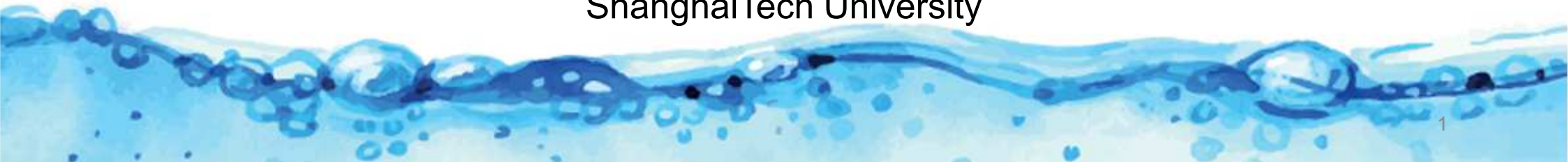


Computer Animation & Physical Simulation

Lecture 7: Rigid-Body Simulation II

XIAOPEI LIU

School of Information Science and Technology
ShanghaiTech University



I. Review of Basic Rigid Body Simulation



Non-constrained Rigid Body Simulation

- Time rate change of the state variable

$$\frac{d}{dt}\mathbf{Y}(t) = \frac{d}{dt} \begin{pmatrix} x(t) \\ R(t) \\ P(t) \\ L(t) \end{pmatrix} = \begin{pmatrix} v(t) \\ \omega(t)^* R(t) \\ F(t) \\ \tau(t) \end{pmatrix}$$

- Computing order

$$F(t) \longrightarrow P(t)$$

$$\tau(t) \longrightarrow L(t)$$

$$P(t) \xrightarrow{P(t) = Mv(t)} v(t) \longrightarrow x(t)$$

$$L(t) \xrightarrow{L(t) = I(t)\omega(t)} \omega(t) \longrightarrow R(t)$$

Non-constrained Rigid Body Simulation

- **Runge–Kutta methods** $y' = f(x, y)$

- Achieve higher accuracy
- Re-evaluate $f(\cdot, \cdot)$ at points intermediate between $(x_n, y(x_n))$ and $(x_{n+1}, y(x_{n+1}))$

$$y_{n+1} = y_n + h\Phi(x_n, y_n; h) ,$$

$$\Phi(x, y; h) = \sum_{r=1}^R c_r k_r ,$$

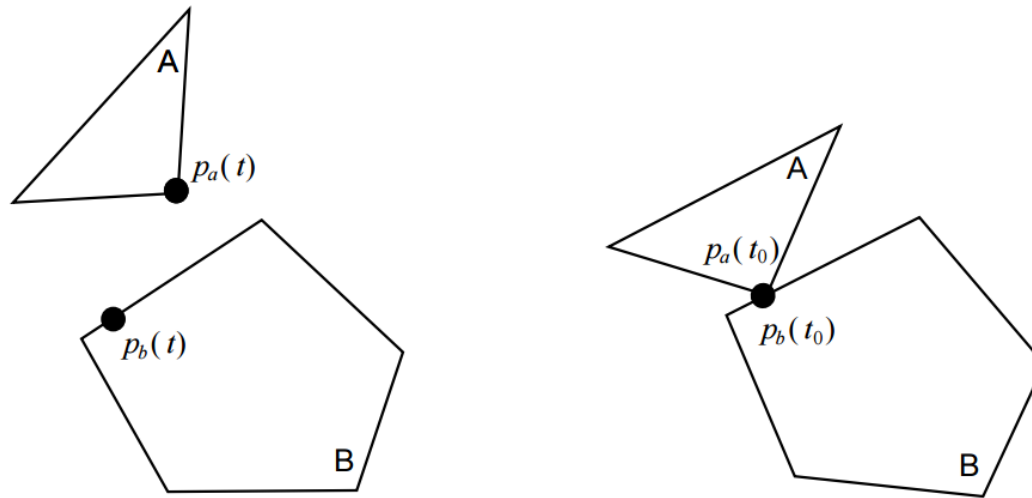
$$k_1 = f(x, y) ,$$

$$k_r = f\left(x + ha_r, y + h \sum_{s=1}^{r-1} b_{rs} k_s\right) , \quad r = 2, \dots, R ,$$

$$a_r = \sum_{s=1}^{r-1} b_{rs} , \quad r = 2, \dots, R .$$

Constrained Rigid Body Simulation

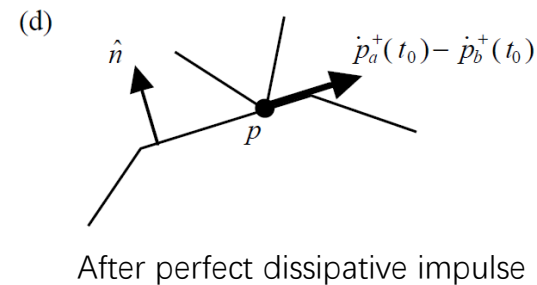
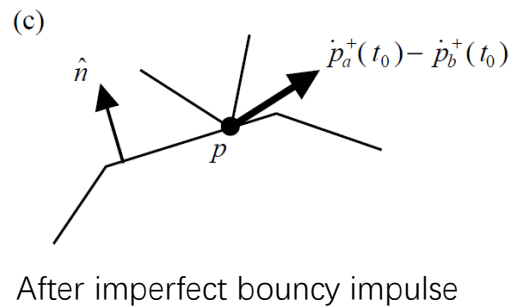
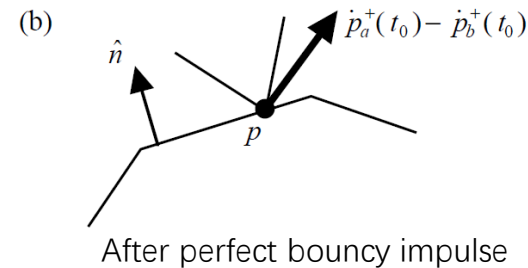
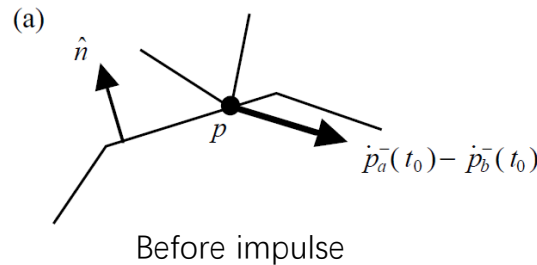
- **Colliding contact**



Constrained Rigid Body Simulation

- **Physical meaning for coefficient of restitution**

- Illustration



Constrained Rigid Body Simulation

- **Colliding contact**

- How to compute the impulse?
- For frictionless bodies, the direction of the impulse will be in the normal direction

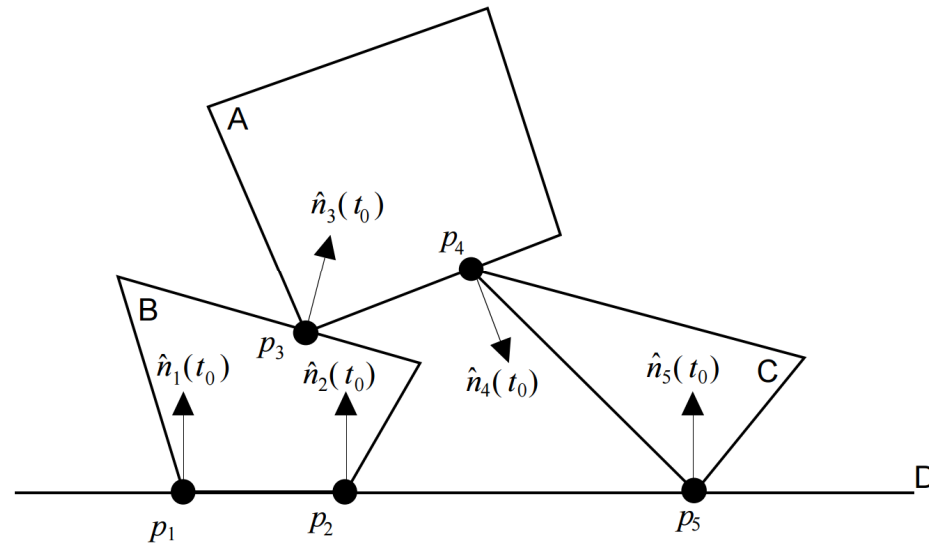
$$J = j\hat{n}(t_0)$$

$$j = \frac{-(1 + \epsilon)v_{rel}^-}{\frac{1}{M_a} + \frac{1}{M_b} + \hat{n}(t_0) \cdot \left(I_a^{-1}(t_0) (r_a \times \hat{n}(t_0)) \right) \times r_a + \hat{n}(t_0) \cdot \left(I_b^{-1}(t_0) (r_b \times \hat{n}(t_0)) \right) \times r_b}$$

$$v_{rel}^+ = -\epsilon v_{rel}^- \quad 0 \leq \epsilon \leq 1$$

Constrained Rigid Body Simulation

- **Resting contact**
 - Computing contact forces

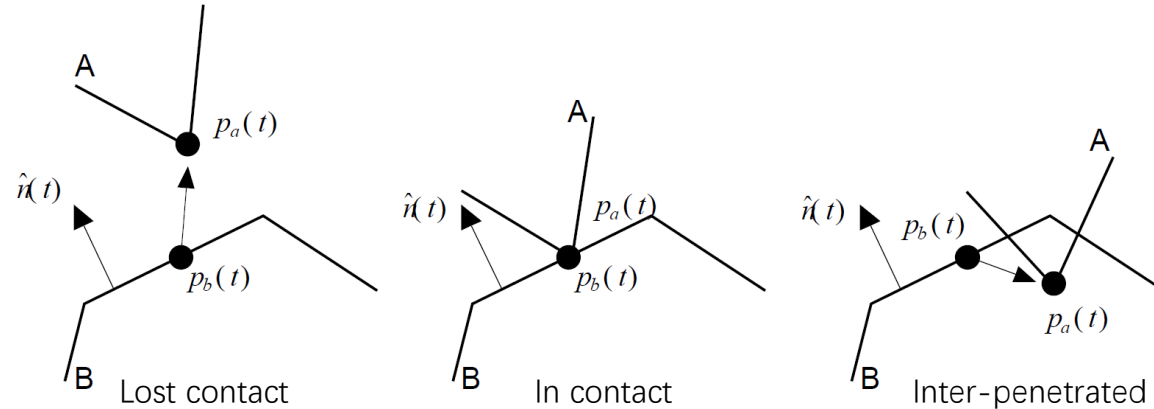


Constrained Rigid Body Simulation

- **Resting contact**

- Preventing inter-penetration
 - Construction of separation distance

$$d_i(t) = \hat{n}_i(t) \cdot (p_a(t) - p_b(t))$$



Constrained Rigid Body Simulation

- **Resting contact**

- Preventing inter-penetration

- Consider $d_i(t_0) = 0$

- We have to keep the two bodies from accelerating towards each other

- Taking the second derivative of $d_i(t_0)$

$$\ddot{d}(t_0) = \hat{n}_i(t_0) \cdot (\ddot{p}_a(t_0) - \ddot{p}_b(t_0)) + 2\dot{\hat{n}}_i(t_0) \cdot (\dot{p}_a(t_0) - \dot{p}_b(t_0))$$

