ← sgnN * arccos(cosN)
16:
17:  for side ∈ [0, 1] do                                                 ▷ Equation (13)
18:    cHorizonCos ← -1
19:    for sample ∈ [0, directionSampleCount) do
20:      s ← sample/directionSampleCount
21:      sTexCoord ← cTexCoord + (-1 + 2 * side) * s * scaling * {ω[0], -ω[1]}          ▷ Flip y due to texture coordinate system
22:      sPosV ← VIEWSPACEPOSFROMDEPTHBUFFER(sTexCoord)
23:      sHorizonV ← NORMALIZE(sPosV - cPosV)
24:      cHorizonCos ← MAX(cHorizonCos, DOT(sHorizonV, viewV))
25:    end for
26:
27:    h[side] ← n + CLAMP((-1 + 2 * side) * arccos(cHorizonCos) - n, - $\pi$ /2,  $\pi$ /2)           ▷ Horizon angle θ
28:    visibility ← visibility + LEN(projNormalV) * (cosN + 2 * h[side] * sin(n) - cos(2 * h[side] - n))/4  ▷ Equation (7)
29:  end for
30: end for
31: visibility ← visibility/sliceCount

```

---

**Algorithm 2** Extension that computes bent normals **b**. Repeated code from ambient occlusion algorithm is omitted.

---

```

1: ...
2: for slice ∈ [0, sliceCount) do
3:   ...
4:   t[0] ← (6 * sin(h[0] - n) - sin(3 * h[0] - n) + 6 * sin(h[1] - n) - sin(3 * h[1] - n) + 16 * sin(n) - 3 * (sin(h[0] + n) + sin(h[1] + n)))/12  ▷ Equations (20) and (21)
5:   t[1] ← (-cos(3 * h[0] - n) - cos(3 * h[1] - n) + 8 * cos(n) - 3 * (cos(h[0] + n) + cos(h[1] + n)))/12
6:   bentNormalL ← {ω[0] * t[0], ω[1] * t[0], -t[1]}                                         ▷ Flip z due to change of handedness
7:   bentNormalV ← bentNormalV + MULT(bentNormalL, ROTFROMTOMATRIX({0, 0, -1}, viewV)) * LEN(projNormalV) ▷ [MH99]
8: end for
9: bentNormalV ← NORMALIZE(bentNormalV)

```

---