- $\ddot{d}_i(t_0) > 0$: contact will break immediately
 - $\ddot{d}_i(t_0) = 0$: contact remains
 - $\ddot{d}_i(t_0) < 0$: must be avoided

II. Mathematical Formulation of Constraints



Resting Contact

- **Expression for three conditions**

- Non-interpenetration

$$\ddot{d}_i(t_0) \geq 0$$

- Repulsiveness

$$f_i \geq 0$$

- Contact breaking

$$f_i \ddot{d}_i(t_0) = 0$$



Resting Contact

- **Computing contact force**

- Force contribution
 - Consider the expression

$$\ddot{d}(t_0) = \hat{n}_i(t_0) \cdot (\ddot{p}_a(t_0) - \ddot{p}_b(t_0)) + 2\dot{\hat{n}}_i(t_0) \cdot (\dot{p}_a(t_0) - \dot{p}_b(t_0))$$

$$\ddot{p}_a(t) = \dot{v}_a(t) + \dot{\omega}_a(t) \times r_a(t) + \omega_a(t) \times (\omega_a(t) \times r_a(t))$$

- Forces contribute to
 - Linear acceleration $\dot{v}_a(t)$
 - Angular acceleration $\dot{\omega}_a(t) = I_a^{-1}(t)\tau_a(t) + I_a^{-1}(t)(L_a(t) \times \omega_a(t))$

Resting Contact

- **Computing contact force**

- Express separating distance acceleration in terms of all associated forces

$$\ddot{d}_i(t_0) = a_{i1}f_1 + a_{i2}f_2 + \cdots + a_{in}f_n + b_i$$

- Write for all contact points

$$\begin{pmatrix} \ddot{d}_1(t_0) \\ \vdots \\ \ddot{d}_n(t_0) \end{pmatrix} = \mathbf{A} \begin{pmatrix} f_1 \\ \vdots \\ f_n \end{pmatrix} + \begin{pmatrix} b_1 \\ \vdots \\ b_n \end{pmatrix}$$

Resting Contact

- **Computing contact force**

- Contribution of a_{ij}
 - Consider linear and angular acceleration

$$\frac{\hat{n}_j(t_0)}{m_a} + \left(I_a^{-1}(t_0) \left((p_j - x_a(t_0)) \times \hat{n}_j(t_0) \right) \right) \times r_a$$

- Contribution of b_i
 - Collect the force independent part

$$\frac{2\dot{\hat{n}}_i(t_0) \cdot (\dot{p}_a(t_0) - \dot{p}_b(t_0))}{m_a} + \left(I_a^{-1}(t_0) \tau_a(t_0) \right) \times r_a + \omega_a(t_0) \times (\omega_a(t_0) \times r_a) + \left(I_a^{-1}(t_0) (\mathbb{L}_a(t_0) \times \omega_a(t_0)) \right) \times r_a$$

Contact Formulation

- **Mathematical Formulation**

$$\mathbf{a} = \mathbf{A}\mathbf{f} + \mathbf{b} \quad a_i \geq 0, \quad f_i \geq 0 \quad \text{and} \quad f_i a_i = 0$$

$$\sum_{i=1}^n f_i a_i = \mathbf{f}^T \mathbf{a} = 0$$

$$\mathbf{A}\mathbf{f} + \mathbf{b} \geq \mathbf{0}, \quad \mathbf{f} \geq \mathbf{0} \quad \text{and} \quad \mathbf{f}^T (\mathbf{A}\mathbf{f} + \mathbf{b}) = 0.$$

Contact Formulation

- **Mathematical Formulation**
 - Linear complementarity problem

$$\mathbf{A}\mathbf{f} + \mathbf{b} \geq \mathbf{0}, \quad \mathbf{f} \geq \mathbf{0} \quad \text{and} \quad \mathbf{f}^T(\mathbf{A}\mathbf{f} + \mathbf{b}) = 0.$$

$$\min_{\mathbf{f}} \mathbf{f}^T(\mathbf{A}\mathbf{f} + \mathbf{b}) \quad \text{subject to} \quad \left\{ \begin{array}{l} \mathbf{A}\mathbf{f} + \mathbf{b} \geq \mathbf{0} \\ \mathbf{f} \geq \mathbf{0} \end{array} \right\}$$

Solution of Linear Complementarity Problem

- **Minimum Map Newton Method**

$$y = ax + b,$$

$$y \geq 0, x \geq 0, \text{ and } xy = 0$$



$$h(x, y) = \min(x, y)$$

$$\min(x, y) = \begin{cases} x & \text{if } x < y \\ y & \text{otherwise} \end{cases}$$

$$h(x^*, y^*) = 0 \quad \text{iff} \quad 0 \leq y^* \perp x^* \geq 0$$

$$\mathbf{H}(\mathbf{x}) = \mathbf{H}(\mathbf{x}, \mathbf{y}) = \begin{bmatrix} h(\mathbf{x}_1, \mathbf{y}_1) \\ \dots \\ h(\mathbf{x}_n, \mathbf{y}_n) \end{bmatrix} = \mathbf{0}$$

Solution of Linear Complementarity Problem

- **Fischer–Newton Method**

$$y = ax + b,$$

$$y \geq 0, x \geq 0, \text{ and } xy = 0$$



$$\phi(x, y) \equiv \sqrt{x^2 + y^2} - x - y$$

$$\phi(x^*, y^*) = 0 \quad \text{iff} \quad 0 \leq y^* \perp x^* \geq 0$$

$$\mathbf{F}(\mathbf{x}) = \mathbf{F}(\mathbf{x}, \mathbf{y}) = \begin{bmatrix} \phi(\mathbf{x}_1, \mathbf{y}_1) \\ \vdots \\ \phi(\mathbf{x}_n, \mathbf{y}_n) \end{bmatrix} = \mathbf{0}$$

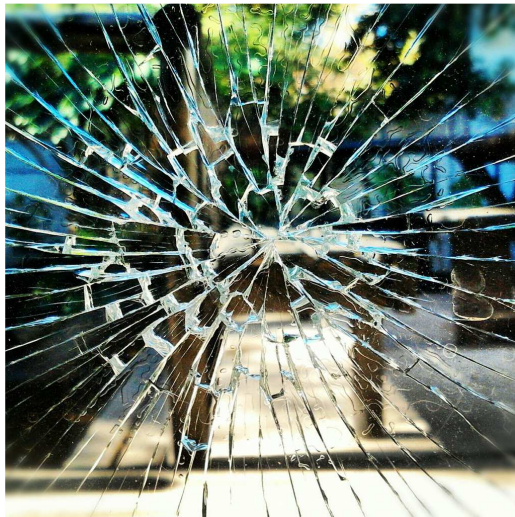
IV. Rigid Body Fracture



What is a fracture?

- **A fracture**

- The separation of an object or material into two or more pieces under the action of stress



How to model fracture?

- **Consider material deformations**

- Even rigid body objects have small deformations
- Deformation causes change of internal stress
- Fracture arises when internal stress exceed the material toughness (strength)

- **Modeling**

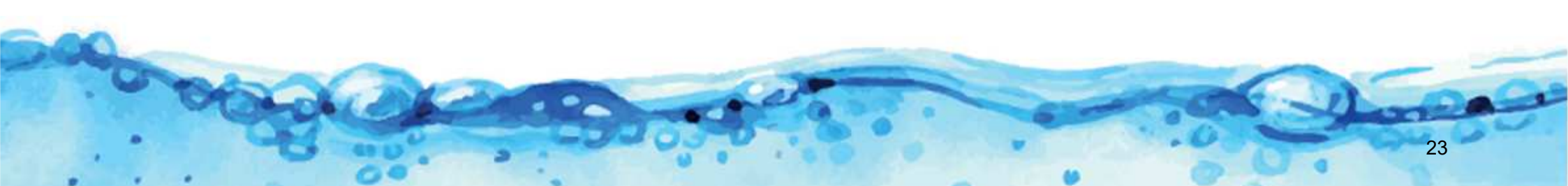
- Computation of internal stress distribution
- Determine the fracture point and fracture geometry



Continuum Mechanics

- **A branch of mechanics**

- Modeled as a continuous mass rather than as discrete particles
- The matter in the body is continuously distributed
- A continuum is a body that can be continually sub-divided into infinitesimal elements
 - Derivatives are available to compute
- Deal with deformable bodies
 - As opposed to ideal rigid bodies
 - Analyzing internal force of rigid bodies should consider deformation (very small)



Material Coordinates

- **Eulerian specification of a field**

- Represented as a function of position \mathbf{x} and time t

$$\mathbf{u}(\mathbf{x}, t)$$

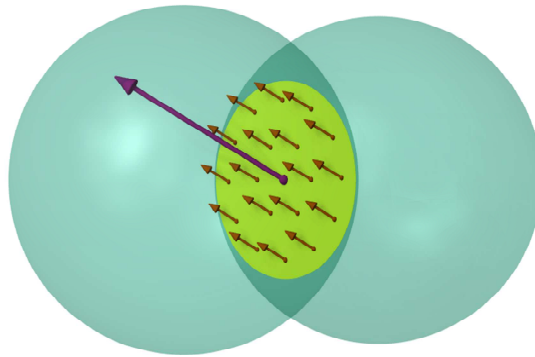
- **Lagrangian (material) specification**

- Particles are followed through time
- Particles are labeled by some (time-independent) vector field \mathbf{x}_0 (material coordinates)

$$\mathbf{u}(\mathbf{X}(\mathbf{x}_0, t), t) = \frac{\partial \mathbf{X}}{\partial t}(\mathbf{x}_0, t)$$

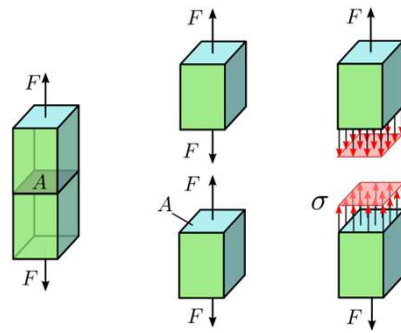
Stress of a Material

- **A physical quantity of material**
 - Internal forces that neighboring particles exert on each other
 - Defined as the force across a "small" boundary per unit area of that boundary

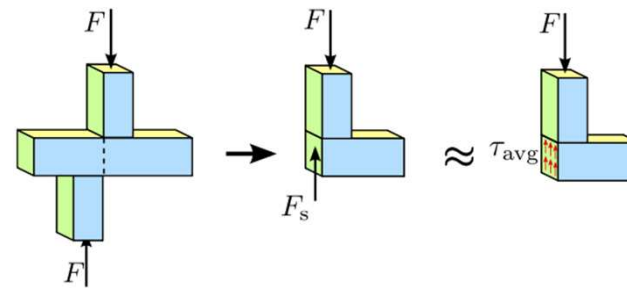


Stress of a Material

- **Stress may be regarded as the sum of two components**
 - Normal stress
 - The stress component perpendicular to the surface (compression or tension)
 - Shear stress
 - The stress component parallel to the surface



Normal stress

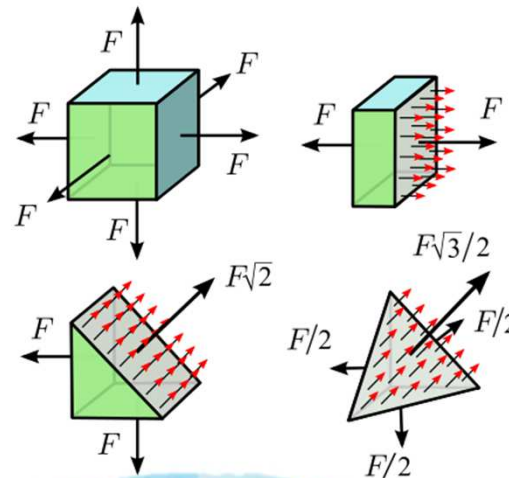


Shear stress

Stress of a Material

- **Isotropic stress**

- A simple type of stress
- Occur when the material body is under equal compression or tension in all directions
- The material is homogeneous



Cauchy Stress Tensor

- **General stress**

- Mechanical bodies experience more than one type of stresses at the same time (combined stress)
- Combined stresses cannot be described by a single vector

- **Cauchy's observation**

- The stress vector across a surface will always be a linear function of the surface's normal

$$T = \sigma(n)$$

$$\sigma(\alpha u + \beta v) = \alpha \sigma(u) + \beta \sigma(v)$$

Cauchy Stress Tensor

- **Definition**

- A 3x3 matrix

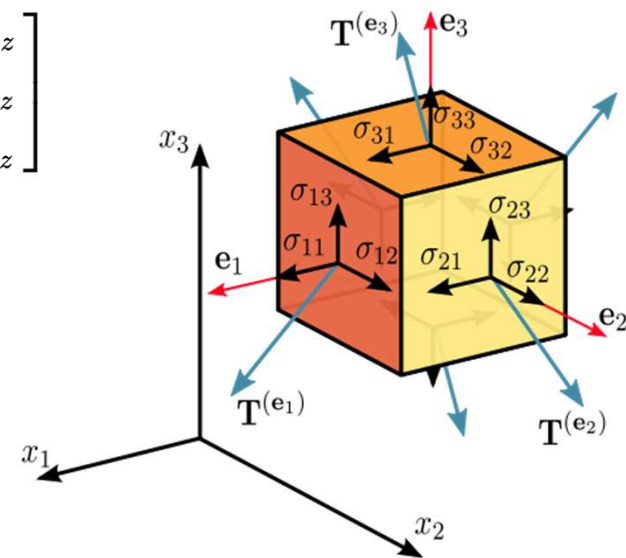
$$\begin{bmatrix} \sigma_{11} & \sigma_{12} & \sigma_{13} \\ \sigma_{21} & \sigma_{22} & \sigma_{23} \\ \sigma_{31} & \sigma_{32} & \sigma_{33} \end{bmatrix}$$

or

$$\begin{bmatrix} \sigma_{xx} & \sigma_{xy} & \sigma_{xz} \\ \sigma_{yx} & \sigma_{yy} & \sigma_{yz} \\ \sigma_{zx} & \sigma_{zy} & \sigma_{zz} \end{bmatrix}$$

$$\mathbf{T} = \mathbf{n} \cdot \boldsymbol{\sigma}$$

$$\begin{bmatrix} T_1 \\ T_2 \\ T_3 \end{bmatrix} = \begin{bmatrix} \sigma_{11} & \sigma_{12} & \sigma_{13} \\ \sigma_{21} & \sigma_{22} & \sigma_{23} \\ \sigma_{31} & \sigma_{32} & \sigma_{33} \end{bmatrix} \begin{bmatrix} n_1 \\ n_2 \\ n_3 \end{bmatrix}$$



Cauchy Stress Tensor

- **Symmetric stress tensor**

- Conservation of angular momentum implies that the stress tensor is symmetric

$$\begin{bmatrix} \sigma_x & \tau_{xy} & \tau_{xz} \\ \tau_{xy} & \sigma_y & \tau_{yz} \\ \tau_{xz} & \tau_{yz} & \sigma_z \end{bmatrix}$$

- Normal stresses

$$\sigma_x, \sigma_y, \sigma_z$$

- Shear stresses

$$\tau_{xy}, \tau_{xz}, \tau_{yz}$$

Deformation of Material

- **Physical view of deformation**

- Transformation of a body from a reference configuration to the current configuration
- A configuration is a set containing the positions of all particles of the body

- **Causes of deformation**

- External loads (usually on the exterior surfaces)
- Body forces (volumetric force within the whole body, e.g., gravity force)



Types of Deformation

- **Elastic deformations**

- Deformations are recovered after the stress field has been removed

- **Irreversible deformation**

- Deformations remain even after stresses have been removed
- Plastic deformation
 - Occurs in material bodies after stresses have attained a certain threshold value (elastic limit or yield stress)



Strain of Material

- **What is a strain**

- A description of deformation in terms of relative displacement of particles
- The relation between stresses and induced strains is expressed by constitutive equations
 - E.g., Hooke's law for linear elastic materials

- **Formulation**

- A general deformation of a body can be expressed in the form

$$\mathbf{x} = \mathbf{F}(\mathbf{X})$$

- \mathbf{X} is the reference position of material points in the body

Strain of Material

- **Formulation**

- Such a measure does not distinguish between rigid body motions and changes in shape of the body

- Mathematical definition

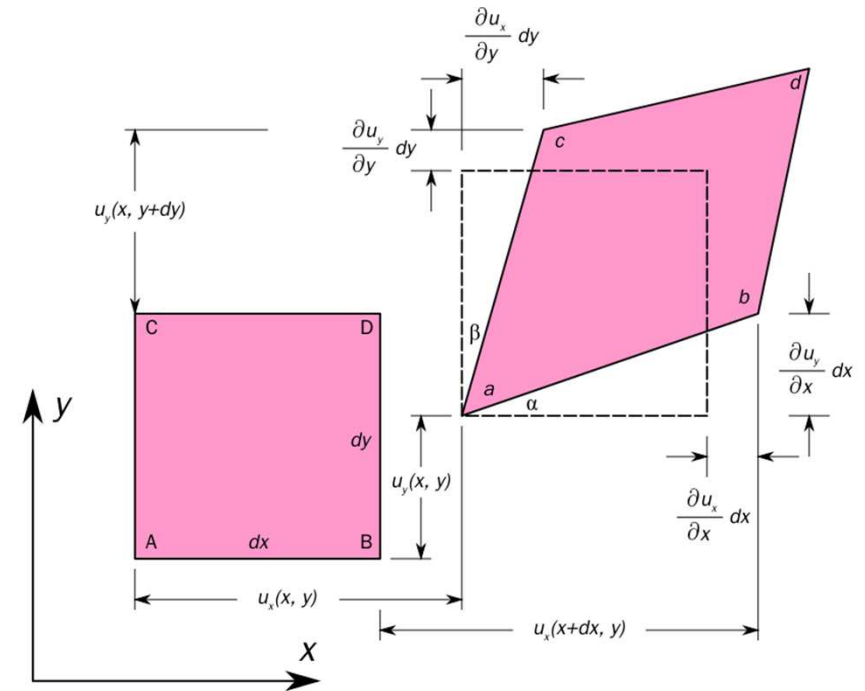
$$\boldsymbol{\epsilon} \doteq \frac{\partial}{\partial \mathbf{X}} (\mathbf{x} - \mathbf{X}) = \mathbf{F}' - \mathbf{I}$$

- Strains measure how much a given deformation differs locally from a rigid-body deformation

Strain of Material

- **Normal and shear strain**

- A normal strain is perpendicular to the face of an element
- A shear strain is parallel to it
- Consistent with normal stress and shear stresses



Strain of Material

- **Strain tensor**

$$\varepsilon_{xx} = \frac{\partial u_x}{\partial x} \quad \varepsilon_{yy} = \frac{\partial u_y}{\partial y} \quad \varepsilon_{zz} = \frac{\partial u_z}{\partial z}$$

$$\gamma_{xy} = \gamma_{yx} = \frac{\partial u_y}{\partial x} + \frac{\partial u_x}{\partial y}$$

$$\gamma_{yz} = \gamma_{zy} = \frac{\partial u_y}{\partial z} + \frac{\partial u_z}{\partial y}$$

$$\gamma_{zx} = \gamma_{xz} = \frac{\partial u_z}{\partial x} + \frac{\partial u_x}{\partial z}$$

$$\boldsymbol{\varepsilon} = \begin{bmatrix} \varepsilon_{xx} & \varepsilon_{xy} & \varepsilon_{xz} \\ \varepsilon_{yx} & \varepsilon_{yy} & \varepsilon_{yz} \\ \varepsilon_{zx} & \varepsilon_{zy} & \varepsilon_{zz} \end{bmatrix} = \begin{bmatrix} \varepsilon_{xx} & \frac{1}{2}\gamma_{xy} & \frac{1}{2}\gamma_{xz} \\ \frac{1}{2}\gamma_{yx} & \varepsilon_{yy} & \frac{1}{2}\gamma_{yz} \\ \frac{1}{2}\gamma_{zx} & \frac{1}{2}\gamma_{zy} & \varepsilon_{zz} \end{bmatrix}$$

Deformation Model

- **Derivation of deformation for fracture**

- Define a set of differential equations
- Describe the aggregate behavior of the material in a continuous fashion
- Use a finite element method to discretize these equations for computer simulation
- Designed to be simple, fast, and suitable for fracture modeling

Continuous Model

- **Primary assumption**

- The scale of the modeled effects is significantly greater than the scale of the material's composition
- Macroscopic fractures can be significantly influenced by effects that occur at small scales
- For graphical simulation, assume that a continuum model is adequate



Continuous Model

- **Define material coordinates**

$$\mathbf{u} = [u, v, w]^T$$

- **Deformation of the material**

$$\mathbf{x}(\mathbf{u}) = [x, y, z]^T$$

- In areas where material exists, $\mathbf{x}(\mathbf{u})$ is continuous
- Except across a finite number of surfaces within the volume that correspond to fractures



Continuous Model

- **Green's strain tensor**

- Measure the local deformation of the material

$$\epsilon_{ij} = \left(\frac{\partial \mathbf{x}}{\partial u_i} \cdot \frac{\partial \mathbf{x}}{\partial u_j} \right) - \delta_{ij}$$

- Kronecker delta

$$\delta_{ij} = \begin{cases} 1 & : i = j \\ 0 & : i \neq j \end{cases}$$

- This strain metric only measures local deformation
- Invariant with respect to rigid body transformations

Continuous Model

- **Strain rate tensor**

- Measure the rate at which the strain is changing
- Take the time derivative of strain

$$\nu_{ij} = \left(\frac{\partial \mathbf{x}}{\partial u_i} \cdot \frac{\partial \dot{\mathbf{x}}}{\partial u_j} \right) + \left(\frac{\partial \dot{\mathbf{x}}}{\partial u_i} \cdot \frac{\partial \mathbf{x}}{\partial u_j} \right)$$

- Strain and strain rate tensors provide the raw information to compute internal elastic and damping forces
- But they do not take into account the properties of the material

Continuous Model

- **Stress tensor**

- Represented as a 3 x 3 symmetric matrix
- Two components
 - Elastic stress due to strain $\sigma^{(\epsilon)}$
 - Viscous stress due to strain rate $\sigma^{(\nu)}$

- Total internal stress

$$\sigma = \sigma^{(\epsilon)} + \sigma^{(\nu)}$$

- Representation of elastic and viscous stresses

$$\sigma_{ij}^{(\epsilon)} = \sum_{k=1}^3 \sum_{l=1}^3 C_{ijkl} \epsilon_{kl}$$

$$\sigma_{ij}^{(\nu)} = \sum_{k=1}^3 \sum_{l=1}^3 D_{ijkl} \nu_{kl}$$

C: a set of the 81 elastic coefficients

D: a set of the 81 damping coefficients

Continuous Model

- **Stress tensor**

- Coefficient reduction

- Symmetric for both elastic and viscous stresses
 - Impose the additional constraint that the material is isotropic
 - **C** reduces to only two independent values (Lamé constants)

$$\sigma_{ij}^{(\epsilon)} = \sum_{k=1}^3 \lambda \epsilon_{kk} \delta_{ij} + 2\mu \epsilon_{ij}$$

Material's rigidity : μ ,
Resistance to changes in volume: λ

Continuous Model

- **Stress tensor**

- Coefficient reduction
 - Similarly, **D** reduces to only two independent values

$$\sigma_{ij}^{(\nu)} = \sum_{k=1}^3 \phi \nu_{kk} \delta_{ij} + 2\psi \nu_{ij}$$

- The coefficients control how quickly the material dissipates internal kinetic energy
- Will not damp motions that are locally rigid
- Has the desirable property of dissipating only internal vibrations

Continuous Model

- **Potential densities**

- Elastic potential density

$$\eta = \frac{1}{2} \sum_{i=1}^3 \sum_{j=1}^3 \sigma_{ij}^{(\epsilon)} \epsilon_{ij}$$

- Damping potential density

$$\kappa = \frac{1}{2} \sum_{i=1}^3 \sum_{j=1}^3 \sigma_{ij}^{(\nu)} \nu_{ij}$$

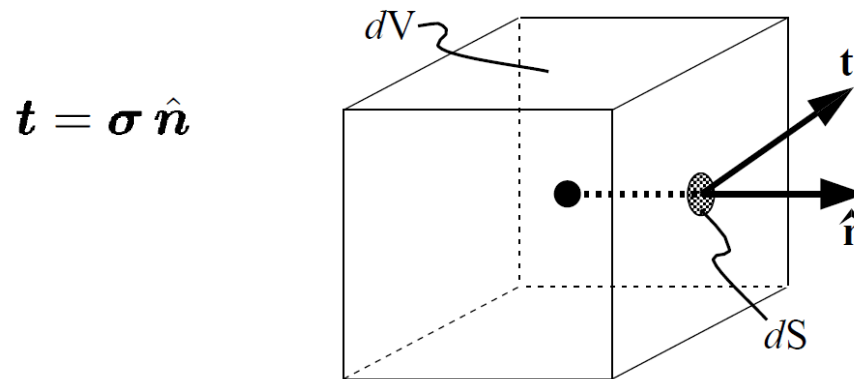
- Integrated over the volume of the material
 - Obtain the total elastic and damping potentials



Continuous Model

- **Internal force**

- The stress can also be used to compute the forces acting internally to the material at any location
- The traction (force per unit area) acting on a face perpendicular to the normal



Finite Element Discretization

- **Discretization for computer simulation**

- Finite difference

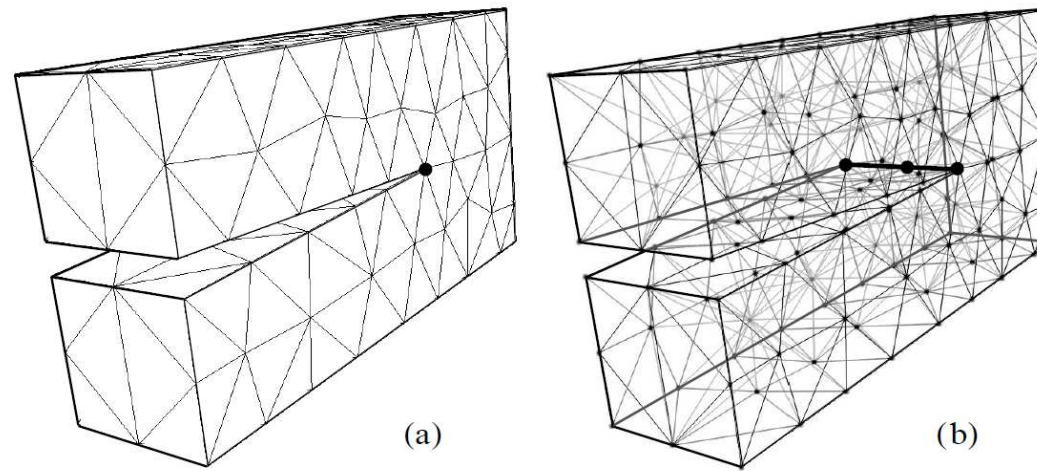
- Well suited for problems with a regular structure but becomes complicated when the structure is irregular

- Finite element

- Partition the domain of the material into distinct sub-domains
 - Within each element, the material is described locally by a function
 - The function is decomposed into a set of shape (basis) functions, each associating one of the nodes
 - The mesh defines a piecewise function over the entire material domain

Finite Element Discretization

- **Tetrahedral mesh representation of objects**

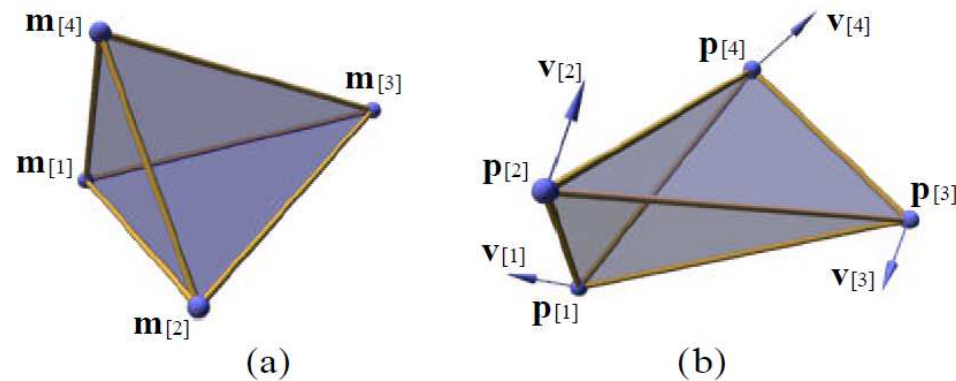


(a) : only the external faces of the tetrahedra are drawn
(b) : the internal structure is shown

Finite Element Discretization

- **Formulation by barycentric coordinates**

- Linear elements



- Each node has
 - A location in the material coordinate system (a)
 - A position and velocity in the world coordinate system (b)

Finite Element Discretization

- **Formulation by barycentric coordinates**

- Linear elements

- Barycentric coordinates provide a natural way to define the linear shape functions
 - Barycentric coordinates definition

$$\mathbf{b} = [b_1, b_2, b_3, b_4]^T$$

- Interpolate the entire region of the element

$$\begin{bmatrix} u \\ 1 \end{bmatrix} = \begin{bmatrix} m_{[1]} & m_{[2]} & m_{[3]} & m_{[4]} \\ 1 & 1 & 1 & 1 \end{bmatrix} \mathbf{b}$$

Finite Element Discretization

- **Formulation by barycentric coordinates**

- Linear elements

- Barycentric coordinates may also be used to interpolate the node's world position and velocity

$$\begin{bmatrix} \mathbf{x} \\ 1 \end{bmatrix} = \begin{bmatrix} \mathbf{p}_{[1]} & \mathbf{p}_{[2]} & \mathbf{p}_{[3]} & \mathbf{p}_{[4]} \\ 1 & 1 & 1 & 1 \end{bmatrix} \mathbf{b}$$

$$\begin{bmatrix} \dot{\mathbf{x}} \\ 1 \end{bmatrix} = \begin{bmatrix} \mathbf{v}_{[1]} & \mathbf{v}_{[2]} & \mathbf{v}_{[3]} & \mathbf{v}_{[4]} \\ 1 & 1 & 1 & 1 \end{bmatrix} \mathbf{b}$$

- Determine the barycentric coordinates of a point within the element

$$\mathbf{b} = \beta \begin{bmatrix} \mathbf{u} \\ 1 \end{bmatrix} \quad \beta = \begin{bmatrix} \mathbf{m}_{[1]} & \mathbf{m}_{[2]} & \mathbf{m}_{[3]} & \mathbf{m}_{[4]} \\ 1 & 1 & 1 & 1 \end{bmatrix}^{-1}$$

Finite Element Discretization


- **Formulation by barycentric coordinates**

- Linear elements

$$\begin{bmatrix} \mathbf{x} \\ 1 \end{bmatrix} = \begin{bmatrix} \mathbf{p}_{[1]} & \mathbf{p}_{[2]} & \mathbf{p}_{[3]} & \mathbf{p}_{[4]} \\ 1 & 1 & 1 & 1 \end{bmatrix} \mathbf{b}$$
$$\begin{bmatrix} \dot{\mathbf{x}} \\ 1 \end{bmatrix} = \begin{bmatrix} \mathbf{v}_{[1]} & \mathbf{v}_{[2]} & \mathbf{v}_{[3]} & \mathbf{v}_{[4]} \\ 1 & 1 & 1 & 1 \end{bmatrix} \mathbf{b}$$

+

$$\mathbf{b} = \boldsymbol{\beta} \begin{bmatrix} \mathbf{u} \\ 1 \end{bmatrix}$$


$$\mathbf{x}(\mathbf{u}) = \mathbf{P} \boldsymbol{\beta} \begin{bmatrix} \mathbf{u} \\ 1 \end{bmatrix}$$
$$\dot{\mathbf{x}}(\mathbf{u}) = \mathbf{V} \boldsymbol{\beta} \begin{bmatrix} \mathbf{u} \\ 1 \end{bmatrix}$$
$$\mathbf{P} = \begin{bmatrix} \mathbf{p}_{[1]} & \mathbf{p}_{[2]} & \mathbf{p}_{[3]} & \mathbf{p}_{[4]} \end{bmatrix}$$
$$\mathbf{V} = \begin{bmatrix} \mathbf{v}_{[1]} & \mathbf{v}_{[2]} & \mathbf{v}_{[3]} & \mathbf{v}_{[4]} \end{bmatrix}$$

Finite Element Discretization

- **Formulation by barycentric coordinates**

- Linear elements

- For non-degenerate elements, β is guaranteed to be non-singular
- Computing the strain and strain rate tensors require

$$\begin{aligned}\frac{\partial \mathbf{x}}{\partial u_i} &= \mathbf{P} \beta \boldsymbol{\delta}_i \\ \frac{\partial \dot{\mathbf{x}}}{\partial u_i} &= \mathbf{V} \beta \boldsymbol{\delta}_i\end{aligned}\quad \boldsymbol{\delta}_i = [\delta_{i1} \ \delta_{i2} \ \delta_{i3} \ 0]^T$$

- Since the element's shape functions are linear, these partials are constant within the element

Finite Element Discretization

- **Computing forces at the node**

- The element will exert elastic and damping forces on its nodes
 - Elastic force of the i-th node
 - Defined as the negative partial of the elastic potential density w.r.t. positions integrated over the volume of the element

$$\eta = \frac{1}{2} \sum_{i=1}^3 \sum_{j=1}^3 \sigma_{ij}^{(\epsilon)} \epsilon_{ij} \quad \epsilon_{ij} = \left(\frac{\partial \mathbf{x}}{\partial u_i} \cdot \frac{\partial \mathbf{x}}{\partial u_j} \right) - \delta_{ij} \quad \frac{\partial \mathbf{x}}{\partial u_i} = \mathbf{P} \boldsymbol{\beta}_i$$

$$\mathbf{f}_{[i]}^{(\epsilon)} = -\frac{\text{vol}}{2} \sum_{j=1}^4 \mathbf{p}_{[j]} \sum_{k=1}^3 \sum_{l=1}^3 \beta_{jl} \beta_{ik} \sigma_{kl}^{(\epsilon)}$$

$$\text{vol} = \frac{1}{6} [(\mathbf{m}_{[2]} - \mathbf{m}_{[1]}) \times (\mathbf{m}_{[3]} - \mathbf{m}_{[1]})] \cdot (\mathbf{m}_{[4]} - \mathbf{m}_{[1]})$$

Finite Element Discretization

- **Computing forces at the node**

- The element will exert elastic and damping forces on its nodes
 - Damping force on the i-th node
 - Defined as the negative partial of the damping potential density w.r.t. velocity integrated over the volume of the element

$$\kappa = \frac{1}{2} \sum_{i=1}^3 \sum_{j=1}^3 \sigma_{ij}^{(\nu)} \nu_{ij} \quad \nu_{ij} = \left(\frac{\partial \mathbf{x}}{\partial u_i} \cdot \frac{\partial \dot{\mathbf{x}}}{\partial u_j} \right) + \left(\frac{\partial \dot{\mathbf{x}}}{\partial u_i} \cdot \frac{\partial \mathbf{x}}{\partial u_j} \right) \quad \frac{\partial \dot{\mathbf{x}}}{\partial u_i} = \mathbf{V} \boldsymbol{\beta} \boldsymbol{\delta}_i$$

$$\mathbf{f}_{[i]}^{(\nu)} = -\frac{\text{vol}}{2} \sum_{j=1}^4 \mathbf{p}_{[j]} \sum_{k=1}^3 \sum_{l=1}^3 \beta_{jl} \beta_{ik} \sigma_{kl}^{(\nu)}$$

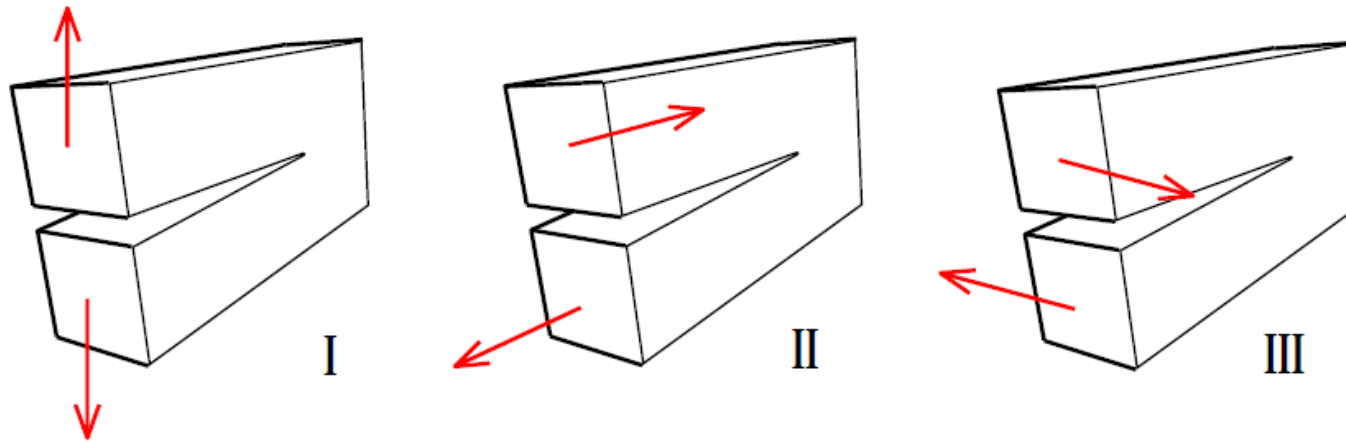
- Total internal force that an element exerts on a node $\mathbf{f}_{[i]}^{\text{el}} = -\frac{\text{vol}}{2} \sum_{j=1}^4 \mathbf{p}_{[j]} \sum_{k=1}^3 \sum_{l=1}^3 \beta_{jl} \beta_{ik} \sigma_{kl}$
- Total force: summing all elements

Fracture Modeling

- **Based on linear elastic fracture mechanics**
 - The region of plasticity near the crack tip is neglected
 - Modeled materials will be brittle
 - Once the material has begun to fail, the fractures will have a strong tendency to propagate across the material
- Three loading modes
 - Opening
 - In-plane shear
 - Out-of-plane shear

Fracture Modeling

- Three loading modes experienced by a crack



Mode I: Opening

Mode II: In-Plane Shear

Mode III: Out-of-Plane Shear

Fracture Modeling

- **Resolve the crack**

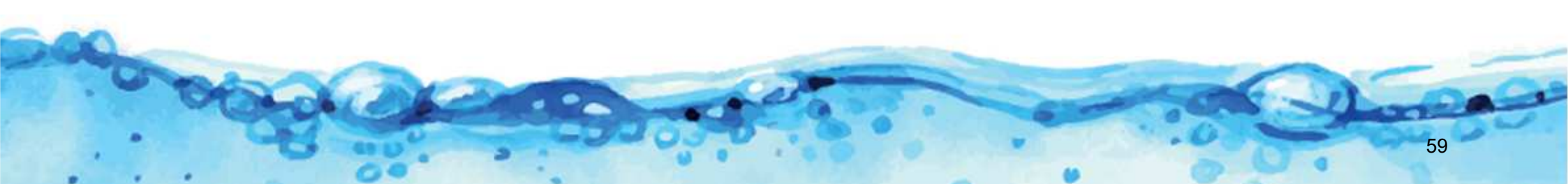
- Analyze the forces acting at the crack tip
 - Tensile forces that are opposed by other tensile forces will cause the crack to continue
 - In a direction perpendicular to the direction of largest tensile load
 - Compressive loads will tend to arrest a crack
 - Perpendicular to the crack
- Use the element nodes to determine where a crack should be initiated



Fracture Modeling

- **Fracture algorithm overview**

- After each time step
 - Resolve the internal forces acting on all nodes
 - Tensile and compressive components
 - At each node
 - The forces are used to form a tensor
 - Describe how internal forces are acting to separate that node
 - If the action is sufficiently large, the node is split into two distinct nodes and a fracture plane is computed
 - All elements attached to the node are divided along the plane



Force Decomposition

- **Forces acting on a node**

- Decomposed by separating the stress tensors into tensile and compressive components
- Eigen analysis of stress tensor
 - Let $v^i(\boldsymbol{\sigma})$ $i \in \{1, 2, 3\}$ be the i -th eigenvalue of $\boldsymbol{\sigma}$
 - Let $\hat{\mathbf{n}}^i(\boldsymbol{\sigma})$ be the corresponding unit length eigenvector
 - Analysis
 - Positive eigenvalues correspond to tensile stresses
 - Negative eigenvalues to compressive stresses

Force Decomposition

- **Decomposition of stress tensor**

- Given a vector \mathbf{a} in \mathbb{R}^3 , we can construct a 3 x 3 symmetric matrix $\mathbf{m}(\mathbf{a})$
 - $|\mathbf{a}|$ as an eigenvalue
 - \mathbf{a} as the corresponding eigenvector
 - The other two eigenvalues equal to zero

$$\mathbf{m}(\mathbf{a}) = \begin{cases} \mathbf{a} \mathbf{a}^T / |\mathbf{a}| & : \mathbf{a} \neq \mathbf{0} \\ \mathbf{0} & : \mathbf{a} = \mathbf{0} \end{cases}$$

Force Decomposition

- **Decomposition of stress tensor**

- The tensile and compressive components can be represented as

$$\boldsymbol{\sigma}^+ = \sum_{i=1}^3 \max(0, v^i(\boldsymbol{\sigma})) \mathbf{m}(\hat{\mathbf{n}}^i(\boldsymbol{\sigma}))$$
$$\boldsymbol{\sigma}^- = \sum_{i=1}^3 \min(0, v^i(\boldsymbol{\sigma})) \mathbf{m}(\hat{\mathbf{n}}^i(\boldsymbol{\sigma}))$$

- Force that an element exerts on a node

- Tensile component $\mathbf{f}_{[i]}^+ = -\frac{\text{vol}}{2} \sum_{j=1}^4 \mathbf{p}_{[j]} \sum_{k=1}^3 \sum_{l=1}^3 \beta_{jl} \beta_{ik} \sigma_{kl}^+$

- Compressive component using $\mathbf{f}_{[i]} = \mathbf{f}_{[i]}^+ + \mathbf{f}_{[i]}^-$

Separation Tensor

- **Separation tensor**

- Used directly to determine whether a fracture should occur at a node
- Formed from the balanced tensile and compressive forces acting at each node

$$\varsigma = \frac{1}{2} \left(-\mathbf{m}(\mathbf{f}^+) + \sum_{\mathbf{f} \in \{\mathbf{f}^+\}} \mathbf{m}(\mathbf{f}) + \mathbf{m}(\mathbf{f}^-) - \sum_{\mathbf{f} \in \{\mathbf{f}^-\}} \mathbf{m}(\mathbf{f}) \right)$$

- Let v^+ be the largest positive eigenvalue, If v^+ is greater than the material toughness τ , the material will fail at the node

Fracture Plane

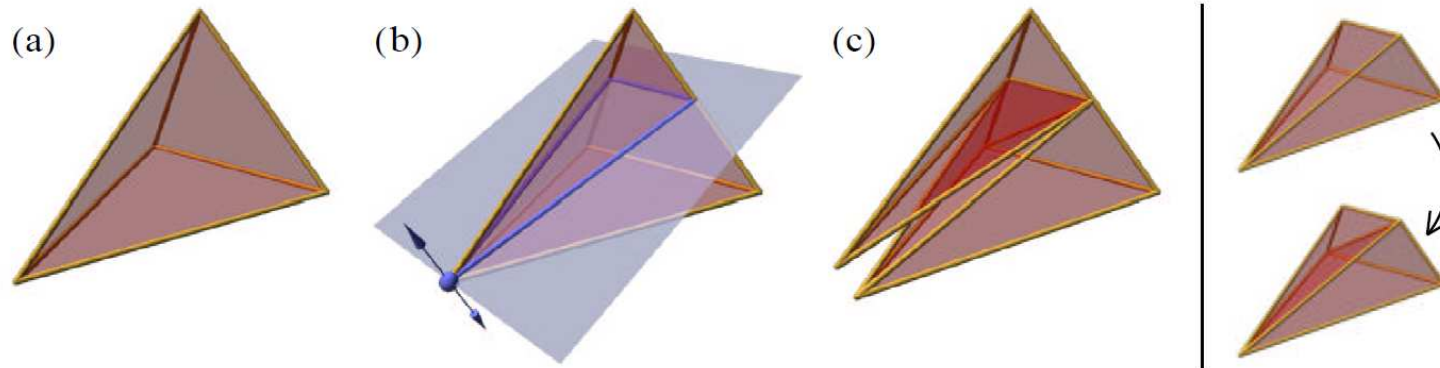
- **Orientation of the fracture plane**

- Perpendicular to the tensile eigen vector
- In case of multiple eigenvalues
 - Multiple fracture planes may be generated
 - First generate the plane for the largest value
 - Remeshing
 - Recompute the new value for separation tensor and proceed the above



Fracture Plane

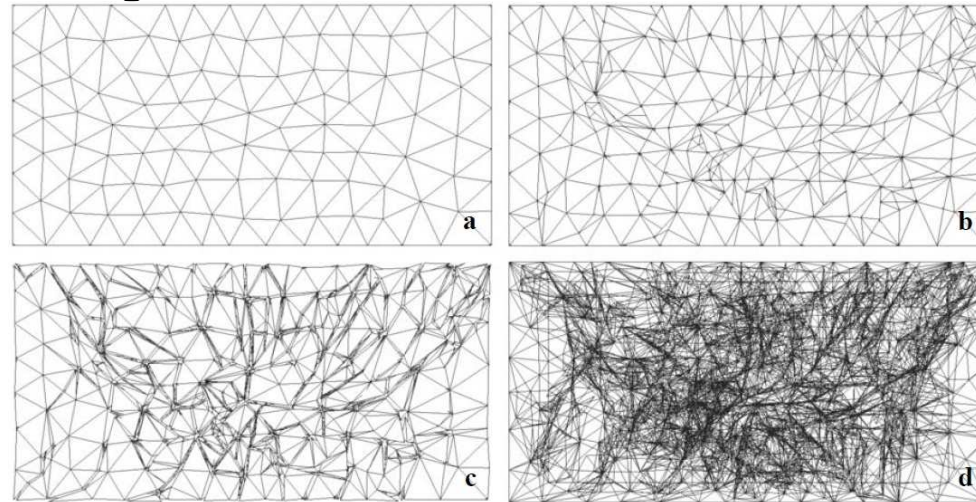
- **An element is split by the fracture plane**



- (a) The initial tetrahedral element
(b) The splitting node and fracture plane are shown in blue
(c) The element is split along the fracture plane into two polyhedra

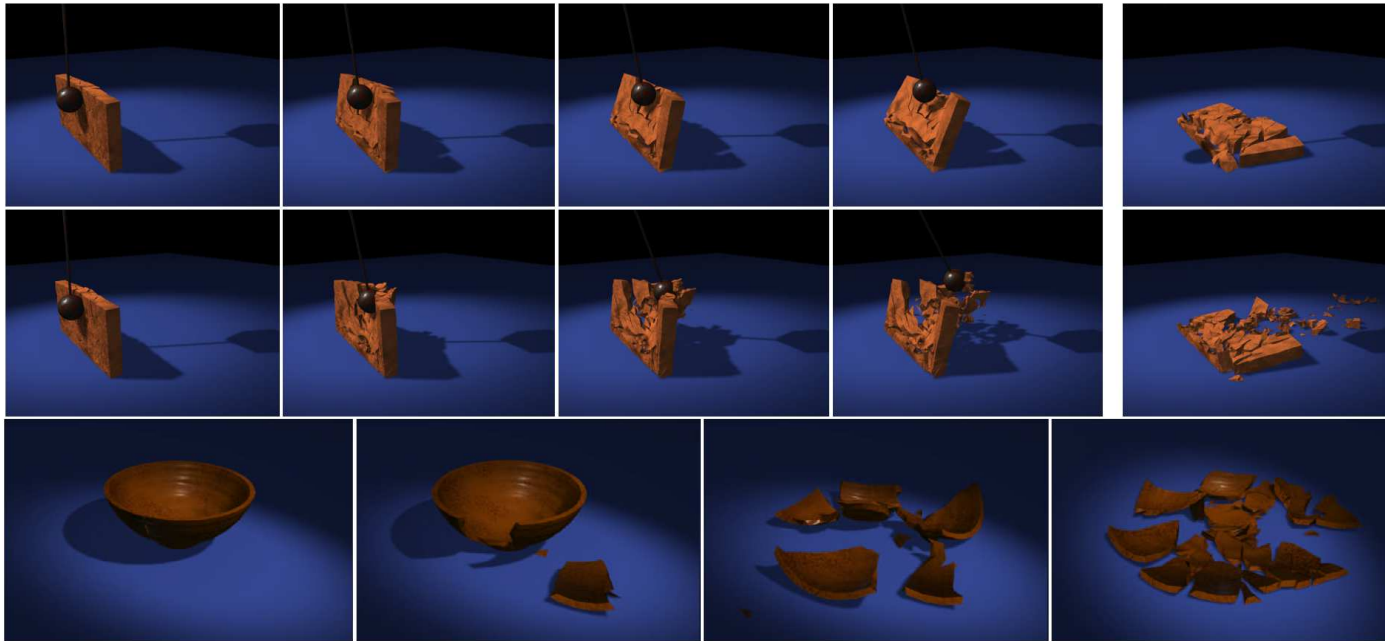
Local Remeshing

- **The mesh must be modified to reflect the new discontinuity**
 - Split elements that intersect the fracture plane
 - Modify neighboring elements to ensure mesh consistency



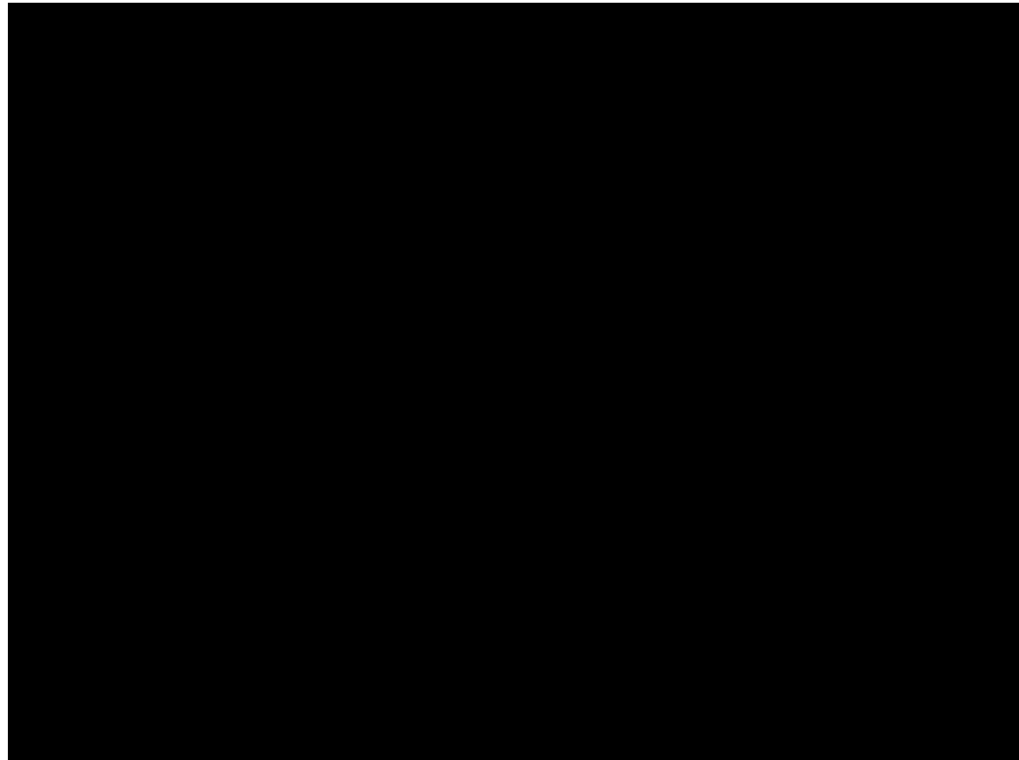
Simulation Results

- Fracture-based rigid body simulation



Simulation Results

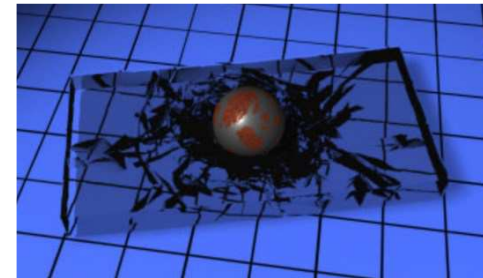
- **Fracture-based rigid body simulation**



Fracture Animation

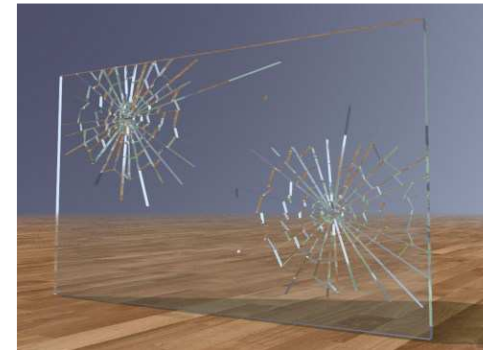
- **Physically simulated fracture**

- Physically accurate
- Stability issue
- Slow in high resolution



- **Pre-defined fracture pattern**

- Easier artistic control
- Fast and robust
- Difficult to create physically plausible details



Adaptive Fracture Refinement

- **Fracture is mesh resolution dependent**

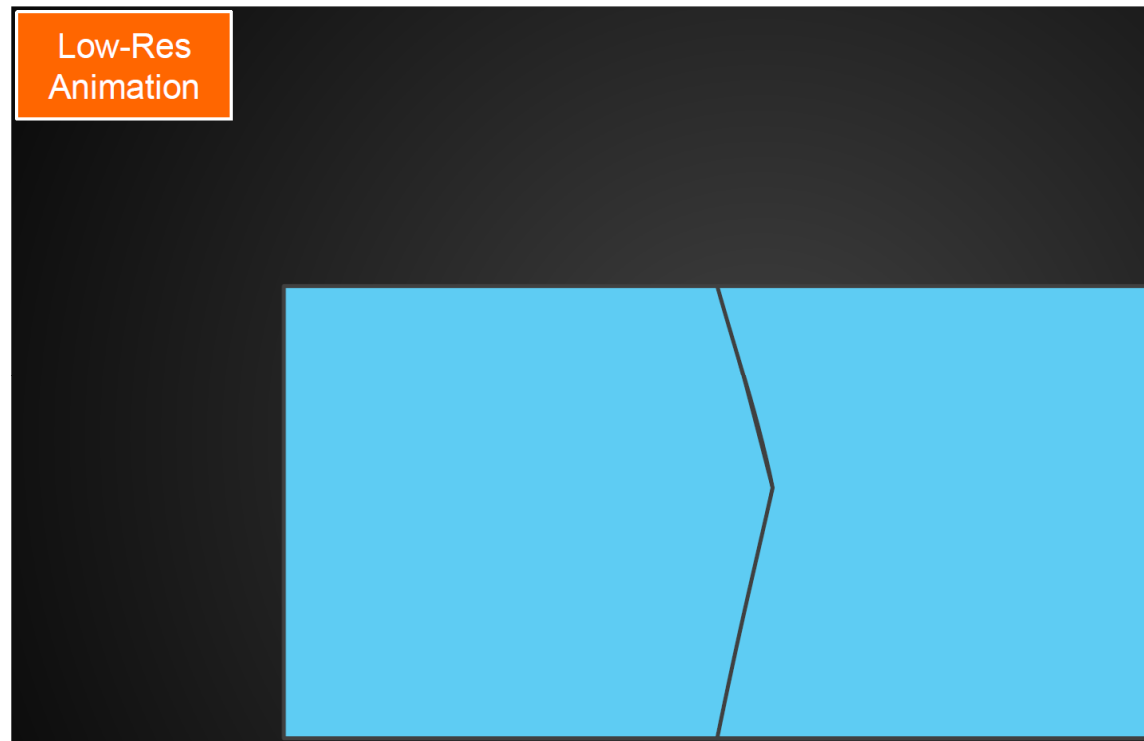
- Low resolution mesh simulation can be dramatically different from high resolution ones
- High resolution simulation
 - Collision detection is computationally costly
 - Fracture computation time is high

- **Adaptive refinement**

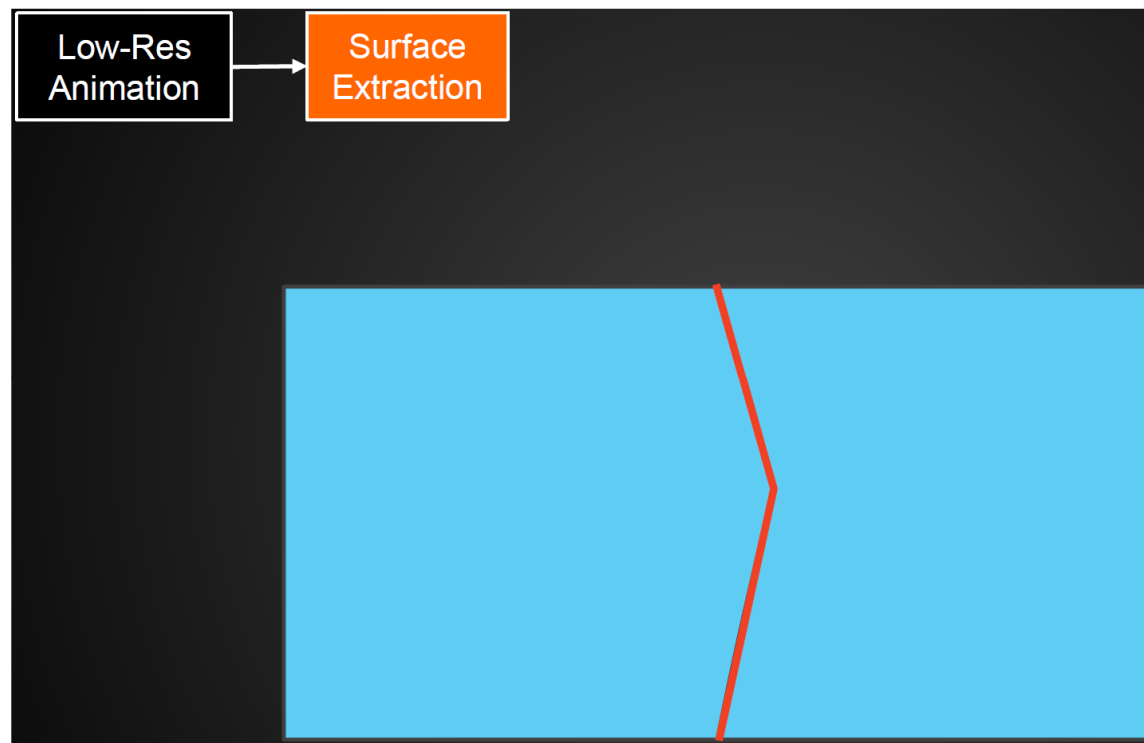
- Given a low resolution fracture animation
 - Adaptively refine the fracture surface within a 3D object



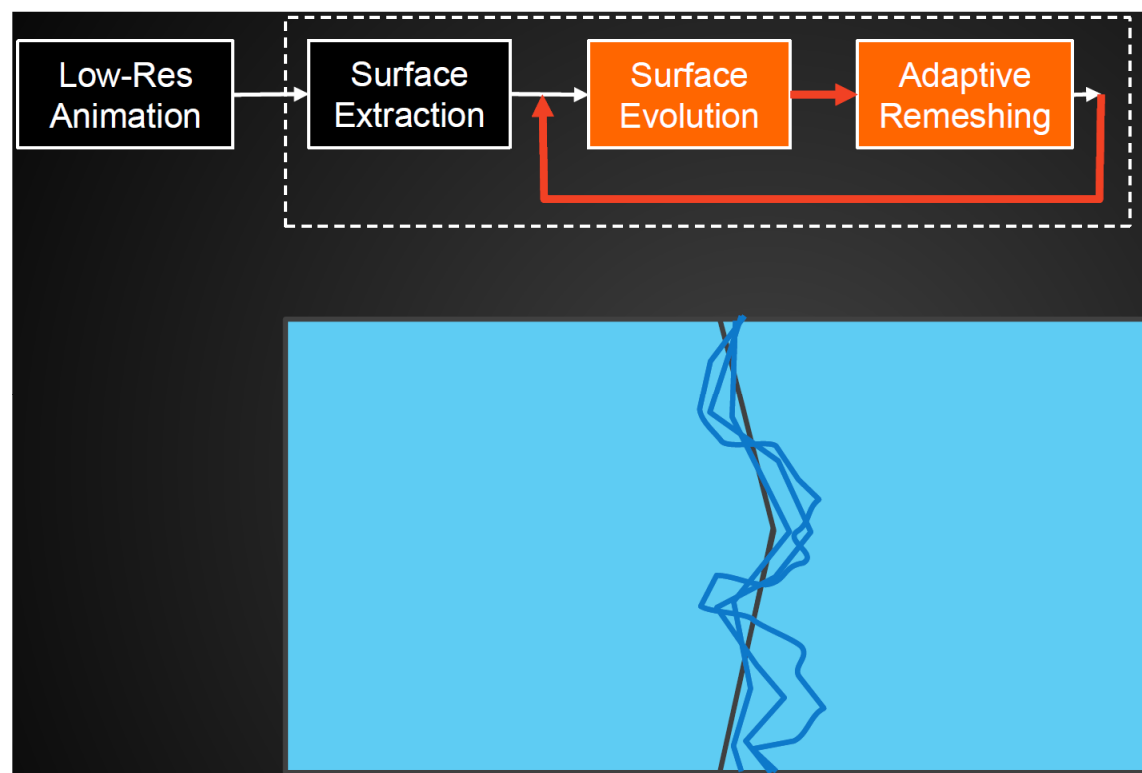
Adaptive Fracture Refinement



Adaptive Fracture Refinement



Adaptive Fracture Refinement



Fracture Surface Evolution

- **Separation field**

- Where the stress is high and material strength is weak
- Vertices should move to the lowest value regions



Gradient Flow & Surface Evolution

- **Evolve surface S to minimize**

$$\mathcal{E}(S) = \int_S \psi(x) ds$$

modified strength field

- Gradient descent for each vertex

$$\frac{d\mathbf{x}_i}{dt} = -\frac{1}{A_i} \sum_{j \in N_i} \left(\int_{S_j} \nabla_{\mathbf{x}_i} \psi(\mathbf{x}) \phi_i(\mathbf{x}) ds - \frac{\mathbf{e}_j^i \times \mathbf{n}_j}{2A_j} \int_{S_j} \psi(\mathbf{x}) ds \right)$$

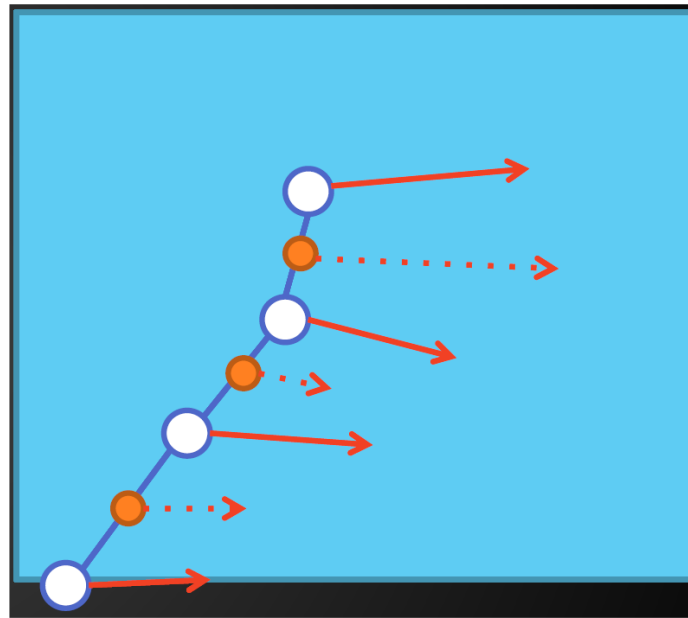
- Approximation

- Separation field varies linearly within triangle plane

$$\frac{d\mathbf{x}_i}{dt} = -\frac{1}{A_i} \sum_{j \in N_i} \left(\frac{1}{3} A_j \nabla \psi(\mathbf{x}_i) - \frac{\mathbf{e}_j^i \times \mathbf{n}_j}{2A_j} \sum_{k \in T_j} \psi(\mathbf{x}_k) \right)$$

Adaptive Fracture Remeshing

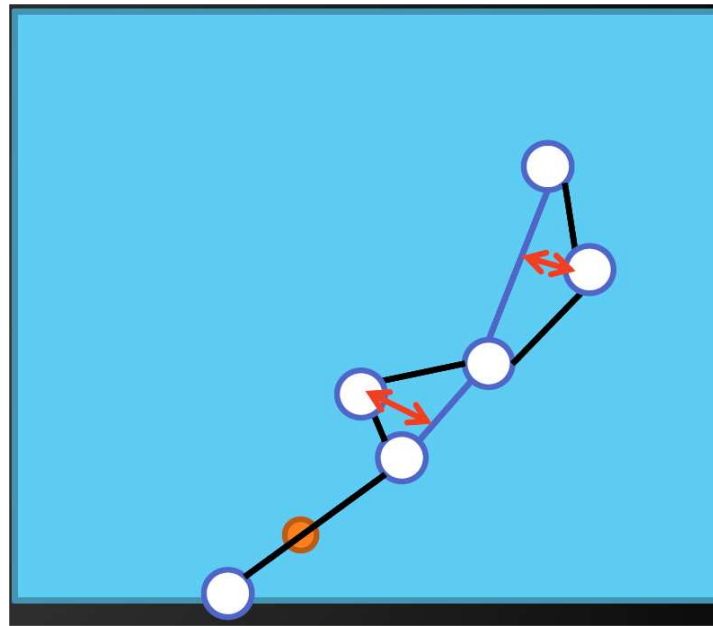
- Random candidate vertices



Adaptive Fracture Remeshing

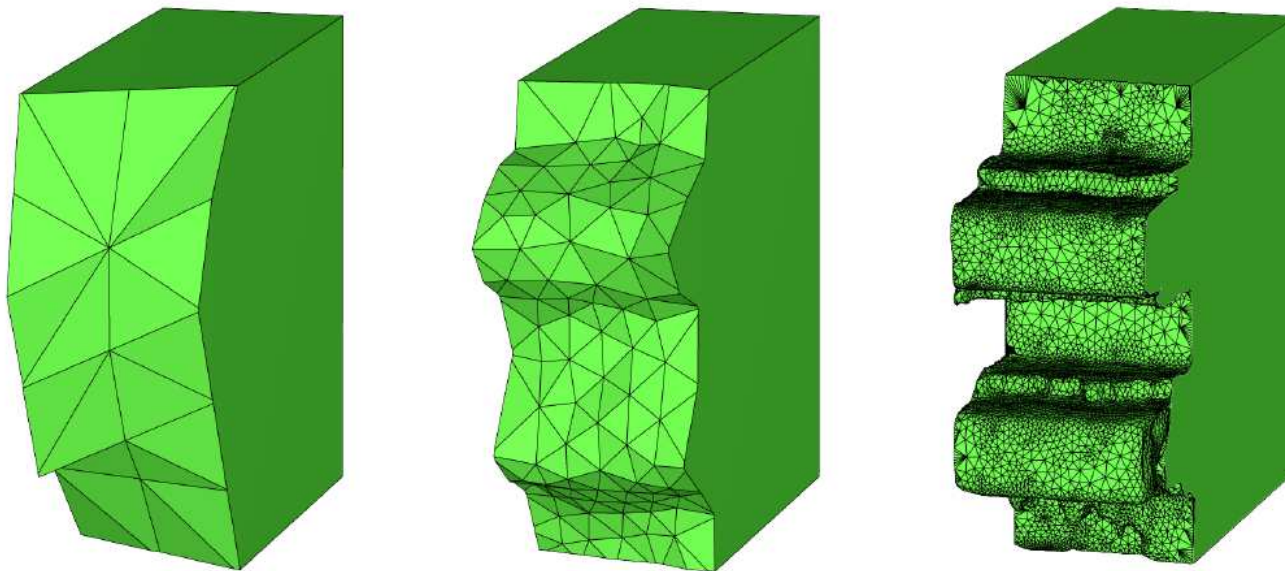
- **Random candidate vertices**

- Select and insert candidates
- Edge flipping optimization



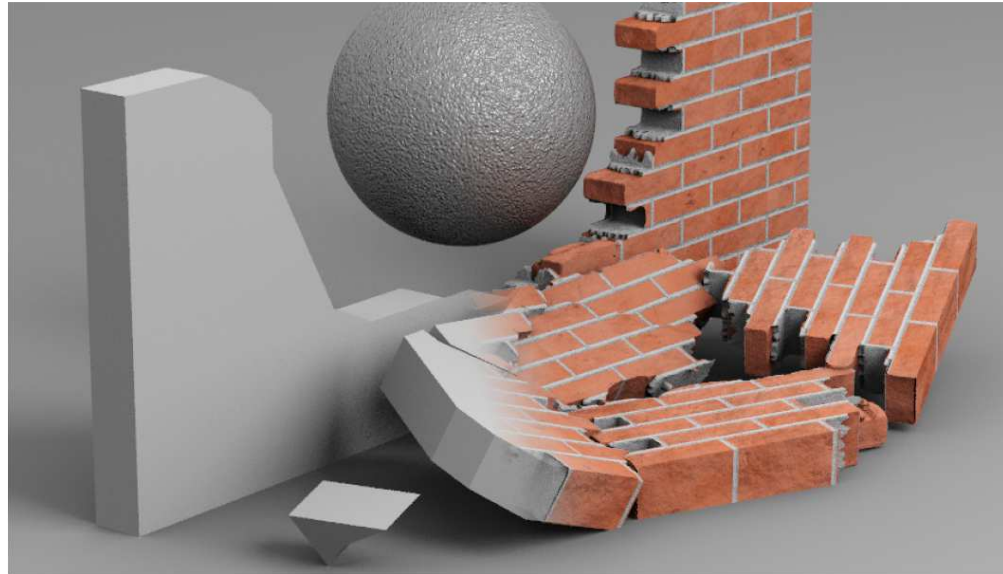
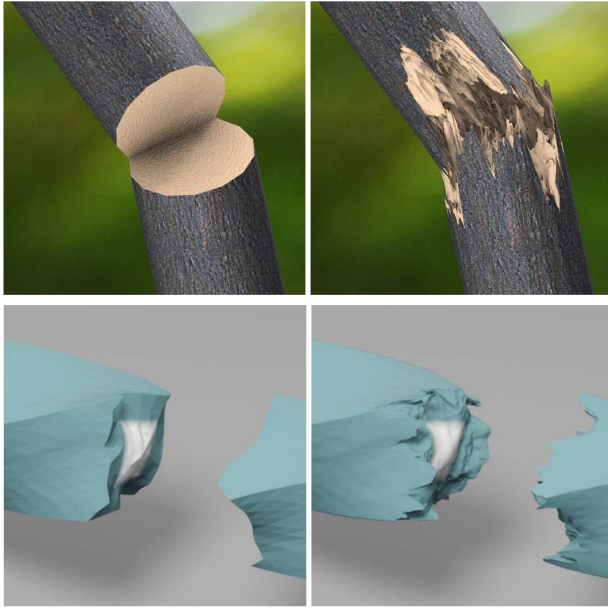
Adaptive Fracture Remeshing

- Iterative remeshing



Results

- **Fine-detail fractures**



Results

Physics-Inspired Adaptive Fracture Refinement

Zhili Chen Miaojun Yao Renguo Feng Huamin Wang
The Ohio State University

SIGGRAPH 2014

Advanced Real-Time Rigid Body Fracture

Real Time Dynamic Fracture
with
Volumetric Approximate Convex Decompositions

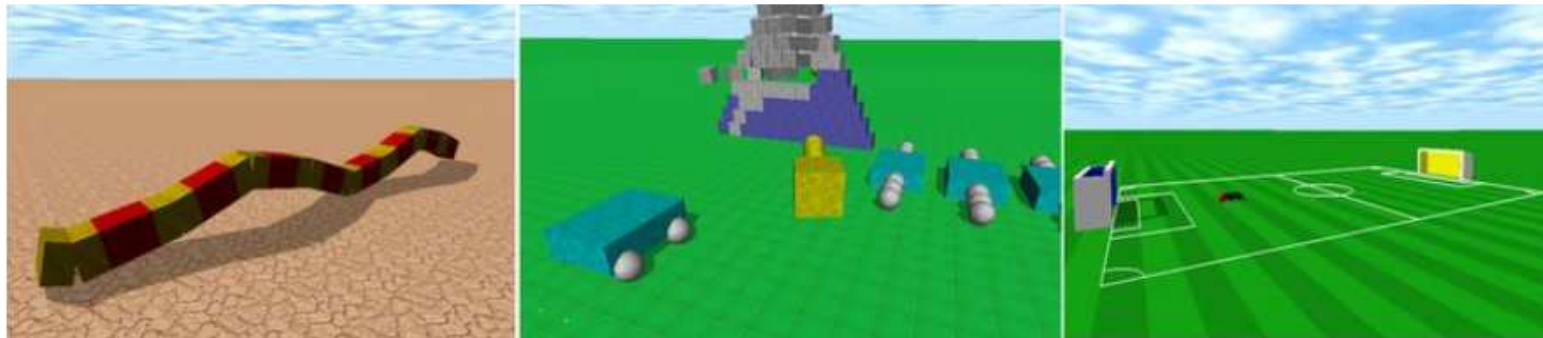
Matthias Müller, Nuttapong Chentanez, Tae-Yong Kim



Rigid Body Simulation Software

- **Open Dynamics Engine (ODE)**

- An open source, high performance library for simulating rigid body dynamics
- <https://www.ode.org/>



Rigid Body Simulation Software

- **Bullet Physics Engine**

- Real-time collision detection and multi-physics simulation for VR, games, visual effects, robotics, machine learning etc
- <https://pybullet.org/wordpress/>
- <https://github.com/bulletphysics/bullet3>



Next Lecture : Soft-Body Simulation – Hair I

